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DNA-templated control of chirality and efficient energy transport in supramolecular DNA architectures with aggregation-induced emission†

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Two conjugates of tetraphenylethylene with D-2'-deoxyuridine (**1b**) and L-2'-deoxyuridine (**1l**) were synthesized to construct new supramolecular DNA-architectures by self-assembly. The non-templated assemblies of **1b** and **1l** show strong aggregation-induced emission and their chirality is exclusively controlled by the configuration of their sugar part. In contrast, the chirality of the DNA-templated assemblies is governed by the configuration of the DNA, and there is no configuration-selective binding of **1b** to D-A₂₀ and **1l** to L-A₂₀. The quantum yield of the assembly of **1b** along the single-stranded DNA A₂₀ is 0.40; approximately every second available binding site on the DNA template is occupied by **1b**. The strong aggregation-induced emission of these DNA architectures can be efficiently quenched and the excitation energy can be transported to Atto dyes at the 5'-terminus. A multistep energy transport "hopping" precedes the final energy transfer to the terminal acceptor. The building block **1b** promotes this energy transport as stepping stones. This was elucidated by reference DNA double strands in which **1b** was covalently incorporated at two distinct sites in the sequences, one near the Atto dye, and one farther away. This new type of completely self-assembled supramolecular DNA architecture is hierarchically ordered and the DNA template controls not only the binding but also the energy transport properties. The high intensity of the aggregation-induced emission and the excellent energy transport properties make these DNA-based materials promising candidates for optoelectronic applications.

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Introduction

Conventional chromophores suffer from aggregation-caused quenching (ACQ), a phenomenon describing the quenched emission at high concentrations. For technical applications, particularly OLEDs, chromophores with strong fluorescence in the aggregated or condensed phase are crucial. The term "aggregation-induced emission" (AIE) was introduced by Tang's group in 2001 describing the emissive behavior of 1,2,3,4,5-pentaphenylsilole in ethanol.¹ Since then, the number of organic chromophores that show AIE is constantly increasing;^{2,3} in particular, tetraphenylethylene (Tpe) has been identified as a lead structure with AIE.⁴ Similar to the aforementioned silol, the propeller-shaped and non-planar Tpe is weakly emissive in dilute solution, but becomes highly emissive in the aggregated or solid state. It is not a distinct type of stacking interaction that increases the emission but rather the restriction of intramolecular rotations blocking the non-radiative decay pathways from the photoexcited state.⁵

The supramolecular organization of chromophores gives access to nanostructured materials with unique optical properties.⁶ The bottom-up approach is a powerful concept to achieve this goal,⁷ and DNA is a unique template with sequence specificity through the recognition by hydrogen bonding to arrange chromophores in a precise way,⁸ and this DNA-like structure also persists in the solid state for optoelectronic applications.⁹ The most precise, covalent incorporation of chromophores into DNA is limited to 5–10 units in a row, until the yields drop.^{10,11} The non-covalent self-assembly of nucleosides along single-stranded DNA is an important alternative, for instance with naphthalenes¹² and porphyrins^{13,14} through binding to T₂₀ or T₄₀. We recently described the sequence-selective assembly of pyrenes,^{15–17} perylenes¹⁸ and nile red^{16,17} as modified 2'-deoxynucleosides along single-stranded DNA templates and mainly ACQ was observed. However, AIE would increase the value of DNA-based materials for optoelectronics. This makes Tpe a promising candidate for DNA-based architectures with AIE.¹⁹ The fluorescence intensity increase of typical DNA staining agents, such as SYBR Green, is a phenomenon related to AIE since rigidification of the chromophore has an impact on the fluorescence in both cases.¹⁹ However, it should be noted here that, molecules, particularly derivatives of Tpe, showing AIE in assemblies are promising candidates for new materials. Achiral Tpe derivatives promote

