



Regular Article

Experimental Physics

Search for CP violation in $D^0 \rightarrow K_S^0 K_S^0$ decays in proton–proton collisions at $\sqrt{s} = 13$ TeV

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Received: 19 May 2024 / Accepted: 13 August 2024
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Abstract A search is reported for charge-parity CP violation in $D^0 \rightarrow K_S^0 K_S^0$ decays, using data collected in proton–proton collisions at $\sqrt{s} = 13$ TeV recorded by the CMS experiment in 2018. The analysis uses a dedicated data set that corresponds to an integrated luminosity of 41.6 fb^{-1} , which consists of about 10 billion events containing a pair of b hadrons, nearly all of which decay to charm hadrons. The flavor of the neutral D meson is determined by the pion charge in the reconstructed decays $D^{*+} \rightarrow D^0 \pi^+$ and $D^{*-} \rightarrow \bar{D}^0 \pi^-$. The CP asymmetry in $D^0 \rightarrow K_S^0 K_S^0$ is measured to be $A_{CP}(K_S^0 K_S^0) = (6.2 \pm 3.0 \pm 0.2 \pm 0.8)\%$, where the three uncertainties represent the statistical uncertainty, the systematic uncertainty, and the uncertainty in the measurement of the CP asymmetry in the $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ decay. This is the first CP asymmetry measurement by CMS in the charm sector as well as the first to utilize a fully hadronic final state.

1 Introduction

The noninvariance of fundamental interactions under the combined charge-parity (CP) transformation is one of the necessary conditions for the generation of the observed baryon asymmetry in the universe [1]. In the standard model (SM), the CP symmetry violation originates from a single phase in the Cabibbo–Kobayashi–Maskawa (CKM) quark mixing matrix [2,3]. Extensive studies of CP violation in weak interaction decays of strange and beauty mesons have been performed by many experiments, with all results to date being consistent with the predictions based on the CKM formalism [4]. However, the magnitude of CP violation in the SM appears to be insufficient to explain the matter–antimatter asymmetry observed in the universe [5–7], suggesting the existence of sources of CP violation beyond the SM. Charmed meson decays are the only meson decays

involving an up-type quark where CP violation can be studied, and are complementary to strange and beauty meson decays. In contrast to the K and B systems, CP violation in charm mesons is severely suppressed by the Glashow–Iliopoulos–Maiani mechanism [8] and by the magnitude of the CKM elements [3]. Given the strong SM suppression, an observation of a significant CP violation in D meson decays may indicate a contribution from new physics, which can be different from those relevant for down-type quark systems. The first observation of CP violation in charm decays was recently reported by the LHCb Collaboration in a measurement of the CP asymmetry (A_{CP}) difference between the $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ decays [9]. However, determining if this (or other) measurements of CP violation is an indication of new physics is hampered by large theoretical uncertainties associated with long-distance contributions and nonperturbative effects [10].

The $D^0 \rightarrow K_S^0 K_S^0$ decay proceeds through the W boson exchange and penguin annihilation Feynman diagrams, some examples of which are shown in Fig. 1, which results in a relatively small branching fraction of $(1.41 \pm 0.05) \times 10^{-4}$ [4]. In this figure and throughout this paper, charge-conjugate states are implied, unless otherwise indicated. Theoretical predictions indicate similar amplitudes and different phases for the two diagrams, which can result in CP violation in this channel as large as a few percent [11–15] and therefore possibly within reach of current experiments.

The CP asymmetry A_{CP} , for the $D^0 \rightarrow K_S^0 K_S^0$ decay, is defined as

$$A_{CP}(K_S^0 K_S^0) = \frac{\Gamma(D^0 \rightarrow K_S^0 K_S^0) - \Gamma(\bar{D}^0 \rightarrow K_S^0 K_S^0)}{\Gamma(D^0 \rightarrow K_S^0 K_S^0) + \Gamma(\bar{D}^0 \rightarrow K_S^0 K_S^0)}. \quad (1)$$

The current world average for the time-integrated CP asymmetry is $A_{CP}(K_S^0 K_S^0) = (-1.9 \pm 1.1)\%$ [4], which is dominated by results from the LHCb [16] and Belle [17] Collaborations.

