

# **Estimation of the time integration error in structural dynamics A comparison of two strategies**

Karl Schweizerhof, Jens Neumann  
Institut für Mechanik  
Universität Karlsruhe

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## Motivation

- Goal:**
- strategy for time step adaptation using time stepping schemes in structural dynamics
  - rigid body dynamics (no effect of spatial discretization)
  - flexible bodies, FE discretized structures

- local + global error estimation needed
- error estimator should work for linear and nonlinear problems  
→ numerical tests
- time step adaptation,  
how to use error estimation efficiently

## General Problem

System of Ordinary Differential Equations in structural dynamics:

here rigid body dynamics:  $M\ddot{\mathbf{u}} + \mathbf{h}(\mathbf{u}, \dot{\mathbf{u}}, t) = \mathbf{0}, \forall t > 0,$  (1)  
 $\mathbf{u} \in \mathbb{R}^n, M \dots$  positive definite

initial conditions:  $\mathbf{u}(t = 0) = \mathbf{u}_0$  and  $\dot{\mathbf{u}}(t = 0) = \dot{\mathbf{u}}_0$

uniqueness of solution of equation (1)

$\mathbf{h}(\mathbf{u}, \dot{\mathbf{u}}, t)$  must satisfy [e.g. Burlisch/Stör]:

$$\|\mathbf{h}(\mathbf{u}_1, \dot{\mathbf{u}}_1, t) - \mathbf{h}(\mathbf{u}_2, \dot{\mathbf{u}}_2, t)\| \leq \mathcal{L} \|\mathbf{u}_1 - \mathbf{u}_2\|, \quad \text{with: } \mathcal{L} > 0,$$

$\mathbf{h} \dots$  continuous, finite

or using mean value theorem:

$$\mathbf{K} = \frac{\partial \mathbf{h}}{\partial \mathbf{u}} \dots \text{continuous, finite, } \mathbf{K} \dots \text{Jacobian of } \mathbf{h}$$

- General:**
- Modal decomposition  $\rightarrow$  numerically high effort required for linear coupled ODE-systems
  - analytical solution for nonlinear ODE's unknown  
 $\rightarrow$  Numerical time integration

**Semidiscretized** equations for structural dynamics (linear case):

$$M\ddot{\mathbf{u}} + C\dot{\mathbf{u}} + K\mathbf{u} = \mathbf{0}, \forall t > 0, \quad \mathbf{u} \in \mathbb{R}^n, \quad (2)$$

initial conditions:  $\mathbf{u}(t = 0) = \mathbf{u}_0, \dot{\mathbf{u}}(t = 0) = \dot{\mathbf{u}}_0$

# Newmark's method derived from Finite Differences (FD-Newm)

Numerical solution of (2) with Finite Differences

→ split time domain  $[0, T]$  into finite number  $N$   
of time intervals  $\Delta t_n = t_{n+1} - t_n, n = 1 \dots N$

Equilibrium at state  $t_{n+1}$ :

$$\mathbf{M}\ddot{\mathbf{u}}_{n+1} + \mathbf{C}\dot{\mathbf{u}}_{n+1} + \mathbf{K}\mathbf{u}_{n+1} = \mathbf{0},$$

Newmark's scheme - derived from FD - with  $\Delta t_n = \Delta t$ :

$$\dot{\mathbf{u}}_{n+1} = \dot{\mathbf{u}}_n + \Delta t(1 - \gamma)\ddot{\mathbf{u}}_n + \Delta t\gamma\ddot{\mathbf{u}}_{n+1},$$

$$\mathbf{u}_{n+1} = \mathbf{u}_n + \Delta t\dot{\mathbf{u}}_n + \Delta t^2(1 - 2\beta)\ddot{\mathbf{u}}_n/2 + \Delta t^2\beta\ddot{\mathbf{u}}_{n+1}$$

Parameters:  $\gamma = 1/2$ , controls accuracy –  $\mathcal{O}(\Delta t^2)$  ,  
 $\beta = 1/4$ , A-stable, not L-stable

A-stable general numerical stability of algorithm

L-stable  $\approx$  damping of 'fast transient' in stiff problems

## Newmark's method as Finite Element formulation

(CG2-Newm)

CG2 – continuous Galerkin formulation with quadratic interpolation for eqn. (1)

- Divide time domain  $[0, T]$  into  $N$  time intervals  $\Delta t_n = t_{n+1} - t_n$ :

$$\sum_{n=0}^N \int_{t_n}^{t_{n+1}} \mathbf{w} \cdot (\mathbf{M}\ddot{\mathbf{u}} + \mathbf{C}\dot{\mathbf{u}} + \mathbf{K}\mathbf{u}) dt = 0, \quad \forall \mathbf{w} \in V$$

$V$  ... function space with inner product,

discrete variational formulation:  $V \rightarrow V_h$  (finite)

- Wood [1990]: Equivalent CG2-formulation for Newmark's method  
with  $\gamma = 2\beta$

Interpolation function:

$$\mathbf{u}_h = \mathbf{u}_n + (t - t_n)\dot{\mathbf{u}}_n + 1/2(t - t_n)^2\ddot{\mathbf{u}}_n \quad \forall t \in [t_n, t_{n+1}]$$

Special weighting function  $\mathbf{w}_h$  (for equivalence):

$$\mathbf{w}_h = \mathbf{W}_h \left[ 1/5 - (t - t_n)/\Delta t + (t - t_n)^2/\Delta t^2 \right] \quad \forall t \in [t_n, t_{n+1}]$$

$$\sum_{n=0}^N \int_{t_n}^{t_{n+1}} \mathbf{w}_h \cdot (\mathbf{M}\ddot{\mathbf{u}}_h + \mathbf{C}\dot{\mathbf{u}}_h + \mathbf{K}\mathbf{u}_h) dt = 0, \quad \forall \mathbf{w}_h \in V_h.$$

.... as  
**Finite Element formulation**  
(CG2-Newm)

$$(4\mathbf{M} + 2\Delta t_n \mathbf{C} + \Delta t_n^2 \mathbf{K}) \ddot{\mathbf{u}}_n = -2(\Delta t_n \mathbf{K} \dot{\mathbf{u}}_n + 2\mathbf{C} \dot{\mathbf{u}}_n + 2\mathbf{K} \mathbf{u}_n)$$

→ Equivalence between schemes - Comparison of Amplification matrix:

$$\begin{pmatrix} \mathbf{u}_{n+1} \\ \dot{\mathbf{u}}_{n+1} \end{pmatrix} = \mathbf{A}_{FD-newm} \begin{pmatrix} \mathbf{u}_n \\ \dot{\mathbf{u}}_n \end{pmatrix} = \mathbf{A}_{CG2-newm} \begin{pmatrix} \mathbf{u}_n \\ \dot{\mathbf{u}}_n \end{pmatrix}$$

→ in general only valid for linear ODE's,

→ techniques for error estimation for Finite Element method can be used

**Advantage:** no computation of 'initial' acceleration for  $\ddot{\mathbf{u}}_0$  necessary as with Finite Difference scheme:  $\ddot{\mathbf{u}}_0 = \mathbf{M}^{-1} (-\mathbf{C} \dot{\mathbf{u}}_0 - \mathbf{K} \mathbf{u}_0)$

## Global and local time integration error

Global error: Error at time  $t_n$ ,

$$e_g(t = t_m) = e(t = t_m) = u(t = t_m) - u_m$$

Local Error: Error in time step  $\Delta t_n$ ,

$$e_l(\Delta t_n) = u(t = t_m) - u_m$$

$$\text{with: } u(t = t_{m-1}) - u_{m-1} = 0$$

$u_m$  in general different for both methods ( $\gamma \neq 2\beta$ )

## Error indicator for global time integration error based on Finite Difference scheme and local error indicator

**Local** error indicator for displacements [Riccius 1998] for  $t \in [t_{m-1}, t_m]$ :

$$e_l(t_m) = \Delta t_m^3 \left(\frac{1}{6} - \beta\right) \ddot{u}_m \approx \Delta t_m^2 \left(\frac{1}{6} - \beta\right) (\ddot{u}_{m-1} - \ddot{u}_m)$$

→ based on difference between Taylor expansion and Newmark solution

**Assumption:**  $e_l(t_m)$  constant in  $t_m/\Delta t_m$  time intervals

Use for **global** error indicator (FD-based error indicator):

$$e_{gh}(FD - newm) = \frac{t_m}{\Delta t_m} e_l(t_m)$$

≡ Multiple of local error  $e_l(\Delta t_m)$  in last time step

## Error estimator for CG2-newm method

**Def.:**  $e_g(t = t_m) = u(t = t_m) - u_m$

Basis: Finite Element formulation of ODE:

$$\int_0^T \mathbf{w}_h \cdot (\mathbf{M}\ddot{\mathbf{u}}_h + \mathbf{C}\dot{\mathbf{u}}_h + \mathbf{K}\mathbf{u}_h) dt = 0, \quad \forall \mathbf{w}_h \in V_h$$

Variational form of ODE:

$$\int_0^T \mathbf{w}_h \cdot (\mathbf{M}\ddot{\mathbf{u}} + \mathbf{C}\dot{\mathbf{u}} + \mathbf{K}\mathbf{u}) dt = 0, \quad \forall \mathbf{w}_h \in V_h$$

→ Galerkin orthogonality of residual

$$\int_0^T \mathbf{w}_h \cdot \mathcal{D}e dt = 0,$$

$$\begin{aligned} \text{with : } \mathcal{D}e &= \mathbf{M}(\ddot{\mathbf{u}} - \ddot{\mathbf{u}}_h) + \mathbf{C}(\dot{\mathbf{u}} - \dot{\mathbf{u}}_h) + \mathbf{K}(\mathbf{u} - \mathbf{u}_h) \\ &= \mathbf{M}\ddot{\mathbf{e}} + \mathbf{C}\dot{\mathbf{e}} + \mathbf{K}\mathbf{e} \\ &= -(\mathbf{M}\ddot{\mathbf{u}}_h + \mathbf{C}\dot{\mathbf{u}}_h + \mathbf{K}\mathbf{u}_h) = \mathbf{R}. \end{aligned}$$

