

# Benchmarking Magnetizabilities with Recent Density Functionals

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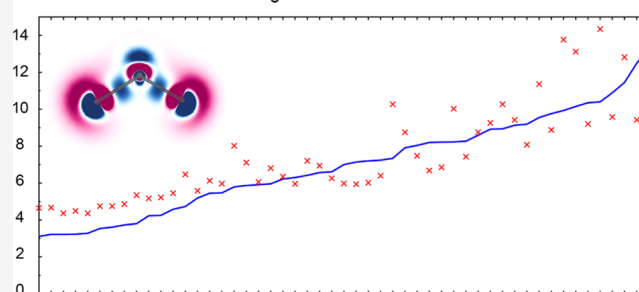
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**ABSTRACT:** We have assessed the accuracy of the magnetic properties of a set of 51 density functional approximations, including both recently published and already established functionals. The accuracy assessment considers a series of 27 small molecules and is based on comparing the predicted magnetizabilities to literature reference values calculated using coupled-cluster theory with full singles and doubles and perturbative triples [CCSD(T)] employing large basis sets. The most accurate magnetizabilities, defined as the smallest mean absolute error, are obtained with the BHandHLYP functional. Three of the six studied Berkeley functionals and the three range-separated Florida functionals also yield accurate magnetizabilities. Also, some older functionals like CAM-B3LYP, KT1, BHLYP (BHandH), B3LYP, and PBE0 perform rather well. In contrast, unsatisfactory performance is generally obtained with Minnesota functionals, which are therefore not recommended for calculations of magnetically induced current density susceptibilities and related magnetic properties such as magnetizabilities and nuclear magnetic shieldings. We also demonstrate that magnetizabilities can be calculated by numerical integration of magnetizability density; we have implemented this approach as a new feature in the gauge-including magnetically induced current (GIMIC) method. Magnetizabilities can be calculated from magnetically induced current density susceptibilities within this approach even when analytical approaches for magnetizabilities as the second derivative of the energy have not been implemented. The magnetizability density can also be visualized, providing additional information that is not otherwise easily accessible on the spatial origin of magnetizabilities.

Mean absolute errors of magnetizabilities and their standard deviations



## 1. INTRODUCTION

Computational methods based on density functional theory (DFT) are commonly used in quantum chemistry because DFT calculations are rather accurate despite their relatively modest computational costs. Older functionals such as the Becke'88–Perdew'86<sup>1,2</sup> (BP86), Becke'88–Lee–Yang–Parr<sup>1,3</sup> (BLYP), and Perdew–Burke–Ernzerhof<sup>4,5</sup> (PBE) functionals at the generalized gradient approximation (GGA) as well as the B3LYP<sup>6</sup> and PBE0<sup>7,8</sup> hybrid functionals are still often employed, even though newer functionals with improved accuracy for energies and electronic properties have been developed.

The accuracy and reliability of various density functional approximations (DFAs) have been assessed in a huge number of applications and benchmark studies.<sup>9–17</sup> It is important to note that functionals that are accurate for energetics may be less suited for calculations of other molecular properties.<sup>16</sup> In specific, the accuracy of magnetic properties calculated within DFAs has been benchmarked by comparing magnetizabilities and nuclear magnetic shieldings to those obtained from coupled-cluster calculations using large basis sets,<sup>18,19</sup> although modern DFAs have been less systematically investigated.<sup>16,20–23</sup> The same also holds for nuclear independent chemical shifts<sup>24–28</sup> and magnetically induced current density susceptibilities,<sup>29–36</sup> which have been studied for a large number of

molecules, but whose accuracy has never been benchmarked properly.

Magnetizabilities are usually calculated as the second derivative of the electronic energy with respect to the external magnetic perturbation<sup>37–41</sup>

$$\xi_{\alpha\beta} = - \left. \frac{\partial^2 E}{\partial B_\alpha \partial B_\beta} \right|_{\mathbf{B}=0} \quad (1)$$

Such analytic implementations for magnetizabilities exist in several quantum chemistry programs. However, as the magnetic interaction energy can also be written as an integral over the magnetic interaction energy density  $\rho^{\mathbf{B}}(\mathbf{r})$  given by the scalar product of the magnetically induced current density  $\mathbf{J}^{\mathbf{B}}(\mathbf{r})$  with the vector potential  $\mathbf{A}^{\mathbf{B}}(\mathbf{r})$  of the external magnetic field  $\mathbf{B}$ <sup>30,31,42–45</sup>

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$$E = \int \rho^{\mathbf{B}}(\mathbf{r}) d^3r = -\frac{1}{2} \int \mathbf{A}^{\mathbf{B}}(\mathbf{r}) \cdot \mathbf{J}^{\mathbf{B}}(\mathbf{r}) d^3r \quad (2)$$

an approach based on quadrature is also possible. As shown in Section 2, the numerical integration approach for the magnetizability provides additional information about its spatial origin that is not available with the analytic approach based on second derivatives: the tensor components of the magnetizability density defined in Section 2 are scalar functions that can be visualized, and the integration approach can be used to provide detailed information about the origin of the corresponding components of the magnetizability tensor. Similar approaches have been used in the literature for studying spatial contributions to nuclear magnetic shielding constants.<sup>46–53</sup>

We will describe our methods for numerical integration of magnetizabilities using the current density susceptibility in Sections 2 and 3. Then, in Section 4, we will list the studied set of density functionals and present the results in Section 5: the functional benchmark is discussed in Section 5.1 and magnetizability densities and spatial contributions to magnetizabilities are analyzed in Section 5.2. The conclusions of the study are summarized in Section 6. Atomic units are used throughout the text unless stated otherwise, and summation over repeated indices is assumed.