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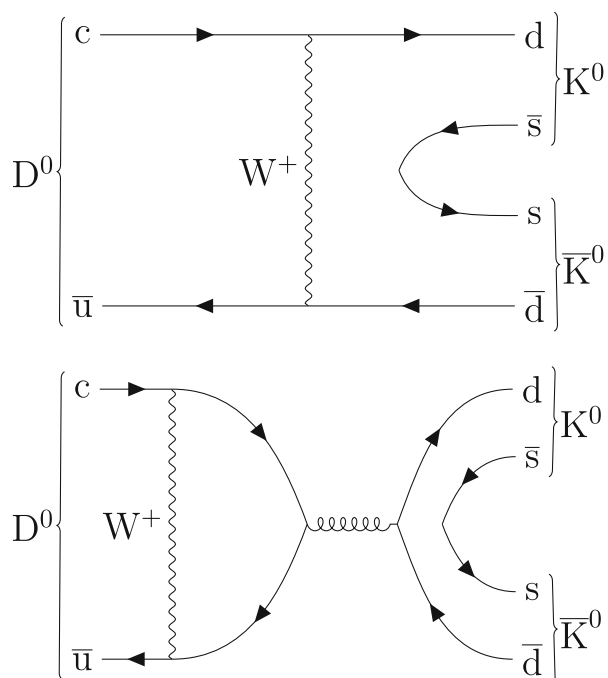


Fig. 1 The decay of neutral charm meson to two neutral kaons: exchange (upper) and penguin annihilation (lower) diagrams

This paper presents the first CP violation measurement by the CMS experiment in the charm sector. The flavor of the neutral D meson is determined from the pion charge found from reconstructing the decays $D^{*+} \rightarrow D^0 \pi^+$ and $D^{*-} \rightarrow \bar{D}^0 \pi^-$. We measure the CP asymmetry difference, ΔA_{CP} , between the signal channel $D^0 \rightarrow K_S^0 K_S^0$ and the reference channel $D^0 \rightarrow K_S^0 \pi^+ \pi^-$. The $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ CP asymmetry has been previously measured [18] and found to be consistent with zero, as expected since this decay is not CKM-suppressed. Therefore, a significant deviation of ΔA_{CP} from 0 would indicate CP violation in the $D^0 \rightarrow K_S^0 K_S^0$ decay.

In proton–proton (pp) collision data, the number of D^{*+} and D^{*-} decays (signal events, N) are measured, where both D^0 and $\bar{D}^0$ are reconstructed in the $K_S^0 K_S^0$ or $K_S^0 \pi^+ \pi^-$ decay modes. The “raw” asymmetry between these numbers, A_{CP}^{raw} (defined in Eq. (2)), is different from the true asymmetry A_{CP} , due to the slightly different production cross sections (σ) of D^{*+} and D^{*-} mesons, as well as to a possible difference in the detection efficiency (ϵ) between D^{*+} and D^{*-} . Because these three asymmetries are all small, the following relation can be written:

$$A_{CP} \approx A_{CP}^{\text{raw}} - A_{CP}^{\text{pro}} - A_{CP}^{\text{det}}, \quad \text{where}$$

$$A_{CP}^{\text{raw}} = \frac{N(D^{*+} \rightarrow D^0 \pi^+) - N(D^{*-} \rightarrow \bar{D}^0 \pi^-)}{N(D^{*+} \rightarrow D^0 \pi^+) + N(D^{*-} \rightarrow \bar{D}^0 \pi^-)},$$

$$A_{CP}^{\text{pro}} = \frac{\sigma_{pp \rightarrow D^{*+} X} - \sigma_{pp \rightarrow D^{*-} X}}{\sigma_{pp \rightarrow D^{*+} X} + \sigma_{pp \rightarrow D^{*-} X}}, \quad \text{and}$$

$$A_{CP}^{\text{det}} = \frac{\epsilon(D^{*+} \rightarrow D^0 \pi^+) - \epsilon(D^{*-} \rightarrow \bar{D}^0 \pi^-)}{\epsilon(D^{*+} \rightarrow D^0 \pi^+) + \epsilon(D^{*-} \rightarrow \bar{D}^0 \pi^-)}, \quad (2)$$

where measuring the difference of A_{CP} between the signal and reference channels, A_{CP}^{pro} ($D^{\pm}$ production asymmetry) and A_{CP}^{det} ($D^{\pm}$ detection asymmetry) cancel out, as they do not depend on the final state ($K_S^0 K_S^0$ or $K_S^0 \pi^+ \pi^-$):

$$\Delta A_{CP} \equiv A_{CP}(K_S^0 K_S^0) - A_{CP}(K_S^0 \pi^+ \pi^-)$$

$$= A_{CP}^{\text{raw}}(K_S^0 K_S^0) - A_{CP}^{\text{raw}}(K_S^0 \pi^+ \pi^-). \quad (3)$$

