

Synchronous detection of longitudinal and transverse bunch signals at a storage ring

Benjamin Kehrer,^{*} Miriam Brosi, Johannes L. Steinmann, Edmund Blomley, Erik Bründermann, Michele Caselle, Stefan Funkner, Nicole Hiller,[†] Michael J. Nasse, Gudrun Niehues, Lorenzo Rota,[‡] Manuel Schedler,[§] Patrik Schönfeldt,^{||} Marcel Schuh, Paul Schütze,[¶] Marc Weber, and Anke-Susanne Müller
Karlsruhe Institute of Technology, Kaiserstraße 12, 76131 Karlsruhe, Germany



(Received 15 August 2018; published 23 October 2018)

To understand and control dynamics in the longitudinal phase space, time-resolved measurements of different bunch parameters are required. For a reconstruction of this phase space, the detector systems have to be synchronized. This reconstruction can be used for example for studies of the microbunching instability which occurs if the interaction of the bunch with its own radiation leads to the formation of substructures on the longitudinal bunch profile. These substructures can grow rapidly—leading to a sawtooth-like behavior of the bunch. At KARA, we use a fast-gated intensified camera for energy spread studies, Schottky diodes for coherent synchrotron radiation studies as well as electro-optical spectral decoding for longitudinal bunch profile measurements. For a synchronization, a synchronization scheme is used which compensates for hardware delays. In this paper, the different experimental setups and their synchronization are discussed and first results of synchronous measurements presented.

DOI: [10.1103/PhysRevAccelBeams.21.102803](https://doi.org/10.1103/PhysRevAccelBeams.21.102803)

I. INTRODUCTION

The investigation of the dynamics of short bunches in a storage ring sets stringent requirements for the diagnostics: Response time and repetition rate of the different detector systems must be sufficient to study a single bunch in a multi-bunch environment on a single-turn basis. The required time scales are given by the rf frequency and the revolution time. In addition, the different detector systems must be synchronized to study simultaneous changes for example in the bunch shape and in the emitted synchrotron radiation.

II. MICROBUNCHING INSTABILITY

Under certain conditions, periodic bursts of coherent synchrotron radiation (CSR) were recorded at different facilities (e.g., NSLS VUV [1], SURF II [2], ALS [3], BESSY-II [4], Diamond [5], SOLEIL [6]). Theoretical

studies of this phenomena gave hints to a self-interaction of the bunch with its own wakefield [7], which can lead to the occurrence of substructures in the longitudinal phase space. The rising amplitude of these substructures can be explained by the CSR impedance [8]. The rapid increase of the amplitude is coupled to an overall increase in projected bunch size. At one point, diffusion and radiation damping outweigh the growth of the substructures and they start to dissolve and damp down, leading to a reduction of the bunch size. The reduced bunch size leads again to a stronger wakefield and subsequently the next CSR instability cycle starts over. Thus, the bunch size is modulated with a repetitive sawtooth pattern in time [9]. Due to the changing bunch profile, the CSR is emitted in bursts and the behavior is referred to as *bursting*.

As the longitudinal phase space is spanned by time and energy, it can be studied by time-resolved measurements of the longitudinal bunch profile and energy spread. The CSR intensity depends on the longitudinal bunch profile, so it can also be used to probe the dynamics. A monitoring of the longitudinal phase space leads to a deeper understanding of the beam dynamics, especially above the bursting threshold and could potentially lead to methods for control. Such a reconstruction requires synchronous measurements of the different bunch properties.

At KARA, we use different diagnostics tools for these measurements which are integrated in a hardware synchronization system to enable simultaneous data acquisition.

^{*}benjamin.kehrer@kit.edu

[†]Present address: PSI, Villigen, Switzerland.

[‡]Present address: SLAC, Menlo Park, CA, USA.

[§]Present address: VARIAN PT, Troisdorf, Germany.

^{||}Present address: DLR-VE, Oldenburg, Germany.

[¶]Present address: DESY, Hamburg, Germany.

III. KARA

KARA (Karlsruhe Research Accelerator), the 2.5 GeV KIT storage ring, can be operated in different modes including a short bunch mode. In this mode, the momentum-compactness factor α_c is lowered by changing the magnet optics [10] to study the microbunching instability [11,12].

At KARA, the revolution time is 368 ns and the bunch spacing of 2 ns is defined by the radio frequency (rf) of 500 MHz. One single measurement should take at most 2 ns, so that only the signal of one bunch is detected. As part of the bunch detection setup, several systems have been installed and commissioned: A fast-gated intensified camera (FGC) to measure the horizontal bunch position and size, fast THz detectors to detect the coherent synchrotron radiation (CSR) as well as an electro-optical bunch profile monitor to determine the longitudinal bunch profile [13–15].