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sensing of DNA by their emission^{20–22} and circularly polarized luminescence.^{23,24} The covalent conjugation of Tpe with DNA²⁵ helps in the visualization of cellular RNA²⁶ and gives DNA-grafted nanosheets,²⁷ and the AIE can be controlled by DNA hybridization.²⁸ To construct new supramolecular DNA-architectures with strong AIE by self-assembly, we present two chiral conjugates of Tpe, **1D** and **1L** (Fig. 1). The configuration of their 2'-deoxyribofuranosides differs to probe the chirality of both the non-templated assembly and the assembly along D- and L-DNA templates by means of optical spectroscopy. These supramolecular architectures are hierarchically ordered and the DNA template controls the binding selectivity, the chirality and the light harvesting through efficient energy transport to the Atto dye attached to the 5'-terminus.

Results and discussion

Non-templated supramolecular assemblies

Both conjugates **1D** and **1L** were synthesized by Sonogashira-coupling as described in the ESI (Schemes S1, S2–S7).[†] Firstly, the formation of assemblies with **1D** and **1L**, respectively, was investigated in aqueous solution without any DNA template. Stock solutions of **1D** and **1L** in THF or DMSO (4.3 mM) were prepared to ensure full solubility of the modified 2'-deoxyuridines. The assembly of **1D** and **1L** was induced by diluting samples (8.8 μ L) of the stock solutions with H₂O (total volume of 1 mL, 37.5 μ M in H₂O with 0.9% THF or DMSO). The first samples of **1D** and **1L** were simply stored for 1 h at room temperature and characterized by optical spectroscopy (Fig. 2). A second set of samples was treated under the typical conditions of DNA annealing, including incubation at 90 °C for 5 min and slow cooling to room temperature (approximately 1 °C min⁻¹). Interestingly, these different preparations yield different types of water-soluble assemblies. Representatively, we discuss the optical properties of the samples that were prepared with THF as cosolvent; the corresponding results with DMSO as cosolvent are similar (see Fig. S22[†]). The fluorescence of the assemblies of **1D** and **1L** prepared at room temperature show maxima at 478

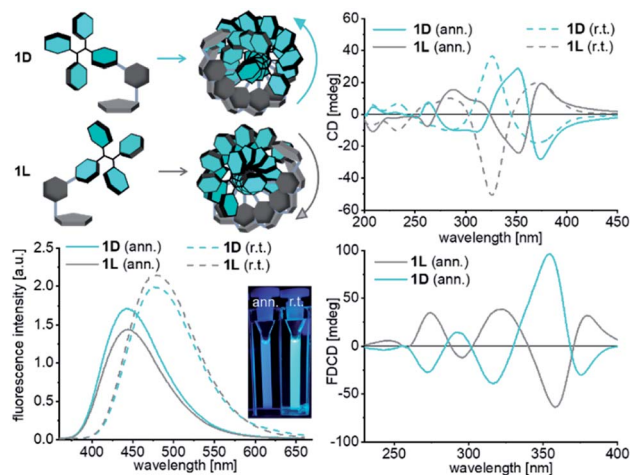


Fig. 2 Fluorescence (left bottom), circular dichroism (CD, right top) and fluorescence-derived circular dichroism (FDCD, bottom right) spectra of the assemblies prepared with **1D** and **1L** in the absence of any DNA template (37.5 μ M **1D/1L**, $\lambda_{\text{exc}} = 341$ nm, H₂O, 0.9% THF, for the UV/Vis absorption see Fig. S21,[†] for the experiments with 0.9% DMSO see Fig. S22[†]) at room temperature (r.t.) or by annealing (ann.) after heating to 90 °C for 5 min.

nm, whereas the annealed assemblies show maxima at 442 nm and reduced intensities. The first type of assembly is dominated by Tpe interactions, whereas the second, the annealed type of assembly is, at least partially, controlled by stacking of the 5-ethynyl-2'-deoxyuridines. Accordingly, the electronic decoupling of Tpe from the 5-ethynyl 2'-deoxyuridine part by a rotational twist along the phenylene bridge causes a blue-shift of fluorescence.

