Rare earth and transition metal based entropy stabilised perovskite type oxides

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ABSTRACT

Multicomponent oxides with perovskite type of structure containing up to 10 different cations in equiatomic amounts have been synthesised for the first time. Out of eleven systems synthesised, only six systems crystallised as single phase perovskite type compounds with random and homogenous cation distribution on the respective sites. The formation of phase pure 10-cationic system, $(Gd_{0.2}La_{0.2}Nd_{0.2}Sm_{0.2}Y_{0.2})(Co_{0.2}Cr_{0.2}Fe_{0.2}Mn_{0.2}Ni_{0.2}Ol_{3.2})$, in contrast to the multiphase mixtures observed in five of the lower entropy systems (containing 6 cations) indicates a possible role of entropy in the stabilisation of a single phase crystal structure. The entropy driven structural stabilisation effect is further supported by the reversible phase transformation, from single phase to multiple phase upon cyclic heat treatment, observed in the $(Gd_{0.2}La_{0.2}Nd_{0.2}Sm_{0.2}Y_{0.2})MnO_3$ system. This type of entropic signature has been observed in rocksalt based high entropy oxide systems. However, it has not been reported before for perovskite based compounds, as shown in this study.

1. Introduction

The discovery of new materials with advanced properties is one of the greatest challenges faced by the scientific community. Research work on high entropy materials was recently extended to oxide systems by Rost et al. [1], where it was observed that high configurational entropy can be effectively used to stabilise single phase rocksalt type systems containing five different cations in equiatomic amount. As per Murty et al. [2], a system can be classified as a high entropy material, if the configurational entropy (ΔS_{config}) of that system is higher than 1.5R $(0.0124 \text{ kJ mol}^{-1} \text{ K}^{-1})$. In case of an oxide systems with cations sharing only one Wyckoff site (like in rocksalt based systems), the ΔS_{config} can only be greater than 1.5R if 5 (or more) cations are present in near equiatomic ratios. Furthermore, ΔS_{config} attains a maximum value when all the different elements are present in equiatomic amount. Although, the chemical bonding in the oxides are very different from the metallic bonds present in high entropy alloys (HEAs) [2,3], the presence of high configurational entropy (1.61R) found in the rocksalt based multicomponent oxide system ((Co_{0.2}Cu_{0.2}Mg_{0.2}Ni_{0.2}Zn_{0.2})O) [1]

cannot be ignored. Additionally, the fact that this system shows a reversible transformation from a multiple phases (at lower temperature) to single phase (at higher temperature) upon cyclic heat treatment confirms the claim that the single phase structure in complex oxides can indeed be stabilised by enhancing the configurational entropy of the system [1,4]. This research on high entropy oxides (HEOs) [1] has opened up a new way to design complex multicomponent oxide cera mics. The functional properties of these new oxides are yet to be fully explored, but the initial studies by Bérardan et al. [5,6] show inter esting and improved properties compared to binary or doped oxide systems, providing further motivation for extensive research in this field.

The field of high entropy oxides was further broadened by the dis covery of multicomponent equiatomic rare earth oxides (ME REOs) [7,8], where it has been shown that up to seven rare earth cations could be successfully incorporated into a phase pure fluorite type crystal lattice. Recently, Jiang et al. [9] observed that entropy stabilisation in oxides can be extended to complex systems like perovskites (ABO₃), where cations share two or more Wyckoff sites. The tendency to form

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secondary phase(s) on addition of more cations (in higher amounts than doping) could be expected to be higher for complex oxides like per ovskites (ABX₃) than for oxides with simpler structure like rocksalt (AX) since the formation of other structures (e.g., ilmenite type) and/or phases of the type $A_xB_yX_z$ (e. g. Ruddlesdon Popper type phases) are often alternatives in systems of higher chemical complexity. In the study by Jiang et al. [9] only the B site was substituted by five different elements, whereas the A site was fixed to one or at most two elements. Hence, in this study more complex perovskites with completely different compositions are studied, where both cation sites are populated with five different elements in equiatomic amounts, i.e., five on the A site, or the B site, and finally five on both the A site and the B site, forming a 10 equiatomic cationic system.

2. Experimental section

2.1. Cations selection

Considering ABO3 type system, three sets of multicomponent equiatomic perovskite oxide ((ME PO)) systems were selected: (i) sys tems where the A site was substituted with 5 rare earth cations in equiatomic amounts ($5A_{0.2} = Gd_{0.2}La_{0.2}Nd_{0.2}Sm_{0.2}Y_{0.2}$), whereas the B site was kept occupied with a single transition metal cation (Co, Cr, Fe, Mn, or Ni) forming five systems: (5A_{0.2})CoO₃, (5A_{0.2})CrO₃, (5A_{0.2})FeO₃, (5A_{0,2})MnO₃ and (5A_{0,2})NiO₃; (ii) systems where the B site was sub stituted with 5 transition metal elements in equiatomic amounts $(5B_{0.2} = Co_{0.2}Cr_{0.2}Fe_{0.2}Mn_{0.2}Ni_{0.2})$ and keeping the A site occupied with a single rare earth cation (Gd, La, Nd, Sm, or Y) forming five systems: (Gd(5B_{0.2})O₃, La(5B_{0.2})O₃, Sm(5B_{0.2})O₃, Nd(5B_{0.2})O₃ and Y $(5B_{0,2})O_3$; and (iii) a system where both the A and B sites were re placed by 5 rare earth and 5 transition metal elements, respectively, in equiatomic amounts, forming 10 cationic (Gd_{0.2}La_{0.2}Nd_{0.2}Sm_{0.2}Y_{0.2}) (Co_{0.2}Cr_{0.2}Fe_{0.2}Mn_{0.2}Ni_{0.2})O₃ system or in an abbreviated form (5A_{0.2}) $(5B_{0,2})O_3$. The cations were chosen following criteria already reported for other HEOs [1,7], i.e., (i) the ionic radii of the cations (on a specific site, A or B) at a specific oxidation state and coordination number should be similar, and (ii) at least one of the perovskite systems should have different prototype structure (e.g. YMnO₃) [10]. An overview of the ionic radii and the space groups of various structures that could be formed using the selected elements are given in the Supplementary information, Tables S1 and S2, respectively.