IV. HORIZONTAL BUNCH SIZE MEASUREMENTS

Measurements of the horizontal bunch size σ_x allow studies of the energy spread σ_δ as they are related by [16]

$$\sigma_x = \sqrt{\beta_x \cdot \epsilon_x + (D_x \cdot \sigma_\delta)^2}$$

in a dispersive section. Therefore, a time-resolved measurement of the horizontal bunch size reveals changes in energy spread. We use a fast-gated intensified camera (FGC). It is located at the Visible Light Diagnostics (VLD) port at KARA which uses incoherent bending radiation from a 5° port of a dipole magnet [17]. The FGC setup consists of a camera and a fast-rotating mirror which sweeps the incoming light over the entrance aperture of the FGC during the image acquisition. In combination with a fast switching of the image intensifier, this allows single turn images of a bunch, even in a multibunch fill [18,19]. The resolution of the setup is confined by a rectangular crotch absorber with a horizontal width of 20 mm which is located 1600 mm downstream of the radiation source point [20]. Taking diffraction effects into account, the resolution is estimated to be $77 \mu\text{m}$ using a wavelength of 400 nm [[21], Eq. (30)].

Figure 1 shows a raw image recorded with the FGC.

Here, the machine was operated without an active feedback system, therefore the bunch is undergoing a synchrotron oscillation ($f_s \approx 11.24$ kHz). The individual spots on this image are the convolution of the charge distribution with the so-called *filament beam spread function* (FBSF). The FBSF can be seen as an extension of the point spread function for a moving point-like source, in this case a moving single electron [22]. We use the software OpTalis [23] to simulate the optical system and the imaging process. Instead of deconvolving the FBSF from the data and fitting a Gaussian curve, we fit a

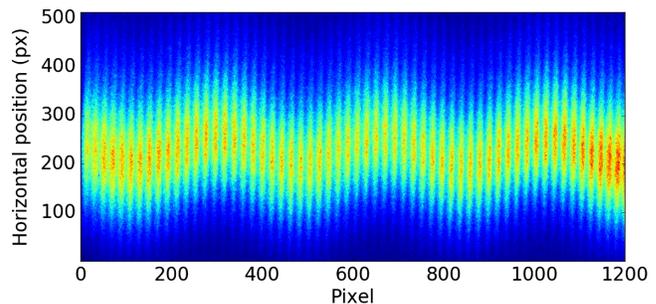


FIG. 1. Raw image taken by the fast-gated intensified camera (FGC). It shows 58 projections of the same bunch for every 14th turn. The optical path of the setup rotates the spots by 90° so their horizontal axis is displayed on the vertical axis of the image.

convolution of the FBSF with a Gaussian curve. Compared to the deconvolution this has the advantage that no estimation of the signal-to-noise ratio is required, in addition, this method is fast and robust. For one spot of an FGC raw image, this analysis is illustrated in Fig. 2.

It can be seen that the fit reproduces the distorted shape of the profile very well, enabling a determination of the horizontal bunch size with single turn resolution.

Experimental studies showed that the horizontal bunch size remains constant below the bursting threshold, while it starts to increase for currents above the threshold [24]. For the equilibrium case below the bursting threshold, the measured bunch size is in good agreement with simulated values using the accelerator toolbox for Matlab (AT) [25] and the lattice model of KARA, see Fig. 3. For the AT model, LOCO fits [26] are used to determine the quadrupole strengths from measurements of the orbit response matrix and the dispersion.

A quantitative determination of the energy spread is not possible due to the unknown contribution of the horizontal emittance to the horizontal bunch size, therefore we are limited to qualitative studies with the horizontal bunch size

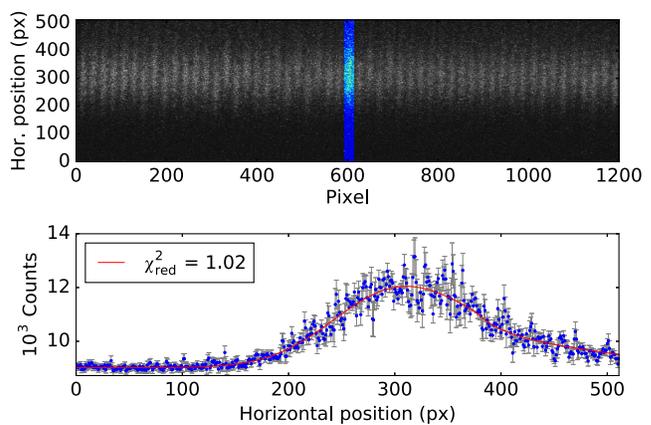


FIG. 2. FGC raw image in grey scale with the analysis range for one spot highlighted in color is shown in the top panel. The corresponding profile and the fit is plotted in the bottom panel.

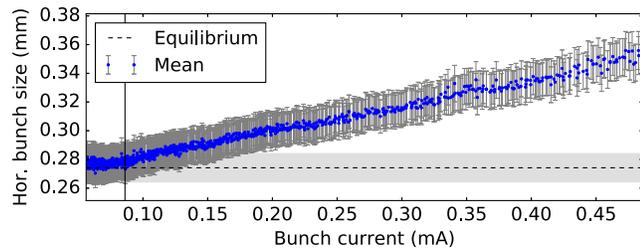


FIG. 3. Mean horizontal bunch size over bunch current. The dashed line and its grey error bar are depicting the simulated equilibrium bunch size at the VLD source point, while the vertical black line illustrates the calculated bursting threshold.

taken as a measure for the energy spread. Besides the studies of the microbunching instability, the system can be used to study different kinds of instabilities, e.g., during the injection or on the energy ramp.