The reference channel was chosen to be as similar as possible in kinematics, topology, and final-state signature to the signal channel, to ensure that the reconstruction efficiency asymmetries cancel in the measured difference of asymmetries, ΔA_{CP} .

The analysis uses proton–proton collisions data recorded by the CMS detector during the CERN LHC Run 2 in 2018, at $\sqrt{s} = 13$ TeV. It utilizes the B parking data set [19,20], collected with a set of single-muon triggers with different minimum thresholds on the muon transverse momentum (p_T) and impact parameter with respect to the beamline. Different triggers were enabled depending on the instantaneous luminosity: as the luminosity decreased, less restrictive triggers were enabled, as allowed by the limited event rate to be processed by the data acquisition system and recorded on tape. The data set contains about 1.2×10^{10} events and corresponds to an integrated luminosity of 41.6 fb^{-1} . More details about this data set can be found in Ref. [19]. These triggers are intended to select events containing a semimuonic decay of a b hadron (or a semimuonic decay of c hadron that originated from a b-hadron decay). Since the trigger requires muons inconsistent with being produced in the primary interaction, most of such muons come from semileptonic decays of beauty hadrons, hence approximately 80% of the events in this sample include b hadrons [19,20]. As beauty hadrons nearly always decay into charm hadrons, this data set also provides a rich sample of charm decays, making it suitable for CP violation studies in the charm sector.

2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter, each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity (η) coverage provided by the barrel and endcap detectors. Muons are measured in gas-ionization detectors embedded in the steel flux-return yoke outside the solenoid. The reconstructed decays used by this analysis contain five pions in the final state. Pions are measured by the silicon tracker whose setup during the 2018 LHC running period, when the data used in this paper were recorded, consisted of 1856 silicon pixel [21] and 15 148 silicon strip detector modules. For non-isolated particles with $|\eta| < 3$ and $1 < p_T < 10$ GeV, the track resolutions are typically 1.5% in p_T and 20–75 μm in the transverse impact parameter [22].

Events of interest are selected using a two-tiered trigger system. The first level, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a fixed latency of 4 μs [23]. The second level, known as the high-level trigger, consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and further reduces the event rate before data storage [24].

A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [25].

3 Simulated event samples

The simulated event samples used in this analysis are generated with PYTHIA 8.230 [26]. The PYTHIA output is interfaced with EVTGEN [27] 1.3.0, which simulates various b and c hadron decays. The underlying event is also modeled with PYTHIA using the CP5 [28] tune. Final-state photon radiation is modeled with PHOTOS 3.61 [29]. Samples with inclusive decays $B^+ \rightarrow D^{*\pm}(\rightarrow D\pi^\pm)X$, $B^0 \rightarrow D^{*\pm}(\rightarrow D\pi^\pm)X$, and prompt $D^{*\pm} \rightarrow D\pi^\pm$ were generated. The events were then passed through a detailed GEANT4-based simulation [30] of the CMS detector, followed by the trigger and reconstruction algorithms identical to those used for the collision data.

4 Reconstruction of charm meson decays

The reconstruction starts with finding $K_S^0 \rightarrow \pi^+\pi^-$ candidates as described in Ref. [31]. The two oppositely-charged

pion tracks are fit to a common vertex that is required to have a χ^2 fit probability $P_{\text{vtx}} > 1\%$. The dipion invariant mass must be within 20 MeV of the world average value of the K_S^0 meson mass [4], corresponding to approximately three times the mass resolution.