## 2. THEORY

The current density  $\mathbf{J}^{\mathbf{B}}(\mathbf{r})$  in eq 2 is formally defined as the real part ( $\mathcal{R}$ ) of the mechanical momentum density

$$\mathbf{J}^{\mathbf{B}}(\mathbf{r}) = -\mathcal{R}[\Psi^*(\mathbf{r})(\mathbf{p} - \mathbf{A}^{\mathbf{B}}(\mathbf{r}))\Psi(\mathbf{r})] \quad (3)$$

where  $\mathbf{p} = -i\nabla$  is the momentum operator. Substituting eq 2 into eq 1 straightforwardly leads to

$$\xi_{\alpha\beta} = \frac{\partial^2}{\partial B_\alpha \partial B_\beta} \frac{1}{2} \int \mathbf{A}^{\mathbf{B}}(\mathbf{r}) \cdot \mathbf{J}^{\mathbf{B}}(\mathbf{r}) d^3r \Big|_{\mathbf{B}=0} \quad (4)$$

The current density susceptibility tensor<sup>29–31</sup> (CDT) is defined as the first derivative of the magnetically induced current density with respect to the components of the external magnetic field in the limit of a vanishing magnetic field<sup>32–35</sup>

$$\mathcal{J}_\gamma^{\mathbf{B}_\beta} = \frac{\partial \mathcal{J}_\gamma^{\mathbf{B}}}{\partial B_\beta} \Big|_{\mathbf{B}=0} \quad (5)$$

The vector potential  $\mathbf{A}^{\mathbf{B}}(\mathbf{r})$  of an external static homogeneous magnetic field is expressed as

$$\mathbf{A}^{\mathbf{B}}(\mathbf{r}) = \frac{1}{2} \mathbf{B} \times (\mathbf{r} - \mathbf{R}_O) \quad (6)$$

where  $\mathbf{R}_O$  is the chosen gauge origin. The  $\alpha\beta$  component of the magnetizability tensor can then be obtained from eqs 4–6 as

$$\xi_{\alpha\beta} = \int \rho_{\alpha\beta}^\xi(\mathbf{r}) d^3r \quad (7)$$

where the magnetizability density is defined as

$$\rho_{\alpha\beta}^\xi(\mathbf{r}) = \frac{1}{2} \sum_{\delta\gamma} \epsilon_{\alpha\delta\gamma} r_\delta \mathcal{J}_\gamma^{\mathbf{B}_\beta}(\mathbf{r}) \quad (8)$$

where  $\epsilon_{\alpha\delta\gamma}$  is the Levi–Civita symbol,  $\alpha, \beta, \gamma$ , and  $\delta$  are one of the Cartesian directions ( $x, y, z$ ), and  $r_\delta$  also denotes one of ( $x, y, z$ ). The components of the magnetizability density tensor  $\rho_{\alpha\beta}^\xi(\mathbf{r})$  are scalar functions that can be visualized to obtain information

about the spatial contributions to the corresponding element of the magnetizability tensor  $\xi_{\alpha\beta}$ .

As the isotropic magnetizability ( $\bar{\xi}$ ) is obtained as the average of the diagonal elements of the magnetizability tensor

$$\bar{\xi} = \frac{1}{3} \text{Tr} \xi = \int \bar{\rho}^\xi(\mathbf{r}) d^3r \quad (9)$$

we introduce the isotropic magnetizability density  $\bar{\rho}^\xi(\mathbf{r})$  defined as

$$\bar{\rho}^\xi(\mathbf{r}) = \frac{1}{3} \text{Tr} \rho^\xi(\mathbf{r}) \quad (10)$$

which yields information about the spatial origin of the isotropic magnetizability, as we demonstrate in Section 5.2.

Although there is freedom with regard to the choice of the gauge origin of  $\mathbf{A}^{\mathbf{B}}(\mathbf{r})$ , the magnetic flux density  $\mathbf{B}$  is uniquely defined via eq 6, because  $\mathbf{B} = \nabla \times (\mathbf{A}(\mathbf{r}) + \nabla f(\mathbf{r}))$  holds for any differentiable scalar function  $f(\mathbf{r})$ . The exact solution of the Schrödinger equation should also be gauge invariant. However, the use of finite one-particle basis sets introduces gauge dependence in quantum chemical calculations of magnetic properties. The CDT can be made gauge origin independent by using gauge-including atomic orbitals (GIAOs), also known as (a.k.a.) London atomic orbitals (LAOs)<sup>32,54,55</sup>

$$\chi_\mu(\mathbf{r}) = e^{-i(\mathbf{B} \times [\mathbf{R}_\mu - \mathbf{R}_O] \cdot \mathbf{r})/2} \chi_\mu^{(0)}(\mathbf{r}) \quad (11)$$

where  $i$  is the imaginary unit and  $\chi_\mu^{(0)}(\mathbf{r})$  is a standard atomic-orbital basis function centered at  $\mathbf{R}_\mu$ . GIAOs eliminate the gauge origin from the expression used for calculating the CDT; the expression we use is given in the Supporting Information (SI). Since the expression for the magnetizability density in eqs 7 and 8 can be computed by quadrature, magnetizabilities can be obtained from the CDT even if the corresponding analytical calculation of magnetizabilities as the second derivative of the energy has not been implemented.

## 3. IMPLEMENTATION

The present implementation is based on the gauge-including magnetically induced current (GIMIC) program<sup>56</sup> and the NUMGRID library,<sup>57</sup> which are both freely available open-source software. Gauge-independent CDTs can be calculated with GIMIC<sup>32–35</sup> using the density matrix, magnetically perturbed density matrices, and information about the basis set.

To evaluate eq 7, a molecular integration grid is first generated from atom-centered grids with the NUMGRID library, as described by Becke.<sup>58</sup> In NUMGRID, the grid weights are scaled according to the Becke partitioning scheme using a Becke hardness of 3;<sup>58</sup> the atom-centered grids are determined by a radial grid generated as suggested by Lindh et al.,<sup>59</sup> and angular grids by Lebedev<sup>60</sup> are used.

Given the quadrature grid, the diagonal elements of the magnetizability tensor are calculated in GIMIC from the Cartesian coordinates of the  $n$  grid points multiplied with the CDT calculated in the grid points. For example, the  $\xi_{xx}$  element of the magnetizability tensor is obtained from eq 7 as

$$\xi_{xx} = \sum_{i=1}^n w_i \rho_{i;xx}^\xi \quad (12)$$

where the  $xx$  component of the magnetizability density tensor at grid point  $i$  with quadrature weight  $w_i$  is

$$\rho_{i;xx}^{\xi} = \frac{1}{2}[(y\mathcal{J}_z^{B_x})_i - (z\mathcal{J}_y^{B_x})_i] \quad (13)$$

where  $(y\mathcal{J}_z^{B_x})_i$  and  $(z\mathcal{J}_y^{B_x})_i$  are the products of the  $z$  and  $y$  components of the CDT calculated in grid point  $i$  with the Cartesian coordinates  $y$  and  $z$  of the grid point, respectively, and the external magnetic field perturbation is along the  $x$ -axis,  $B_x$ . The  $\xi_{yy}$  and  $\xi_{zz}$  elements are obtained analogously.