2.2. Synthesis

ME POs powders were synthesised using nebulised spray pyrolysis (NSP) [11]. The illustration of the experimental setup can be found elsewhere [12]. The nitrate salts of Gd, La, Nd, Sm, Y Co, Cr, Fe, Mn and Ni (Gd(NO₃)₃·6H₂O, Sigma Aldrich, 99.9%; La(NO₃)₃·6H₂O, Sigma Al drich, 99.9%; Nd(NO₃)₃·6H₂O, Sigma Aldrich, 99.9%; Sm(NO₃)₃·6H₂O, ABCR, 99.9%; Y(NO₃)₃·6H₂O, ABCR, 99.9%; Co(NO₃)₂·6H₂O, Sigma Aldrich, 99.9%; Cr(NO₃)₃·9H₂O, Sigma Aldrich, 99.9%; Fe(NO₃)₃·9H₂O, Sigma Aldrich, 99.9%; Mn(NO₃)₂·4H₂O, ChemPure, 98.5%; Ni (NO₃)₂·6H₂O, Sigma Aldrich, 99.9%) were used as precursors. A water based solution containing appropriate equiatomic cation combinations $(0.1 \text{ mol}.\text{L}^{-1})$ was continuously delivered (~ $0.12 \text{ L}\cdot\text{h}^{-1}$) into the piezo driven nebuliser. The generated mist containing fine droplets of the precursor solution was transported by flowing oxygen (5 standard L·min⁻¹) into the hot wall reactor where the particles were formed at a temperature of 1050 °C and a pressure of 900 mbar. The as synthesised particles were collected using a filter based collector and calcined at 1200 °C for 2 h in air atmosphere with a heating/cooling rate of $10 \degree \text{C·min}^{-1}$.

2.3. Characterisation

The microstructure and elemental composition of the synthesised

powders were studied using a Philips XL30 FEG scanning electron mi croscope (SEM) equipped with an energy dispersive spectroscopy (EDS) system (Amatek, USA). Particle size distribution was obtained by measuring around 300 particles and the average particle size was ob tained from a log normal fit [13].

For the validation of the elemental composition of the powders, selected samples have been analysed in triplicate by inductively cou pled plasma mass spectrometry (ICP MS) using an Agilent 7500cx ICP MS (Agilent, Santa Clara, USA); detailed analytical conditions are given in the Supplementary information (Table S3). For the measurements 10 mg of the synthesised powders were first dissolved in 2 ml HNO₃ (supra quality, 68%, Carl Roth, Karlsruhe, Germany) and 8 ml MilliQ deionised water (18.2 mΩ·cm). In a second step, 10 µL of the sample solution was further diluted with 9.68 mL MilliQ water. Additionally, 10 µL of internal standard (10 mg L⁻¹ scandium and holmium, respec tively) was added to correct the time dependent sensitivity variation of the ICP MS system followed by an addition of 300 µL of nitric acid (68%, supra quality) to prohibit metal ion sorption onto the tube sur face. For quantification a calibration with 0, 5, 25, 100 and 500 µg L⁻¹ of all analysed elements were prepared under similar conditions.

Room temperature X ray diffraction (XRD) patterns of the powders were recorded using a Bruker D8 diffractometer with Bragg Brentano geometry using Cu Ka radiation with a Ni filter and a VANTEC detector having a fixed divergence slit (0.3°) . A step size of 0.015° and a col lection time of 4s per step at 30 kV and 40 mA over the diffraction angle (2 θ) range between 10 and 90° were used. For in situ high tem perature XRD (HT XRD) measurements a temperature stage (Anton Paar HTK 1200, Bruker) was used to collect data in a temperature range between 30 °C and 1000 °C with a heating and cooling rate of 30 °C·min⁻¹. The step size and collection time were set to 0.015° and 2 s, respectively. Rietveld analysis (TOPAS 5, Bruker) [14] of the XRD patterns was performed to determine the structure and phase compo sition of powders. Perovskites phases could be refined using a standard orthorhombic perovskite type structure of GdFeO₃ (*Pbmn*, ICSD 23823). The structural files were modified according to the number and the type of elements present in the refined system. The instrumental intensity distribution for the XRD data was determined using a reference scan of LaB₆ (NIST 660a). Thermal displacement parameters were constrained to be the same for all atoms.

The TEM specimen was prepared by directly dispersing the powders onto a gold grid coated with holey carbon film. It was examined with a FEI Titan 80 300 electron microscope (FEI, Eindhoven, the Netherlands) equipped with CEOS image spherical aberration corrector, Fischione model 3000 high angle annular dark field (HAADF) scanning transmission electron microscopy (STEM) detector, EDAX SUTW energy dispersive X ray spectroscopy (EDX) detector and Gatan Tridiem image filter. The microscope was operated at an accelerating voltage of 300 kV in TEM mode for HRTEM and in nanoprobe mode for STEM imaging, EDX and electron energy loss spectroscopy (EELS) spectra acquisition.