In the signal channel, two $K_S^0 \rightarrow \pi^+\pi^-$ candidates are each fit again with kinematic constraints to the K_S^0 mass, and subsequently, the K_S^0 candidates are fitted as two virtual tracks to a common vertex, assumed to be the $D^0 \rightarrow K_S^0 K_S^0$ decay vertex. The $K_S^0 K_S^0$ invariant mass is required to be between 1.77 and 1.95 GeV, and the vertex fit probability must exceed 1%. Both K_S^0 decay vertices have to be displaced in three-dimensional (3D) space by at least one standard deviation (s.d.) from the fitted $K_S^0 K_S^0$ vertex, and the corresponding pointing angle (the angle between the particle momentum and the vector joining the production vertex with its decay vertex) for each K_S^0 candidate is required to be less than 90° .

In the reference channel, after the single K_S^0 selection, two additional high-purity [32] and opposite-sign tracks with $p_T > 0.6$ GeV (and at least one of them with $p_T > 0.7$ GeV) are selected. The $K_S^0 \pi^+ \pi^-$ combination is then fit to a common vertex, assumed to be the D^0 decay vertex, which must have a $P_{\text{vtx}} > 5\%$, and an invariant mass between 1.823 and 1.908 GeV, assuming charged-pion mass [4] for both tracks, which corresponds to approximately twice the mass resolution.

The primary vertex (PV) is selected from the reconstructed pp interaction vertices as the one with the smallest pointing angle of the D^0 candidate. After the D^0 reconstruction, an additional track is added to form the $D^{*+} \rightarrow D^0 \pi^+$ or $D^{*-} \rightarrow \bar{D}^0 \pi^-$ candidates. A two-object vertex fit is performed to reconstruct the $D^{*\pm}$ decay vertex, which is required to have a $P_{\text{vtx}} > 1\%$. The D^{*+} candidate invariant mass is determined from the refitted pion and D^0 four-momenta and then corrected by subtracting the difference between the reconstructed D^0 candidate mass and the world-average D^0 mass, to remove the effect of the D^0 detector mass resolution. The candidates are rejected if they are compatible with an incorrect decay topology that assumes negligible decay time of any of the K_S^0 candidates.

5 Final selection criteria

A mixture of different triggers with varying thresholds in the data set makes it challenging to properly model the kinematic distributions of charm mesons in the simulation. Therefore, an optimization of the selection criteria is done using the experimental data directly. A two-dimensional (2D) fit to the distribution of $m(D\pi^\pm)$ vs $m(K_S^0 K_S^0)$, similar to the one

Table 1 Optimized selection criteria in the signal channel $K_S^0 K_S^0$. The requirements on the K_S^0 candidates in the third and fourth lines are given first for the K_S^0 with larger p_T , then for the K_S^0 with lower p_T

Variable	Requirement
$ \eta $ of the tagging pion from $D^{*\pm} \rightarrow D\pi^\pm$	<1.2
p_T of the tagging pion from $D^{*\pm} \rightarrow D\pi^\pm$	>0.35 GeV
$p_T(K_S^0)$	>2.2 and >1.0 GeV
K_S^0 vertex displacement significance from the D^0 vertex in xyz	>7 and >9
D^0 vertex displacement significance from the PV in xy	>2
D^0 vertex displacement significance from the PV in xyz	>9
$P_{\text{vtx}}(D\pi^\pm)$	$>5\%$
$P_{\text{vtx}}(K_S^0 K_S^0)$	$>1\%$
$P_{\text{vtx}}(\pi^+ \pi^-)$ for $K_S^0 \rightarrow \pi^+ \pi^-$	$>1\%$
Angle between D^0 momentum and displacement from PV in xyz	<0.205 rad
Angle between D^0 momentum and displacement from PV in xy	<0.237 rad
Angle between D^0 momentum and displacement from beamline in xy	<0.237 rad

described below, is performed for the data (with the D^{*+} and D^{*-} samples merged) while the selection criteria are varied. The variables used in the optimization include the candidate p_T , η , P_{vtx} , distances between production and decay vertices for K_S^0 and D^0 candidates divided by their corresponding uncertainties, and corresponding pointing angles. The optimal criteria were chosen as those which result in the smallest relative uncertainty on the fitted signal yield. Cross-validation was used to ensure there is no bias due to statistical fluctuations in the data, via randomly splitting the data into six equal sub-samples, finding optimal criteria using five of them and applying them to the last part. The procedure is repeated six times (each time leaving out a different part of the full data set) and results in six almost identical sets of selection criteria. The average value for each selection is taken as the final selection criteria, presented in Table 1.

