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Analysis of Current Harmonics in Inverter-Driven Interior Permanent Magnet Synchronous Machines

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<<Permanent magnet motor>>, <<Converter machine interactions>>, <<Harmonics>>, <<Modulation strategy>>

Abstract

This paper presents the analysis of current harmonics in an inverter-driven interior permanent magnet machine for different modulation methods. In contrast to other approaches, the influence of the saliency $(L_d \neq L_q)$ on the harmonics is being considered. By use of the rms-value of the currents, analytically closed equations can be deduced and tedious Fourier analyses of single harmonics can be avoided. The derivation of the analytic equations is described and the results are verified by numerical simulation as well as by experimental results.

Introduction

The influence of the modulation method on the current waveform in converter systems with high pulse rate has been investigated in different publications (e.g. [1] - [4]). There the calculation of the current harmonics has been done with a three-phase symmetric load at the converter output, which is represented by three equal inductances (fig. 1). However, in the case of magnetic anisotropic permanent magnet synchronous machines the inductance in each of the three output phases differs from the others and variates depending on the rotor position γ of the machine.



Figure 1: Voltage source inverter with symmetric load

Additionally, the phase angle between the stator voltage of an interior permanent magnet synchronous machine (IPMSM) and the rotor position γ depends on the stator current, which makes the values of the inductances at the inverter output also dependent on the operation point of the machine.

This paper presents the estimation of the current harmonics in a drive system consisting of a voltage source inverter (VSI, fig. 1) and an IPMSM depending on the modulation depth and the phase angle ϕ_U of the stator voltage for different modulation methods. Verification of the developed formulas is performed by numerical simulation of a drive system. Measurements on a machine test bench have also been accomplished in order to validate the derived formulas.

Calculation of current harmonics

In order to estimate the current harmonics, first the machine model and the switching states of the inverter are being described. Thereon, the current waveform within one half pulse period of the inverter will be examined. Subsequently, the current harmonics related to one full period of the inverter output voltages u_1, u_2, u_3 will be evaluated.

Machine model

Since the inductances of the IPMSM are constant in the rotor oriented dq-reference frame (fig. 2a), this frame will be used for the estimation of the current waveform. Thus the IPMSM can be described by the stator voltage equations and the equations for the flux linkages:

$$u_{\rm d} = R_{\rm s} \cdot i_{\rm d} + \frac{\mathrm{d}\psi_{\rm d}}{\mathrm{d}t} - \omega \cdot \psi_{\rm q} \qquad \qquad u_{\rm q} = R_{\rm s} \cdot i_{\rm q} + \frac{\mathrm{d}\psi_{\rm q}}{\mathrm{d}t} + \omega \cdot \psi_{\rm d} \qquad (1)$$

$$\psi_{\rm d} = L_{\rm d} \cdot i_{\rm d} + \psi_{\rm PM} \qquad \qquad \psi_{\rm q} = L_{\rm q} \cdot i_{\rm q} \tag{2}$$

The functions for the current component waveforms can be obtained by solving eq. (1) and (2) for the current derivatives. Assuming the inductances L_d and L_q and the speed ω as constant and neglecting the stator resistance R_s , this results in a system of ordinary differential equations:

$$\frac{\mathrm{d}i_{\mathrm{d}}}{\mathrm{d}t} = \frac{u_{\mathrm{d}} + \omega \cdot L_{\mathrm{q}} \cdot i_{\mathrm{q}}(t)}{L_{\mathrm{d}}} \qquad \qquad \frac{\mathrm{d}i_{\mathrm{q}}}{\mathrm{d}t} = \frac{u_{\mathrm{q}} - \omega \cdot (L_{\mathrm{d}} \cdot i_{\mathrm{d}}(t) + \psi_{\mathrm{PM}})}{L_{\mathrm{q}}} \tag{3}$$