**Idea:** Use adjoint ODE to equation (2), Bangerth[1997], Maute[2001]

$$\mathcal{D}^* \mathbf{z} = \mathbf{M}\dot{\mathbf{z}} - \mathbf{C}\mathbf{z} + \mathbf{K}\mathbf{z} = \mathbf{J}, \quad \forall t < t_m \quad (\text{Backward analysis})$$

with initial conditions at  $t = t_m$ :  $\mathbf{z}(t = t_m) = \mathbf{z}_m$  and  $\dot{\mathbf{z}}(t = t_m) = \dot{\mathbf{z}}_m$

$\mathbf{J}$  ... functional, specified later (used to select error quantity)



## ... Global Error Estimator CGZ-Newton (contd.)

use  $e$  as test function for adjoint ODE:

$$\int_0^T e (M\ddot{z} - C\dot{z} + Kz) dt = \int_0^T e J dt$$

by partial integration:

$$\int_0^T e (M\ddot{z} - C\dot{z} + Kz) dt = \int_0^T z (M\ddot{e} + C\dot{e} + Ke) dt$$

$\hat{=}$  using  $z$  as test function for  $\mathcal{D}e = R$ :

$$\int_0^T z \cdot \mathcal{D}e dt = \int_0^T z R dt = \int_0^T e J dt$$

→ analogy to Betti's Principle from structural mechanics – here in time

### Choice for $J$ :

$J$  ... Dirac-Distribution – select quantity of interest for error:

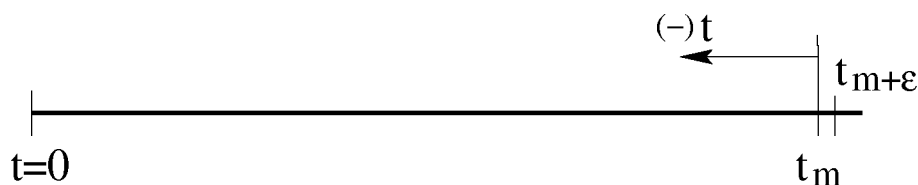
$$J = \delta(t = t_m) \frac{e(t = t_m)}{|e(t = t_m)|} \rightarrow \int_0^T e \cdot J dt = |e(t = t_m)| = e_g(t = t_m)$$

Corresponding initial conditions of adjoint ODE using linear momentum:

$$M\dot{z}(t_m + \epsilon) - M\dot{z}(t_m) = \int_{t_m}^{t_m + \epsilon} J dt = \frac{e(t = t_m)}{|e(t = t_m)|},$$

Approximation for  $e$  :  $e(t = t_m) \approx e_l(t = t_m)$ , (from numerical tests),

$$z(t_m) = 0, \quad \dot{z}(t_m + \epsilon) = 0, \quad \rightarrow \quad \dot{z}(t_m) = -M^{-1} \frac{e_l(t_m)}{|e_l(t_m)|}.$$



→ all quantities in  $z$  selected ,

*Alternative:* one single dof with linear momentum

## ... global Error estimator CG2-newm (contd.)

### Numerical effort:

- a) same numerical effort required for computation of  $z_h$  as for computation of  $u_h$  in linear case  
→ numerical solution of adjoint ODE, e.g. with same FE method

b) Computation of estimated error for  $e_g(t = t_m)$ : 
$$e_{gh} = \int_0^{t_m} \mathbf{R} z_h dt$$

- Requires:**
- storage of  $u_h, \dot{u}_h, \ddot{u}_h$  at each time step,
  - evaluation of  $\mathbf{R}(u_h, \dot{u}_h, \ddot{u}_h)$  and  $z_h$  at integration points, e.g. 3 Gauß points

## ... global Error estimator CG2-newm (contd.)

*Same estimator for nonlinear problems?*

Nonlinear ODE:

$$M\ddot{\mathbf{u}} + \mathbf{h}(\mathbf{u}, \dot{\mathbf{u}}, t) = \mathbf{0}, \quad \forall t > 0, \quad \mathbf{u} \in \mathbb{R}^n$$

e.g. solution with CG2-method and Newton-Raphson scheme

→ linearization in each time step – iterative solution

→ Forward nonlinear analysis

**Problem:** duality principle only valid for linear operators:

$$\int_0^T \mathbf{z} \cdot \mathcal{D}\mathbf{e} dt = \int_0^T \mathbf{e} \cdot \mathcal{D}^*\mathbf{z} dt$$

For  $\mathcal{D}^*\mathbf{z}$  linearized differential operator of primal equation used

→ total differential:

$$d\mathbf{h} = \frac{\partial \mathbf{h}}{\partial \dot{\mathbf{u}}} d\dot{\mathbf{u}} + \frac{\partial \mathbf{h}}{\partial \mathbf{u}} d\mathbf{u} = \mathbf{C}(\dot{\mathbf{u}}, \mathbf{u}) d\dot{\mathbf{u}} + \mathbf{K}(\dot{\mathbf{u}}, \mathbf{u}) d\mathbf{u}$$

Corresponding linearized adjoint ODE:

$$M\ddot{\mathbf{z}} - \left( \frac{\partial \mathbf{h}}{\partial \dot{\mathbf{u}}} \right)^T \dot{\mathbf{z}} + \left( \frac{\partial \mathbf{h}}{\partial \mathbf{u}} \right)^T \mathbf{z} = \mathbf{J}$$

**Thus:** backward linear analysis with operators linearized  
at current time of primal solution

## ... global Error estimator CG2-newm (contd.)

**Algorithm** to compute  $e_{gh}(CG2-newm)$  at  $t_m$  for nonl. problems:

- i) compute numerical solution  $\mathbf{u}_h$  at discrete time states  $t_0, \dots, t_m$   
 $\rightarrow \mathbf{u}_h, \dot{\mathbf{u}}_h, \ddot{\mathbf{u}}_h$
- ii) compute  $\mathbf{z}_h$  with **linearized** adjoint ODE  
with  $\mathbf{K} = \mathbf{K}(t)$  and  $\mathbf{C} = \mathbf{C}(t)$
- iii) compute time integral  $\int_0^{t_m} \mathbf{R} \cdot \mathbf{z}_h dt,$ 
  - $\mathbf{z}_h$  with quadratic interpolation
  - Gauß integration with e.g. 4 points due to nonlinearity

## Numerical Examples

→ investigation of error estimators and indicators

**Single degree of freedom:** linear problem

$0.25\ddot{u} + 0.9u = 0$ ,  $\forall t > 0$  with initial conditions:  $u_0 = 1.0$ ,  $\dot{u}_0 = 0.0$

- Numerical integration using CG2,
- Variation of  $\Delta t = 0.05/0.1/0.2$

→ Global time integration error-  $e_g$  at  $t_m = 1.0$ :

$\Delta t$	$e_{gh}(CG2 - newm)$	$e_{gh}(FD - newm)$	$e_g$ (exact)	$e_l(t = 1.0)$
0.05	1.3454e-3	1.3669e-3	1.3463e-3	6.834e-5
0.1	5.3533e-3	5.5221e-3	5.3674e-3	5.552e-4
0.2	2.0968e-2	2.2217e-2	2.1189e-2	4.443e-3

**Results:** - efficiency index  $\eta = \frac{e_{gh}}{e_g(exact)} \approx 1$ ,

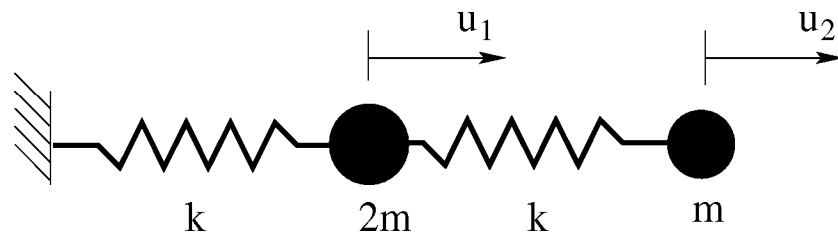
- order of local error:  $e_l = C_l \Delta t^3$ ,

- order of global error:  $e_g(t = t_m) = C_g \Delta t^2$

Reason: order reduction due to transport of local error

Hairer, Norsett, Wanner [1991]: min. order reduction  $\mathcal{O}(\Delta t)$

**Two degree of freedom system:** linear problem



$$\begin{bmatrix} 400 & 0 \\ 0 & 200 \end{bmatrix} \begin{bmatrix} \ddot{u}_1 \\ \ddot{u}_2 \end{bmatrix} + \begin{bmatrix} 200 & -100 \\ -100 & 100 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix} = \mathbf{0} \quad \text{for } t > 0,$$

with:  $M\ddot{u} + Ku = \mathbf{0}$

Initial conditions:  $\begin{bmatrix} u_1^0 \\ u_2^0 \end{bmatrix} = \begin{bmatrix} 0.5 \\ 1.0 \end{bmatrix}$  and  $\begin{bmatrix} \dot{u}_1^0 \\ \dot{u}_2^0 \end{bmatrix} = \begin{bmatrix} 0.0 \\ 0.0 \end{bmatrix}$

- Numerical integration - Newmark, Time step size:  $\Delta t = 0.05$

**Goal:** • global displacement error:

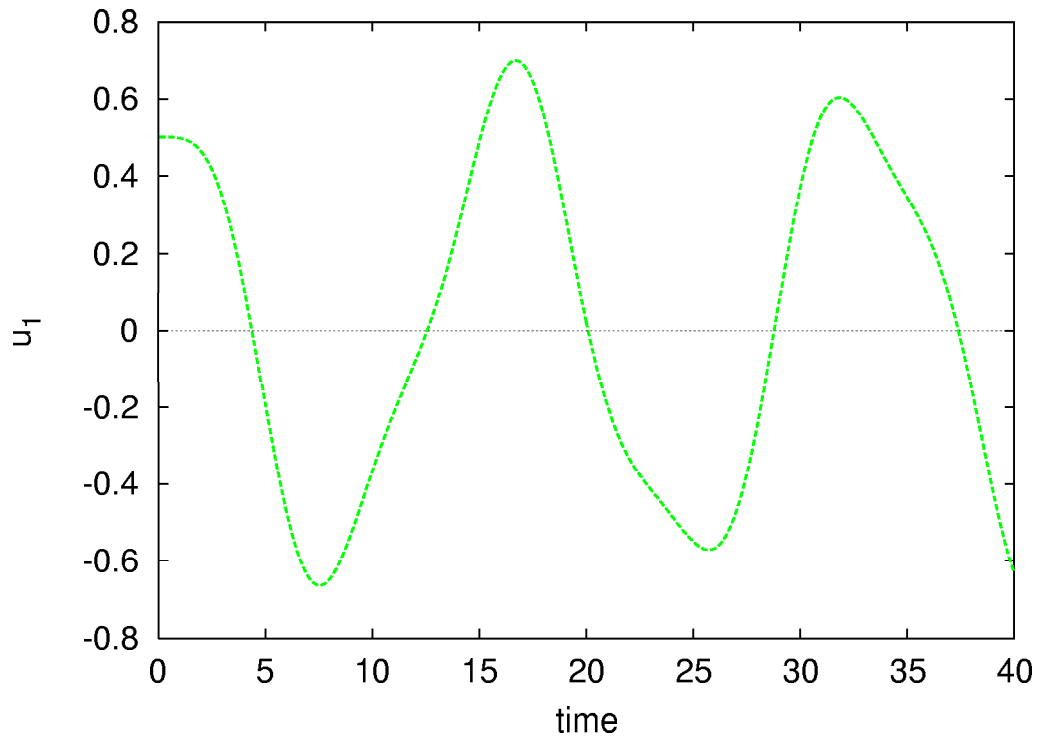
$$e_g(t = t_m) = \left[ (u_1 - u_{h,1})_{t=t_m}^2 + (u_2 - u_{h,2})_{t=t_m}^2 \right]^{1/2}$$