V. COHERENT SYNCHROTRON RADIATION

The investigation of the CSR dynamics includes THz detectors that are fast enough to resolve the emitted synchrotron radiation for each bunch in a multi-bunch environment. We use commercially available room temperature zero-bias Schottky diode detectors with different sensitivity ranges [11,12]. To read out these detectors we use the data acquisition (DAQ) system KAPTURE [27,28]. KAPTURE is a fast FPGA-based system with four input channels that can be used to sample the signal of one detector with four sample points per bunch and turn. Alternatively, it is possible to sample up to four detectors in parallel with a bunch-by-bunch sampling rate of 500 MHz. For the measurements presented here, we configured KAPTURE to digitize the peak intensity of multiple different detectors simultaneously.

Figure 4 illustrates this multidetector operation mode. It shows the signals recorded by an avalanche photo diode (APD) sensitive to incoherent synchrotron radiation in the visible range and a Schottky diode sensitive in the THz range. The APD is insensitive to arrival time oscillations caused by synchrotron oscillation, as the duration of the APD pulse is longer than the amplitude of the synchrotron oscillation and KAPTURE samples at a fixed phase relative to the rf. Thus, the observed oscillation of the APD signal is due to an intensity fluctuation. The bunch oscillates horizontally around the focal point of the optics and the limited aperture of the beam line optics transfers this into an intensity oscillation. This signal will play a vital role for the calibration of the synchronization in Sec. VII.

The combination of KAPTURE and Schottky diodes is used for detailed and systematic studies of the CSR emission during the microbunching instability. The analysis of the fluctuation of the CSR intensity allows for example a fast and precise mapping of the bursting threshold [11]. Using four Schottky diodes with different

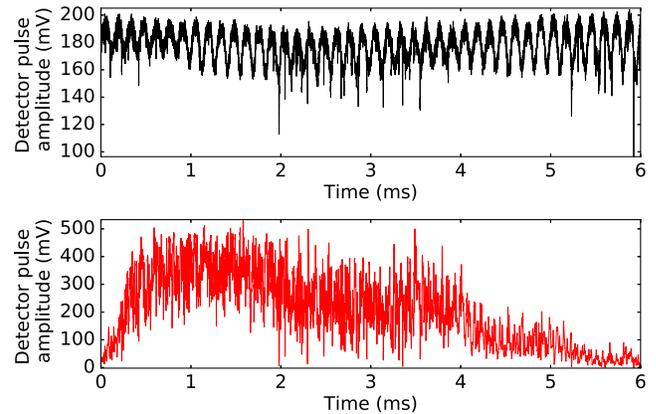


FIG. 4. Data recorded by KAPTURE using two different detectors in parallel. The top panel shows the signal from an APD in the visible range while the bottom panel shows the turn-by-turn THz signal recorded using a narrow-band Schottky detector [29]. The THz signal has the sawtooth-like pattern induced by the bursting and the APD signal is modulated by the synchrotron frequency.

frequency bands we can configure the KAPTURE system into a bunch-by-bunch spectrometer for CSR studies [12].

VI. LONGITUDINAL BUNCH PROFILE

The technique of electro-optical spectral decoding (EOSD) [30] is used to determine the longitudinal bunch profile and arrival time [31]. This is done by inserting an electro-optical crystal into the vacuum pipe to sample the electric near-field of the bunch [32]. This field turns the crystal birefringent. A laser pulse, chirped on the ps-time scale, is directed through the crystal and the induced birefringence turns the initial linear polarization into an elliptical. Thus the bunch profile is imprinted into the polarization of the laser pulse. A crossed polarizer transfers this into an intensity modulation and, due to the unique, ideally linear, correlation between time and wavelength in the chirped laser pulse, the bunch profile can be determined by recording the modulated laser pulse spectrum. Measuring in the near field has the advantage that there is no frequency cut-off of the electric field by the beam-line.

The spectrometer used for this measurement is based on the KALYPSO system (KARlsruhe Linear arraY detector for MHz-rePetition rate SpectrOscopy [33,34]). KALYPSO consists of a 256-pixel line array that can be read out with a turn-by-turn frame rate of 2.7 MHz.

An example recorded by this system is shown in Fig. 5. It shows color-coded longitudinal bunch profiles (in units of ps) as a function of time. These measurements were taken with an early development version of the KALYPSO system, which manifests in pixel errors and other artifacts. They are visible for example as dark horizontal lines between -10 and -20 ps. At the beginning, the bunch is undergoing a synchrotron oscillation as also here the

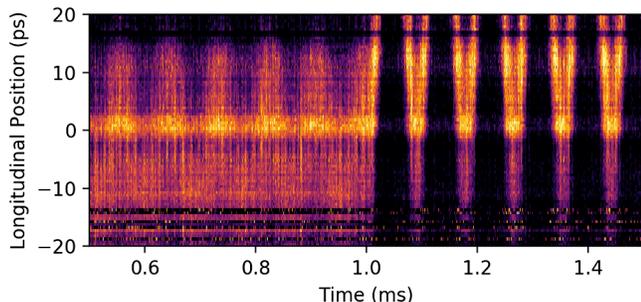


FIG. 5. Longitudinal bunch profiles recorded with the KALYPSO-based electro-optical spectral decoding setup. Each image column corresponds to one bunch profile. Synchrotron oscillations can be seen during the entire time range. The abrupt change of its amplitude at 1 ms is due to an rf phase step that leads to an oscillation around the new phase. The upper part of this oscillation is not visible due to the finite size of the line array and the length of the chirped laser pulse chosen to achieve a better temporal resolution.

feedback system was switched off. At 1 ms, this oscillation abruptly changes its amplitude due to a triggered step in the rf phase at this point in time. For these measurements, the EO system was adjusted such that the phase step has the maximum possible effect on the signal. So, the bunch shape is not reproduced optimally. Instead, the frequency of the forced synchrotron oscillation after the rf phase step can be seen clearly. This enables to take the system into account for the detector synchronization.