Solving this system of ordinary differential equations, the functions for the current components become

$$i_{\rm d}(t) = \frac{\left[\left(L_{\rm d} \cdot i_{\rm d,0} + \psi_{\rm PM} \right) \cdot \omega - u_{\rm q} \right] \cdot \cos\left(\omega t\right)}{\omega \cdot L_{\rm d}} + \frac{\left(\omega \cdot L_{\rm q} \cdot i_{\rm q,0} + u_{\rm d} \right) \sin\left(\omega t\right) + u_{\rm q} - \omega \cdot \psi_{\rm PM}}{\omega \cdot L_{\rm d}} \tag{4}$$

$$i_{q}(t) = \frac{\left[-\left(L_{d} \cdot i_{d,0} + \psi_{PM}\right) \cdot \omega + u_{q}\right] \cdot \sin\left(\omega t\right)}{\omega \cdot L_{q}} + \frac{\left(u_{d} + \omega \cdot L_{q} \cdot i_{q,0}\right) \cdot \cos\left(\omega t\right) - u_{d}}{\omega \cdot L_{q}}$$
(5)

with the initial values $i_d(t = 0) = i_{d,0}$ and $i_q(t = 0) = i_{q,0}$.

In case of a high pulse rate, the variation of the rotor position angle $\Delta \gamma = \Delta \omega t$ during one half pulse period can also be neglected. Therefore the small-angle approximation $\sin(\omega t) \approx \omega t$ and $\cos(\omega t) \approx 1$ can be used. This simplifies eq. (4) and (5) to

$$i_{\rm d}(t) = \frac{u_{\rm d} + \omega \cdot L_{\rm q} \cdot i_{\rm q,0}}{L_{\rm d}} \cdot t + i_{\rm d,0} \qquad \qquad i_{\rm q}(t) = \frac{u_{\rm q} - \omega \cdot \left(L_{\rm d} \cdot i_{\rm d,0} + \psi_{\rm PM}\right)}{L_{\rm q}} \cdot t + i_{\rm q,0} \tag{6}$$

The voltage components u_d and u_q in eq. (6) depend on the discrete switching states of the inverter.

Switching states of the inverter

The converter voltage space vector is defined as

$$\underline{u}_{s} = \frac{2}{3} \cdot \left(u_{1} + \underline{a} \cdot u_{2} + \underline{a}^{2} \cdot u_{3} \right) \qquad \text{with} \qquad \underline{a} = -\frac{1}{2} + j \cdot \frac{\sqrt{3}}{2} \tag{7}$$

Assuming the location of the desired voltage space vector in sector I in fig. 2b, the discrete converter voltage space vectors used during the pulse period are

$$\underline{u}_1 = \frac{2}{3} \cdot U_{dc} \qquad \qquad \underline{u}_2 = \frac{2}{3} \cdot U_{dc} \cdot \left(\frac{1}{2} + j\frac{\sqrt{3}}{2}\right) \tag{8}$$



Figure 2: (a): Stator voltage space vector in the rotary dq-frame and its location in the stationary $\alpha\beta$ -frame. (b): Switching states and possible output voltage of the inverter.



Figure 3: Exemplary switching state sequence in sector I (a) and exemplary current waveform at $\gamma = 0$ (b).

Additionally, one or two of the freewheeling states 7 and 8 can be inserted in order to achieve the desired mean value of the voltage space vector during the pulse period T_p . For sinusoidal phase modulation, the switching state sequence in sector I is shown in fig. 3a.

The effect on the current components in the dq-frame can be obtained by transforming the discrete voltage space vectors from the stationary to the rotating frame. The resulting voltage components are

$$u_{d,1} = \Re\left\{\underline{u}_{\alpha\beta,1} \cdot e^{-j \cdot \gamma}\right\} = \frac{2}{3} \cdot U_{dc} \cdot \cos\gamma$$
(9)

$$u_{d,2} = \Re\left\{\underline{u}_{\alpha\beta,2} \cdot e^{-j \cdot \gamma}\right\} = \frac{1}{3} \cdot U_{dc} \cdot \left(\cos\gamma + \sqrt{3} \cdot \sin\gamma\right)$$
(10)

$$u_{\rm d,7} = u_{\rm d,8} = 0 \tag{11}$$

$$u_{q,1} = \Im\left\{\underline{u}_{\alpha\beta,1} \cdot e^{-j \cdot \gamma}\right\} = -\frac{2}{3} \cdot U_{dc} \cdot \sin\gamma$$
(12)

$$u_{q,2} = \Im \left\{ \underline{u}_{\alpha\beta,2} \cdot e^{-j \cdot \gamma} \right\} = \frac{1}{3} \cdot U_{dc} \cdot \left(\sqrt{3} \cdot \cos \gamma - \sin \gamma \right)$$
(13)

$$u_{q,7} = u_{q,8} = 0 \tag{14}$$

Fig. 3b shows exemplary waveforms of $i_d(t)$ and $i_q(t)$ within one pulse period assuming $\gamma = 0$.