- Initial conditions – adjoint problem:

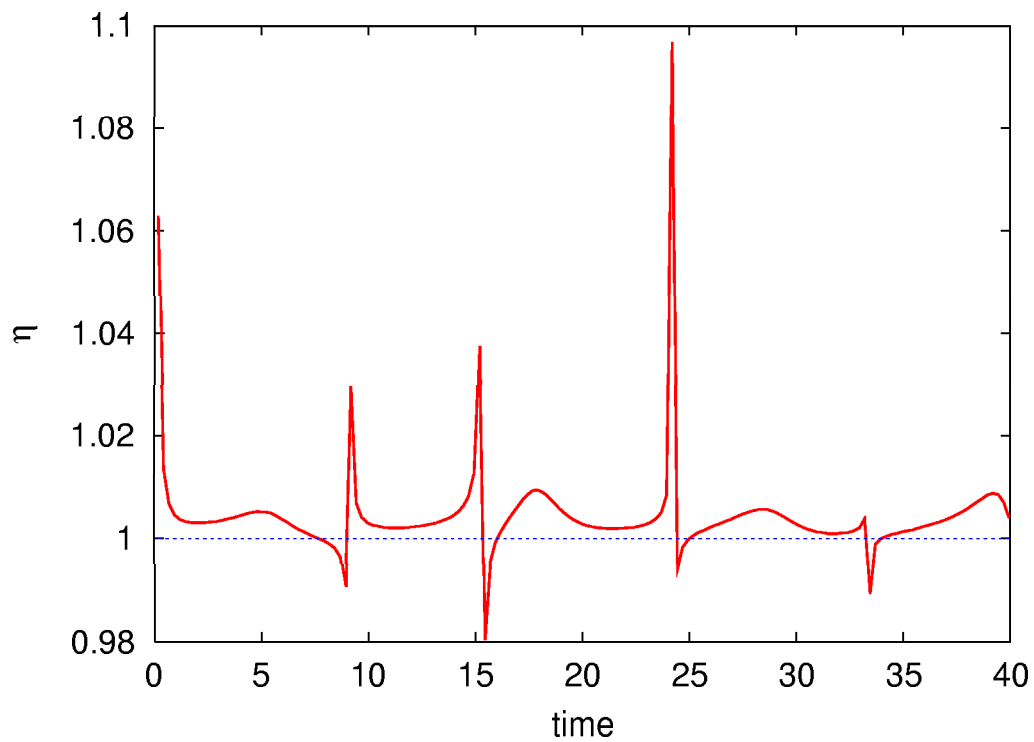
$$z(t = t_m) = \mathbf{0}, \quad \dot{z}(t = t_m) = -\mathbf{M}^{-1} \frac{\mathbf{e}_l(t_m)}{|\mathbf{e}_l(t_m)|}$$

## Two degree of freedom system:

Results: •  $u_1(t)$  - displacement



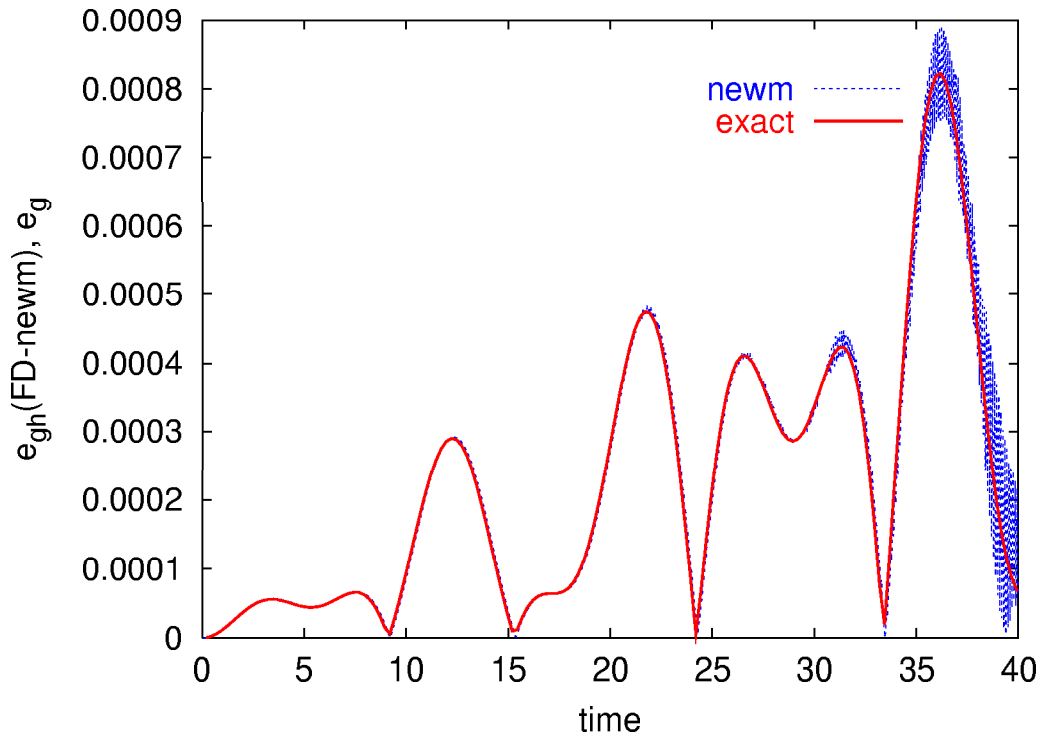
• efficiency index  $\eta(t)$  for  $e_{gh}(CG2 - newm)$



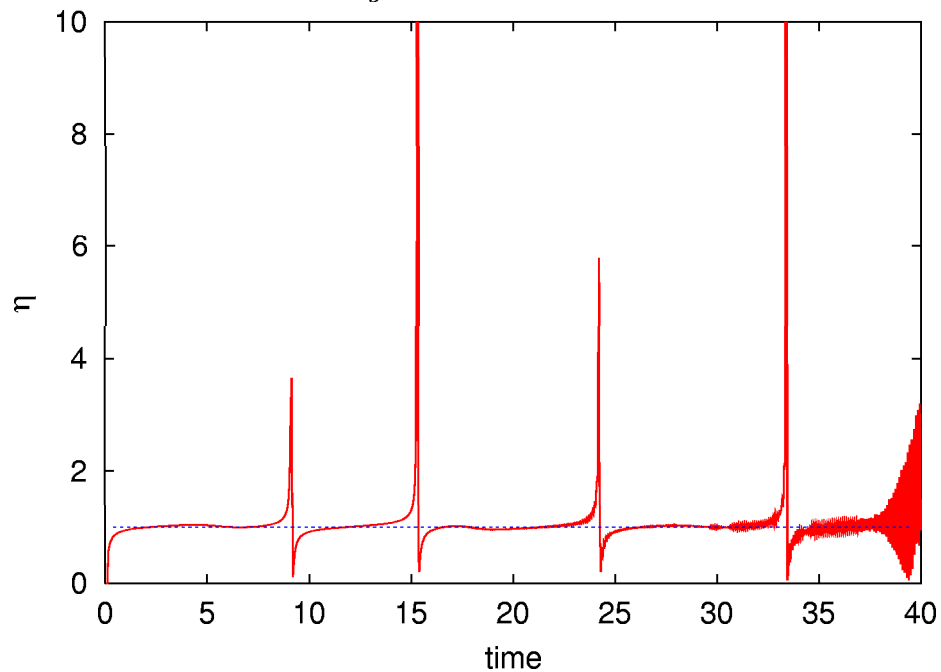
$$\eta = \frac{e_{gh}}{e_g}, \approx 1$$

## Two degree of freedom system: contd.

Results: • Estimated error  $e_{gh}(FD - newm)$  compared to exact error



• efficiency index  $\eta(t)$  for  $e_{gh}(FD - newm)$



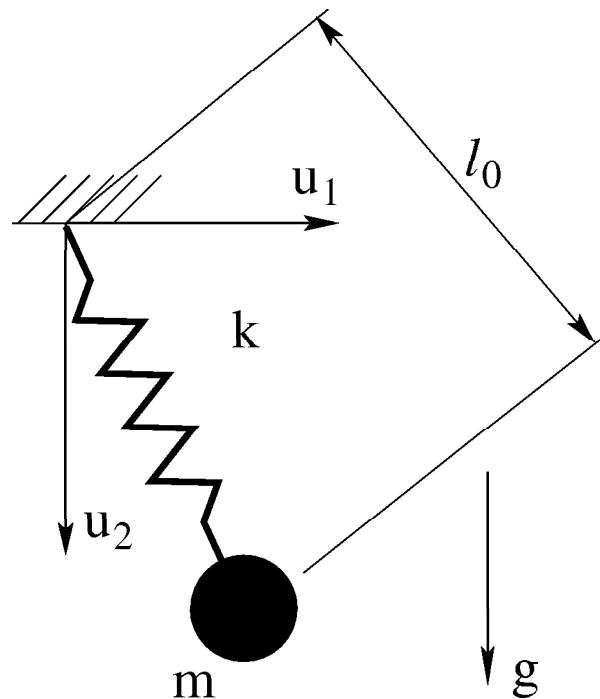
growing oscillations of  $\eta = \frac{e_{gh}}{e_g}$ ,