Generally, the EOSD system can also be set up to allow measurements of the bunch profile with a sub-ps resolution, that allows to resolve substructures on the longitudinal bunch profile [35]. By using KALYPSO this setup allows turn-by-turn measurements [[33], Fig. 3].

VII. DETECTOR SYNCHRONIZATION

A straightforward way to measure synchronously the different bunch parameters is to feed the signals from the different detectors into one common DAQ system, e.g., an oscilloscope. This is suitable if all devices are located close to each other at a storage ring with compact design (e.g., [2]). In our case, this is not directly possible for several reasons. The first point is the location of the detector systems at different positions around the storage ring. In addition, at these systems dedicated post-processing of the raw-data is required. The FGC and the EOSD setup delivers 2D images while the KAPTURE system provides—in the multidetector mode—one data point per bunch, turn, and detector. The acquisition has to be aligned temporally to enable correlation studies between derived parameters like CSR intensity, energy spread and bunch length.

To provide a common origin for the time axis with single-turn resolution, we need to synchronize the acquisition. The timing system [36]—based on one event generator (EVG)

and several event receivers (EVR)—is used to generate a synchronized measurement trigger.

Furthermore, the intrinsic hardware delays have to be taken into account. KAPTURE and KALYPSO sample continuously, independent of the coupled detectors. The acquisition trigger starts the data storage into memory and therefore an instantaneous synchronization of the recording is achieved.

For the FGC this is different: Due to the measurement controls and the inertia of the rotating mirror, a certain advance time is needed before the first sample can be recorded. To protect the sensor from overexposure, the rotating mirror has a hold position in which the light is not on the sensor. Thus the light spot has to be driven first onto the sensor to be recorded. This also ensures the measurement is taken in the regime of linear mirror motion. The synchronization schemes takes these delays into account and compensates for them by triggering the FGC the required time before other devices (see [[37], Fig. 2]).

The test of the synchronization requires a signal with a signature on the longitudinal and horizontal plane of the beam that all detectors are sensitive to. We used the low-level rf (LLRF) system to trigger a step in the rf phase. This sudden change leads to the onset of a synchrotron oscillation around the new longitudinal equilibrium, which all different measurement stations are sensitive to. For the calibration of KAPTURE, we used the visible light APD as a reference. This convenient reference is insensitive to fluctuations on the bunch profile that can occur following a sudden rf phase-step, which can affect the CSR signal recorded by Schottky detectors.

Figure 6 shows a successful synchronization event. The onset of the synchrotron oscillation is recorded by all detectors at the same time. While the synchrotron oscillation on the FGC and the APD are in phase and both in the horizontal plane, the one from KALYPSO is phase-shifted by $\pi/2$. This can be explained by the dynamics in longitudinal phase space. Here, the synchrotron oscillation corresponds to a rotation of the longitudinal electron phase space density. The projection onto the position and energy axes are separated by $\pi/2$. The FGC and the APD are sensitive to the energy because their source points are located in dispersive sections. KALYPSO, in contrast, measures the arrival time.

The APD signal recorded with KAPTURE shows voltage dips at the highest values and maximum bunch deflection. The origin of the dips are unclear, but are likely due to intensity cutoff from the finite aperture of the optical beam path or jumps in the timing trigger as KAPTURE samples at a fixed phase. Nevertheless, the main feature, a strong synchrotron oscillation, can be seen and used to determine the phase and the frequency of the electron beam motion.

As discussed before, the EOSD setup was optimized for temporal resolution. Due to the finite size of the line array and the length of the chirped laser pulse, this leads to a

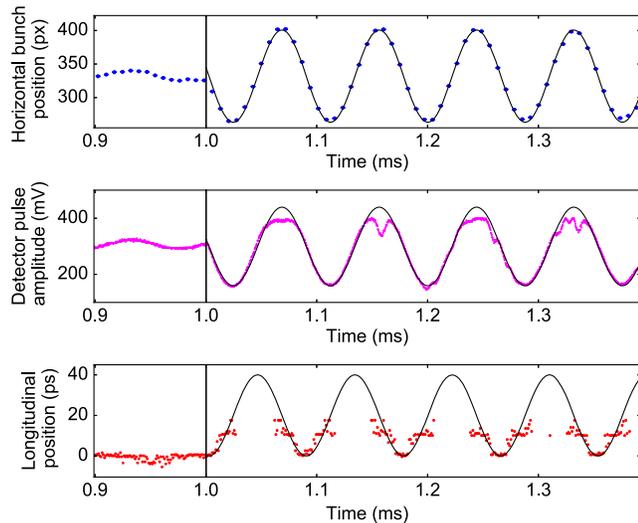


FIG. 6. Signature of the triggered rf phase-step seen by the different systems. Top: Horizontal bunch position from the FGC. Center: Signal of the incoherent synchrotron radiation recorded using an optical APD and KAPTURE. Bottom: Arrival time of the bunch recorded with KALYPSO. It is defined by the position of the maximum electro-optical modulation. The vertical line in all plots marks the rf phase step triggering the onset of a synchrotron oscillation. As a guide to the eye the plots contain a sinusoidal curve at the synchrotron frequency.

cutoff of the upper part of the signal showing the longitudinal bunch position.

