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The Role of Thermodiffusion and Dimensionality in the Formation of Cellular Instabilities in Hydrogen Flames

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Motivation

Thermodiffusive instabilities

• A common numerical approach to study thermodiffusive instabilities in laminar flame propagation until flame fingers form



• Hydrogen – air at $\phi = 0.4$ and atmospheric conditions

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Thermodiffusive instabilities

- The previous simulations were all performed in 2D
- Motivation of this work:
 - What is the impact on cell formation when simulations are performed in 3D?
 - What is the impact of thermodiffusion (Soret effect)?
- Aspects of these were studied by other authors before:
 - Theoretical work by Matalon
 - Grcar et al. (2017): <u>https://doi.org/10.1016/j.proci.2008.06.075</u>
 - Schlup & Blanquart (2017): <u>https://doi.org/10.1080/13647830.2017.1398350</u>
 - X. Wen et al. (2024): https://doi.org/10.1016/j.combustflame.2024.113497

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Numerical Setup

Numerical Setup

- Hydrogen-air at $\phi = 0.4$ and atmospheric conditions
- Laminar inflow
- Initial perturbation of the flame front
- Simulation done in 2D/3D & with/without Soret diffusion
- Mechanism by Li et al. with 9 species and 19 reactions

Table 1: Properties of reference laminar 1D flames.

ϕ	Soret	s_L^0 (m/s)	$\delta_{ m th}^{ m 0}(m mm)$	au (ms)	$Y_{\rm H_2,iso}$
0.4	no	0.239	0.599	2.510	0.00138
0.4	yes	0.227	0.613	2.703	0.00128

wave-transmissive boundary for pressure, 1 atm

Cell formation

- $t/\tau = \mathbf{0}$ $t/\tau = 3.5$ $t/\tau = 1.4$ $t/\tau = 2.1$ $t/\tau = 2.8$ 1400 1200 1000 × $t/\tau = 4.5$ $t/\tau = 5.2$ $t/\tau = 6$ $t/\tau = 7$ $t/\tau = 8.4$ 800 600
- First, linear growth regime, then decay into secondary structures (non-linear regime)

• 3D Simulation with Soret diffusion. Iso-surface of $Y_{H_{2,iso}}$ colored by temperature.

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Flame propagation

Flame propagation in 2D/3D and with/without Soret diffusion



• Faster propagation in 3D, and also faster propagation with Soret diffusion

Effect on flame structure

• In 3D, curvatures are higher, leading to a stronger focusing/defocusing of diffusion fluxes



 This leads to local heat release rates larger by up to 80% and more extreme changes in local curvatures (shown in the paper)

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2D, w/o Soret

2D, Soret

3D, Soret

Trajectories of progress variable $c = Y_{H_2O}$





Diffusive fluxes at positively curved flame segments

$$\boldsymbol{j}_k = -\rho D_k \nabla Y_k - D_k^{\mathrm{th}} \frac{\nabla T}{T}$$



Due to non-unity Lewis
 number/preferential diffusion,
 locations of species and
 temperature minima/maxima
 do not coincide



At positively curved flame segments, thermodiffusion flux of oxygen opposes that of the Fickian diffusion





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Effect on global flame properties

- For the 3D case, the flame surface area grows more rapidly
- However, the time instance at which transition from primary to secondary structures occurs, is the same as in 2D
- · Without Soret diffusion, the transition occurs later
- Higher efficiency factors in 3D and with Soret diffusion
- After secondary structures developed, quasi-steady state (also observed by other authors)

$$I = (s_T / s_L^0) / (A_t / A_0)$$

$$s_T = -\frac{1}{\rho_0 A_0 (Y_{H_2}^0 - Y_{H_2}^b)} \int \dot{\omega}_{H_2} \, \mathrm{d}V$$





Effect of Soret diffusion as function of perturbation wavelengths

 Cell growth depends on perturbation wavelengths



 Soret diffusion increases efficiency factors depending on initial curvature



Conclusions

- Even at the same cross-section in 2D, a 3D flame exhibits larger curvatures and thus heat release rates exceeding by 80% at the positively curved flame segments, resulting in faster flame propagation
- However, the time instance when the collapse happens is nearly the same
- Soret diffusion leads to lower laminar flame speeds but larger local propagation speeds due to opposed diffusive fluxes at positively curved flame segments
- Collapse into secondary cells occurs later without Soret diffusion
- Effect of Soret diffusion on efficiency factors for early cell formation is a function of perturbation wavelength.



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Thank you!



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