→ higher efficiency index at lower values

→ appears usable for wide region



## Spring pendulum – nonlinear problem



$$\begin{bmatrix} m\ddot{u}_1 + ku_1 \frac{\sqrt{u_1^2 + u_2^2} - l_0}{\sqrt{u_1^2 + u_2^2}} = 0 \\ m\ddot{u}_2 + ku_2 \frac{\sqrt{u_1^2 + u_2^2} - l_0}{\sqrt{u_1^2 + u_2^2}} = mg \end{bmatrix}$$

for  $t > 0$

initial conditions:

$$u_1^0 = 10^{-11}$$

$$u_2^0 = 1.5$$

$$\dot{u}_1^0 = 0.0$$

$$\dot{u}_2^0 = 0.0$$

→ dominant vertical initial condition,  
minimal horizontal initial deflection

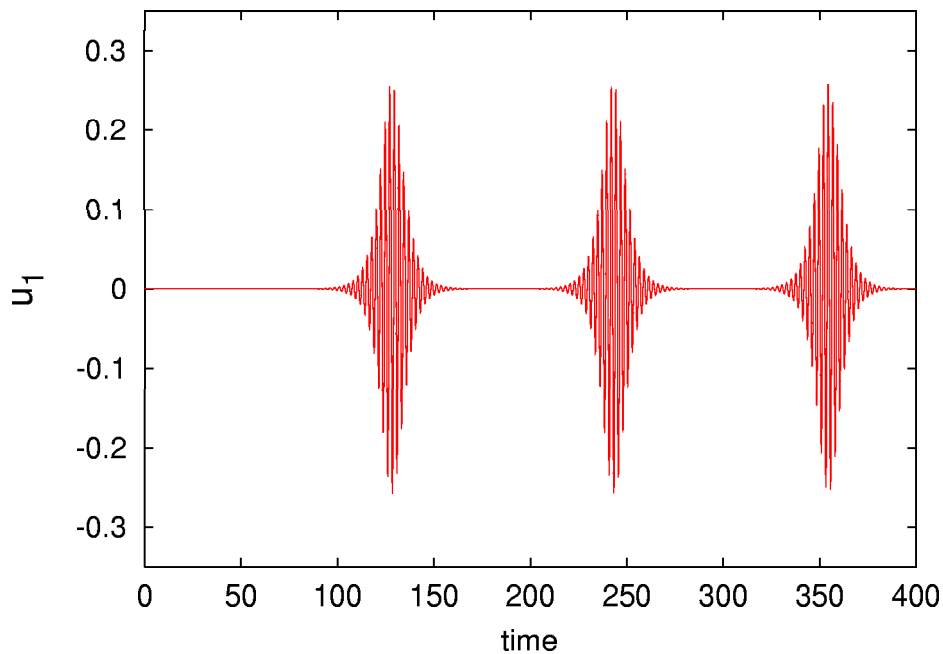
Parameters:  $k = 38.5$ ,  $m = 1.4$ ,  $l_0 = 1.0$

Numerical time integration CG2,  $\Delta t = 0.05$ ,

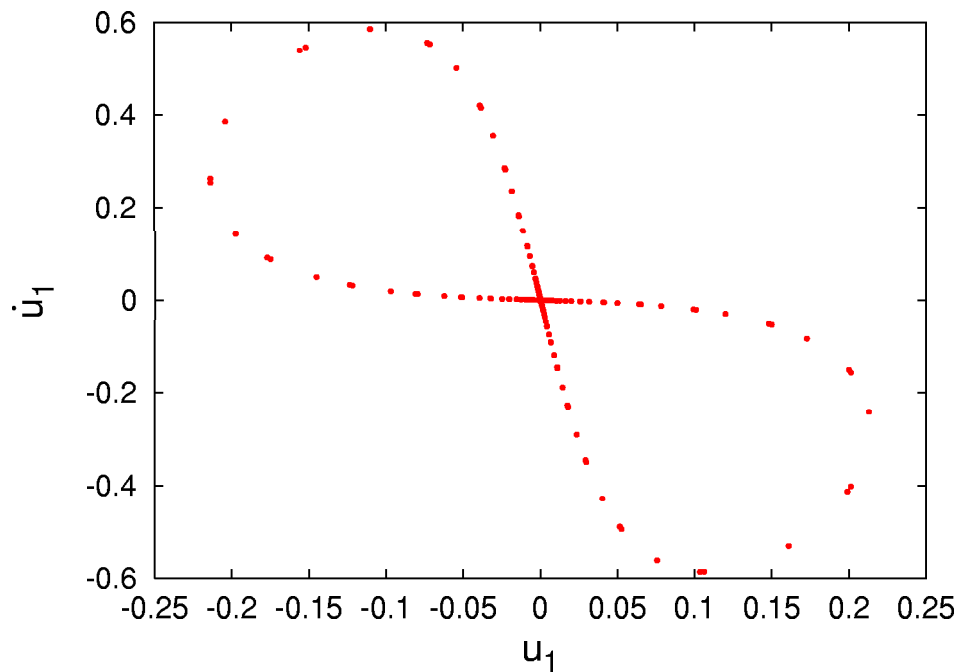
Reference ('exact') solution with  $\Delta t(ref) = 0.000625$

## Spring pendulum – contd.

Results: •  $u_1(t)$  horizontal displacement



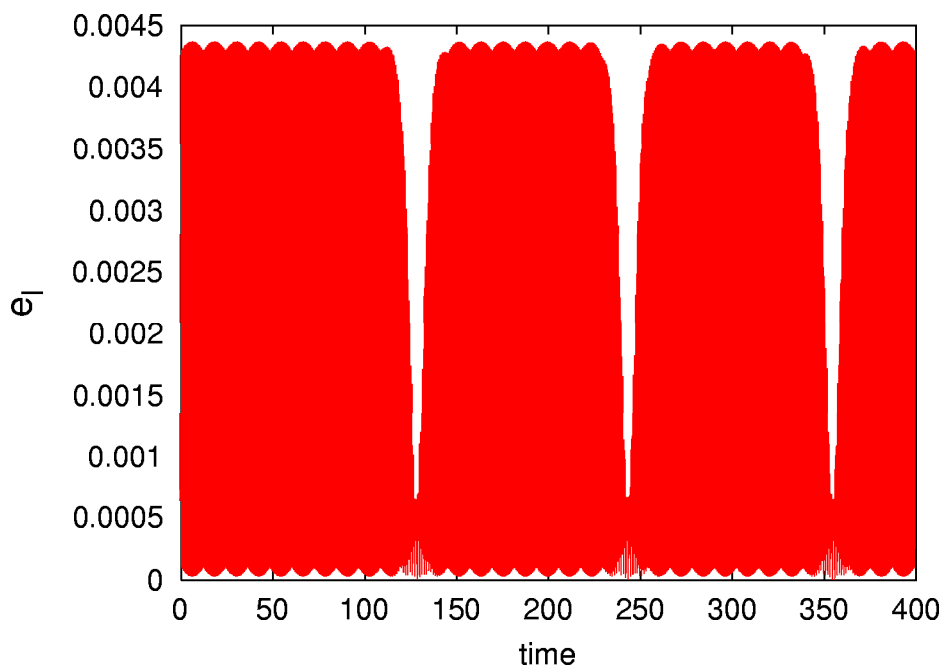
• Poincaré map  $u_1 - \dot{u}_1$  for  $u_2 - (l_0 + mg/k) = 0$ ,  $\dot{u}_2 > 0$



→ Poincaré map shows quasi-periodic solution,  
Behaviour: perturbation of primal vibration in vertical direction  
with and  $u_1 = \dot{u}_1 = 0$

## Spring pendulum – contd.

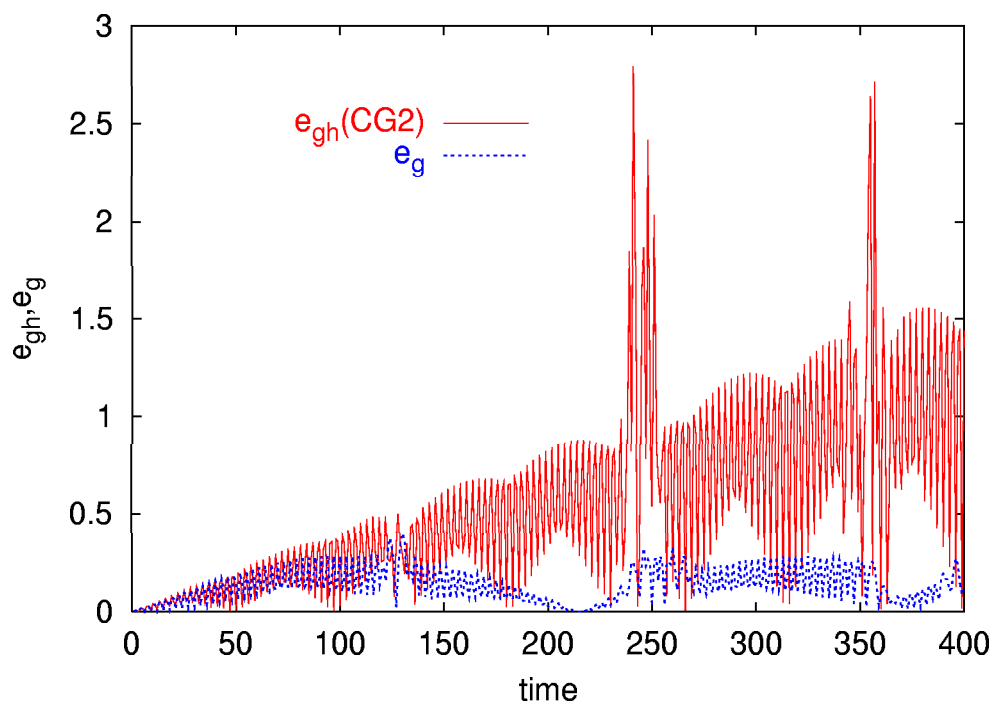
Results: • Estimated local time integration error  $e_l$  (from FD considerations)