3. Results and discussion

3.1. Microstructure and chemical composition

The as synthesised as well as the calcined multicomponent equia tomic perovskite based oxide (ME PO) powders are mostly black in colour except for $(5A_{0.2})CrO_3$, $(5A_{0.2})FeO_3$, and $(5A_{0.2})NiO_3$, which are light green, brown and olive green, respectively (Fig. 1a). The mor phology of the calcined powders studied by SEM (Fig. 1b) show sphe rical hollow agglomerates of particles having a mean particle size of 180 nm (inset in Fig. 1b).

The actual compositions of the systems could be the determined for 5 of the 11 compounds by ICP MS (Supplementary information Table S4). The other 6 compounds could not be dissolved completely in a concentrated acid. The obtained overall compositions confirm that all the elements in those 5 ME POs are present in the desired



Fig. 1. (a) Photographs of several ME-PO systems showing variation in powder colour: 1. (5A_{0.2})CrO₃, 2. (5A_{0.2})FeO₃, 3. (5A_{0.2})NiO₃, 4. (5A_{0.2})(5B_{0.2}O₃. (b) SEM micrograph of (5A_{0.2}) (5AB_{0.2})O₃ powder representing a typical morphology of ME-POs systems produced by NSP method. Inset represents log-normal particle size distribution with an average particle size of 180 nm.

stoichiometry. The EDS analysis (see Supplementary information Table S5) performed on the systems are in good agreement with the ICP MS results, which confirm the correct overall compositions for the systems which could not be analysed using ICP MS.

3.2. Phase composition and structural analysis

The XRD patterns of the as synthesised powders, Supplementary information Fig. S1(a d), show that only one of the ME PO systems, (5A_{0.2})CoO₃, completely crystallises as a single phase solid solution directly under the synthesis conditions used, whereas in some of the systems ((5A_{0.2})CrO₃, (5A_{0.2})FeO₃, La(5B_{0.2})O₃, Nd(5B_{0.2})O₃) several phases are formed. In all the other systems ((5A_{0.2})MnO₃, (5A_{0.2})NiO₃, Gd(5B_{0.2})O₃, Y(5B_{0.2})O₃ and (5A_{0.2})(5B_{0.2})O₃) a large amount of amorphous phase is observed. One of the reasons for the presence of more than one phase in some of the systems and incomplete crystal lisation in the others can be the low residence time (around 150 ms) of the particles in the NSP reactor. The residence time plays an important role in diffusion and crystallisation processes. Hence, all as synthesised powders were subjected to a further calcination step at 1200 °C for 2 h in air and only the calcined powder are used for further characterisation and discussion. The corresponding XRD patterns are shown in Fig. 2(a d) and the summary of their detailed structural analysis are given in Table 1.

The Rietveld analysis shows that all of the systems except $(5A_{0.2})$ NiO₃ crystallise in orthorhombic perovskite type structure, but with different structural parameters and degree of secondary phases. The Rietveld refinements of all the systems (along with the Bragg peak position of the secondary phases for the multiphase mixtures) are provided in the Supplementary information (Fig. S2 S11). The average crystallite size of the perovskite phases for the different ME POs ranges from 75 125 nm where $(5A_{0.2})COO_3$ has the smallest and Sm $(5B_{0.2})O_3$ has the largest crystallite size. Minor amount of non perovskite phases (up to a maximum of 5 wt.%) are observed in three of the systems: $(5A_{0.2})CrO_3$, Sm $(5B_{0.2})O_3$ and Y $(5B_{0.2})O_3$.

In the $(5A_{0.2})MnO_3$ system no additional non perovskite type phases are found. However, strong tails at lower diffraction angles are ob served for all the reflections in the XRD pattern for this system. The pattern of this phase cannot be refined using a single perovskite type phase, which is seen by the misfit (Fig. 3a) of the (020) reflection at 31.1° (20). This reflection shows a misfit due to its strong asymmetry, which indicates a distribution of perovskites phases with similar lattice parameters, mutually resulting from a non isotropic composition.



Fig. 2. XRD patterns of ME-PO powders: (a) $(5A_{0.2})BO_3$, B = Co, Cr, Fe, Mn, or Ni, (b) $A(5B_{0.2})O_3$, A = Gd, La, Nd, Sm, or Y, (c) $(5A_{0.2})(5B_{0.2})O_3$, and (d) simulated orthorhombic perovskite (*Pbnm*) pattern (ICSD 23823).

Table 1

Structural data for all (ME-PO) systems obtained from Rietveld refinement: space group (SG) along with prototype structures of secondary phases, phase fraction (*f* in wt.%), lattice parameters (*a*, *b*, *c* in Å, and β in °) and goodness of fit (GoF).