Chirality is an essential feature of these supramolecular structures and therefore the CD spectra were recorded. Remarkably, the CD spectra of assemblies of **1D** and **1L** from both types of preparation show strong mirror bisignate signals in the absorption range between 305 nm and 450 nm of the Tpe chromophore. Interestingly, the D-configured conjugate **1D** yields an assembly with left-handed helicity according to this bisignate signal,²⁹ whereas **1L** yields assemblies with right-handed chirality. The CD signals of the annealed assemblies of **1D** and **1L** look similar in the range between 305 nm and 400 nm and their crossing points are shifted from 350 nm to 370 nm. The absorbance range between 200 nm and 305 nm, which is the typical absorption range for the 5-ethynyl-uracil part of the conjugates, shows another bisignate signal. Here, the annealed assembly formed with **1D** shows left-handed helicity whereas the annealed assembly formed with **1L** shows right-handed chirality. Their crossing point is at 270 nm, which is near the absorbance maximum of 5-ethynyl-2'-deoxyuridine. This supports our aforementioned hypothesis that the annealed assemblies are controlled by the nucleoside stacking. The chirality order may be unexpected, but, in fact, it is not. For instance, the D-configured Nile red-modified 2'-deoxyuridines yielded assemblies with left-handed chirality, too.³⁰ The CD spectra show very clearly that the chirality of the assemblies of **1D** and **1L** is exclusively controlled by the configuration of the

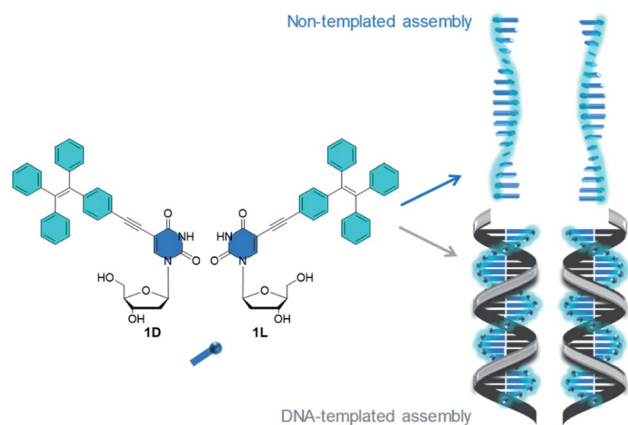


Fig. 1 Tetraphenylethylene-modified 2'-deoxyuridines **1D** and **1L**. Illustration of the non-templated assembly and DNA-templated assembly of **1D** and **1L**.



attached 2'-deoxyribofuranosides. The fluorescence detected circular dichroism (FD CD) combines conventional CD with the high sensitivity of fluorescence spectroscopy. It allows investigating the chiral behavior of fluorescent chromophores in a non-emissive chiral architecture, such as DNA.³¹ The detected FD CD signals are mirror bisignate signals and again show that the chirality of the assemblies of **1D** and **1L** is controlled by the different configurations of the attached chiral 2'-deoxyribofuranosides.

DNA-templated supramolecular architectures

The next level of hierarchically ordered supramolecular architectures is the structural control of the self-assembly of **1D** and **1L** in aqueous solution by DNA templates. Therefore, we followed our established protocol to prepare such DNA-based supramolecular architectures.^{32,33} We used **1D** and **1L** as building blocks and probed their assembly along D- and L-configured DNA templates with the sequences A₂₀ and T₂₀ by means of optical spectroscopy. Small volumes of 2'-deoxynucleoside stock solutions in DMSO were added to an aqueous solution of the DNA templates (1 mL, 1.25 μM). The concentrations of both stock solutions were prepared sufficiently high that not more than 0.9% DMSO as cosolvent was added to the final samples in water. This amount of cosolvent is typically tolerated by the DNA conformation. Both, **1D** and **1L**, were added in a 1.5-fold excess to promote the occupancy of the available binding sites along the templates. 250 mM NaCl is known to stabilize double-helical DNA and thus expected to also support the formation of the non-covalent supramolecular DNA architectures with **1D** and **1L**. Both conjugates are nearly insoluble in this high salt aqueous solution unless they are bound to single stranded DNA and therefore remain in the solution. The excess and unbound **1D** or **1L** was removed by short centrifugation (3 min @ 16 000g). The supernatant contained the DNA-templated assembly and was