#### 4. COMPUTATIONAL METHODS

Calculations are performed for the set of 28 molecules studied in ref 18 that also provides our molecular structures and the CCSD(T) reference values: AlF, C<sub>2</sub>H<sub>4</sub>, C<sub>3</sub>H<sub>4</sub>, CH<sub>2</sub>O, CH<sub>3</sub>F, CH<sub>4</sub>, CO, FCCH, FCN, H<sub>2</sub>C<sub>2</sub>O, H<sub>2</sub>O, H<sub>2</sub>S, H<sub>4</sub>C<sub>2</sub>O, HCN, HCP, HF, HFCO, HOF, LiF, LiH, N<sub>2</sub>, N<sub>2</sub>O, NH<sub>3</sub>, O<sub>3</sub>, OCS, OF<sub>2</sub>, PN, and SO<sub>2</sub>. However, as in ref 18, O<sub>3</sub> was omitted from the analysis since it is an outlier and due to the fact that the reliability of the CCSD(T) level of theory is not guaranteed for this system: the perturbative triples correction to the magnetizability of O<sub>3</sub> is  $-46.2 \times 10^{-30}$  J/T<sup>2</sup>, indicating that the CCSD(T) result might still have large error bars.<sup>18</sup> The results of this work thus only pertain to the 27 other molecules, as in ref 18.

Electronic structure calculations were performed with Hartree–Fock (HF) and the functionals listed in Tables 1 and 2 using TURBOMOLE 7.5.<sup>110</sup> Several rungs of Jacob's ladder were considered when choosing the functionals listed in Tables 1 and 2: local density approximations (LDAs), generalized gradient approximations (GGAs), and meta-GGAs (mGGAs). Several kinds of functionals are also included: (pure) density functional approximations, global hybrid (GH) functionals with a constant amount of HF exchange, and range-separated (RS) hybrids with a given amount of HF exchange in the short range (SR) and the long range (LR). As can be seen in Tables 1 and 2, the evaluated functionals consist of 1 pure LDA, 8 pure GGAs, 8 global hybrid GGAs, 10 range-separated hybrid GGAs, 12 mGGAs, 8 global hybrid mGGAs, and 4 range-separated mGGAs, in addition to HF.

The Dunning aug-cc-pCVQZ basis set<sup>111–115</sup> (with aug-cc-pVQZ on the hydrogen atoms) and benchmark quality integration grids were employed in all calculations. Universal auxiliary basis sets<sup>116</sup> were used with the resolution-of-the-identity approximation for the Coulomb interaction in all TURBOMOLE calculations. All density functionals were evaluated in TURBOMOLE with LIBXC,<sup>117</sup> except the calculations with the recently published CAMh-B3LYP functional for which XCFun was used.<sup>118</sup> Magnetizabilities were subsequently evaluated with GIMIC by numerical integration of eq 7. The data necessary for evaluating the CDT in GIMIC were obtained from TURBOMOLE calculations of nuclear magnetic resonance (NMR) shielding constants employing GIAOs.<sup>54,55,110,119,120</sup>

Although response calculations are not possible at the moment in the presence of the non-local correlation kernel used in  $\omega$ B97X-V, B97M-V, and  $\omega$ B97M-V, we have estimated the importance of the van der Waals (vdW) effects on the magnetic properties by comparing magnetizabilities obtained with orbitals optimized with and without the vdW term in the case of SO<sub>2</sub>. The magnetizability obtained with the vdW-optimized orbitals differed by only  $0.4 \times 10^{-30}$  J/T<sup>2</sup> (0.14%) from that obtained from a calculation where the vdW term was omitted in the orbital optimization. Thus, the vdW term appears to have very little influence on magnetizabilities, as is already

**Table 1. Functionals at the Local Density Approximation (LDA) and the Generalized Gradient Approximation (GGA) Considered in This Work<sup>f</sup>**

functional	hybrid	type	notes	LIBXC ID <sup>a</sup>	references
LDA		LDA		1 + 7	61–63
BLYP		GGA		106 + 131	1, 3, and 64
BP86		GGA		106 + 132	1 and 2
CHACHIYO		GGA		298 + 309	65 and 66
KT1		GGA		167	67
KT2		GGA		146	67
KT3		GGA	PySCF data used	587	68
N12		GGA		82 + 80	69
PBE		GGA		101 + 130	4 and 5
B3LYP	GH	GGA	20% HF	402	6
revB3LYP <sup>b</sup>	GH	GGA	20% HF	454	70
B97-2	GH	GGA	21% HF	410	71
B97-3	GH	GGA	26.9% HF	414	72
BHLYP <sup>c</sup>	GH	GGA	50% HF	435	61, 62, and 73
BHandHLYP <sup>d</sup>	GH	GGA	50% HF	436	1 and 73
PBE0	GH	GGA	25% HF	406	7 and 8
QTP-17	GH	GGA	62% HF	416	74
N12-SX	RS	GGA	25% SR, 0% LR	81 + 79	75
CAM-B3LYP	RS	GGA	19% SR, 65% LR	433	76
CAMh-B3LYP <sup>e</sup>	RS	GGA	19% SR, 50% LR		77
CAM-QTP-00	RS	GGA	54% SR, 91% LR	490	78
CAM-QTP-01	RS	GGA	23% SR, 100% LR	482	79
CAM-QTP-02	RS	GGA	28% SR, 100% LR	491	80
$\omega$ B97	RS	GGA	0% SR, 100% LR	463	81
$\omega$ B97X	RS	GGA	15.8% SR, 100% LR	464	81
$\omega$ B97X-D	RS	GGA	22.2% SR, 100% LR	471	82
$\omega$ B97X-V	RS	GGA	16.7% SR, 100% LR	531	83

<sup>a</sup>Two numbers indicate the exchange and correlation functionals, respectively. A single number indicates an exchange–correlation functional. <sup>b</sup>Revised version. <sup>c</sup>Following King et al. in refs 84–86, BHLYP is defined as 50% LDA exchange, 50% HF exchange, and 100% LYP correlation. It is sometimes also known as BHandH, which is its keyword in Gaussian. <sup>d</sup>BHandHLYP is 50% Becke'88 exchange, 50% HF exchange, and 100% LYP correlation. <sup>e</sup>CAMh-B3LYP is defined using the XCFun library with  $\alpha = 0.19$ ,  $\beta = 0.31$ , and  $\mu = 0.33$ . <sup>f</sup>GH stands for global hybrid and RS for range-separated hybrid. The amount of Hartree–Fock (HF) exchange or exact exchange in the short range (SR) and the long range (LR) is also given.

well-known in the literature for other properties.<sup>121</sup> The vdW term was therefore not included in the calculations using the  $\omega$ B97X-V, B97M-V, and  $\omega$ B97M-V functionals in this study.