Minima of estimated local error at  $t \approx 125/240/360$

→ not usable for global error estimation

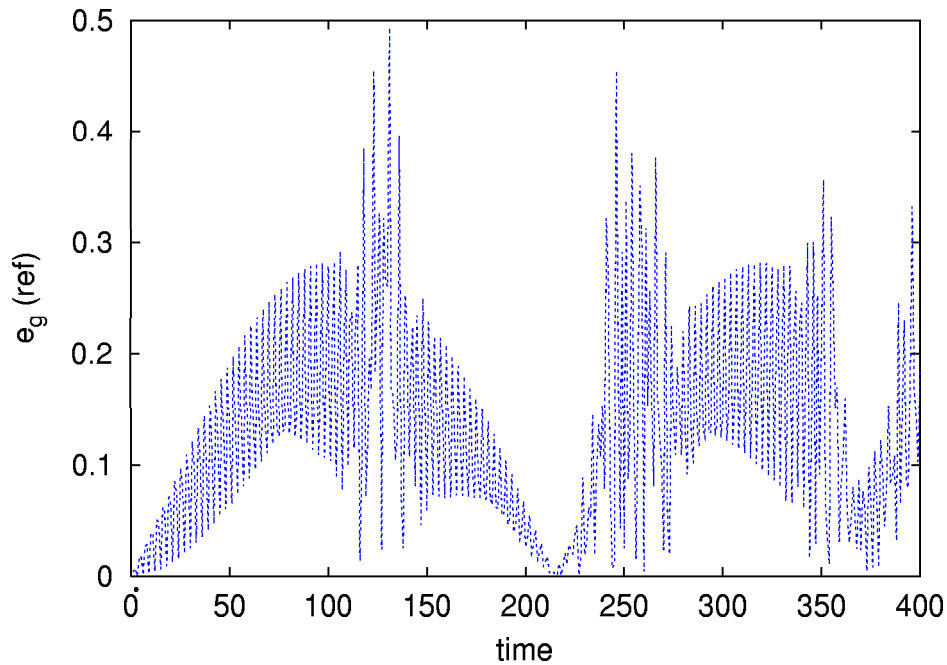
• Estimated global time integration error  $e_{gh}(cg2)$  and exact error  $e_g$



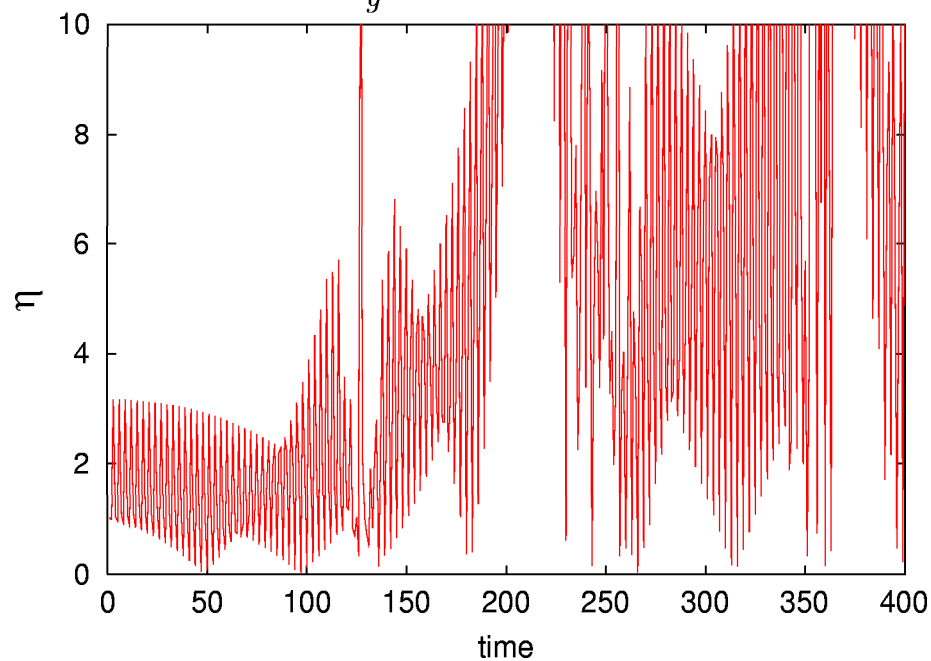
'Bifurcation' from longitudinal motion to horizontal motion at  $t \approx 125/240/360$ ,  
→ Maxima of estimated global error at  $t \approx 125/240/360$ ,

## Spring pendulum – contd.

**Results:** • Global time integration error using reference solution,  $\Delta t(ref) = 0.000625 \rightarrow$  'exact' error



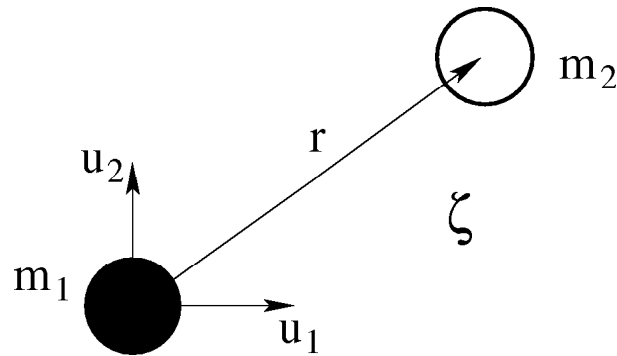
• Efficiency index  $\eta = \frac{e_{gh}}{e_g}$



- Maxima of 'exact' global time integration error at  $t \approx 125/240/360$ ,
- error somehow captured by error estimator only up to first 'bifurcation'
- But:**  $t > 100: \eta \gg 1 \rightarrow$  huge overestimation of global time integration error

## Two body interaction nonlinear problem

e.g. celestial mechanics



initial conditions:

$$\left[ \begin{array}{l} m\ddot{u}_1 + \zeta \frac{u_1}{(u_1^2 + u_2^2)^{3/2}} = 0 \\ m\ddot{u}_2 + \zeta \frac{u_2}{(u_1^2 + u_2^2)^{3/2}} = 0 \end{array} \right] \text{ for } t > 0$$
$$\begin{array}{l} u_1^0 = 0.4 \\ u_2^0 = 0.0 \\ \dot{u}_1^0 = 0.0 \\ \dot{u}_2^0 = 2.0 \end{array}$$

Parameters:  $\zeta = 1.0$  (gravitational constant),  $m = \frac{m_1 m_2}{m_1 + m_2} = 1$

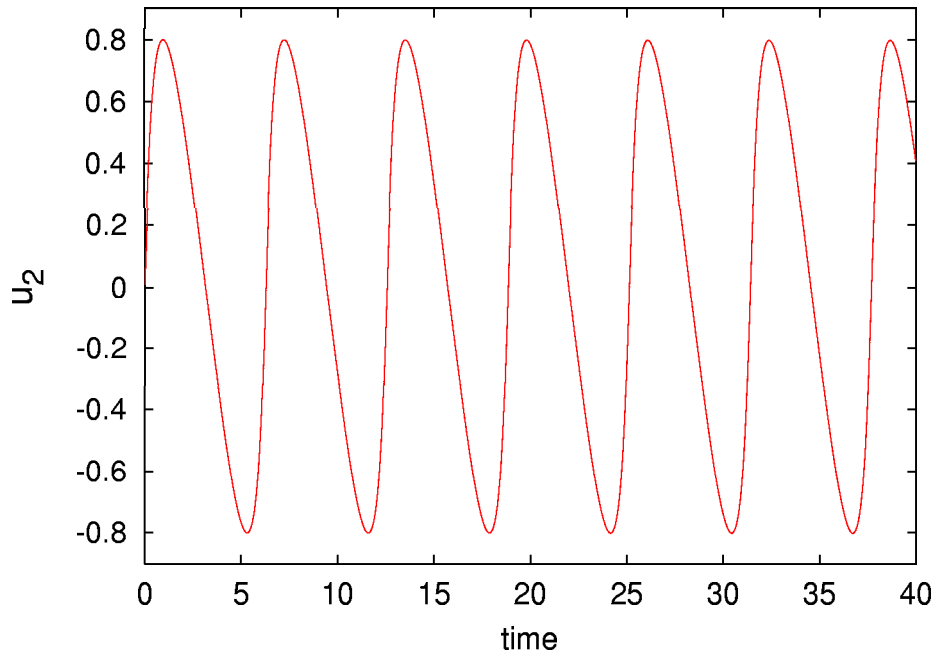
• Numerical time integration CG2-newm,  $\Delta t = 0.05$ ,

• Exact solution:  $u_1(t) = \cos(\tau) - 0.6$ ,

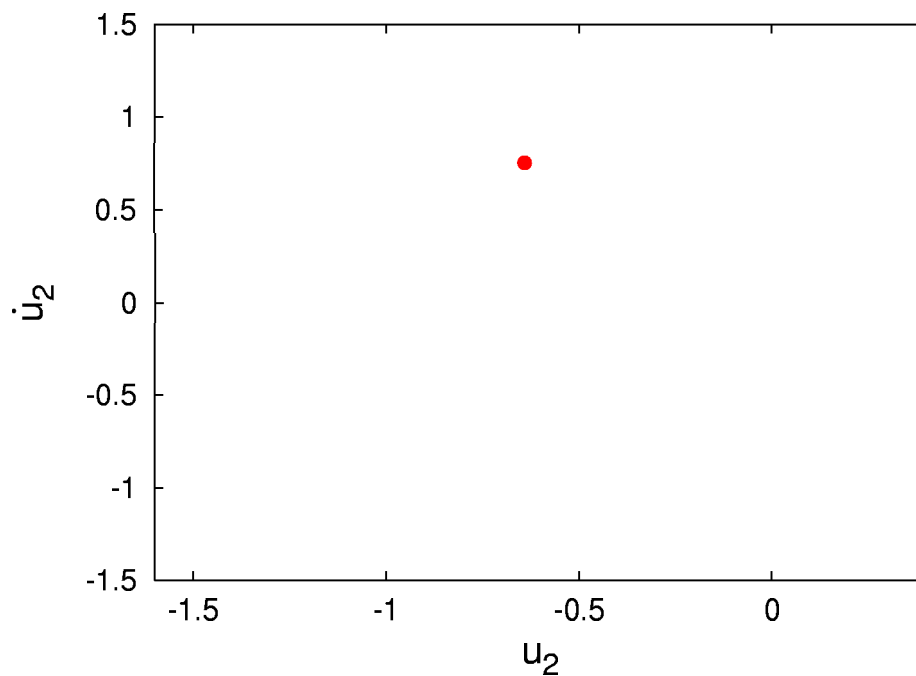
$$u_2(t) = 0.8 \sin(\tau), \quad \text{with: } t = \tau - 0.6 \sin(\tau)$$