investigated by optical spectroscopy (Fig. 3). The absorbances in the range between 300 nm and 400 nm of the assemblies of **1D** with D-A₂₀ and **1L** with L-A₂₀ as templates are significantly higher than those with D-T₂₀ and L-T₂₀. This result reveals preferred binding of **1D** and **1L** to D-A₂₀ and L-A₂₀, respectively, which matches the rules of canonical base-pairing, because both modified 2'-deoxyridines are complementary to 2'-deoxyadenosines in the DNA template A₂₀, but not complementary to T₂₀. This selectivity is not as pronounced as for the Watson-Crick base pairing in unmodified double-stranded DNA because "mismatched" binding of **1D** and **1L** to the wrong templates D-T₂₀ and L-T₂₀ occurs to a certain extent. There is, however, no configuration-selective recognition by the DNA templates. The absorbances of the assemblies of **1D** with D-A₂₀ and L-A₂₀ are equally high (Fig. 3), which is also the case for the assemblies of **1L** with both D-A₂₀ and L-A₂₀ (Fig. S23†). The absorbances can be used to calculate the occupancy fraction *f* which is the number of occupied binding sites (by **1D** or **1L**) divided by the number of available binding sites on the template. For the assemblies of **1D** and **1L** along D-A₂₀ and L-A₂₀, *f* lies in the range of 0.55 ± 0.1; approximately 11 out of 20 available binding sites of the DNA templates are occupied. Obviously, the four phenyl groups of Tpe are sterically hindering and prevent a complete occupancy of all binding sites, as previously observed for planar chromophores, such as Nile red.³² The fluorescence gives additional support for the selective binding of the chromophore-nucleosides. Since Tpe shows AIE and unbound chromophore conjugates were removed by centrifugation, the fluorescence can be clearly attributed to the DNA-templated assemblies of **1D** and **1L**. The assembly of **1D** along D-A₂₀ (and L-A₂₀) shows the strongest fluorescence intensity with a maximum at 492 nm and a quantum yield of Φ = 0.40. The assemblies of **1L** along L-A₂₀ (and D-A₂₀) show similarly strong fluorescence intensities (Fig. S23†), whereas the assemblies of **1L** and **1D** along the wrong templates L-T₂₀ and D-T₂₀ show significantly weaker fluorescence intensities. These results track well with the absorbance differences and the idea of selective binding of **1L** and **1D** to D-A₂₀ and L-A₂₀.

Although configuration selectivity was not observed for the binding of **1D** and **1L** to the DNA templates, the chirality of the formed helical assemblies was probed by CD spectroscopy. The CD spectra of the DNA-templated assemblies of **1D** and **1L** show generally weaker signals than those of the non-templated assemblies, and there are only signals of 2'-deoxynucleosides in the absorbance range between 200 nm and 300 nm (dU in the Tpe-conjugates and A or T in the templates). In contrast to the non-templated assemblies of **1D** and **1L**, as discussed above, the chirality in the templated assemblies is not controlled by the configuration of 2'-deoxyribofuranosides in **1D** and **1L**, but instead by the configuration of the DNA templates. The CD signal of all four DNA templates shows the slightly different helical preorganized conformations of the single strands (Fig. S24†) which are not changed by the binding of **1D** (Fig. 3) or **1L** (Fig. S23†). According to the characteristic bisignate signals with a zero crossing at approximately 280 nm, assembling **1D** with D-A₂₀ and D-T₂₀ yields right-handed helicity, whereas with L-A₂₀ and L-T₂₀ it gives left-handed helicity in accordance with the