Current harmonics during one half pulse period

During the following considerations, the mean value of the inverter output voltage within one half pulse period will be assumed to be equal to the stator voltage of the machine. The inner voltage components

$$e_{d} = -\omega \cdot L_{q} \cdot i_{q,0} \qquad \qquad e_{q} = \omega \cdot \left(L_{d} \cdot i_{d,0} + \psi_{PM} \right) \tag{15}$$

are approximately constant during $\frac{T_p}{2}$ for high pulse rates [1]. They can be expressed by use of the modulation depth *M*, which is defined as (cf. [2])

$$M = \frac{2 \cdot |\underline{u}_{\rm s}|}{U_{\rm dc}} \tag{16}$$

For steady-state operation the stator voltage is equal to the inner voltage. The angle φ_U of the stator voltage space vector in the dq-frame (fig. 2a) then depends on the inner voltage components:

The components of the inner voltage can be expressed in terms of the modulation depth by

$$e_{\rm d} = \frac{1}{2} \cdot M \cdot U_{\rm dc} \cdot \cos \varphi_{\rm U} \qquad \qquad e_{\rm q} = \frac{1}{2} \cdot M \cdot U_{\rm dc} \cdot \sin \varphi_{\rm U} \qquad (18)$$

The current during each of the switching states $x \in \{1, 2, 7, 8\}$ in one half pulse period then is

$$i_{d}(t) = \frac{u_{d,x} - \frac{1}{2} \cdot M \cdot U_{dc} \cdot \cos \varphi_{U}}{L_{d}} \cdot t + i_{d,k-1} \qquad i_{q}(t) = \frac{u_{q,x} - \frac{1}{2} \cdot M \cdot U_{dc} \cdot \sin \varphi_{U}}{L_{q}} \cdot t + i_{q,k-1}$$
(19)

with $i_{d/q,k-1}$ being the value of $i_{d/q}$ at the moment of changing the switching state. To calculate the current harmonics, the mean values of i_d and i_q during one half pulse period are not of interest. Therefore the initial current component values $i_{d,0}$ and $i_{q,0}$ at t = 0 can be set to zero.

To get the square of the rms-value of the harmonic current during one half pulse period, the squares of the two currents are integrated to:

$$\Delta i_{\rm s}^2 = \frac{2}{T_{\rm p}} \int_0^{\frac{T_{\rm p}}{2}} \frac{i_{\rm d}^2(t) + i_{\rm q}^2(t)}{2} \,\mathrm{d}t \tag{20}$$

Since the voltage space vector applied to the machine needs to be equal to the stator voltage of the machine $\underline{u}_s = e_d + j e_q$, the turn-on times of the different switching states for sinusoidal pulse width modulation (SPWM) can be estimated directly from the stator voltage (cf. fig. 3a):

$$\frac{t_7}{T_{\rm p}/2} = \frac{1}{2} \cdot \left[1 - M \cdot \cos\left(\varphi_{\rm U} + \gamma\right)\right] \tag{21}$$

$$\frac{t_1}{T_{\rm p}/2} = \frac{1}{2} \cdot \left[1 - M \cdot \cos\left(\varphi_{\rm U} + \gamma - \frac{2\pi}{3}\right) \right] - \frac{t_7}{T_{\rm p}/2}$$
(22)

$$\frac{t_2}{T_p/2} = \frac{1}{2} \cdot \left[1 - M \cdot \cos\left(\varphi_U + \gamma - \frac{4\pi}{3}\right) \right] - \frac{t_7}{T_p/2} - \frac{t_1}{T_p/2}$$
(23)

$$\frac{t_8}{T_p/2} = 1 - \frac{t_7}{T_p/2} - \frac{t_1}{T_p/2} - \frac{t_2}{T_p/2}$$
(24)

Eq. (21) to (24) are only valid for sector I of the voltage plane in fig. 2b, so the rotor position angle has to be set to the corresponding interval

$$\gamma \in \left\{-\phi_{\rm U}\dots - \phi_{\rm U} + \frac{\pi}{3}\right\} \tag{25}$$

With the estimation of the turn-on times of the switching states, the square of the harmonic current rms value Δi_s^2 within one half pulse period is determined.