ME-POs	SG	<i>f</i> [wt.%]	a [Å]	b [Å]	c [Å]	β [°]	GoF
$5A_{0,2} = Gd_{0,2}La_{0,2}Nd_{0,2}$	2Sm _{0.2} Y _{0.2}						
(5A _{0.2})CoO ₃	Pbnm	100	5.2964(2)	5.3655(2)	7.5128(3)	-	1.16
(5A _{0.2})CrO ₃	Pbnm	97.7(6)	5.3799(7)	5.5035(7)	7.6614(1)	-	1.74
	$Ia\overline{3}$ (Gd ₂ O ₃)	2.3(6)	10.729(0)				
(5A0.2)FeO3	Pbnm	100	5.4094(3)	5.5936(2)	7.7259(2)	-	1.63
(5A _{0.2})MnO ₃	Pbnm	42.3(1)	5.3658(4)	5.7396(4)	7.5352(5)	-	1.43
	Pbnm	40.9(1)	5.3670(4)	5.7637(5)	7.5201(7)	-	
	Pbnm	16.8(5)	5.3633(2)	5.6916(2)	7.5652(3)	-	
(5A _{0.2})NiO ₃	$C2/m (Sm_2O_3)^a$	60.4(3)	14.259(4)	3.6241(9)	8.8706(2)	100.6(0)	1.27
	$Fm\overline{3}m$ (NiO)	33.1(3)	4.1791(1)	-	-	-	
	I4/mmm (La ₂ NiO ₄)	6.5(1)	3.8465(3)	-	12.5330(2)	-	
$5B_{0,2} = Co_{0,2}Cr_{0,2}Fe_{0,2}I$	Mn _{0.2} Ni _{0.2}						
Gd(5B _{0,2})O ₃	Pbnm	100	5.2958(6)	5.5303(6)	7.5708(9)	-	1.25
La(5B _{0.2})O ₃	Pbnm	100	5.4656(6)	5.5101(5)	7.7424(9)	-	1.19
Nd(5B _{0.2})O ₃	Pbnm	100	5.4087(6)	5.4760(6)	7.6718(8)	-	1.47
Sm(5B _{0.2})O ₃	Pbnm	98.3(3)	5.3523(6)	5.5024(6)	7.6178(9)	-	1.42
	$C2/m (Sm_2O_3)^a$	1.7(3)	14.236(1)	3.6400(3)	8.7290(1)	100.4(1)	
Y(5B _{0.2})O ₃	Pbnm	94.7(1)	5.2255(6)	5.5361(7)	7.4977(1)	-	1.44
	$Ia\overline{3}$ (Y ₂ O ₃)	3.2(7)	10.597(1)	-	-	-	
	$Fm\overline{3}m$ (NiO)	2.1(1)	4.2021(1)	-	-	-	
$(5A_{0.2})(5B_{0.2})O_3$	Pbnm	100	5.3614(7)	5.5012(7)	7.6309(1)	-	1.45

^a In (5A_{0.2})NiO₃ and Sm(5B_{0.2})O₃, the prototype structure used for the refinement of the C2/m phase is Sm₂O₃, however the composition of the phase is most likely to be (5A_{0.2})₂O₃.

Therefore, the flexibility of the model needs to be increased to describe the observed pattern, and it was found that a series of three orthor hombic perovskite phases with very close lattice parameters (same shape and positional parameters for all perovskite phases, see Table 1) can be used to approximately model the pattern (Fig. 3b). For this 3 phase model an increase of the *a* axis with decrease of the *b* and *c* axes for the different phase fractions was found. The fit can be further im proved by choosing a smoother distribution of *n* perovskite phases, by adding a constraint in such a way that the lattice parameters $(a,b,c)_i$ of phase fraction *i* are calculated according to :

$$(a, b, c)_i = (a, b, c)_{i=1} + [(a, b, c)_{i=n} - (a, b, c)_{i=1}]^* (i/n)$$
(1)

in addition to constraining the phases to the same shape and positional parameters (i.e., independent of the number of *n*, only 6 lattice para meters, $(a,b,c)_{i-1}$ and $(a,b,c)_{i-n}$, are being refined, which mimics the width of the distribution of the phases, with additional *n* scaling parameters). The phase distribution model for a fit with n = 8 is given in the Supplementary information, Table S6. Again, this agrees well with a segregation of rare earth cations based on their size. Fig. 3(a c) compares Rietveld fits of the $(5A_{0.2})MnO_3$ XRD pattern by either con sidering a single perovskite (a) or three perovskite (b) or an n = 8 perovskite phases (c). In the $(5A_{0.2})NiO_3$ system, a major monoclinic phase (Gd₂O₃ type) instead of an orthorhombic perovskite phase, with higher amount of secondary rocksalt (NiO type) and tetragonal (La₂NiO₄ type) phases, are observed. The major Gd₂O₃ type phase is most likely a solid solution containing several of the rare earth cations in it.

From the XRD patterns (Fig. 2a c), it can be clearly observed $(2\theta = 20^{\circ} 30^{\circ})$ that the intensities of the superstructure reflections varies in different ME POs. This intensity variation is strongly dependent on the degree of deviation from the cubic perovskite structure, which indeed strongly depends on the lattice parameters (Table 1) of the perovskite phases. The pictorial representation of distortion from the ideal cubic lattice, based on the structural information obtained from Rietveld refinement, is shown in the Supplementary information (Figs. S12(a) and (b)). The Goldschmidt's tolerance factor (*t*) is a well known theoretical parameter used for explaining this type of deviation from cubic symmetry and is calculated from the ionic radii [15] based on the following equation [16]:

$$t = \frac{r_A + r_0}{\sqrt{2(r_B + r_0)}}$$
(2)

where r_A and r_B are the ionic radii of the cation at A site and B site (see Table S1 in the Supplementary information), respectively, and r_o is the radius of the oxygen ion. In case of multiple cations at a specific site an average of the ionic radii is considered.

