

Thoughts of an Isovoluminetric Thermal Desorption Experiment

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- Experiment as automotive application





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Simple analytical Solutions

Without re-diffusion, complete outgasing:

$$\underbrace{p_{r}}_{t \to \infty} = \underbrace{p_{end}}_{switching off} + \frac{V_{sa} k_{s,sa} \sqrt{p_{load} R T}}{\left(\frac{V_{c}}{chamber} - \frac{V_{sa}}{sample}\right)^{2}}$$

With phase equilibrium, mass conservation (number of "hydrogens" in atomic interstitial and molecular gaseous state constant) and <u>non</u> interacting confinement condition:

$$0 = \frac{2(V_c - V_{sa})}{RT} \left(\sqrt{p_r}\right)^2 + V_s k_{s,sa} \underbrace{\sqrt{p_r}}_{="x"} - \left(V_s k_{s,sa} \sqrt{p_{end}} + \frac{2p_{end}(V_c - V_{sa})}{RT}\right)$$

$$p_{T}_{t \to \infty} = \left(\frac{-1 + \left(\frac{8(V_{c} - V_{sa})\sqrt{p_{load}}}{R T V_{s} k_{s,sa}} + \frac{16 p_{end}(V_{c} - V_{sa})^{2}}{\left(R T V_{sa} k_{s,sa}\right)^{2}}\right)}{\left(\frac{V_{sa} k_{s,sa} R T}{4 (V_{c} - V_{sa})}\right)^{-1}} \right)^{2}$$

3.: Structure of differential equations:



$$c(0 \le r \le r_s, z = \pm z_s, \forall t) = k_{s,sa} \sqrt{p(t)}$$

$$c(r = r_s, |z| \le z_s, \forall t) = k_{s,sa} \sqrt{p(t)}$$

$$d(1)(r \le r_{co}, z = -z_{ci}, \forall t) = k_{s,cu} \sqrt{p(t)}$$

$$d(2)(r = r_{ci}, -z_{ci} \le z \le z_{ci}, \forall t) = k_{s,cu} \sqrt{p(t)}$$

$$d(2)(r_{ci} \le r \le r_{co}, z = \pm z_{ci}, \forall t) = k_{s,cu} \sqrt{p(t)}$$

$$d(3)(r \le r_{co}, z = z_{ci}, \forall t) = k_{s,cu} \sqrt{p(t)}$$



3. Structure of differential equations



$$\frac{dm}{dt} = -2 \pi D_{sa} \int_{0}^{r_{s}} \frac{2}{symmetric} r \frac{\partial}{\partial z} c(r, z = z_{s}, t) dr$$

$$- \underbrace{2 \pi r_{s} D_{sa} \int_{-z_{s}}^{z_{s}} \frac{\partial}{\partial r} c(r = r_{s}, z, t) dz}_{superficies surface of specimen}$$

$$- \underbrace{2 \pi D_{cu} \int_{0}^{r_{co}} r \frac{\partial}{\partial z} d(1)(r, z = -z_{ci}, t) dr}_{circular area of body 1}$$

$$- 2 \pi r_{ci} D_{cu} \int_{-z_{ci}}^{z_{ci}} \frac{\partial}{\partial r} d(2)(r = r_{ci}, z, t) dz$$

$$- \underbrace{2 \pi D_{cu} \int_{0}^{r_{co}} r \frac{\partial}{\partial z} d(3)(r, z = z_{ci}, t) dr}_{circular area of body 3}$$

$$p(t) = p_{start} + \underbrace{k_{v}}_{RT_{abs}/v_{gas}} \int_{t_{1}+t_{2}}^{t} \underbrace{0.5}_{gaseous\leftrightarrow} \frac{dm}{dt} dt$$

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4. Outlook to analytical solution



Die Lösung für die Diffusionsgleichung mit Zylinderkoordinaten (φ -unabhängig)

$$\frac{dc}{dt} = D_{sa} \left\{ \frac{\partial^2 c}{\partial r^2} + \frac{1}{r} \frac{\partial c}{\partial r} + \frac{\partial^2 c}{\partial z^2} \right\}$$

mit Randbedingung $K_{s,sa}\sqrt{p}$ (Konstante) für

$$r=r_s, |z|\leq z_s$$

und

 $0 \le r \le r_s, \ z = \pm z_s$

ist

$$c(r,z,t) = \sum_{k=1}^{\infty} \sum_{n=0}^{\infty} B_{kn} e^{-\mu_{kn}t} \cos\left(\frac{n\pi z}{2z_s}\right) J_0\left(\sqrt{\frac{\mu_{kn}}{D_{sa}}} - \left(\frac{n\pi}{2z_s}\right)^2 r\right) + K_{s,sa}\sqrt{p},$$

wobei

$$B_{kn} = \frac{\int_{-z_s}^{z_s} \int_{0}^{r_s} \left(u(r, z, 0) - K_{s, sa} \sqrt{p} \right) \cos\left(\frac{n\pi z}{2z_s}\right) J_0 \left(\sqrt{\frac{\mu_{kn}}{D_{sa}} - \left(\frac{n\pi}{2z_s}\right)^2 r} \right) r dr dz}{\int_{-z_s}^{z_s} \int_{0}^{r_s} \cos^2\left(\frac{n\pi z}{2z_s}\right) J_0 \left(\sqrt{\frac{\mu_{kn}}{D_{sa}} - \left(\frac{n\pi}{2z_s}\right)^2 r} \right)^2 r dr dz}$$

und

$$\mu_{kn} = D_{sa} \left\{ \left(\frac{\alpha_k}{r_s} \right)^2 + \left(\frac{n\pi}{2z_s} \right)^2 \right\}$$

mit α_k : *k*-te Nullstelle von Besselfunktion J_0 . Eine Symmetrie für *z*-Achse wurde vorausgesetzt. Um Konstante B_{kn} zu berechnen, muss u(r, z, 0) durch eine gegebene Funktion (Anfangsbedingung) ersetzt werden.



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Comparison to analytical solution





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Main goal of the first one is the determination of transport parameters of hydrogen in structural metallic materials used for components in fusion power stations:

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