VIII. SYNCHRONOUS MEASUREMENTS

Synchronization of the detector systems allows time-resolved studies of the beam dynamics. In the following, we limit our studies to energy spread and the CSR emission during the microbunching instability.

Figures 7 and 8 show two synchronous measurements of the horizontal bunch for the energy spread (top) and the corresponding CSR emission (bottom) recorded with a broadband Schottky diode detector (50 GHz–2 THz). In both cases, the two curves have the same modulation period. This behavior was predicted by simulations [9] and also observed experimentally at SURF III [2].

In Fig. 7, the CSR intensity decreases until it reaches a constant level between two bursts. After the blow-up at the onset of a burst, the energy spread decreases until it reaches the same lower limit where it is blown-up again. At the same time, the CSR burst starts to arise.

Figure 8 shows a different case using different machine settings and a lower bunch current. Here, the CSR intensity is not constant between two bursts, but increases slightly while the energy spread is still decreasing. This can be explained by the coupling of bunch length and energy spread in the longitudinal phase space. If the bunch gets shorter, it emits more CSR—regardless of substructures. The sensitivity to these bunch length fluctuations is

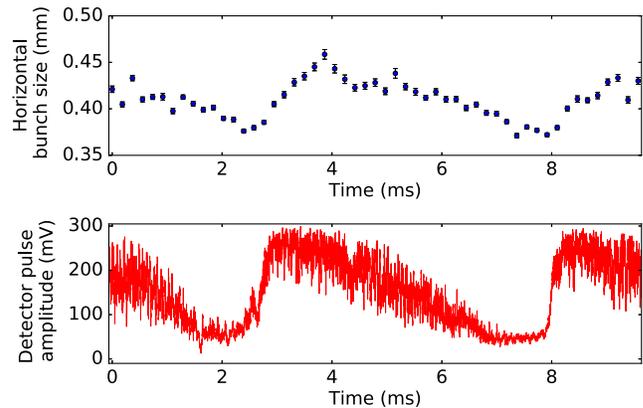


FIG. 7. Horizontal bunch size as measure for the energy spread (top) and the synchronously measured CSR signal sampled using a Schottky diode (bottom). The energy spread shows the same modulation pattern as the CSR (raw data published in [[38], Fig. 10]). Beam parameter: f_s : 8.1 kHz, V_{RF} : 1500 kV, I_{bunch} : 0.88 mA.

dependent on the frequency band of the detector and possible shielding of the CSR due to apertures in the beam line. Thus, we do not observe it for longer bunches.

For both cases, the bunch contracts until an instability threshold is reached. At this point, the charge density inside the bunch is high enough and the bunch becomes unstable. Numerical simulations of the longitudinal phase space using the Vlasov-Fokker-Planck solver Inovesa [40] showed, that at this point the amplitude of the substructures starts to grow rapidly [41].

To study these effects in more detail, the FGC can be configured to record shorter time ranges with a higher temporal resolution. For the third example discussed here, the gate separation was set to 24 turns and a time range of 500 μ s was recorded. In this case an additional narrow-band Schottky diode was used synchronously. The result is

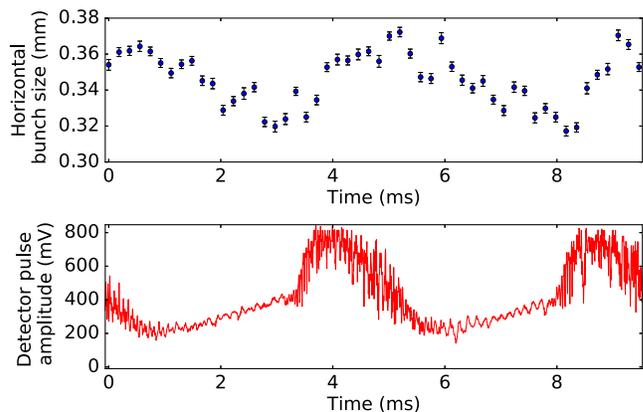


FIG. 8. Horizontal bunch size (top) and the synchronously measured CSR from a broadband Schottky diode (raw data published in [[39], Fig. 2]). Beam parameter: f_s : 6.7 kHz, V_{RF} : 1500 kV, I_{bunch} : 0.25 mA.