In an ideal case, for the cubic structure *t* should be equal to 1. For t < 1, i.e., for systems with smaller A site cation or bigger B site ca tions, orthorhombic or rhombohedral structures are preferred, whereas tetragonal or hexagonal structures are preferably formed for t > 1. In case of ME POs orthorhombic perovskite phases are observed as the value of *t* is lower than 1 (Table 2) and the deviation from the cubic symmetry becomes more prominent with adding (or substituting smaller cation) on the A site or bigger on the B site. Another parameter which is often used to explain the deviation from the pseudo cubic symmetry is the degree of metric distortion (ε) which can be defined as [17]:

$$\varepsilon = \left(\frac{1}{3}\right) \cdot \left[\left\{ \frac{a_{norm} - a_{ps.cubic}}{a_{ps.cubic}} \right\}^2 + \left\{ \frac{b_{norm} - a_{ps.cubic}}{a_{ps.cubic}} \right\}^2 + \left\{ \frac{c_{norm} - a_{ps.cubic}}{a_{ps.cubic}} \right\}^2 \right]^{0.5}$$
(3)

where $a_{ps.cubic}$ is the pseudo cubic lattice parameter, and a_{norm} , b_{norm} , c_{norm} are the normalised values of pseudo cubic lattice parameters taking into account the orientation of the lower symmetry cell. These values are calculated from the following equations:

$$a_{ps.cubic} = \sqrt[3]{V} \tag{4}$$

$$a_{norm} = \frac{a}{\sqrt{2\sqrt[3]{V}}}, \ b_{norm} = \frac{b}{\sqrt{2\sqrt[3]{V}}}, \ c_{norm} = \frac{c}{2\sqrt[3]{V}}$$
(5)

where V is the volume per ABO₃ unit, i.e.,

$$V = \frac{a^* b^* c}{4} \tag{6}$$

with *a*, *b*, *c* being the original lattice parameter of the orthorhombic (*Pbnm*) perovskites obtained from the Rietveld analysis (Table 1). In case of ε , the extent of deviation from the ideality follows a similar trend as the theoretical parameter *t* (Table 2). With decrease in the cationic radii at the A site, or increase in the cationic radii at the B site



Fig. 3. XRD pattern together with Rietveld fit of (5A_{0.2})MnO₃, where either a single (a), 3 (b) or 8 (c) perovskite phases are used to refine the structure. The zoomed in section (b, d and f) shows that a good fit can be achieved with a strongly constrained distribution model of 8 perovskite phases.

Table 2

Overview of the Goldschmidt's factor (t), metric distortion (ε), and effective coordination number (ECoN) for A- and B-site elements in ME-PO systems obtained from ionic radii, lattice parameters of phase with perovskite structure and structural model, respectively.

ME-POs	t	ε [%]	ECoN(A)	ECoN(B)					
$5A_{0.2} = Gd_{0.2}La_{0.2}Nd_{0.2}Sm_{0.2}Y_{0.2}$									
(5A _{0.2})CoO ₃	0.932	0.344	7.41	6.00					
(5A _{0.2})CrO ₃	0.927	0.547	7.24	6.00					
(5A _{0.2})FeO ₃	0.918	0.580	7.12	5.99					
(5A _{0.2})MnO ₃	0.914	2.299	6.12	5.48					
(5A _{0.2})NiO ₃ ^a	0.963	-	-	-					
$5B_{0.2} = Co_{0.2}Cr_{0.2}Fe_{0.2}Mn_{0.2}Ni_{0.2}$									
Gd(5B _{0.2})O ₃	0.900	1.609	7.29	5.98					
La(5B _{0.2})O ₃	0.963	0.198	10.09	5.99					
Sm(5B _{0.2})O ₃	0.921	0.685	7.59	5.99					
Nd(5B _{0.2})O ₃	0.932	0.438	7.14	5.97					
Y(5B _{0.2})O ₃	0.880	1.425	6.59	5.98					
$(5A_{0.2})(5B_{0.2})O_3$	0.919	0.633	7.43	5.99					

^a Perovskite phase was not formed in this system.

(see Table S2 for the ionic radii of the cations given in the Supple mentary information), ε deviates strongly from ideality ($\varepsilon > 0$) and hence, a lowering in degree of symmetry or a higher magnitude of

distortion is observed. The tilting and distortion of the BO₆ octahedra in the perovskite oxides is known to be the genesis of several of their important functional properties like magnetic, ferroelectric, etc. [18]. The distortion in perovskite structure from the ideal cubic symmetry is strongly governed by the shape or relative orientation of these BO₆ polyhedra [19]. There can be three different cases: (i) cations dis placement inside the octahedra mostly observed for smaller B cations (often leading to polar structures), (ii) tilting of the octahedra mostly observed for smaller A cations which are too small to be accommodated within a cuboctahedral cage site (often leading to orthorhombic or rhombohedral structures), and (iii) distortion of the BO₆ polyhedra often seen for compounds with Jahn Teller active cations (e.g., Cu^{2+} , Mn³⁺) [18,20]. Fig. S12(c e) in the Supplementary information illus trates the increasing tilt of the BO₆ octahedra in ME POs based on the size of the A cations, and it shows that the tilting of the octahedra follows a similar trend as t.