Figure 4: (a): Current harmonics for SPWM. (b): Current harmonics for magnetic isotropic machines for use of different modulation methods

RMS value of harmonic currents related to one full period of the fundamental

Calculation of the mean value of the particular harmonic currents related to the pulse periods within sector I leads to the square of the resulting harmonic current rms value:

$$\Delta I^2 = \frac{1}{m} \cdot \sum_{m} \Delta i_s^2 \tag{26}$$

Due to the symmetries of the output characteristics of the inverter, this is equivalent to the mean value for an entire fundamental [2]. Assuming high pulse rate in relation to the fundamental, eq. (26) can be approximated by integration:

$$\Delta I^2 \approx \frac{3}{\pi} \cdot \int_{-\varphi_{\rm U}}^{-\varphi_{\rm U} + \pi/3} \Delta i_{\rm s}^2 \,\mathrm{d}\gamma \tag{27}$$

Evaluating ΔI^2 leads to the result for sinusoidal modulation (SPWM):

$$\Delta I_{\text{SPWM}}^{2} = \left(\frac{T_{\text{P}}}{2}\right)^{2} \cdot \left(\frac{U_{\text{dc}}}{l L_{\text{d}}}\right)^{2} \cdot M^{2} \cdot \left[\frac{1}{128}\left(1 + (l^{2} - 1)\left(\cos\left(\varphi_{\text{U}}\right)\right)^{2}\right) \cdot M^{2} - \frac{\sqrt{3}\left(12\left(\cos\left(\varphi_{\text{U}}\right)\right)^{2}\left(l^{2} - 1\right) + 11 - l^{2}\right)}{360 \pi} \cdot M + \frac{1}{96}\left(1 + (l^{2} - 1)\left(\cos\left(\varphi_{\text{U}}\right)\right)^{2}\right)\right]$$
(28)

with the saliency factor

$$l = \frac{L_{\rm q}}{L_{\rm d}} \tag{29}$$

The result for $I_{\rm B} = \Delta I_{\rm SPWM} / \left(\frac{T_{\rm P}}{2} \cdot \frac{U_{\rm dc}}{L_{\rm d}}\right)$ for a machine with $L_{\rm d} = 0.3$ mH and $L_{\rm q} = 1.5$ mH is shown in fig. 4a. In eq. (28) as well as in fig. 4a the dependency of ΔI on the angle $\varphi_{\rm U}$ of the stator voltage can clearly be seen. The minimum of the current harmonics for constant modulation depth appears at $\varphi_{\rm U} = \frac{\pi}{2}$ and $\varphi_{\rm U} = \frac{3\pi}{2}$, which corresponds to the direction of $L_{\rm q}$. The highest value of the current harmonics occur at $\varphi_{\rm U} = 0$ and $\varphi_{\rm U} = \pi$, as the voltage space vector is located in the direction of the smaller inductance $L_{\rm d}$. However, at modulation depths near $M \approx 1$ the current harmonics become higher for $\varphi_{\rm U} = \frac{\pi}{2}$ and $\varphi_{\rm U} = \frac{3\pi}{2}$ than for $\varphi_{\rm U} = 0$ and $\varphi_{\rm U} = \pi$.



Figure 5: Reference functions. (a): 3rdPWM. (b): SYPWM. (c): DPWM1. (d): DPWM2.

In case of a magnetic isotropic PMSM, we have $L_d = L_q$ and therefore l = 1. Eq. (28) then simplifies to

$$\Delta I_{\rm SPWM}^2 = \left(\frac{T_{\rm P}}{2}\right)^2 \cdot \left(\frac{U_{\rm dc}}{L_{\rm d}}\right)^2 \cdot M^2 \cdot \left(\frac{1}{128} \cdot M^2 - \frac{\sqrt{3}}{36\pi} \cdot M + \frac{1}{96}\right) \tag{30}$$

This corresponds to the red line in fig. 4b. There is no dependency of the current harmonics on the angle φ_U for machines with $L_d = L_q$. Fig. 4b also shows the values of the current harmonics for other modulation methods, which will be discussed in the next section. In case of $L_d = L_q$ they do not depend on the angle φ_U either.