The accuracy of the numerical integration in GIMIC was assessed by comparing the TURBOMOLE/GIMIC magnetizability data to analytical values from PySCF,<sup>122</sup> in which LIBXC<sup>117</sup> was also used to evaluate the density functionals. Since PySCF does not currently support magnetizability calculations with mGGA functionals or range-separated functionals, further calculations were undertaken with Gaussian

**Table 2. Meta-GGA Functionals (mGGA) Considered in This Work<sup>d</sup>**

functional	hybrid	type	notes	LIBXC ID <sup>a</sup>	references
B97M-V		mGGA		254	87
M06-L		mGGA		449 + 235	88
revM06-L <sup>b</sup>		mGGA		293 + 294	89
M11-L		mGGA		226 + 75	90
MN12-L		mGGA		227 + 74	91
MN15-L		mGGA		268 + 269	92
TASK		mGGA		707 + 13	93 and 94
MVS		mGGA		257 + 83	95 and 96
SCAN		mGGA		263 + 267	97
rSCAN <sup>c</sup>		mGGA		493 + 494	98
TPSS		mGGA		457	99 and 100
revTPSS <sup>b</sup>		mGGA		212 + 241	96 and 101
TPSSh	GH	mGGA	10% HF	457	102
revTPSSh <sup>b</sup>	GH	mGGA	10% HF	458	96, 101, and 102
M06	GH	mGGA	27% HF	449 + 235	103
revM06 <sup>b</sup>	GH	mGGA	40.4% HF	305 + 306	104
M06-2X	GH	mGGA	54% HF	450 + 236	103
M08-HX	GH	mGGA	52.2% HF	295 + 78	105
M08-SO	GH	mGGA	56.8% HF	296 + 77	105
MN15	GH	mGGA	44% HF	268 + 269	106
M11	RS	mGGA	42.8% SR, 100% LR	297 + 76	107
revM11 <sup>b</sup>	RS	mGGA	22.5% SR, 100% LR	304 + 172	108
MN12-SX	RS	mGGA	25% SR, 0% LR	248 + 73	75
$\omega$ B97M-V	RS	mGGA	15% SR, 100% LR	531	109

<sup>a</sup>Two numbers indicate the exchange and correlation functionals, respectively. A single number indicates an exchange–correlation functional. <sup>b</sup>Revised version. <sup>c</sup>Regularized version. <sup>d</sup>The notation is the same as in Table 1.

16.<sup>123</sup> The analytical magnetizabilities from PySCF and Gaussian were found to be in perfect agreement for the studied LDA and GGA functionals available in both codes (LDA, BP86, PBE, PBE0, BLYP, B3LYP, and BHandLYP). A comparison of the data from PySCF to the GIMIC data revealed the numerically integrated magnetizabilities to be accurate, as the magnetizabilities agreed within  $0.5 \times 10^{-30} \text{ J/T}^2$  for all molecules using the B3LYP, B97-2, B97-3, BLYP, BP86, KT1, KT2, LDA, PBE, and PBE0 functionals; the small discrepancy may arise from the use of the resolution-of-identity approximation<sup>124</sup> in TURBOMOLE or from the numerical integration of the magnetizability density. A comparison of the raw data for BP86 and B3LYP is given in the SI.

The magnetizabilities calculated with Gaussian and TURBOMOLE using the meta-GGA functionals were found to differ. The discrepancies between the magnetizabilities obtained with the two programs are due to the use of different approaches to handle the gauge invariance of the kinetic energy density in meta-GGAs, which are described in refs 125 and 126 for Gaussian and TURBOMOLE, respectively. We found the TURBOMOLE data to be significantly closer to the CCSD(T) reference values.

Finally, since we found the implementation of the KT3 functional in LIBXC version 5.0.0 used by TURBOMOLE to be flawed, the KT3 results in this study are based on calculations with PySCF with a corrected version of LIBXC.