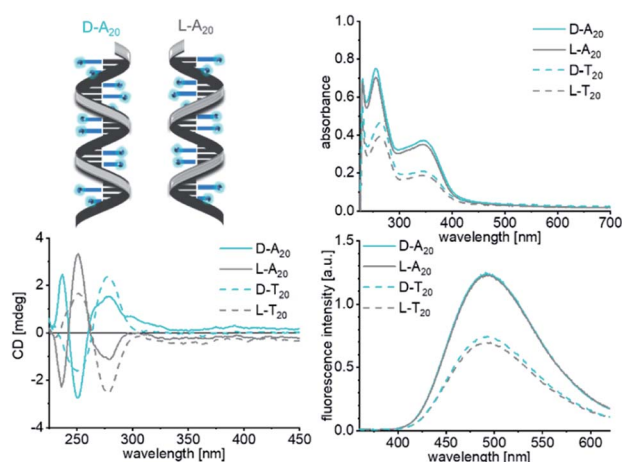


Fig. 3 UV/Vis absorbance (top right), fluorescence (bottom right) and circular dichroism (bottom left) spectra of the DNA-templated assemblies of **1D** with D-A₂₀, L-A₂₀, D-T₂₀ and L-T₂₀ (1.25 μM DNA in water + 0.9% DMSO, 250 mM NaCl, λ_{exc} = 341 nm, and supernatant after centrifugation 3 min @ 16 000g). The corresponding spectra of **1L** are shown Fig. S23†



expected chirality for such DNA-like helical structures. There is no measurable CD above 305 nm and thus no information on the ordered chirality of the Tpe chromophores along the DNA templates. Only approximately every second binding site of the DNA templates is occupied by **1D** or **1L**, preventing the chromophores from CD-active excitonic interactions. Based on the determined general helicity of the DNA-like architectures as determined by their CD signals between 200 nm and 300 nm, we assume that the Tpe chromophores are also helically arranged.

The next step towards DNA-based light harvesting systems is to transfer the excitation energy from the aggregation-emissive Tpe to an appropriate energy acceptor. For the preparation of such energy transfer systems, we focused solely on the components with *D*-configuration (**1D** and *D*-**A**₂₀) due to the commercial availability of appropriate dye–DNA conjugates. But in principle, we assume that the following experiments would also work with the components with *L*-configuration. The DNA template *D*-**A**₂₀ bears the Atto565 and the Atto633 dye conjugated to the 5'-termini that serve as energy acceptors in the supramolecular assemblies of **1D**. The preparation was carried out following the previous protocol including centrifugation to remove unbound **1D** by precipitation from the solution. The fluorescence of the Tpe chromophores in the assemblies along the **A**₂₀–Atto565 and **A**₂₀–Atto633 templates at 492 nm is strongly reduced compared to that in the assembly of **1D** along *D*-**A**₂₀. As a result of energy transfer, the strong fluorescence signals of the Atto dyes appear at 598 nm and 660 nm, respectively. When the donor **1D**, as part of the assembled DNA architectures, is excited at 389 nm, the quantum yields in the fluorescence range of the donor (400 nm–550 nm, **1D** with *D*-**A**₂₀, without acceptor: $\Phi_D = 0.400$; **1D** with **A**₂₀–Atto565: $\Phi_{DA} = 0.056$; and **1D** with **A**₂₀–Atto633: $\Phi_{DA} = 0.079$) and the quantum yields in the fluorescence range of the acceptor (550–675 nm, **A**₂₀–Atto565: $\Phi_A = 0.575$ and **A**₂₀–Atto633: $\Phi_A = 0.366$) were determined. This allows the calculation of the energy transfer efficiencies according to $E = 1 - \frac{\Phi_{DA}}{\Phi_A}$ which gives remarkable values of $E = 86\%$ for the DNA architecture with *D*-**A**₂₀–Atto565 and $E = 80\%$ for that with *D*-**A**₂₀–Atto633. The fluorescence lifetime of **1D** along the template **A**₂₀ at 492 nm is 3.78 ns when excited at 366 nm. It is shortened to 2.11 ns in the DNA architecture with **A**₂₀–Atto565 and to 2.59 ns with **A**₂₀–Atto633, which is caused by the energy transfer. The lifetimes of the energy acceptors are 2.78 ns at 598 nm (Atto565), and 3.41 ns at 660 nm (Atto633) in the DNA architectures with **1D** when excited at the donor wavelength. We tried to measure the lifetimes of **A**₂₀–Atto565 and **A**₂₀–Atto633 without **1D** at the donor excitation wavelength (366 nm) too, but the fluorescence intensities were too low for any lifetime determination. As an alternative, we measured the steady-state fluorescence of **A**₂₀–Atto565 and **A**₂₀–Atto633 in the absence of **1D**, yielding only very small intensities when excited at the donor wavelength (341 nm, Fig. 4, top right), which rules out direct excitation of the acceptors and gives clear evidence for the efficient energy migration from **1D** to the Atto dyes in these assembled DNA architectures.