Utilization of different modulation methods

The described method of estimating current harmonics can easily be applied to different modulation methods. As the duration of the two active switching states t_2 and t_3 needs to be the same for each modulation method, only the turn-on times of the freewheeling states need to be adjusted.

For magnetic isotropic machines, the optimal modulation method in order to reduce the harmonic current rms value is realized by adding a third harmonic (3rdPWM) with an amplitude of $\frac{M_3}{M_1} = \frac{1}{4}$ [2]. The resulting waveforms of the reference functions can be seen in fig. 5a.

The turn-on times of the freewheeling-states for this modulation method result to

$$\frac{t_{7,3rdPWM}}{T_{p}/2} = \frac{1}{2} \left[1 - M \cdot \cos\left(\varphi_{\rm U} + \gamma\right) + \frac{M}{4} \cdot \cos\left(3 \cdot \varphi_{\rm U} + \gamma\right) \right]$$
(31)

$$\frac{t_{8,3rdPWM}}{T_p/2} = 1 - \frac{t_{7,3rdPWM}}{T_p/2} - \frac{t_1}{T_p/2} - \frac{t_2}{T_p/2}$$
(32)

Applying these turn-on times to eq. (20) and solving eq. (27) yields the solution for 3rdPWM:

$$\Delta I_{3rdPWM}^{2} = \left(\frac{T_{P}}{2}\right)^{2} \cdot \left(\frac{U_{dc}}{l L_{d}}\right)^{2} \cdot M^{2} \cdot \left[\frac{7}{1024} \frac{\left(1 + (l^{2} - 1) \left(\cos\left(\varphi_{U}\right)\right)^{2}\right) \left(-\frac{3\sqrt{3}}{4} + \pi\right)}{\pi} \cdot M^{2} - \frac{\sqrt{3} \left(12 \left(\cos\left(\varphi_{U}\right)\right)^{2} \left(l^{2} - 1\right) + 11 - l^{2}\right)}{360 \pi} \cdot M + \frac{1}{96} \left(1 + (l^{2} - 1) \left(\cos\left(\varphi_{U}\right)\right)^{2}\right)\right]$$
(33)



Figure 6: Current harmonics $I_{\rm B} = \Delta I / \left(\frac{T_{\rm P}}{2} \frac{U_{\rm dc}}{L_{\rm q}}\right)$ for SYPWM (a) and DPWM2 (b).



Figure 7: Current harmonics $I_{\rm B} = \Delta I / \left(\frac{T_{\rm P}}{2} \frac{U_{\rm dc}}{l L_{\rm q}}\right)$ for use of different modulation methods. (a): $\varphi_{\rm U} = \pi$. (b): $\varphi_{\rm U} = \frac{\pi}{2}$.

In case of distributing the remaining time $\frac{T_p}{2} - t_1 - t_2$ symmetrical at the beginning and at the end of the half pulse period (SYPWM, fig. 5b), which represents a suboptimal approximation for 3rdPWM [2], the duration of the two freewheeling-states resolves to

$$\frac{t_{7,\text{SYPWM}}}{T_{\text{p}}/2} = \frac{t_{8,\text{SYPWM}}}{T_{\text{p}}/2} = \frac{1}{2} \cdot \left(1 - \frac{t_1}{T_{\text{p}}/2} - \frac{t_2}{T_{\text{p}}/2}\right)$$
(34)

Applying the turn-on times in eq. (34) to Δi_s leads to the result for SYPWM:

$$\Delta I_{\text{SYPWM}}^{2} = \left(\frac{T_{\text{P}}}{2}\right)^{2} \cdot \left(\frac{U_{\text{dc}}}{l L_{\text{d}}}\right)^{2} \cdot M^{2} \cdot \left[\frac{3}{256} \frac{\left(1 + \left(l^{2} - 1\right)\left(\cos\left(\varphi_{\text{U}}\right)\right)^{2}\right)\left(-\frac{3\sqrt{3}}{4} + \pi\right)}{\pi} \cdot M^{2} - \frac{\sqrt{3}\left(12\left(\cos\left(\varphi_{\text{U}}\right)\right)^{2}\left(l^{2} - 1\right) + 11 - l^{2}\right)}{360 \pi} \cdot M + \frac{1}{96}\left(1 + \left(l^{2} - 1\right)\left(\cos\left(\varphi_{\text{U}}\right)\right)^{2}\right)\right]$$
(35)

Fig. 6a shows $I_{\rm B} = \Delta I_{\rm SYPWM} / \left(\frac{T_{\rm P}}{2} \cdot \frac{U_{\rm dc}}{l L_{\rm d}}\right)$ with $L_{\rm d} = 0.3$ mH and $L_{\rm q} = 1.5$ mH. The results for $\Delta I_{\rm 3rdPWM}$ do not differ significantly from those of $\Delta I_{\rm SYPWM}$.