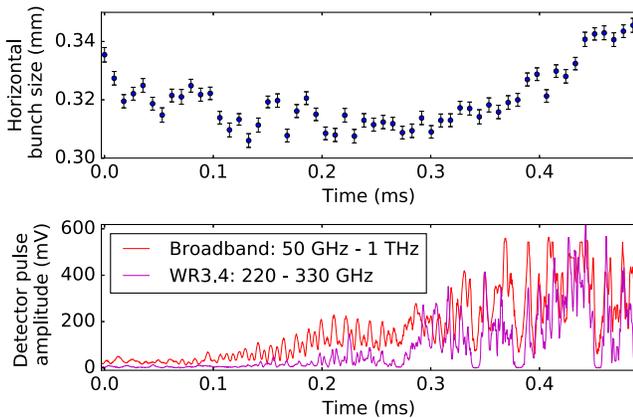


FIG. 9. Horizontal bunch size recorded with a gate separation of 24 turns (top) and the corresponding CSR intensity (bottom) sampled with two Schottky diodes with different bandwidths. The error bars for the spot size are large compared to the ones in Fig. 8 due to the lower amplification of the microchannel plate leading to less counts on the CCD (raw data published in [[39], Fig. 3]). Beam parameter: f_s : 13.4 kHz, V_{RF} : 1500 kV, I_{bunch} : 1.57 mA.

illustrated in Fig. 9. At the beginning, the energy spread given by the horizontal bunch size (top panel) is slightly decreasing due to damping effects while the CSR intensity on the broadband Schottky diode is constant, for the narrow-band diode it is almost zero. Around $t = 0.1$ ms, the first changes can be seen when the CSR intensity on the broadband Schottky diode starts to increase, while for the narrow-band this increase starts approximately 0.1 ms later. The earlier increase of the signal on the broadband Schottky diode is due to its lower limit in the frequency sensitivity [12].

While the CSR intensity starts to increase, the energy spread is still decreasing due to damping effects. At 0.27 ms, it reaches a lower limit triggering the onset of another burst.

These examples clearly illustrate the diagnostic capability of synchronized detector systems. It allows to fully exploit the potential of the different, synchronous systems. While with KAPTURE in the multidetector mode it is possible to study phase differences between the different frequency bands of the CSR [12], these phase studies can now be extended to take multiple detector systems into account.

IX. SUMMARY

The analysis of fast rising instabilities, like the microbunching instability, benefits from turn-by-turn detector systems. To resolve complimentary features in 6-D phase space, several detector systems which observe different bunch parameters need to be synchronized. At KARA, several independent detector systems with a single-turn resolution are combined in a distributed sensor network.

The energy spread is investigated using an FGC, while the CSR intensity is measured using a KAPTURE system and Schottky diodes. For studies of the longitudinal bunch profile an EOSD setup is used in the near-field combined with KALYPSO. The integration of these systems into a hardware synchronization scheme opens up new diagnostics opportunities. While the individual setups allowed an insight into the dynamics of the microbunching instability, successful synchronization is a first step toward a multi-dimensional analysis of the longitudinal phase space. We experimentally showed that the energy spread and the CSR undergo a modulation with the same period length. In addition, detailed studies revealed a certain phase difference between the CSR in different frequency ranges and energy spread.

In the future we will improve and upgrade the individual detector systems while keeping the capability for synchronized recording. A fast line array optimized for the visible frequency range will be combined with KALYPSO and replace the FGC setup. This will allow us to continuously stream horizontal bunch profiles at each revolution. It will also be used at the EOSD setup with a line array, optimized for the laser frequency and improving the signal-to-noise ratio. Further, the next version of KALYPSO will support a streaming rate of up to 10 Mfps and up to 2048 pixels [34]. A new version of KAPTURE, KAPTURE2 will provide 8 simultaneous readout paths, allowing more accurate peak reconstruction and the measurement of the arrival time [28]. Alternatively, the amplitudes of up to 8 different detectors can be recorded. With narrow-band detectors, sensitive in different frequency ranges, this allows recording single-shot spectra [12].

ACKNOWLEDGMENTS

We would like to thank A. Plech, Y.-L. Mathis and M. Süpfle for their support. This work has been supported by the Initiative and Networking Fund of the Helmholtz Association under Contract No. VH-NG-320 and by the BMBF under Contracts No. 05K10VKC and No. 05K13VKA. M. Brosi, P. Schönfeldt, and J.L. Steinmann acknowledge the support of the Helmholtz International Research School for Teratronics (HIRST) and E. Blomley the support of the Karlsruhe School of Elementary Particle and Astroparticle Physics (KSETA).

- [1] G. Carr, S. Kramer, J. Murphy, R. Lobo, and D. Tanner, Observation of coherent synchrotron radiation from the NSLS VUV ring, *Nucl. Instrum. Methods Phys. Res., Sect. A* **463**, 387 (2001).
- [2] U. Arp, G. T. Fraser, A. R. Hight Walker, T. B. Lucatorto, K. K. Lehmann, K. Harkay, N. Sereno, and K.-J. Kim, Spontaneous coherent microwave emission and the sawtooth instability in a compact storage ring, *Phys. Rev. ST Accel. Beams* **4**, 054401 (2001).