## 5. RESULTS

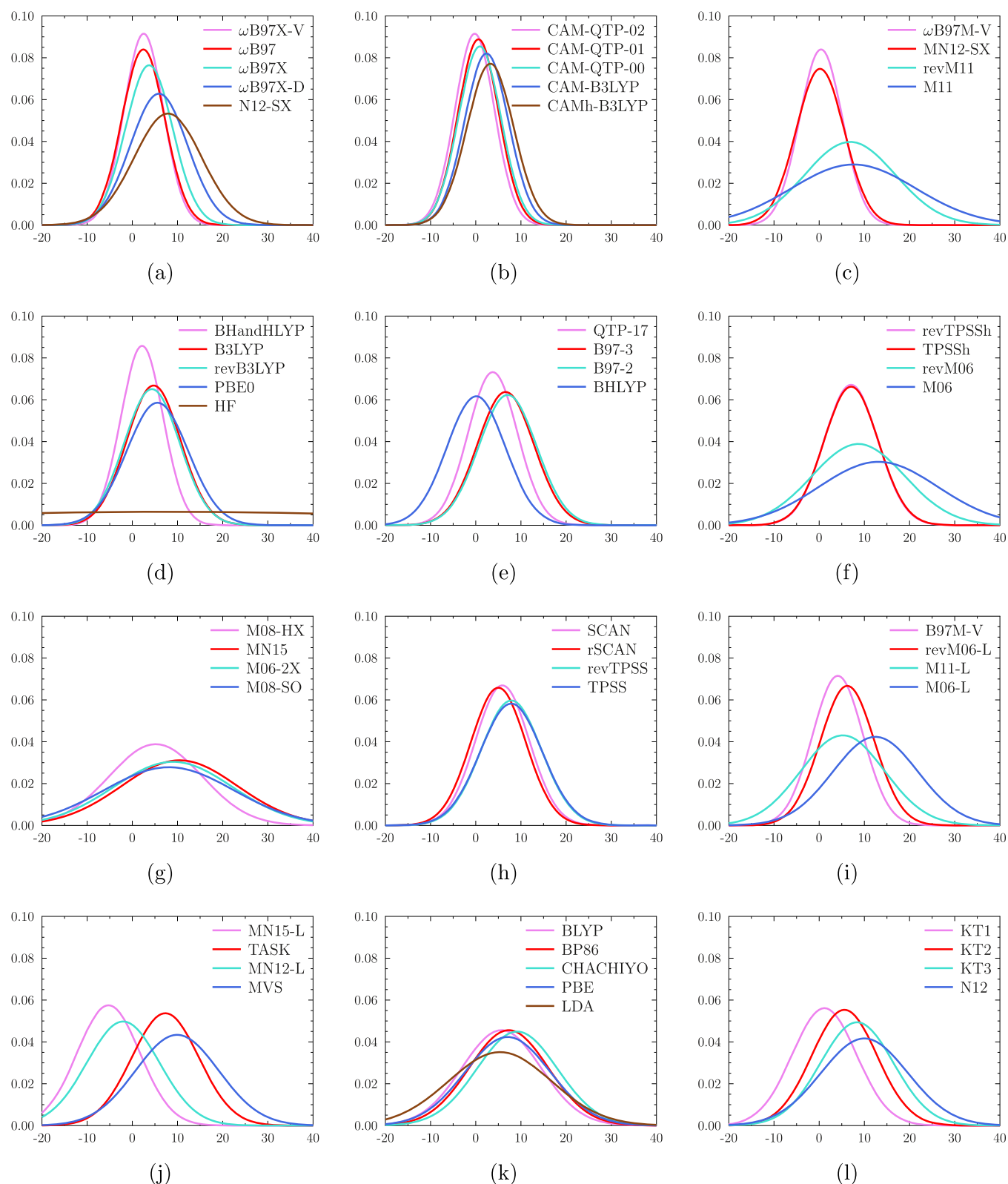
**5.1. Functional Benchmark.** The deviations of the DFT magnetizabilities from the CCSD(T) reference values of ref 18 are visualized as ideal normal distributions (NDs) in Figure 1. The visualization shows the idealized distribution of the error in the magnetizability for each functional, based on the computed mean errors (ME) and standard deviation of the error (STD) given in Table 3. The raw data on the magnetizabilities and the differences from the CCSD(T) reference are available in the SI. Although the error distributions in Figure 1 are instructive, we will employ mean absolute errors (MAEs) to rank the functionals studied in this work in a simple, unambiguous fashion. The MAEs are also given in Table 3.

Examination of the data in Table 3 shows that range-separated (RS) functionals generally yield accurate magnetizabilities. Judged by the mean absolute error, the best performance is obtained with the BHandHLYP GH functional. BHandHLYP is followed by 10 RS functionals, which have much sharper distributions than the rest of the studied functionals. The best performing RS functionals are three of the six Berkeley RS functionals ( $\omega$ B97X-V,  $\omega$ B97, and  $\omega$ B97M-V) and the three RS functionals from the University of Florida's Quantum Theory Project (QTP) CAM-QTP-00, CAM-QTP-01, and CAM-QTP-02. Five of these functionals have 100% long-range (LR) HF exchange, while the CAM-QTP-00 functional has 91% LR HF exchange. The two other RS Berkeley functionals with 100% LR exchange are ranked 11th ( $\omega$ B97X) and 21st ( $\omega$ B97X-D) among the studied functionals. The NDs of the studied RS GGA functionals are shown in Figure 1a,b, whereas the NDs of the studied RS mGGA functionals are shown in Figure 1c.

The CAM-B3LYP (65% LR HF exchange) and CAMh-B3LYP (50% LR HF exchange) functionals are among the top 10 functionals (ranked 8th and 10th, respectively). CAM-B3LYP was designed for the accurate description of charge transfer excitations in a dipeptide model,<sup>76</sup> while CAMh-B3LYP functional is aimed at excitation energies of biochromophores.<sup>77</sup>

The best Minnesota functional, MN12-SX, is ranked 9th. MN12-SX is a highly parameterized functional with 58 parameters that is known to require the use of extremely accurate integration grids.<sup>13</sup> Furthermore, since MN12-SX is an RS functional with HF exchange only in the short range (SR), it may have problems modeling magnetic properties of antiaromatic molecules sustaining strong ring currents in the paratropic (nonclassical) direction.<sup>127–129</sup> We illustrate this with calculations on the strongly antiaromatic tetraoxa isophlorin molecule in the Supporting Information: MN12-SX yields a magnetizability that is 4 times larger than the local second-order Møller–Plesset perturbation theory (LMP2) reference value, while the magnetizabilities from BHandHLYP and CAM-B3LYP are in good agreement with LMP2. The MN12-SX functional ranked 32nd is also an RS functional with 0% LR exchange. The RS Minnesota functionals with 100% LR HF exchange (M11 and revM11) have large MAEs of  $9.93 \times 10^{-30} \text{ J/T}^2$  and  $8.87 \times 10^{-30} \text{ J/T}^2$  and are ranked 44th and 35th, respectively.

The best global hybrid (GH) functional is BHandHLYP, which is ranked first among all functionals of this study, as was already mentioned above. Among GHs, BHandHLYP is followed by QTP-17, which is ranked 12th. Old and established GH functionals like BHandH, B3LYP, and PBE0 perform almost as well as QTP-17 and are ranked 13th, 16th, and 20th, respectively. The performance of revB3LYP is



**Figure 1.** Normal distributions (ND) representing the errors in the magnetizabilities for the 27 benchmark reproduced by the studied functionals, obtained by plotting the data presented in Table 3. The curves are ordered in each figure by increasing standard deviation. The NDs of RS functionals are shown in (a)–(c). The NDs of the GH functionals are shown in (d)–(g). The NDs of the mGGA functionals are shown in (h)–(j). The NDs of the LDA and GGA functionals are shown in (k) and (l).