For discontinuous modulation methods (DPWM), different distributions of the freewheeling states are commonly used [5]. The reference functions for two of them are shown in fig. 5c and 5d.

Because of a distinction of cases in the turn-on times for DPWM1 within sector I and the resulting length in the expression for the current harmonics, DPWM1 will not be further investigated within this paper. In the case of DPWM2 the relative turn-on times are

$$\frac{t_{7,\text{DPWM2}}}{T_{\text{p}}/2} = 0 \qquad \qquad \frac{t_{8,\text{DPWM2}}}{T_{\text{p}}/2} = 1 - \frac{t_1}{T_{\text{p}}/2} - \frac{t_2}{T_{\text{p}}/2} \tag{36}$$

These lead to the result for the current harmonics for DPWM2:

$$\Delta I_{\rm DPWM}^{2} = \left(\frac{T_{\rm P}}{2}\right)^{2} \left(\frac{U_{\rm dc}}{l L_{\rm d}}\right)^{2} M^{2} \cdot \left[\left(\frac{3\left(1 + (l^{2} - 1)\left(\cos\left(\varphi_{\rm U}\right)\right)^{2}\right)\left(3\sqrt{3} + 8\pi\right)}{1024 \pi} - \frac{27(l^{2} - 1)\cos\left(\varphi_{\rm U}\right)\sin\left(\varphi_{\rm U}\right)}{1024 \pi}\right) M^{2} - \left(\frac{\sqrt{3}\left(183\left(\cos\left(\varphi_{\rm U}\right)\right)^{2}\left(l^{2} - 1\right) + 179 - 4l^{2}\right)}{1440 \pi} + \frac{45\left(l^{2} - 1\right)\cos\left(\varphi_{\rm U}\right)\sin\left(\varphi_{\rm U}\right)}{1440 \pi}\right) M + \frac{1}{24}\left(1 + (l^{2} - 1)\left(\cos\left(\varphi_{\rm U}\right)\right)^{2}\right) \right]$$
(37)

The corresponding surface can be seen in fig. 6b.

Fig. 7 shows the results of the calculated current harmonics for SPWM, SYPWM and DPWM2 for stator voltage angles $\varphi_U = \pi$ and $\varphi_U = \frac{\pi}{2}$. Since the pulse period T_p for each of the three modulation methods has been considered as constant, for discontinuous modulation there are less switching operations within one pulse period. At an angle $\varphi_U = \pi$ this causes higher current harmonics for DPWM, especially at lower modulation depths. With higher modulation depths, the difference becomes smaller.

As can be seen in fig. 7b, the impact of less switching operations within one pulse period decreases as φ_U changes towards the direction of the d-axis. For $\varphi_U = \frac{\pi}{2}$ the resulting current harmonics for all three investigated modulation methods are almost equal.

Correlation to the operating point of the IPMSM

The inner voltage components of the machine are dependent on the operation point of the machine. They correspond to the speed of the machine and the current reference values $i_d^* = i_{d,0}$ and $i_q^* = i_{q,0}$ of the torque control. For IPMSM, i_d^* and i_q^* can be determined with look-up-tables of the trajectory of maximum torque per ampere (MTPA) for lower speeds and along the voltage constraint for the field-weakening area [6]. By use of eq. (18) the modulation depth *M* and the voltage phase angle φ_U can be determined for every operation point.

Simulation results

Simulations have been made in Matlab/Simulink for a high-saliency IPMSM with constant inductances L_d and L_q . Simulation parameters are shown in table I.