- [3] J. M. Byrd, W. P. Leemans, A. Loftsdottir, B. Marcellis, M. C. Martin, W. R. McKinney, F. Sannibale, T. Scarvie, and C. Steier, Observation of Broadband Self-Amplified Spontaneous Coherent Terahertz Synchrotron Radiation in a Storage Ring, *Phys. Rev. Lett.* **89**, 224801 (2002).
- [4] M. Abo-Bakr, J. Feikes, K. Holldack, P. Kuske, and G. Wüstefeld, in *Proceedings of the 2003 Particle Accelerator Conference, Portland, OR* (IEEE, New York, 2003), p. 3023.
- [5] W. Shields, R. Bartolini, G. Boorman, P. Karataev, A. Lyapin, J. Puntree, and G. Rehm, Microbunch Instability Observations from a THz Detector at Diamond Light Source, *J. Phys. Conf. Ser.* **357**, 012037 (2012).
- [6] C. Evain, J. Barros, A. Loulergue, M. A. Tordeux, R. Nagaoka, M. Labat, L. Cassinari, G. Creff, L. Manceron, J. B. Brubach, P. Roy, and M. E. Couprie, Spatio-temporal dynamics of relativistic electron bunches during the micro-bunching instability in storage rings, *Europhys. Lett.* **98**, 40006 (2012).
- [7] G. Stupakov and S. Heifets, Beam instability and micro-bunching due to coherent synchrotron radiation, *Phys. Rev. ST Accel. Beams* **5**, 054402 (2002).
- [8] M. Venturini and R. Warnock, Bursts of Coherent Synchrotron Radiation in Electron Storage Rings: A Dynamical Model, *Phys. Rev. Lett.* **89**, 224802 (2002).
- [9] R. Warnock, Study of bunch instabilities by the nonlinear Vlasov–Fokker–Planck equation, *Nucl. Instrum. Methods Phys. Res., Sect. A* **561**, 186 (2006).
- [10] A.-S. Müller *et al.*, Experimental Aspects of CSR in the ANKA Storage Ring, ICFA Beam Dynamics Newsletter **57**, 154 (2012), https://icfa-usa.jlab.org/archive/newsletter/icfa_bd_nl_57.pdf.
- [11] M. Brosi, J. L. Steinmann, E. Blomley, E. Bründermann, M. Caselle, N. Hiller, B. Kehrer, Y.-L. Mathis, M. J. Nasse, L. Rota, M. Schedler, P. Schönfeldt, M. Schuh, M. Schwarz, M. Weber, and A.-S. Müller, Fast mapping of terahertz bursting thresholds and characteristics at synchrotron light sources, *Phys. Rev. Accel. Beams* **19**, 110701 (2016).
- [12] J. L. Steinmann, M. Brosi, E. Bründermann, M. Caselle, B. Kehrer, L. Rota, P. Schönfeldt, M. Schuh, M. Siegel, M. Weber, and A.-S. Müller, Continuous bunch-by-bunch spectroscopic investigation of the micro-bunching instability, [arXiv:1710.09568](https://arxiv.org/abs/1710.09568).
- [13] A. Fisher, M. Petree, R. Kraus, Y.-S. Au, B. Chan, T. Meyer, and R. Webber, Turn-by-turn imaging of the transverse beam profile in PEP-II, *AIP Conf. Proc.* **868**, 303 (2006).
- [14] H.-W. Hübers, A. Semenov, K. Holldack, U. Schade, G. Wüstefeld, and G. Gol'tsman, Time domain analysis of coherent terahertz synchrotron radiation, *Appl. Phys. Lett.* **87**, 184103 (2005).
- [15] I. Wilke, A. M. MacLeod, W. A. Gillespie, G. Berden, G. M. H. Knippels, and A. F. G. van der Meer, Single-Shot Electron-Beam Bunch Length Measurements, *Phys. Rev. Lett.* **88**, 124801 (2002).
- [16] H. Wiedemann, *Particle Accelerator Physics* (Springer Science & Business Media, New York, 2007).
- [17] B. Kehrer, A. Borysenko, E. Hertle, N. Hiller, M. Holz, A.-S. Müller, P. Schönfeldt, and P. Schütze, in *Proceedings of IPAC'15* (JACoW, Geneva, Switzerland, 2015), p. 866.
- [18] P. Schütze, A. Borysenko, E. Hertle, N. Hiller, B. Kehrer, A.-S. Müller, and P. Schönfeldt, in *Proceedings of IPAC'15* (JACoW, Geneva, Switzerland, 2015), p. 872.
- [19] P. Schütze, *Transversale Strahldynamik bei der Erzeugung kohärenter Synchrotronstrahlung*, BestMasters (Springer, New York, 2017).
- [20] N. Hiller, A. Hofmann, E. Huttel, V. Judin, B. Kehrer, M. Klein, S. Marsching, and A.-S. Müller, in *Proceedings of the 2nd International Particle Accelerator Conference, San Sebastián, Spain* (EPS-AG, Spain, 2011), p. 2951.
- [21] A. Hofmann and F. Méot, Optical resolution of beam cross-section measurements by means of synchrotron radiation, *Nucl. Instrum. Methods Phys. Res.* **203**, 483 (1982).
- [22] Å. Andersson, V. Schlott, M. Rohrer, A. Streun, and O. Chubar, in *Proceedings of the 10th European Particle Accelerator Conference, Edinburgh, Scotland, 2006* (EPS-AG, Edinburgh, 2006), p. 1223.
- [23] OpTaliX Pro 8.82, www.optenso.com.
- [24] K. L. F. Bane, K. Oide, and M. Zobov, in *1st CARE-HHH-APD Workshop on Beam Dynamics in Future Hadron Colliders and Rapidly Cycling High-Intensity Synchrotrons* (CERN, Geneva, Switzerland, 2005).
- [25] B. Nash, N. Carmignani, L. Farvacque, S. Liuzzo, T. Perron, P. Raimondi, R. Versteegen, and S. White, in *Proceedings of IPAC'15* (JACoW, Geneva, Switzerland, 2015), p. 113.
- [26] J. Safranek, G. Portmann, and X. Huang, Theme section - LOCO, ICFA Beam Dynamics Newsletter **44**, 43 (2009), https://icfa-usa.jlab.org/archive/newsletter/icfa_bd_nl_44.pdf.
- [27] M. Caselle, M. Brosi, S. Chilingaryan, T. Dritschler, V. Judin, A. Kopmann, A.-S. Müller, J. Raasch, N. Smale, J. Steinmann, M. Vogelgesang, S. Wuensch, M. Siegel, and M. Weber, in *Real Time Conference (RT), 2014 19th IEEE-NPSS* (IEEE, New York, 2014), p. 1.
- [28] M. Caselle, L. Ardila Perez, M. Balzer, A. Kopmann, L. Rota, M. Weber, M. Brosi, J. Steinmann, E. Bründermann, and A.-S. Müller, A high-speed DAQ framework for future high-level trigger and event building clusters, *J. Instrum.* **12**, C01040 (2017).
- [29] Virginia Diodes, Inc., <http://vadiodes.com/index.php/en/>.
- [30] Z. Jiang and X.-C. Zhang, Electro-optic measurement of THz field pulses with a chirped optical beam, *Appl. Phys. Lett.* **72**, 1945 (1998).
- [31] E. Roussel, S. Bielawski, A. Borysenko, C. Evain, N. Hiller, A.-S. Müller, P. Schönfeldt, J. L. Steinmann, and C. Sz waj, in *Proceedings of IBIC'15* (JACoW, Geneva, Switzerland, 2016), p. 33.
- [32] N. Hiller, A. Borysenko, E. Hertle, V. Judin, B. Kehrer, S. Marsching, A. Müller, and M. Nasse, in *Proceedings of IPAC'14* (JACoW, Geneva, Switzerland, 2014).
- [33] L. Rota, M. Balzer, M. Caselle, M. Weber, G. Niehues, P. Schönfeldt, M. Nasse, A.-S. Müller, N. Hiller, A. Mozzanica, C. Gerth, B. Steffen, D. Makowski, and A. Mielczarek, in *Proceedings of IBIC'16* (JACoW, Geneva, Switzerland, 2016), p. 740.
- [34] M. Caselle, L. Rota, M. Balzer, M. Brosi, S. Funkner, B. Kehrer, M. Nasse, G. Niehuse, M. Patil, P. Schönfeldt, M. Schuh, J. Steinmann, M. Yan, E. Bründermann, M. Weber, A.-S. Müller, G. Borghi, M. Boscardin, and S. Ronchin, in