## Two body interaction – contd.

Results: •  $u_2(t)$  – vertical displacement



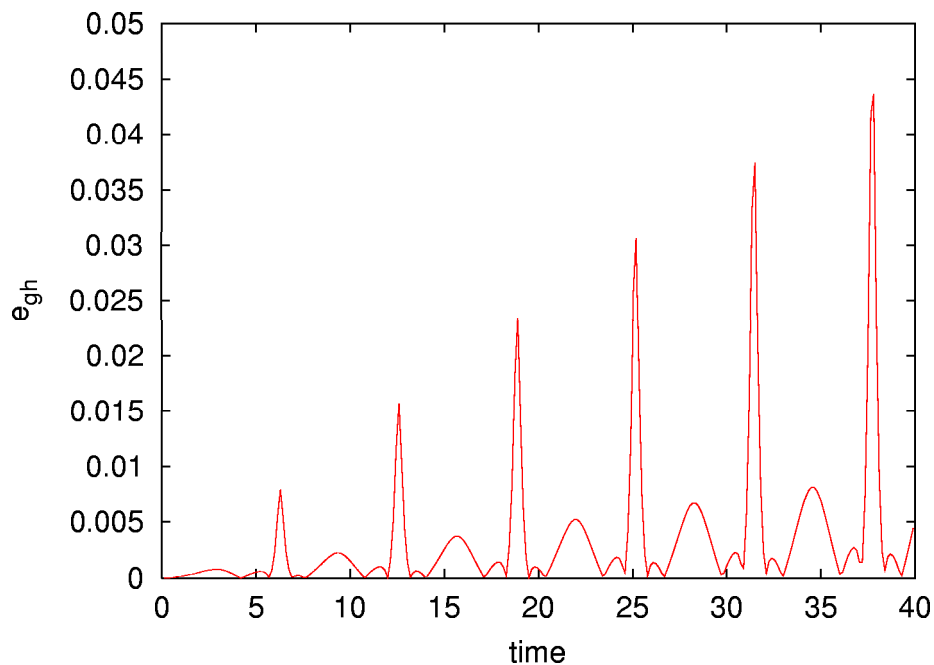
• Poincaré map  $u_2 - \dot{u}_2$  at  $u_1 = 0, \dot{u}_1 > 0$



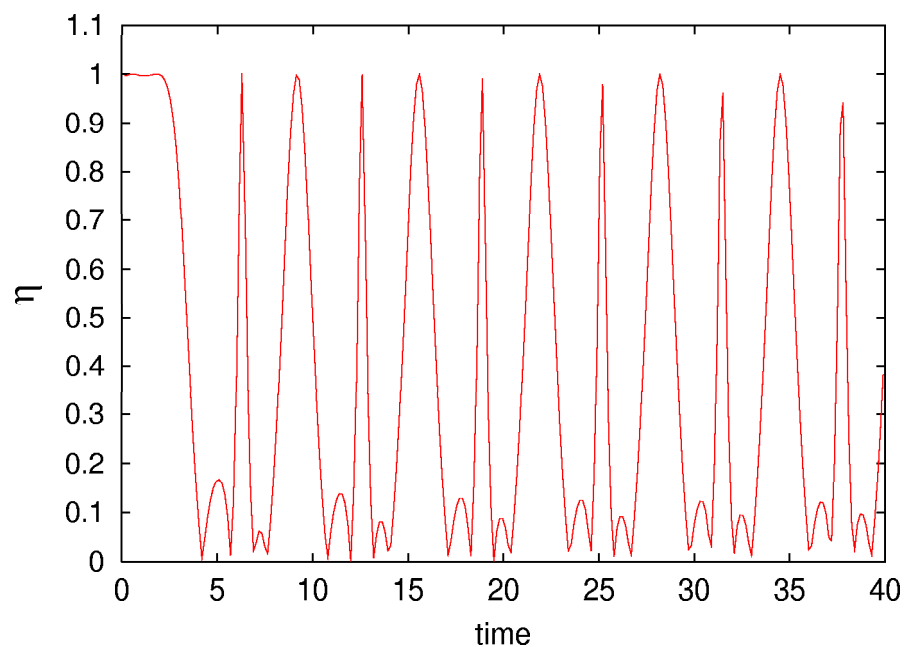
→ periodic solution

## Two body interaction – contd.

Results: • Estimated global time integration error  $e_{gh}(cg2 - newm)$  based on adjoint problem



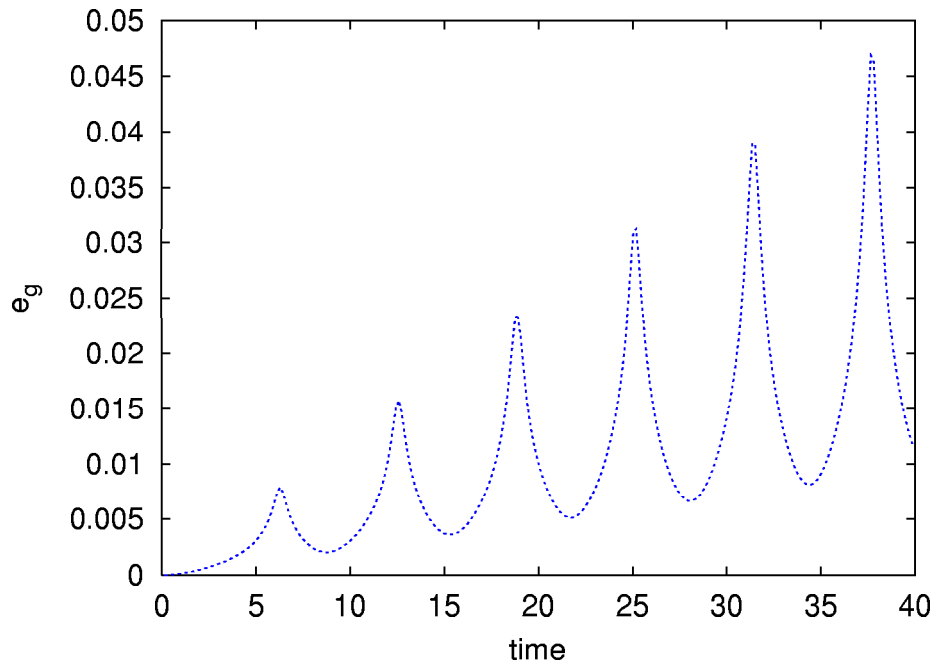
• Efficiency index  $\eta = \frac{e_{gh}}{e_g}$



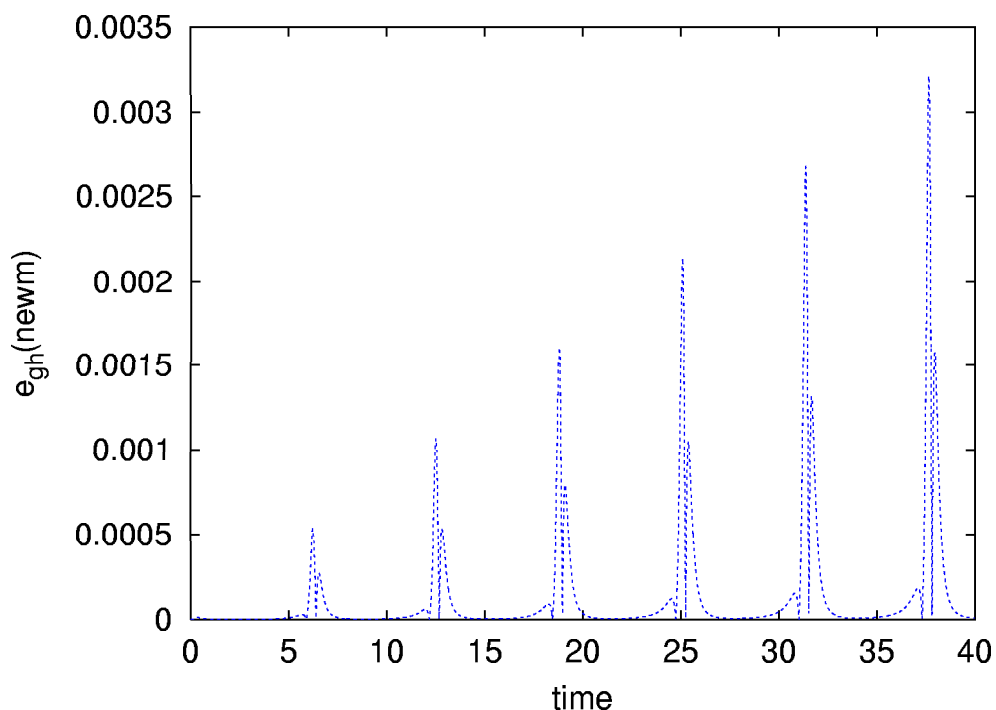
→ 'smooth' periodic solution →  $\eta \approx 1$  at maxima of error ,

## Two body interaction – contd.

- Results:** • Exact global time integration error  $e_g$  for Finite Element scheme (CG2-newm)



- Estimated global time integration error based on FD-version of Newmark  $e_{gh}(FD - newm)$

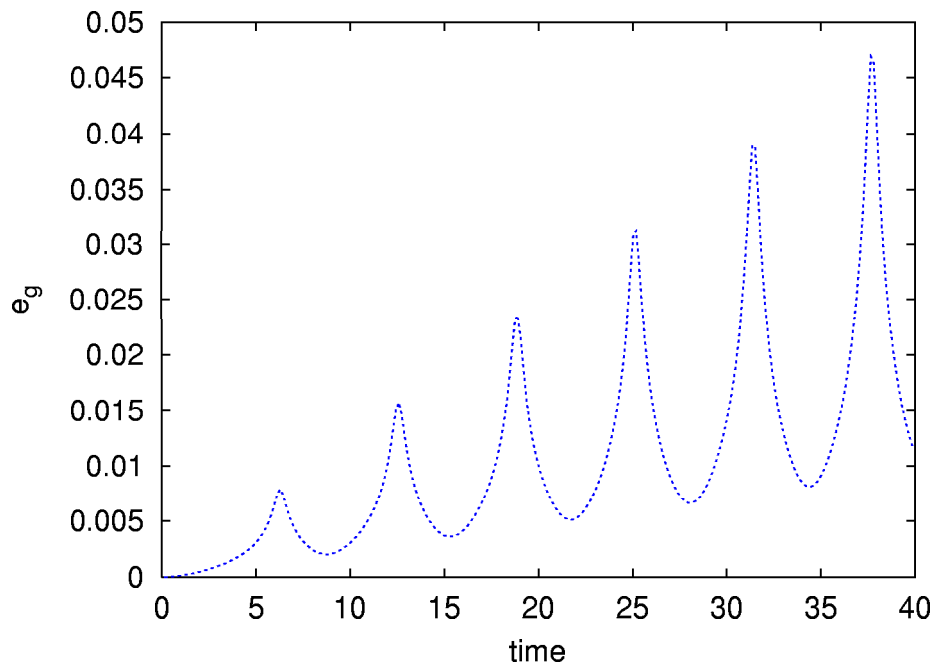


→ times of maxima of  $e_{gh}(FD - newm)$  correspond to times of maxima in  $e_g$  and  $e_{gh}(cg2)$