DC link U_{dc}	300 V	Flux linkage Ψ_{PM}	0.065 Vs
Pulse period T_p	100 µs	Inductance <i>L</i> _d	0.35 mH
pole pairs p	3	Inductance L _q	1.5 mH

Table I: Simulation para	ameters
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The desired torque value of T = 80 Nm has been generated with the current reference values shown in fig. 8a. They correspond to the MTPA-curve for lower speeds, whereas at higher speed the current reference values have been chosen for operation at maximum modulation depth. The corresponding modulation depth M and voltage angle φ_U can be seen in fig. 8b. During MTPA-operation φ_U stays constant, while M increases proportional to the speed of the machine. As the maximum modulation depth is reached, Mstays constant, while φ_U increases towards values $> \pi$.



Figure 8: Current components (a) and modulation depth and voltage angle (b) for a torque of T = 80 Nm.



Figure 9: Simulation and analytic results of $\Delta I/I_1$ for two different modulation methods.

Fig. 9 shows the results of the simulation for SYPWM and DPWM2. The effect of $L_d \neq L_q$ on the current harmonics can clearly be seen: In the field weakening area, with constant modulation depth M, the current harmonics do not stay constant. They vary with increasing speed because of the shift of φ_U . The rms values of the current harmonics have also been calculated with the derived formulas. As can be seen, the analytic values comply very well to the harmonics in simulation. Therefore, the derived formulas for the current harmonics apply to machines assuming constant inductance values L_d and L_q .

Experimental results

Measurements of current harmonics have been made with an IPMSM with parameters similar to those in tab. I. The current distortions have been measured with a *Norma D6000*. During measurements, the angle of the voltage space vector has been constant, while the modulation depth has been varied by increasing the speed of the machine. With current components $i_d = -100 \text{ A}$, $i_q = 50 \text{ A}$ for $\varphi_U \approx \pi$ and $i_d = -100 \text{ A}$, $i_q = 10 \text{ A}$ for $\varphi_U \approx \frac{\pi}{2}$ the desired angles of the voltage space vector were realized. The results for the total harmonic distortion, which is being defined as

$$THD = \frac{\sqrt{I_{\rm rms}^2 - I_1^2}}{I_1}$$
(38)

can be seen in fig. 10, together with a 4th order polynomial interpolation through the setpoints. The dashed lines in fig. 10 show the results of eq. (35) and (37).

There is a significant offset between the analytic and the experimental results in the THD-value for both modulation methods. This is due to the following cases:

- The IPMSM cannot be driven in an open-loop control. Therefore closed-loop current control is being used, which itself causes additional harmonics in the phase currents.
- Saturation effects of the inductances might have a significant impact on the current distortion. Although the evaluation of eq. (35) and (37) has been done with inductance values which have been derived from parameter measurements of the machine, their values still might differ from those that are essential for the current harmonics.



Figure 10: Measured harmonic distortion for use of different modulation methods. (a): $\phi_U \approx \pi$. (b): $\phi_U \approx \frac{\pi}{2}$.

• During measurements, the voltage vector angle φ_U might slightly differ from the desired one, which also influences the current distortion. This could be the reason for the measured THD values decreasing at higher modulation depths at $\varphi_U \approx \frac{\pi}{2}$, although the calculated estimated values are monotonously increasing with the modulation depth.

A closer investigation of these issues will be carried out in future work.

However, the shape of the measured THD values with constant φ_U and varying modulation depth and especially the relative comparison between the two illustrated modulation methods clearly corresponds to the shape of the calculated values for the rms value of the harmonic currents. At $\varphi_U \approx \pi$, the difference between the distortion with use of SYPWM and DPWM is significantly large, especially at modulation depths M < 0.8. As the modulation depth becomes larger, the difference decreases. For $\varphi_U \approx \frac{\pi}{2}$, the measured current distortions for both SYPWM and DPWM are almost equal, although the number of switching operations within one pulse period at DPWM is reduced by $\frac{1}{3}$ compared to SYPWM.

Conclusion and future work

Analysation of the current distortion in interior permanent magnet synchronous machines for different modulation methods has been presented in this paper. Analytic equations for estimation of harmonics have been derived and verified by numerical simulation. The results might be used to optimize the drive system by choosing the optimal modulation method and using the ideal switching frequency with respect to the operation point in order to minimize the total losses.

Measurements on a machine test bench have shown that the qualitative results for the current distortions can be used for a comparison of different modulation methods regarding the resulting current harmonics. The absolute values do not comply to the experimental results. The reasons for these discrepancies will be subject to further investigations.

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