Similar selection criteria are applied to the reference channel, to minimize the differences in kinematic distributions between the signal and the reference channels: the only adjustment is that the scalar sum of the p_T of the two pions that are not from the K_S^0 decay in the reference channel must exceed 1 GeV and the single K_S^0 candidate in the reference channel must satisfy the requirements applied to the high- p_T K_S^0 candidate in the signal channel.

6 A_{CP} measurement: reference channel

The signal and reference channels are found to have consistent η and ϕ distributions, but slightly different $p_T(D^{*\pm})$ ones, and thus the detection and production asymmetries may not cancel out fully in the ΔA_{CP} measurement. In order to suppress this effect, the reference channel data are reweighted

to match the $p_T(D^{*\pm})$ distribution found in the signal channel, before splitting the samples by the pion charge.

To extract the raw CP asymmetry, a simultaneous binned extended maximum likelihood fit is performed on the invariant mass distributions $m(D\pi^\pm)$ of weighted D^{*+} and D^{*-} candidates. The signal in $x = m(D\pi^\pm)$ is fitted with the S_U Johnson transformation of the normal distribution [33], with the shape parameters shared between the D^{*+} and D^{*-} components while the signal yields are independent. The background is modeled with a modified threshold function $(x - x_0)^\alpha (1 + ax)$, where x_0 is the threshold value equal to the sum of the masses of D^0 and π^+ , and α and a are floated in the fit and they are not shared between the D^{*+} and the D^{*-} background model. The results of the fit to the $m(D\pi^\pm)$ distributions are presented in Fig. 2 and Table 2. The measured raw asymmetry is $A_{CP}^{\text{raw}}(K_S^0 \pi^+ \pi^-) = (0.78 \pm 0.10)\%$, where the uncertainty is statistical only and accounts for the correlations found in the simultaneous fit.

7 A_{CP} measurement: signal channel

To reduce the statistical uncertainty arising from the signal channel yield, which is the dominant uncertainty in the analysis, the signal extraction is performed using a 2D unbinned maximum likelihood fit performed simultaneously on the D^{*+} and D^{*-} samples to the distribution of $m(D\pi^\pm)$ vs. $m(K_S^0 K_S^0)$. The fit function consists of the following components:

- $D^0 \times D^{*+}$, the signal component;
- $D^0 \times bkg$, for events containing genuine D^0 and background pion combinations;
- $bkg \times bkg$, for the background in both dimensions,

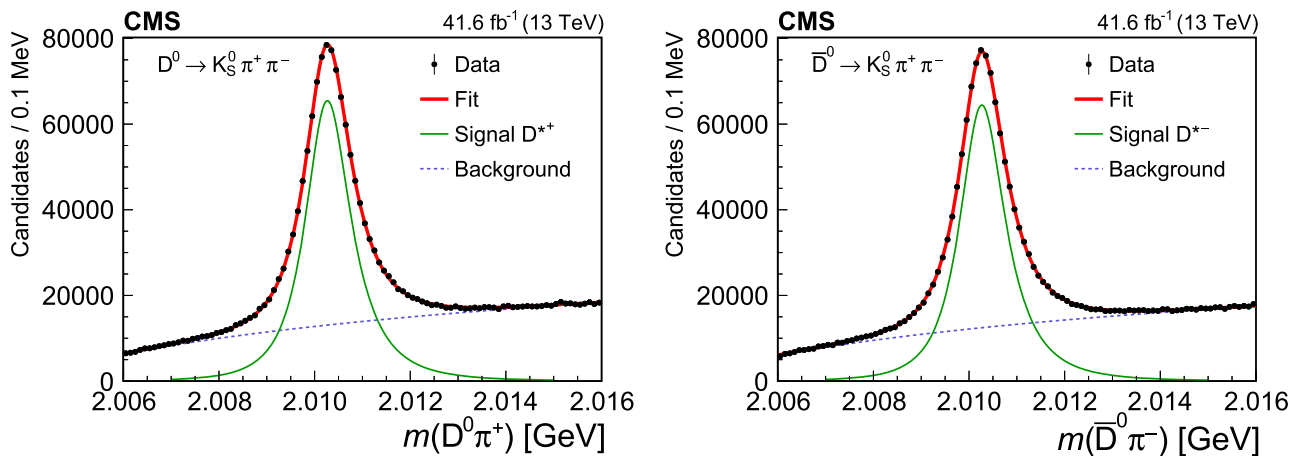


Fig. 2 The $D^0 \pi^+$ (left) and $\bar{D}^0 \pi^-$ (right) invariant mass distributions for the $K_S^0 \pi^+ \pi^-$ channel, with the result of the fit to both distributions