A strong increase in the tilting is observed accompanied by the in crease of ε from 0.198 in La(5B_{0.2})O₃ to 0.633 in (5A_{0.2})(5B_{0.2})O₃, followed by 1.609 in Gd(5B_{0.2})O₃. The lowering of symmetry due to the structural distortion is often found to be accompanied by a lowering of the coordination number of the A site cations [17,21,22]. Hence, the effective coordination number (ECoN) for both the A and B site

elements are calculated following the formulation proposed by Hoppe [23] and are listed in Table 2. The details of the ECoN calculation to gether with the refined distances of the A O, as well as B O are provided in the Supplementary information (see Table S7 and related discussion). From the calculated ECoN values, it can be seen that for all ME POs the B site cations are almost ideally octahedrally coordinated independent of the type of cations present in the A or B sites. However, a large deviation from the ideal 12 fold coordination of the A site cations is observed for all systems. The ECoN(A) also follows the similar trend as t, ε and the degree of tilting of the BO₆ octahedra, i.e., a gradual de crease (from the ideal 12 fold coordination) on addition or substitution of the smaller cation at A site or larger cation on the B site. The origin of secondary phases observed in some of the systems can be strongly related to structural distortions in the lattices. In case of the (5A_{0.2}) MnO₃ system, the strong Jahn Teller distortion can be one of the plausible reasons for the high metric distortion (2.299) and the low ECoN(A) (6.12) and ECoN(B) (5.48) value. From the ε vs. t, ECoN(A) vs. t and ECoN(B) vs. t plots (Fig. 4(a), (b) and (c), respectively) it can be observed that the system which deviates most from the otherwise fairly linear behaviour is the (5A_{0.2})MnO₃ system. This indicates a strong enthalpic contribution due the possible Jahn Teller distortion, which in turn can be prohibitive for the formation of a single phase perovskite structure in (5A_{0,2})MnO₃ at lower temperature (see Sec. 3.4 for the high temperature studies of this system). In $Y(5B_{0,2})O_3$, one possible reason for the formation of multiple phases can be the small size of Y³⁺ and its inclination towards a lower coordination state. For the other systems the following reasons can be possible for the formation of the multi phase mixtures: (i) The chemical nature of the B site with most stable oxidation states other than 3 + is found to prohibit the formation of a single phase in the case of a multicomponent equiatomic A site con figuration, for example in case of (5A_{0.2})NiO₃ the difficulty to stabilise Ni³⁺ (under the synthesis conditions used) could plausibly explain the formation of high amount bunsenite (NiO) phase along with of major monoclinic type non perovskite phase. (ii) Even for compositions which allow for the formation of perovskite type phase only, segregation of larger from smaller A site cations along with the Jahn Teller effect can be prohibitive for the formation of a single phase, as seen in $(5A_{0,2})$ MnO₃. Therefore, these above mentioned factors can increase the en thalpy of formation of a single phase compound and hence, favour the formation of non single phase (either with non perovskite type im purities or multi perovskite type phase) compositions. Thus, con sidering all these factors and the fact that several of the heptanary ME POs are not phase pure makes the chemically complex single phase 10 cationic system interesting.

3.3. Chemical homogeneity

A detailed TEM study was carried out on the 10 cationic system. On all regions of the sample investigated by TEM, perfect crystallinity has been observed, evident from the sharp lattice fringes shown in the HRTEM micrographs (Fig. 5a and c). The fast Fourier transformations (FFTs) of the HRTEM micrographs are shown in Fig. 5b and d, re spectively. The FTTs can be indexed to a single orthorhombic per ovskite (*Pbnm*) structure, which is in agreement with XRD findings. To substantiate the local homogeneity of the element distribution in the 10 cationic system, STEM EELS/EDX spectra imaging was conducted (Supplementary information Fig. S13). Due to the heavily overlapped signals in the EDX spectra caused by the presence of multiple cations (i.e., ten), the elemental distribution maps were mostly extracted from the EELS using the M4.5 edges for Gd, Nd and Sm, only M5 edge for La (since La M₄ overlaps with Ni L₃) and the L_{2.3} edges for Co, Cr, Fe and Mn. The Y and Ni maps were constructed from EDX Y L and Ni K sig nals, respectively. These maps represent the distribution of the elements qualitatively and show that all the elements are uniformly distributed without any significant fluctuations down to nanometer scale. This confirms the local homogeneity in the studied 10 cationic system. Fig. 6



Fig. 4. Metric distortion (ϵ), EcoN(A) and ECoN(B) are plotted with respect to the tolerance factor (t) in diagram (a), (b) and (c) respectively. A strong deviation of the (5A_{0.2}) MnO₃ system from the otherwise fairly linear behaviour can be observed.

shows a HAADF STEM image for $(5A_{0.2})(5B_{0.2})O_3$ system along the [010] zone axis. The atomic layers of A site containing heavy rare earth cations can be easily distinguished from the light transition metal ca tions on the B site by the Z contrast (here Z is average atomic number for site A and B, respectively). The atomic columns at the A site show similar contrast indicating a random occupancy of the rare earth cations (Gd, La, Nd, Sm and Y). Similar conclusion can be drawn for the B site cations confirming the random distribution of Co, Cr, Fe, Mn and Ni, as well.

3.4. Possible reasons for the stabilisation of single phase

In Section 3.2 it has been mentioned that factors like ion size mediated segregation and affinity of an element towards particular oxidation state (and co ordination number) can increase the enthalpy of mixing (ΔH_{mix}) for formation of the single phase. Hence, these factors



Fig. 5. HR-TEM micrographs and corresponding FFTs, respectively for $(5A_{0.2})(5B_{0.2})O_3$ system along the [010] (a, b) and $[02\overline{1}]$ (c, d) zone axes.