When control experiments with the “wrong” templates **T**₂₀ were performed, **T**₂₀–Atto565 and **T**₂₀–Atto633 show

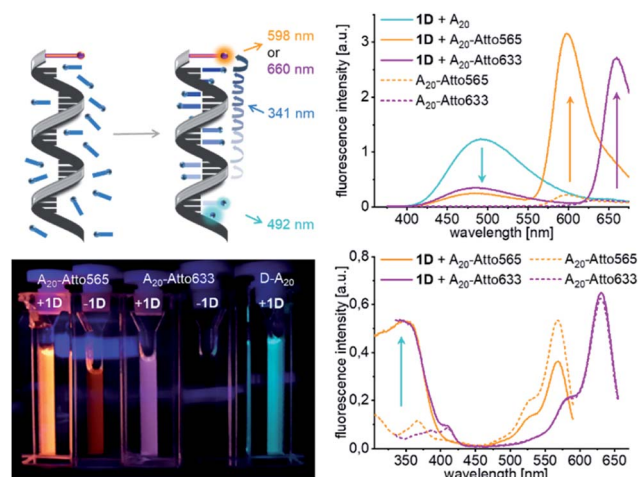


Fig. 4 Formation of supramolecular light-harvesting DNA architectures with **1D** and *D*-**A**₂₀ templates that are modified at the 5'-terminus with the Atto565 or the Atto633 dye; fluorescence (top right), 1.25 μ M DNA in water + 0.9% DMSO, 250 mM NaCl, and $\lambda_{exc} = 341$ nm), excitation spectra (bottom right), detected at an emission wavelength of $\lambda_{em} = 598$ or 660 nm and images of cuvettes under the UV lamp (bottom left). For the UV/Vis absorption see Fig. S25.†

significantly less energy transfer since very few molecules of **1D** are bound to these “wrong” DNA templates according to the low absorbance at 345 nm after removal of unbound **1D** by centrifugation (Fig. S26†). As mentioned above, the assembly of **1D** along *D*-**A**₂₀ (without any attached Atto dye) has a quantum yield of $\Phi_F = 0.40$. Accordingly, the emission of the **1D** assemblies along **A**₂₀–Atto565 and **A**₂₀–Atto633 as templates is quenched by 86% and 80%, respectively. The excitation spectra clearly show that the emission of the Atto dyes originates from the excited **1D** units. If we assume an occupancy fraction of $f = 0.55$ for the 20 binding sites available on these templates, the emission of approximately 9 out of 11 DNA-bound molecules of **1D** is quenched. Based on a regular stacking distance of 3.4 Å, the farthest **1D** engaged in energy transfer would be approximately 58 Å away from the Atto dye. The efficient quenching over such long distances is likely not the result of one-step energy transfer processes from the individual Tpe that is assembled as a donor at different distances to the 5'-terminal Atto dye as the acceptor. In contrast, the efficient quenching indicates a step-wise energy transport between the Tpe molecules taking place before the final energy transfer to the Atto acceptor dyes occurs.