- Proceedings of IBIC'17* (JACoW, Geneva, Switzerland, 2017), p. 12.
- [35] S. Funkner, E. Blomley, E. Bründermann, M. Caselle, N. Hiller, M. J. Nasse, G. Niehues, L. Rota, P. Schönfeldt, S. Walther, M. Weber, and A.-S. Müller, High throughput data streaming of individual longitudinal electron bunch profiles in a storage ring with single-shot electro-optical sampling, [arXiv:1809.07530](https://arxiv.org/abs/1809.07530).
- [36] A. Hofmann, I. Birkel, M. Fitterer, S. Hillenbrand, N. Hiller, E. Huttel, V. Judin, M. Klein, S. Marsching, A.-S. Müller, N. Smale, K. Sonnad, and P. Tavares, in *Proceedings of the International Particle Accelerator Conference, Kyoto, Japan* (ICR, Kyoto, 2010), p. 924.
- [37] B. Kehrer, E. Blomley, M. Brosi, E. Bründermann, N. Hiller, A.-S. Müller, M. J. Nasse, M. Schedler, P. Schönfeldt, M. Schuh, P. Schütze, N. Smale, and J. L. Steinmann, in *Proceedings of IPAC'16* (JACoW, Geneva, Switzerland, 2016), p. 109.
- [38] J. Steinmann, M. Brosi, E. Bründermann, M. Caselle, S. Funkner, B. Kehrer, M. Nasse, G. Niehues, L. Rota, P. Schönfeldt, M. Siegel, M. Weber, and A.-S. Müller, in *Proceedings of the 60th ICFA Advanced Beam Dynamics Workshop on Future Light Sources* (JACoW, Geneva, Switzerland, 2018).
- [39] B. Kehrer, E. Blomley, M. Brosi, E. Bründermann, N. Hiller, A.-S. Müller, M. Nasse, M. Schedler, M. Schuh, M. Schwarz, P. Schönfeldt, P. Schütze, N. Smale, and J. Steinmann, in *Proceedings of IPAC'17* (JACoW, Geneva, Switzerland, 2017) pp. 53–56.
- [40] P. Schönfeldt, M. Brosi, M. Schwarz, J. L. Steinmann, and A.-S. Müller, Parallelized Vlasov-Fokker-Planck solver for desktop personal computers, *Phys. Rev. Accel. Beams* **20**, 030704 (2017).
- [41] T. Boltz, Comprehensive analysis of micro-structure dynamics in longitudinal electron bunch profiles, Master thesis, Karlsruhe Institute of Technology, 2017.