Contributions to the Joint DFG-RSF Project - Generation of Ultra-Short Microwave Pulses -

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Introduction

In a joined DFG-RSF project of the Institute of Applied Physics (IAP-RAS) and the Institute for Pulsed Power and Microwave Technology (IHM-KIT), the generation of a periodic sequence of powerful short RF pulses is studied [1]. Such powerful pulses in the millimeter and sub-millimeter wavelength range can be useful for a large number of fundamental problems and practical applications, including diagnostics of plasmas, photochemistry, biophysics, new locating systems, and spectroscopy of various media.

A periodic sequence of powerful pulses can be generated in a feedback loop consisting of an amplifier and a saturable absorber [1]. The saturable absorber acts as a nonlinear filter which transmits high intensity signals while signals with low intensity are absorbed. In such a feedback loop, the periodic signal is generated by the mechanism of passive mode-locking. In the millimeter wavelength range, gyro-devices as the gyro-traveling-wave-tube (gyro-TWT) are well suited for the realization of both devices, the amplifier and the nonlinear absorber.

Components of a feedback loop

Amplifier and absorber can be realized as gyro-TWTs with helical interaction region [2]. These devices provide a series of advantages compared to classical gyro-TWTs based on cylindrical interaction circuits with dielectric losses. The helical gyro-TWTs can operate at the 2\textsuperscript{nd} cyclotron harmonic and therefore, the required magnetic field is reduced by the factor of 2 compared to classical gyro-TWTs operating at the fundamental cyclotron harmonic. Further advantages are that helical gyro-TWTs offer a broad frequency bandwidth up to 20\%, which is important for the generation of ultra-short pulses in the sub-nanosecond range, and that they are not sensitive to electron velocity spread. Also, it is possible to operate helical them with a single window for the input and output of the RF signal [2]. In the single input/output window operation it is possible to couple even high-power signals into the device, which is an essential condition in the feedback loop. In addition, the single input/output window simplifies the coupling of the amplifier and the absorber in a feedback loop.

Amplifier and absorber will operate in different operating regimes. The amplifier will run in a regime optimal for the maximal amplification of ultra-short pulses [3], while the absorber will run in the so-called Kompfner dip regime [4].

Simulations

Full-wave PIC simulations including the interaction between electrons and electromagnetic waves are...
Fig. 2. Dependence of the transmission coefficient on the power of the incident signal for a helical gyro-TWT operating in the Kompfner dip regime.

performed with the advanced simulation program PICLas [6], developed by the Institute of Aerodynamics and Gas Dynamics (IAG) at the University of Stuttgart. The required high-order hexahedral meshes are generated with the mesh-generator HOPR [7], also developed by the IAG.

In the following, simulations of an absorber as suggested in [8] for a feedback loop at a center frequency of 30 GHz are performed. The suggested feedback loop is designed to provide pulses with a peak-power up to 400 kW and a width of only 0.25 ns by a separation of 5-10 ns [8]. The absorber has a mean waveguide radius of 0.4 cm with a corrugation of amplitude 0.1 cm and a period of 1.0 cm. The electron beam has a beam voltage of 54 kV, a current of 2.4 A and a pitch factor of 0.53.

Fig. 3. RF-power of a Gaussian pulse propagating along the interaction region of a helical gyro-TWT operating in the Kompfner dip regime, for a low power input signal a), and a high power signal b).

In figure 2, the simulated transmission coefficient for short Gaussian pulses with different peak powers is shown. While weak signals are almost completely absorbed, high power signals are only weakly attenuated. Therefore, the simulations proof the qualification of gyro-TWTs operating in the Kompfner dip regime as saturable absorbers. From figure 3, the saturable absorber effect can be understood in detail. In figure 3 a), the propagation of a 10 kW Gaussian RF pulse along the absorber is shown. The absorption of almost the complete pulse can be clearly seen. In figure 3 b), the propagation of a high-power input pulse (200 kW) along the same absorber is shown. It can be seen that first a maximum fraction of 30 kW of the input power is absorbed by the electron beam. The transfer of the RF power to the kinetic energy of the electrons causes a shift from the Kompfner dip regime to the amplification regime. Thus, the RF pulse gains a fraction of its previous emitted energy back. A comparison of the input and output signal shows that the high-power pulse lost only a small fraction of its power. Therefore, the device exhibits properties of a saturable absorber.

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References