In order to probe the distance dependence of the energy transfer between **1D** and the Atto dyes, we incorporated **1D** as a modified nucleotide into synthetic single-stranded DNA. The optical properties of this covalent modification were compared with the non-covalent supramolecular assemblies (Fig. 5). The phosphoramidite as a DNA building block was synthesized by standard procedures and used for automated oligonucleotide synthesis on a solid phase as described in the ESI (Fig. S8–S12).† The synthetic oligonucleotides **TPE1a**, **TPE1b**, **TPE2a** and **TPE2b** were purified by semi-preparative HPLC and were checked by MALDI-TOF mass spectrometry. **TPE1a** bears a single **1D** modification in the middle of the DNA sequence, and **TPE1b**, nearer to the 5'-terminus.



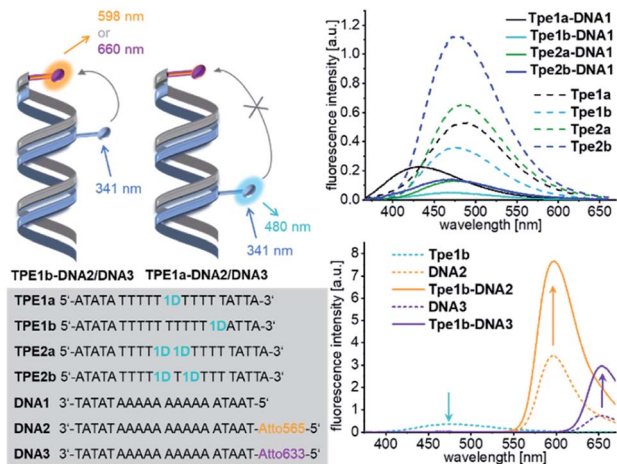


Fig. 5 Sequences of DNA strands with covalent **1D** modification **TPE1a** to **TPE2b** and complementary counterstrands **DNA1** to **DNA3** (bottom left), forming DNA hybrids, e.g. **TPE1-DNA1**; fluorescence of hybrids with unmodified counterstrand **DNA1** (top right) and with Atto-counterstrands **DNA2/DNA3** (bottom right), $c(\text{DNA}) = c(\text{cs}) = 2.5 \mu\text{M}$, 250 mM NaCl, 10 mM NaPi, and $\lambda_{\text{exc}} = 341 \text{ nm}$. For UV/Vis absorption see Fig. S27.†

The fluorescence of these single-stranded oligonucleotides show maxima at 489 nm. The fluorescence intensity is significantly reduced and the maximum is shifted to 430 nm in the double stranded hybrid **TPE1a-DNA1**, but not in **TPE1b-DNA1**, when the modified single strands were annealed with the complementary and unmodified counterstrand **DNA1**. This was similarly observed in the annealed assemblies of **1D** (see above). We explain this observation again by stacking the 5-ethynyl-uracil part in a DNA-like double-helical assembly which induces a rotational twist to the phenylene group and decouples the Tpe chromophore from the ethynyl-nucleoside part. This stacking of the core part of the DNA architectures is induced by annealing with complementary counter strand **DNA1** with strong effects on the fluorescence – less fluorescence intensity and blue-shift. The conformation of the Tpe chromophore is obviously influenced also by the base on the 3'-side of the **1D** modification, since the fluorescence of the DNA hybrid **TPE1b-DNA1** is only quenched, but not blue-shifted. **TPE1a** and **TPE1b** were subsequently annealed with **DNA2** bearing the Atto565 dye or with **DNA3** bearing the Atto633 dye as acceptor dyes. The TPE emission of the hybrids **TPE1a-DNA2** and **TPE1a-DNA3** is not quenched at all and there is no observable energy transfer. Obviously, the distance of 34 Å is too long for energy transfer. In comparison, the hybrids **TPE1b-DNA2** and **TPE1b-DNA3** show 48% and 46% quenching of the TPE emission and a detectable energy transfer to the Atto dyes over a distance of 17 Å. In comparison, in the non-covalent DNA architectures **1D**_{20-D-A20}-Atto565 and **1D**_{20-D-A20}-Atto, the pronounced fluorescence quenching involves an energy transfer distance of up to 58 Å, as discussed above. The typical sigmoidal Förster dependence of the energy transfer efficiency on the distance with a characteristic Förster radius cannot be applied to DNA architectures with chromophores that are attached with short and rigid