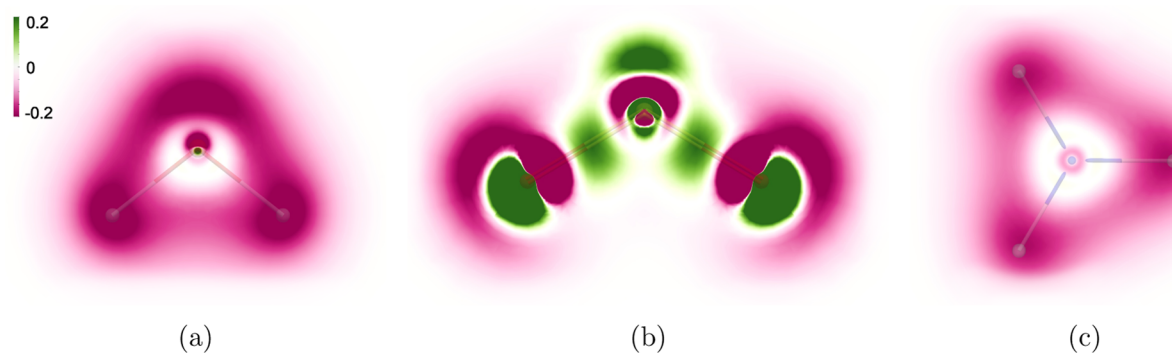
practically the same as that of B3LYP; the same holds for revTPSSh and TPSSh. The other established GH functionals like B97-2, B97-3, and TPSSh and newer ones like revTPSSh

and M08-HX are found in the beginning of the second half of the ranking list, whereas M08-SO, M06, revM06, M06-2X, MN15,

**Table 3.** Mean Absolute Errors (MAEs), Mean Errors (MEs), and Standard Deviations (STDs) for the Magnetizabilities of the 27 Studied Molecules in Units of  $10^{-30}$  J/T<sup>2</sup> from the CCSD(T) Reference with the Studied Functionals<sup>a</sup>

rank	functional	MAE	ME	STD	rank	functional	MAE	ME	STD
1	BHandHLYP	3.11	2.15	4.65	27	revTPSSh	7.14	7.05	5.94
2	CAM-QTP-00	3.22	0.88	4.67	28	TPSSh	7.20	7.07	6.02
3	$\omega$ B97X-V	3.22	2.51	4.36	29	B97-2	7.24	7.07	6.40
4	CAM-QTP-01	3.23	0.59	4.49	30	M08-HX	7.34	5.17	10.27
5	CAM-QTP-02	3.28	-0.23	4.36	31	BLYP	7.91	5.69	8.75
6	$\omega$ B97	3.54	2.44	4.75	32	N12-SX	8.04	7.89	7.48
7	$\omega$ B97M-V	3.61	0.41	4.75	33	revTPSS	8.20	7.86	6.68
8	CAM-B3LYP	3.73	2.38	4.86	34	TPSS	8.22	7.85	6.85
9	MN12-SX	3.80	0.22	5.34	35	revM11	8.23	6.83	10.03
10	CAMh-B3LYP	4.23	3.22	5.17	36	TASK	8.27	7.31	7.43
11	$\omega$ B97X	4.25	3.71	5.22	37	BP86	8.59	7.30	8.75
12	QTP-17	4.58	3.77	5.45	38	M11-L	8.92	5.20	9.26
13	BHLYP	4.73	0.10	6.47	39	revM06	8.94	8.67	10.27
14	B97M-V	5.19	4.13	5.58	40	PBE	9.13	7.07	9.42
15	revB3LYP	5.45	4.34	6.13	41	KT3	9.19	8.38	8.08
16	B3LYP	5.47	4.72	5.97	42	LDA	9.55	5.37	11.36
17	MN12-L	5.79	-2.03	8.02	43	CHACHIYO	9.76	9.17	8.88
18	KT1	5.87	1.15	7.11	44	M11	9.93	7.61	13.77
19	rSCAN	5.91	5.00	6.06	45	M06-2X	10.15	9.01	13.12
20	PBE0	5.96	5.56	6.81	46	MVS	10.35	9.92	9.20
21	$\omega$ B97X-D	6.22	5.89	6.35	47	M08-SO	10.40	8.09	14.34
22	SCAN	6.30	5.89	5.96	48	N12	10.89	10.01	9.58
23	KT2	6.42	5.58	7.21	49	MN15	11.45	10.45	12.82
24	MN15-L	6.57	-5.27	6.94	50	M06-L	12.49	12.45	9.42
25	B97-3	6.61	6.61	6.26	51	M06	13.34	13.11	13.16
26	revM06-L	7.00	6.23	5.98	52	HF	18.40	7.48	61.81

<sup>a</sup>The functionals are ordered in increasing MAE.



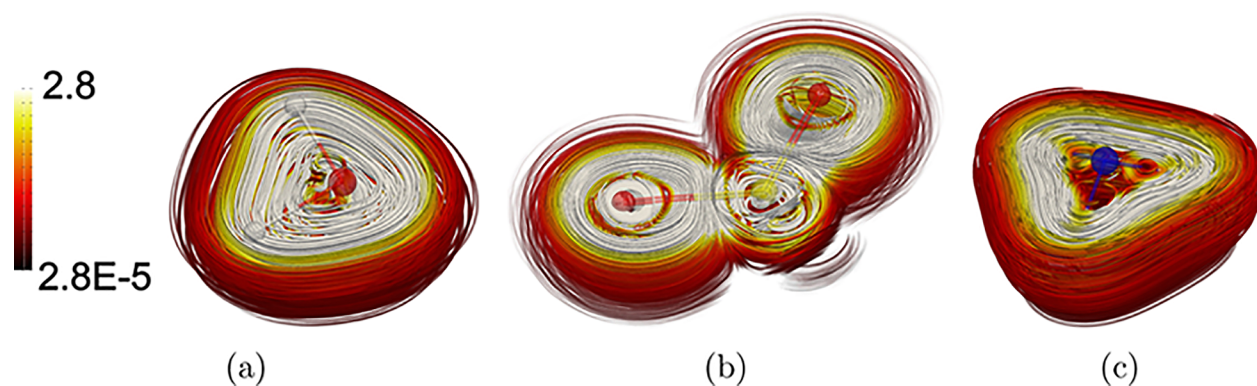
**Figure 2.** Visualization of the isotropic magnetizability density  $\bar{\rho}^E(\mathbf{r})$  (eq 10) shown in the molecular plane of H<sub>2</sub>O (a) and SO<sub>2</sub> (b) as well as in the plane formed by the hydrogen atoms of NH<sub>3</sub> (c), positioned  $0.06 a_0$  away from the N atom toward the hydrogen atoms. Negative contributions are shown in pink and positive ones in green. The gauge origin  $\mathbf{R}_G$  is  $(0, 0, 0) a_0$ .

and M06 are ranked between 39th and 51st. The NDs of the GH functionals are compared in Figure 1d–g.