Fig. 6. HAADF-STEM image taken along [010] zone axis, where the atomic layers containing heavier atoms (Gd, La, Nd, Sm and Y) and lighter atoms (Co, Cr, Fe, Mn and Ni), respectively can be clearly distinguished.

stabilise the multiphase mixtures. Thus considering these facts, the high configurational entropy in these ME PO systems can be one of the plausible reasons for the formation of the single phase, especially in case of the 10 cationic system. In case of multi principal elemental al loys, better known as the high entropy alloys (HEAs), the high config urational entropy is often regarded as the reason for the formation of single phase structures in chemically complex systems. Similar entropic behaviour in stabilisation of a single phase was also observed in rock salt based high entropy oxides (HEOs) by Rost et al. [1]. In case of the ME POs the configurational entropy (ΔS_{config}) can be calculated using the following formula:

$$\Delta S_{config} = -R \left[\left(\sum_{a=1}^{n} x_a ln x_a \right)_{A-site} + \left(\sum_{b=1}^{n} x_b ln x_b \right)_{B-site} + 3 \left(\sum_{c=1}^{n} x_c ln x_c \right)_{O-site} \right]$$
(7)

where x_a , x_b and x_c are the mole fraction of the ions present in the A site, B site and O site (or the anion site which is oxygen in this case), respectively. For the ΔS_{config} calculations, all the ME POs are considered to be stoichiometric, i.e., the possible oxygen defects that can be present are not taken into account. For 10 cationic system, (5A_{0.2})(5B_{0.2})O₃, S_{config} is 3.22R (0.0268 kJ mol⁻¹ K⁻¹) and is twice that of the other systems, 1.61R (0.0134 kJ mol⁻¹ K⁻¹) in which only one of the sites, either A or B sites, are substituted by 5 equiatomic cations, (5A_{0.2})BO₃ or A(5B_{0.2})O₃. The competing factors of enthalpy and entropy are generally described within the approximation of an ideal solid solution. Independent of the systems (whether it is an alloy or an oxide), whenever a non identical species is added to an existing sub lattice, two thermodynamically opposing factor come into existence; one is the change of ΔH_{mix} which hinders the stability of the solid solution and the other is the entropic $(\Delta S_{mix} / \Delta S_{config})$ gain from the mixing which sup ports the formation of the solid solution. The existence of the high ΔS_{mix} (or ΔS_{config}) (calculated from Eq. 7) due to the presence of several ele ments in the ME PO systems cannot be ignored, independent of whether it actually plays a role in the structural stabilisation or not. In case of rocksalt based HEOs [1], a reversible transformation from multiphase



Fig. 7. Full range (a) and magnified portion (b) of the in-situ HTXRD pattern of (5A_{0.2})MnO₃ at different temperatures. A reversible transformation from multiple phases to a single phase perovskite occurs around 900 °C.

to single phase structure was observed upon cyclic heat treatment. At higher temperature the $T\Delta S$ term becomes dominant over the ΔH term, hence decreasing the free energy (as $\Delta G_{mix} = \Delta H_{mix}$ $T\Delta S_{mix}$).

Thus, in order to investigate the role of entropy in the structural stabilisation, in situ high temperature XRD (HTXRD) measurements were performed on the ME POs. For the $(5A_{0.2})MnO_3$ system, where a series of perovskite phases (see Fig. 3) are observed, in situ HTXRD was performed at various temperatures (Fig. 7a). A gradual transformation from a multiphase to a single phase accompanied by distinct changes in the XRD pattern is observed with increasing temperature (Fig. 7a and b). A table with the details of the lattice parameters and phase com position of the powder at each temperature has been provided in the Supplementary information (see Table S8).

The complete transformation to single perovskite type phase in $(5A_{0.2})$ MnO₃ is found to occur around 900 °C. The Rietveld refinement for the same has been shown in the Fig. 8. The single phase, thus obtained is preserved even at higher temperature 1000 °C, however, a complete re version, i.e., a series of perovskite phases is observed while cooling back to

room temperature (Fig. 7). This behaviour is very similar to the one ob served in case of the rocksalt based HEOs [1], where entropic factors were found to be the reason for the stabilisation of single phase compounds at elevated temperature. Hence, a similar entropic stabilisation is likely to be true in case of the $(5A_{0,2})MnO_3$ system as well, since at higher temperature the $T\Delta S_{mix}$ term compensates the high ΔH_{mix} which is indicated to be especially high for the (5A_{0.2})MnO₃ system owing to the possible Jahn Teller effect (see previous Section 3.2). HTXRD patterns of the 10 cationic system show that the single phase is preserved all along (till 1000 °C) and no phase transformation is observed at lower temperature (Supplementary information Fig. S12). This is explained as follows: (i) the entropy of mixing can be large enough compared to the enthalpy of mixing, i.e., $\Delta S_{mix} > > \Delta H_{mix}$, making ΔG_{mix} independent of ΔH_{mix} , even at lower temperatures, or (ii) even though the enthalpic penalty is larger than $T\Delta S_{mix}$ at lower temperature, the kinetics plays a limiting role (owing to the sluggish diffusion due to presence of several different cations). How ever, considering the facts that high configurational entropy is inherently present in these systems, one of the heptanary systems show reversible



Fig. 8. In-situ HTXRD pattern together with Rietveld fit of $(5A_{0.2})MnO_3$ at 900 °C. The pattern can be well fitted using only a single perovskite phase compared to the fit at room temperature where a minimum of three perovksite phases are needed to fit the pattern and to describe the peak shape due to demixing of A-site cations adequately.

phase transformation as observed in rocksalt based HEOs [1,4], and the 10 cationic system is phase pure (in spite of its high chemical complexity) points towards the dominant role of entropy in structural stabilisation in the ME POs. HTXRD measurements were done on all the other systems as well (see Supplementary information Figs. S15(a d) and S16(a d)). In case of the multiple phase systems ($(5A_{0.2})CrO_3$, $(5A_{0.2})NiO_3$, Sm($5B_{0.2})O_3$ and Y($5B_{0.2})O_3$) single phase could not be achieved even at higher tempera tures. This further proves the fact that in these systems the enthalpic pe nalties needed to stabilise the single phase is very high. Hence, supports the fact that the 10 cationic system is most likely entropy stabilised.