linkers.^{34–36} Taken together, it implies that a multistep homo-energy transport between the assembled Tpe chromophores precedes the final energy transfer to the Atto dyes as acceptors in the supramolecular structures. Such homo-energy transport cannot be promoted by the unmodified A–T pairs, separating the Tpe from the Atto dyes in the hybrids **TPE1a-DNA2** and **TPE1a-DNA3**. This result underscores the significant role of the DNA templates that control not only the direction of the energy transport but also significantly improve the energy transport properties by their building blocks within the supramolecular DNA architectures, such that the light is harvested much better. The strands **TPE2a** and **TPE2b** were modified twice with the nucleotide **1D**. The hybridization with the counter strand **DNA1** affects the fluorescence in the same manner like in case of **TPE1b**. **TPE2a** or **TPE2b** was annealed with **DNA2** modified with the Atto565 dye or with **DNA3** modified with the Atto633 dye. In both cases, there was no energy transfer observable between **1D** and the Atto dyes in these double-stranded DNA hybrids.

Conclusions

The two new conjugates **1D** and **1L** with 2'-deoxyuridine as the recognition unit for 2'-deoxyadenosine in single-stranded DNA templates and with Tpe as a fluorescent chromophore can be applied to prepare a new type of supramolecular DNA architecture by self-assembly. Most importantly, all assemblies, including the non-templated and the DNA-templated ones, show significant AIE. The quantum yield of the assembly of **1D** along **A**₂₀ is 0.40; approximately 55% of the available binding sites on the DNA template are occupied by **1D**. The two conjugates **1D** and **1L** differ by the configurations of 2'-deoxyribofuranosides to probe the chirality of both their non-templated assembly and their templated assembly along **D**- and **L**-configured DNA templates by means of optical spectroscopy. The chirality of the non-templated assemblies of **1D** and **1L** is exclusively controlled by the configuration of their sugar parts. In contrast, the chirality of the templated assemblies is governed by the configuration of the sugar part of the DNA templates, and there is no configuration-selective binding of **1D** to **D**-**A**₂₀ and **1L** to **L**-**A**₂₀. These supramolecular DNA architectures are hierarchically ordered and the DNA template controls not only the chirality, but also the binding selectivity and the energy transport properties. Moreover, the strong AIE of these DNA architectures with **1D** and **1L** can be efficiently quenched and the excitation energy can be transported to Atto dyes that were attached to the 5'-termini of the DNA templates. The building blocks of the self-assembled DNA architectures participate themselves in the energy transport along the DNA template; a multistep energy “hopping” transport precedes the final energy transfer to the terminal Atto dyes. This was elucidated by the reference DNA double strands in which we incorporated **1D** covalently at two distinct sites in the sequences, one near the Atto dye (4 base pairs, *ca.* 17 Å), and one farther away (9 base pairs, 34 Å). Only the first DNA construct showed energy transfer properties. The high intensity of the AIE of this new type of supramolecular DNA architecture and the excellent energy transport properties efficiently harvest light and make



these DNA-based materials promising candidates for optoelectronic applications.

Author contributions

HU performed the experiments and wrote parts of the manuscript. HAW wrote the manuscript and supervised the research project.

Conflicts of interest

There are no conflicts to declare.

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