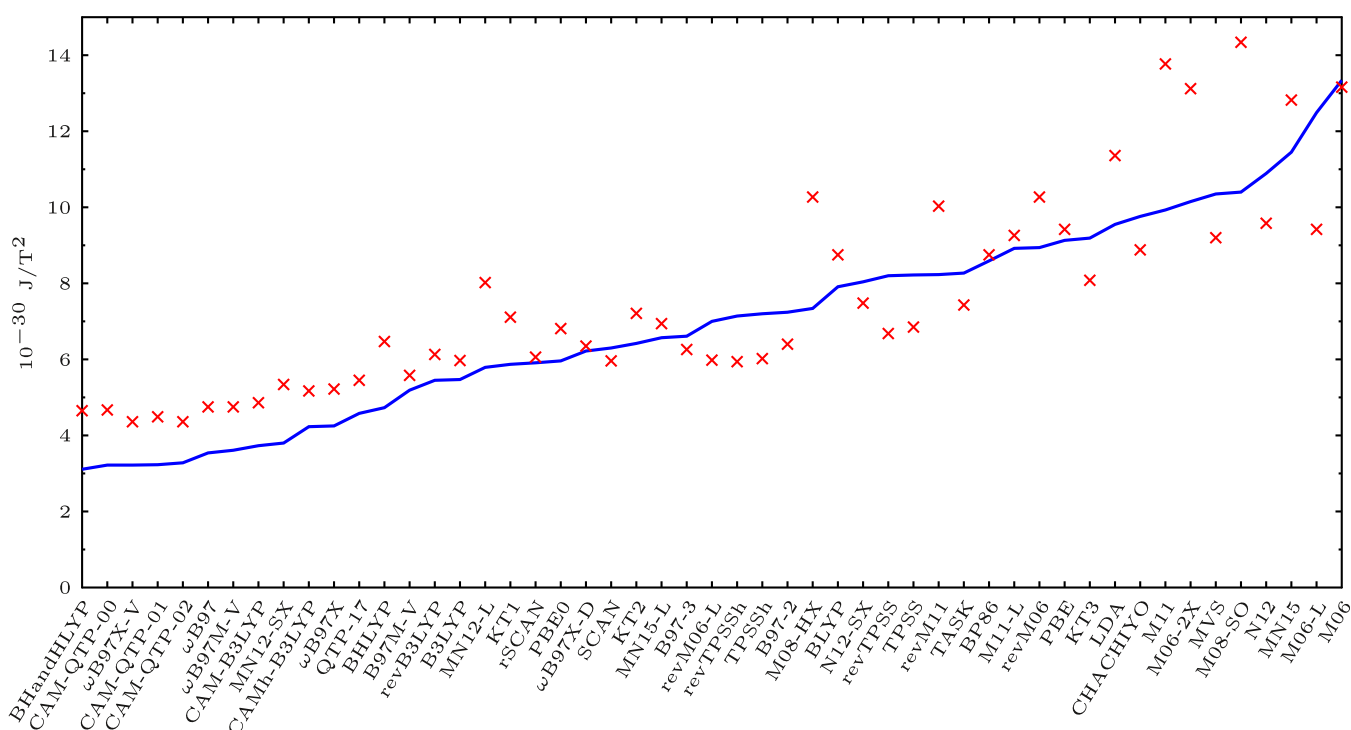
B97M-V, at the 14th place, is the best pure mGGA functional. The rSCAN and SCAN functionals are ranked 19th and 22th, respectively, whereas revTPSS and TPSS appear at positions 33 and 34, respectively. The pure mGGA functionals of the Minnesota series are ranked 17th (MN12-L), 24th (MN15-L), 26th (revM06-L), and 50th (M06-L). The performance of the Minnesota pure mGGA functionals, excluding M06-L, is about the same as that of TASK and the other mGGA functionals. The magnetizabilities calculated with the revised M06-L (revM06-L) functional are more accurate than those with M06-L. The MVS mGGA functional is ranked 46th. The NDs for the mGGA functionals are shown in Figure 1h–j.

The magnetizabilities calculated with several of the Minnesota functionals are inaccurate. Seven of the eight worst performing functionals (M11, M06-2X, MVS, M08-SO, N12, MN15, M06-L, and M06) in Table 3 are Minnesota functionals. Five other Minnesota functionals are also ranked in the lower half, placing 30th (M08-HX), 32th (N12-SX), 35th (revM11), 38th (M11-L), and 39th (revM06).

The KT1 and KT2 functionals are the best GGA functionals, ranking 18th and 23rd, respectively; both KT1 and KT2 have been optimized for NMR shieldings.<sup>67</sup> The older commonly used GGAs i.e., BLYP, BP86, and PBE are ranked 31st, 37th, and 40th, respectively, which is only slightly better than KT3 ranked 41st and LDA ranked 42nd. The CHACHIYO and N12 functionals, which are newer GGAs, are ranked 43rd and 48th,



**Figure 3.** Streamline representation of the CDT (eq 5) of H<sub>2</sub>O (a), SO<sub>2</sub> (b), and NH<sub>3</sub> (c). The CDT is calculated with the magnetic field perpendicular to the molecular plane of H<sub>2</sub>O and SO<sub>2</sub> as well as with it along the symmetry axis of NH<sub>3</sub>. The color scale represents the strength of the CDT in nAT<sup>-1</sup>a<sub>0</sub><sup>-2</sup>.



**Figure 4.** Mean absolute errors (the blue solid line) as well as the errors' standard deviations (red crosses) of the magnetizabilities in 10<sup>-30</sup> J/T<sup>2</sup> of the 27 studied molecules obtained with the 51 functionals compared to the CCSD(T) reference.

respectively. The NDs of the GGA functionals and the LDA are shown in Figure 1k,l.

The magnetizabilities calculated at the HF level are significantly less accurate and have a much larger MAE-STD than those obtained at the DFT levels, and we cannot recommend the use of HF for magnetic properties.