It is also interesting that in spite of the presence of so many elements in ME POs (especially for the 10 cationic system) the orthorhombic distortion of the perovskite phase can be maintained, even at high temperatures (see Supplementary information Fig. S12 and Table S9), instead of formation of a highly symmetric cubic aristotype structure, which is known to be formed for many multinary systems. The tendency to maintain a high symmetry is also observed in the rocksalt type en tropy stabilised oxide [1] and the fluorite type multicomponent rare earth oxides [7], where distortion and ordering variants are well known for both systems [24]. Hence, the tilting of the octahedra can be as sumed to have an enthalpic contribution (similar to the Jahn Teller contribution in Cu enriched HEOs) [25] in stabilisation of the lower symmetry perovskite phase. This distortion is known to result from the lowering of the A site coordination, and apparently this enthalpic contribution is so strong that the adoption of the cubic aristotype is energetically costly. This is further confirmed by a HT study of the ME POs system with the lowest deviation from the cubic symmetry (i.e. lowest value of strain ε , see Table 2), La(5B_{0.2})O₃. This phase can also not be transformed to a cubic structure, even at higher temperatures. However, in case of the $La(5B_{0.2})O_3$, a reversible structural transition from an orthorhombic (*Pbnm*) phase to a trigonal (R3ch) phase is ob served (Fig. 9). The transformation starts around 125 °C but a dis tinctive change in the in situ XRD patterns of $La(5B_{0,2})O_3$ can only be observed above 140 °C. Complete transformation to a trigonal structure takes place around 300 °C, where the volume of both the phase frac tions, i.e., the trigonal cell and the orthorhombic cell become equal (see Supplementary information Figs. S17 and S18 and Table S10). The trigonal structure formed is preserved even at higher temperatures 1050 °C. This explains the fact that the octahedral tilting even in this systems (where the t = 0.963 and $\varepsilon = 0.198$) plays a crucial role in prohibiting the system to attain a cubic structure.

4. Conclusions

Based on the results obtained from the combine XRD, ICP MS, TEM and STEM/EDX EELS studies, three major conclusions can be drawn:

- i Out of the eleven different ME PO systems synthesised, six of the systems can be stabilised as a single phase perovskite type compound at ambient temperature. Even the 10 cationic ((Gd_{0.2}La_{0.2}Nd_{0.2}Sm_{0.2}Y_{0.2})(Co_{0.2}Cr_{0.2}Fe_{0.2}Mn_{0.2}Ni_{0.2})O₃) system is phase pure along with random and homogenous cation distribution on the respective sites, in spite of its high chemical complexity for which strong stabilisation of distinct binary combinations could be expected.
- ii The fact that five of the heptanary systems are multiphase mixtures whereas the 10 cationic system is phase pure indicates the possible role of unfavourable enthalpic interactions of some of the lower entropy (heptanary) systems which cannot be overcome by the en tropic contribution. This indicates a possible role of entropy in the structural stabilisation.
- iii The signature of entropy driven structural stabilisation effect in these perovskites systems is strongly supported by the fact that the $(Gd_{0.2}La_{0.2}Nd_{0.2}Sm_{0.2}Y_{0.2})MnO_3$ system shows a reversible trans formation from a multiphase to single phase upon cyclic heat treatment. At higher temperature the single phase solid solution is favoured over the multiphase mixtures as the entropic $(T\Delta S_{mix})$ term dominates the free energy landscape. This is a typical entropy sta bilisation behaviour, which has been reported only for the rocksalt based multicomponent equiatomic oxides. Hence, this study extends the concept of entropy driven structural stabilisation to multi component perovskite type oxides.
- iv Large metric distortion from the cubic symmetry along with strong tilting of the BO_6 polyhedra and lowering of the effective co ordination number of A site cation is observed for the systems with smaller cations on the A site and/or larger B site cations. The en thalpic factors like cation size adopted metric distortion, affinity of certain elements towards a particular oxidation state and Jahn Teller distortion which favour an ordered scenario can possibly make the formation of a cubic aristotype structure in ME POs, en ergetically expensive.

Future studies will focus on the deeper understanding of the local structure of these systems along with the functional properties for their complete exploitation.



Fig. 9. Full range (a) and magnified portion (b) of the in-situ HTXRD pattern of $La(5B_{0.2})$ O₃ at different temperatures. Reversible structural transition from an orthorhombic to trigonal phase starts around 125 °C and complete transition to a phase pure trigonal phase observed above 300 °C.

Conflicts of interest

There are no conflicts of interest to declare.

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Appendix A. Supplementary data

Supplementary material related to this article can be found, in the online version, at doi:https://doi.org/10.1016/j.jeurceramsoc.2017.12. 058.

Supplementary information (Electronic annex): Ionic radii of the constituent cations, space groups of the possible ternary oxide combinations, elemental compositions from EDS, operating parameters and analytical conditions of ICP MS measurements, Rietveld refinements along with structural parameters of the different phases, the cation oxygen bond distances, description regarding the ECoN calculation, high temperature XRD patterns and related discussion are provided in the supplementary information.

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