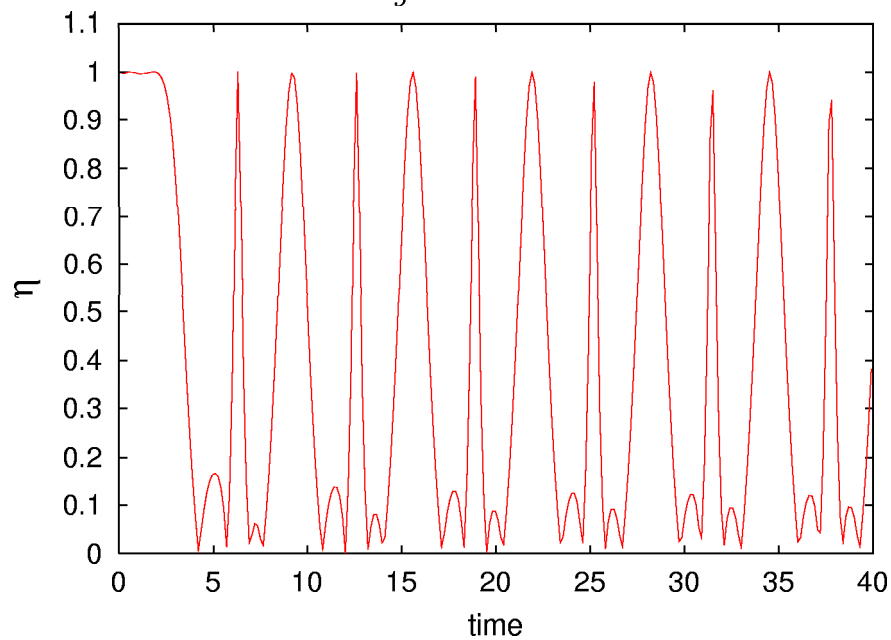


## Two body interaction – contd.

- Results: • Exact global time integration error  $e_g$   
for Finite Element scheme (CG2-newm)



- Efficiency index  $\eta = \frac{e_{gh}}{e_g}$



## Time step adaptation scheme

Based on combination of estimator for global and local time integration error

Global estimator:  $e_{gh}(CG2 - newm) = \int_0^{t_m} \mathbf{R} \mathbf{z}_h dt$ , Local indicator:  $e_l$

**Idea** • perform multiple analysis

- assume upper/lower tolerance for time step (local tolerance)
- adjust local time step size
- compute global error, check maximum of analysis

**Basis analysis** • compute with constant time step, determine max. displ.

- set global tolerance  $gtol =$  fraction of max. displacement
- local tolerances  $upper\ ltol = \frac{gtol}{N}$ ,  $lower\ ltol = \frac{upper\ ltol}{1...10}$

## Adaptive error control procedure

a) check local error  $e_l$

modify  $\Delta t_m$ , if

$e_l < lower\ ltol$  or

$e_l > upper\ ltol$  then  $\rightarrow \Delta t_m(new) = \Delta t_m \sqrt[3]{\frac{ltol}{e_l}}$

b) check global error, compute  $emax = \max_{t \in (0, T)} e_{gh}$

modify  $lower\ ltol$  and  $upper\ ltol$  if

$emax > gtol$  then

$$upper\ ltol(new) = \left( \frac{gtol}{emax} \right)^{3/2} upper\ ltol(old),$$

$$lower\ ltol(new) = \left( \frac{gtol}{emax} \right)^{3/2} lower\ ltol(old),$$

$$\rightarrow \Delta t_0(new) = \left( \frac{gtol}{emax} \right)^{1/2} \Delta t_0(old)$$

*Idea:* global error is reduced quadratically, order of local error is  $\mathcal{O}(\Delta t)^3$

if  $\Delta t(new) = 0.5 \Delta t(old)$

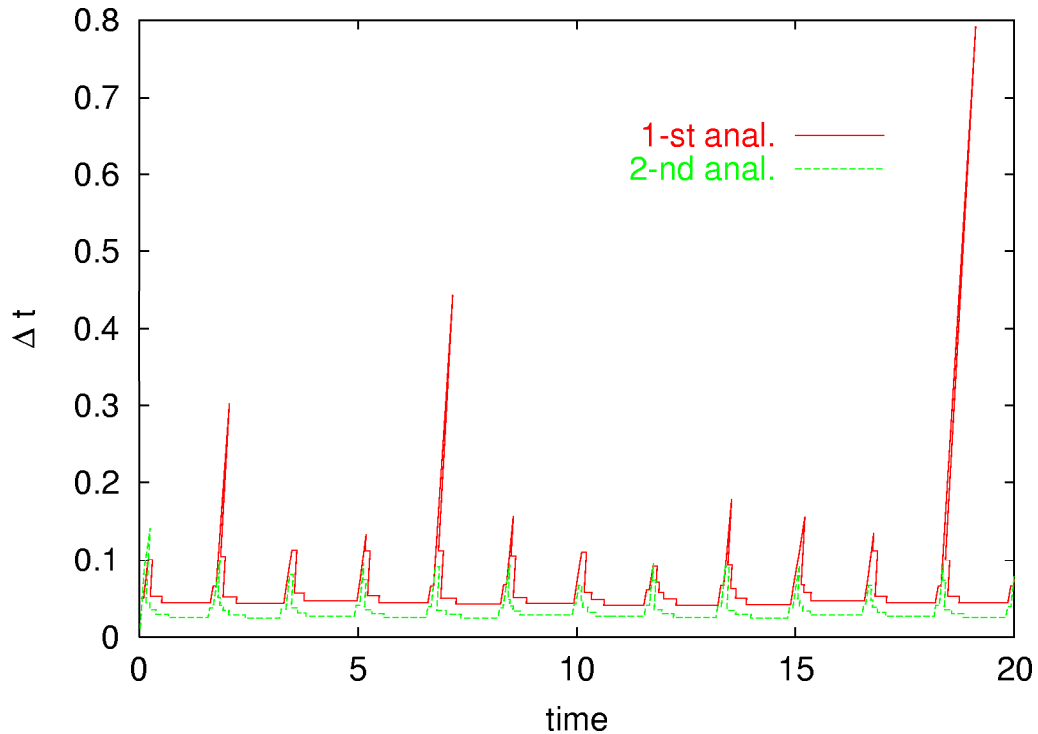
→ **Restart** at  $t = 0$  if  $emax > gtol$

## Time step adaptation

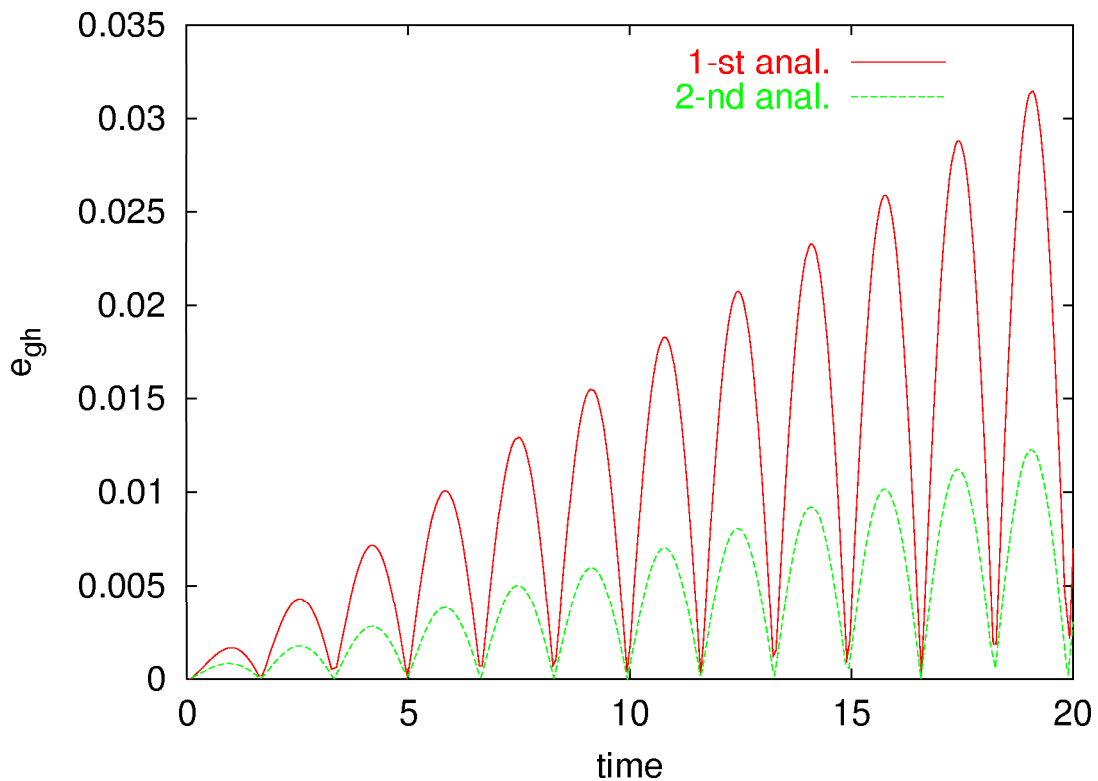
**SDOF** as before, starting tolerances:

$$\text{lower } ltol = 1.0 \cdot 10^{-5}, \text{ upper } ltol = 5.0 \cdot 10^{-5}, gtol = 0.015$$