**5.2. Magnetizability Densities.** Spatial contributions to the magnetizability densities, i.e., the integrand in eq 7, are illustrated for H<sub>2</sub>O, NH<sub>3</sub>, and SO<sub>2</sub> in Figure 2, with Figure 3 showing the corresponding CDTs. The magnetizability densities are calculated with the gauge origin of the external magnetic field ( $\mathbf{R}_0$ ) at  $(x, y, z) = (0, 0, 0)$ . In the calculations on H<sub>2</sub>O and SO<sub>2</sub>, the magnetic field perturbation is perpendicular to the molecular plane, while for NH<sub>3</sub>, the perturbation is parallel to the C<sub>3</sub> symmetry axis. In the case of H<sub>2</sub>O, the current density flux around the whole molecule (Figure 3) leads to the ring-shaped contribution shown in Figure 2. The magnetic field

along the symmetry axis of NH<sub>3</sub> also results in a current density flux around the molecule at the hydrogen atoms (Figure 3), giving rise to a similar ring-shaped contribution shown in Figure 2.

The isotropic magnetizability density of SO<sub>2</sub> shown in Figure 2 has positive (green) and negative (pink) values. Calculations of the CDT show that the oxygen atoms sustain a strong diatropic atomic CDT that flows around the atom, whereas the atomic CDT of the sulfur atom is much weaker (Figure 3). The p-orbital shaped contributions to the magnetizability density of SO<sub>2</sub> around the oxygen atoms in Figure 2 originate from the atomic CDTs. The patterns of the CDT of H<sub>2</sub>O and SO<sub>2</sub> lead to the different magnetizability densities shown in Figure 2a,b, respectively. The positive magnetizability densities in H<sub>2</sub>O and NH<sub>3</sub> are extremely localized close to the atomic nuclei, also because of the vortices of the atomic CDT.

The magnetizability density depends on the gauge origin of the vector potential of the external magnetic field, even though the magnetizability is independent of the gauge origin.<sup>43</sup> The magnetizability densities for H<sub>2</sub>O, NH<sub>3</sub>, and SO<sub>2</sub> calculated with the gauge origin at  $\mathbf{R}_O = (1, 1, 1) a_0$  are shown in the SI. The contribution of the choice of the gauge origin to the magnetizability computed from eq 7 vanishes when the CDT fulfills the charge conservation condition<sup>29</sup>

$$\int \mathcal{J}_\alpha^{B_\beta}(\mathbf{r}) d^3r = 0 \quad (14)$$

Calculating the magnetizability for NH<sub>3</sub> with a gauge origin set to  $\mathbf{R}_O = (100, 100, 100) a_0$  yielded a value that differs by 0.32% from the one computed for  $\mathbf{R}_O = (0, 0, 0)$ . When the gauge origin is set to  $\mathbf{R}_O = (1, 1, 1) a_0$ , the deviation is 2 orders of magnitude smaller because the change in the magnetizability depends linearly on the relative position of the gauge origin. The magnetizabilities of H<sub>2</sub>O and SO<sub>2</sub> also change by only 0.46 and 0.03% when moving the gauge origin from  $(0, 0, 0) a_0$  to  $(100, 100, 100) a_0$ , respectively, showing that charge conservation is practically fulfilled in our calculations. All other positions than  $(0, 0, 0)$  for the gauge origin lead to a small, spurious CDT contribution to the magnetizability density.

The GIAO ansatz modifies the atomic orbitals leading to a magnetic response of an external magnetic field that is correct to the first order for the one-center problem.<sup>30,130</sup> Even though GIAOs do not guarantee that the integral condition for the charge conservation of the CDT is fulfilled,<sup>131</sup> the basis set convergence is faster and the leakage of the CDT is much smaller when GIAOs are used.<sup>32</sup>

## 6. CONCLUSIONS

We have calculated magnetizabilities for a series of small molecules using both recently published density functionals, as well as older, established density functionals. The accuracy of the magnetizabilities predicted by the various density functional approximations has been assessed by comparison to coupled-cluster calculations with singles and doubles and perturbative triples [CCSD(T)] reported by Lutnæs et al.<sup>18</sup> Our results are summarized graphically in Figure 4: the top functionals afford both small mean absolute errors and standard deviations, but the same is not true for all recently suggested functionals.

Numerical methods for calculating magnetizabilities based on the quadrature of the magnetizability density have been implemented. We have shown that this method allows studies of spatial contributions to the magnetizabilities by visualization of the magnetizability density. The method has been employed to calculate magnetizabilities from magnetically induced current density susceptibilities, which were obtained from TURBO-MOLE calculations of nuclear magnetic shielding constants. Thus, magnetizabilities can be calculated in this way with TURBOMOLE even though analytical methods to calculate magnetizabilities as the second derivative of the energy are not yet available in this program. Further information about spatial contributions to the magnetizability could be obtained in the present approach by studying atomic contributions and investigating the positive and negative parts of the integrands separately in analogy to our recent work on nuclear magnetic shieldings in ref 53, which may be studied in the future work.

Our calculations show that the most accurate magnetizabilities (judged by the smallest MAE) for the studied database are obtained with BHandHLYP, which is an old global hybrid with 50% HF exchange and 50% B88 exchange. The

calculations also show that the modern range-separated functionals with 100% long-range HF exchange developed by Head-Gordon and co-workers and by Bartlett and co-workers yield accurate magnetizabilities for the database. Calculations with other range-separated functionals like CAM-B3LYP and CAMh-B3LYP as well as with global hybrid functionals like QTP-17, BHandHLYP a.k.a. BHandH, B3LYP, and PBE0 yield relatively accurate magnetizabilities for the studied molecules. Meta-GGA functionals are found to yield somewhat better magnetizabilities than GGA and LDA functionals.

However, functionals developed by Truhlar and co-workers do not appear to be well-aimed for calculations of magnetizabilities and other magnetic properties that involve magnetically induced current densities. Magnetizabilities calculated using the popular M06-2X functional are found to be unreliable, and we do not recommend the use of the M06-2X functional in calculations of nuclear magnetic shieldings, magnetizabilities, ring-current strengths, and other magnetic properties that depend on magnetically induced current density susceptibilities. Previous studies have also suggested that the M06-2X functional sometimes underestimates magnetizabilities and ring-current strengths.<sup>128,129,132</sup> Revised versions of Minnesota functionals have been studied in this work and found to yield somewhat more accurate magnetizabilities than the original parameterizations. However, the revised versions also still appear on the second half of the ranking list.

## ■ ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.jctc.0c01190>.

Calculated magnetizabilities for the studied molecules and functionals as well as the difference between the calculated magnetizabilities and the CCSD(T) reference values taken from ref 18; magnetizability densities calculated with the gauge origin  $\mathbf{R}_O = (1, 1, 1) a_0$ ; and additional calculations on tetraoxa isophlorin are discussed (PDF)

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### Notes

The authors declare no competing financial interest.

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