**Results:** •  $\Delta t(t)$  - time step size



• Estimated global error  $e_{gh}(CG2 - newm)$



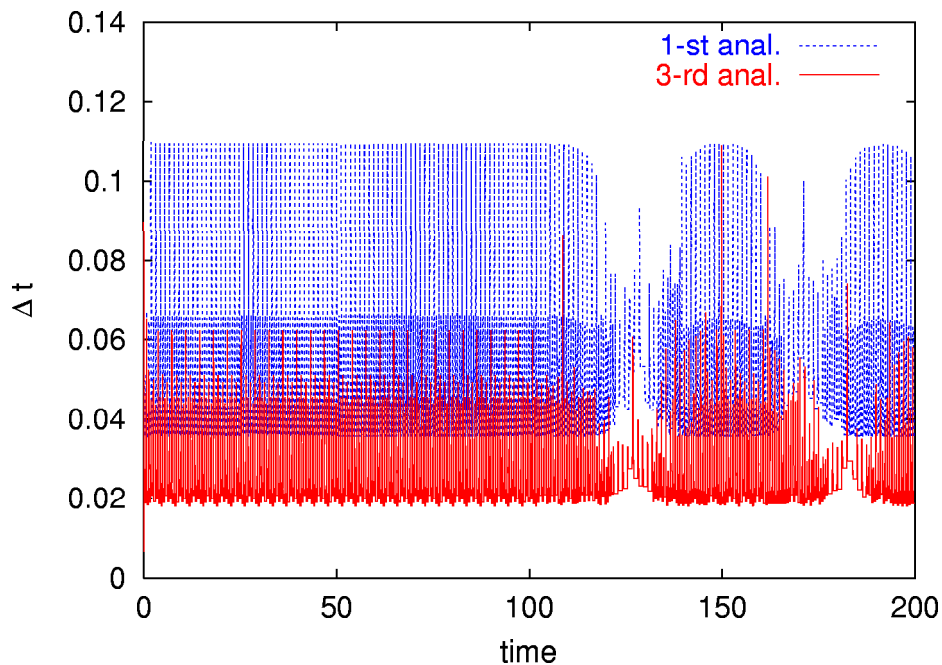
→ in 2-nd analysis  $e_{max} < gtol$

## Time step adaptation

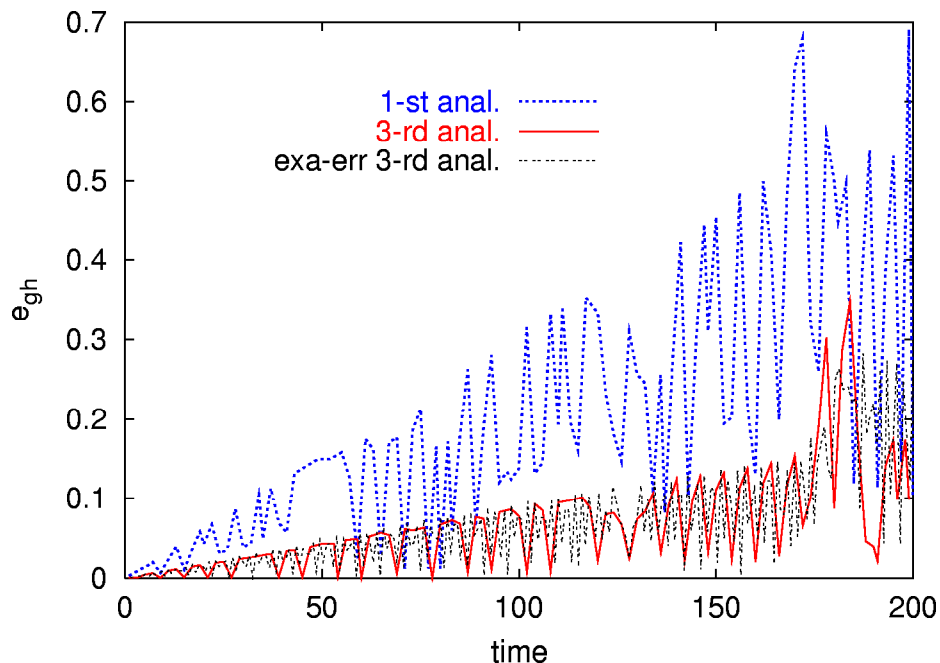
**Spring pendulum:** starting tolerances:

$$\text{lower tol} = 2.0 \cdot 10^{-5}, \text{ upper tol} = 1.0 \cdot 10^{-4}, \text{ gtol} = 0.6$$

**Results:** •  $\Delta t(t)$  - time step size



• Estimated global error  $e_{gh}(CG2 - newm)$ ; Exact error (exa-err)



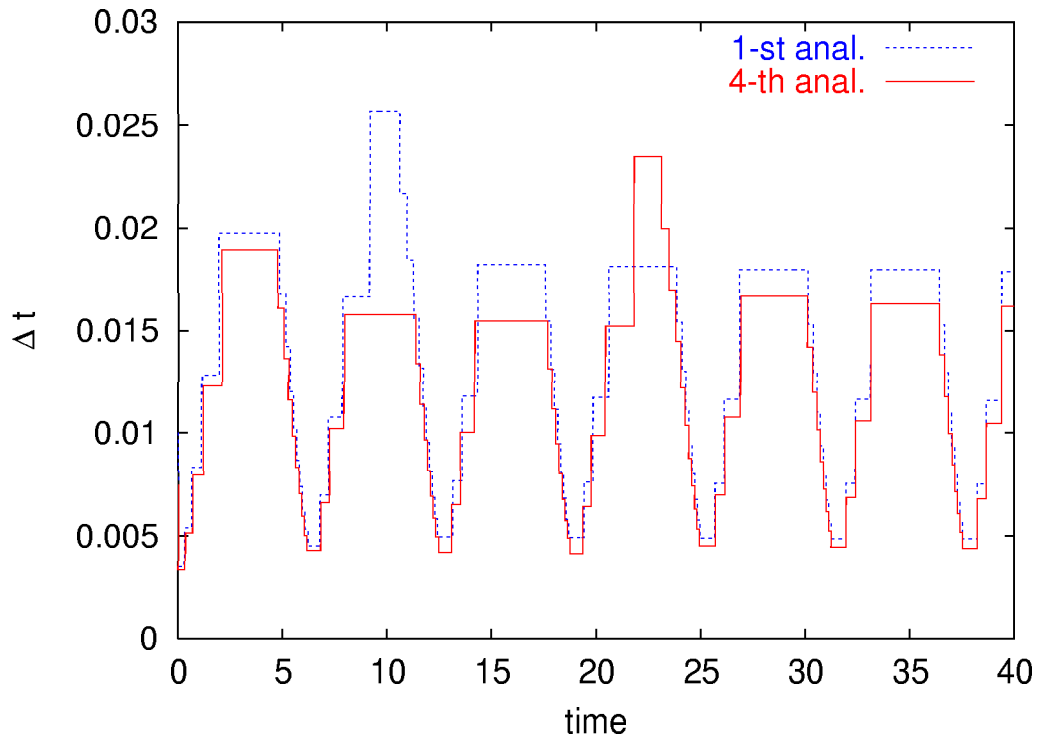
- adap. scheme reduces error for complete time range, in 3-rd anal.  $emax < gtol$
- better results for time before 'bifurcation' (at  $t \approx 125$ ),
- after  $t \approx 125$  reasonable results, but not 'perfect' error reduction

## Time step adaptation

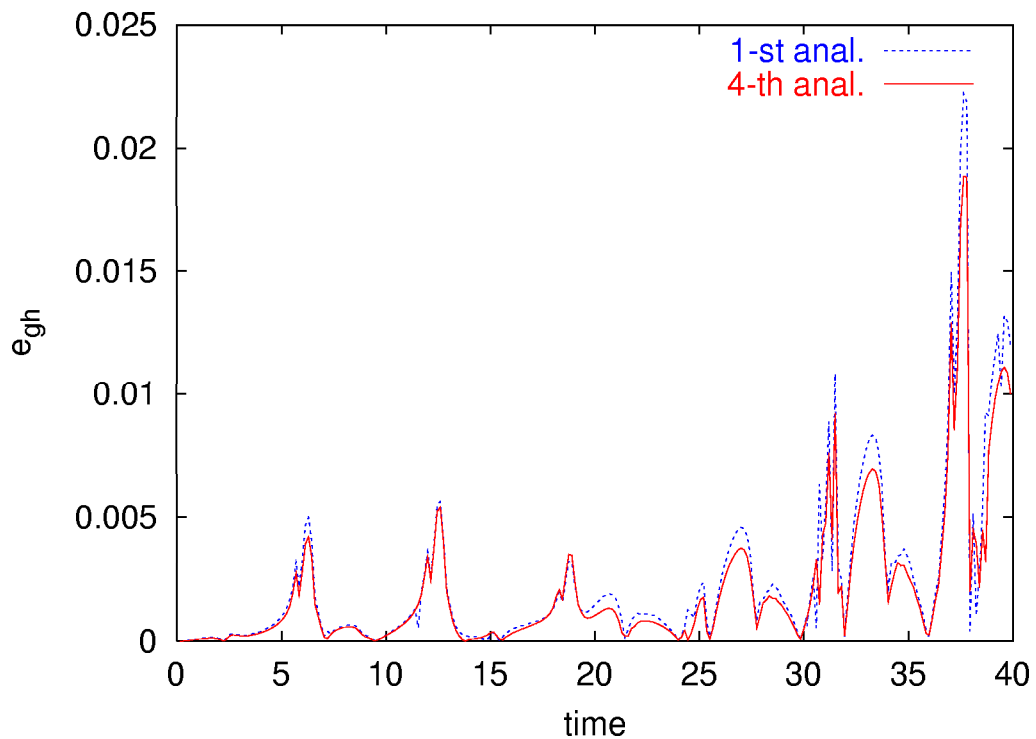
**Two body interaction:** starting tolerances:

$$\text{lower } ltol = 5.0 \cdot 10^{-8}, \text{ upper } ltol = 3.0 \cdot 10^{-7}, gtol = 0.02$$

**Results:** •  $\Delta t(t)$  - time step size



• Estimated global error  $e_{gh}(CG2 - newm)$



→ in 2-nd analysis  $emax < gtol$

## Conclusions

- effectivity and efficiency of local and global error indicator / estimator for structural dynamic problems shown
- simple local and FD-based global indicator for linear problems reasonably applicable
- FE-based global estimator
  - high efficiency for linear and mildly nonlinear problems
  - less efficient for highly nonlinear problems
- Time step adaptation scheme based on combination of global FE estimator and local FD indicator works reasonably well for linear and nonlinear problems