Table 2 Results of the fit to the selected $D^{*+} \rightarrow D^0 \pi^+$ and $D^{*-} \rightarrow \bar{D}^0 \pi^-$ candidates, where $D^0 (\bar{D}^0) \rightarrow K_S^0 \pi^+ \pi^-$. The $D^{*\pm}$ signal yields N given in the second column are used in the evaluation of A_{CP}^{raw} . The uncertainties are statistical only

Decay	N	χ^2 with 100 bins
$D^{*+} \rightarrow D^0 \pi^+$	$944\,800 \pm 3\,500$	78
$D^{*-} \rightarrow \bar{D}^0 \pi^-$	$930\,150 \pm 3\,400$	93

where each component is a product of two one-dimensional (1D) functions. For the $D^{*\pm}$ signal, the Johnson function is used with all signal shape parameters fixed to those found in the fit to the reference channel. This approach is verified to be reasonable using simulated event samples. The D^0 signal is modeled with a sum of two Johnson functions, all parameters of which are fixed to values determined from the simulation, except for a single free parameter that is used to scale the width. The background in $x = m(D\pi^\pm)$ is modeled with the same function as in the reference channel. The background in $y = m(K_S^0 K_S^0)$ is described with an exponential function $\exp(\beta y)$, where β is floating in the fit, plus a Gaussian function with free parameters to describe the partially-reconstructed background from the $D_s^\pm \rightarrow K_S^0 K_S^0 \pi^\pm$ decay producing an excess at about 1.83 GeV in the $K_S^0 K_S^0$ invariant mass distribution.

The projections of the data and the 2D fit on both axes are shown in Fig. 3; additional projections in sub-ranges are shown in Appendix A. The fit results are listed in Table 3. The measured raw asymmetry is $A_{CP}^{\text{raw}}(K_S^0 K_S^0) = (7.1 \pm 3.0)\%$ and in combination with the results of Sect. 6 the A_{CP} difference is measured to be $\Delta A_{CP} = (6.3 \pm 3.0)\%$, where the uncertainty is statistical only and accounts for the correlations found in the simultaneous fit.

8 Systematic uncertainties

The measured difference in the asymmetries is largely insensitive to many systematic uncertainties that would affect a measurement of A_{CP} in a single channel, such as the difficult-to-measure production and detection asymmetries that would need a dedicated calibration procedure.

Uncertainties related to the choice of the signal and background models are calculated separately using alternative models and assessing the observed variations in ΔA_{CP} .

In the baseline approach, the signal in the $m(D\pi^\pm)$ invariant mass distribution is modeled with the Johnson function [33]. As an alternative, we use a Johnson+Gaussian function with a common mean. Another alternative is a sum of two Crystal Ball functions [34]. For each case, the reference channel is fit as a first step, then the obtained shape parameters are fixed in the 2D fit to the signal channel. Other components of the 2D fit remain unchanged from the baseline fit. The largest deviation in ΔA_{CP} from the baseline value is taken as a systematic uncertainty.

The baseline signal function for the $m(K_S^0 K_S^0)$ invariant mass distribution is a sum of two Johnson functions. As an alternative, we use a Johnson+Gaussian function or a sum of two Crystal Ball functions. These variations have no effect on the fit of the reference channel, just on that of the signal channel. The largest deviation in ΔA_{CP} from the baseline value is taken as a systematic uncertainty.

The baseline background model in the $x = m(D\pi^\pm)$ distribution is $(x - x_0)^\alpha (1 + ax)$. An alternative background model is obtained by changing the function multiplying the threshold function from a linear polynomial to an exponential function. The baseline background model in the $m(K_S^0 K_S^0)$ distribution is an exponential function and an exponential multiplied by a linear polynomial is used as an alternative.

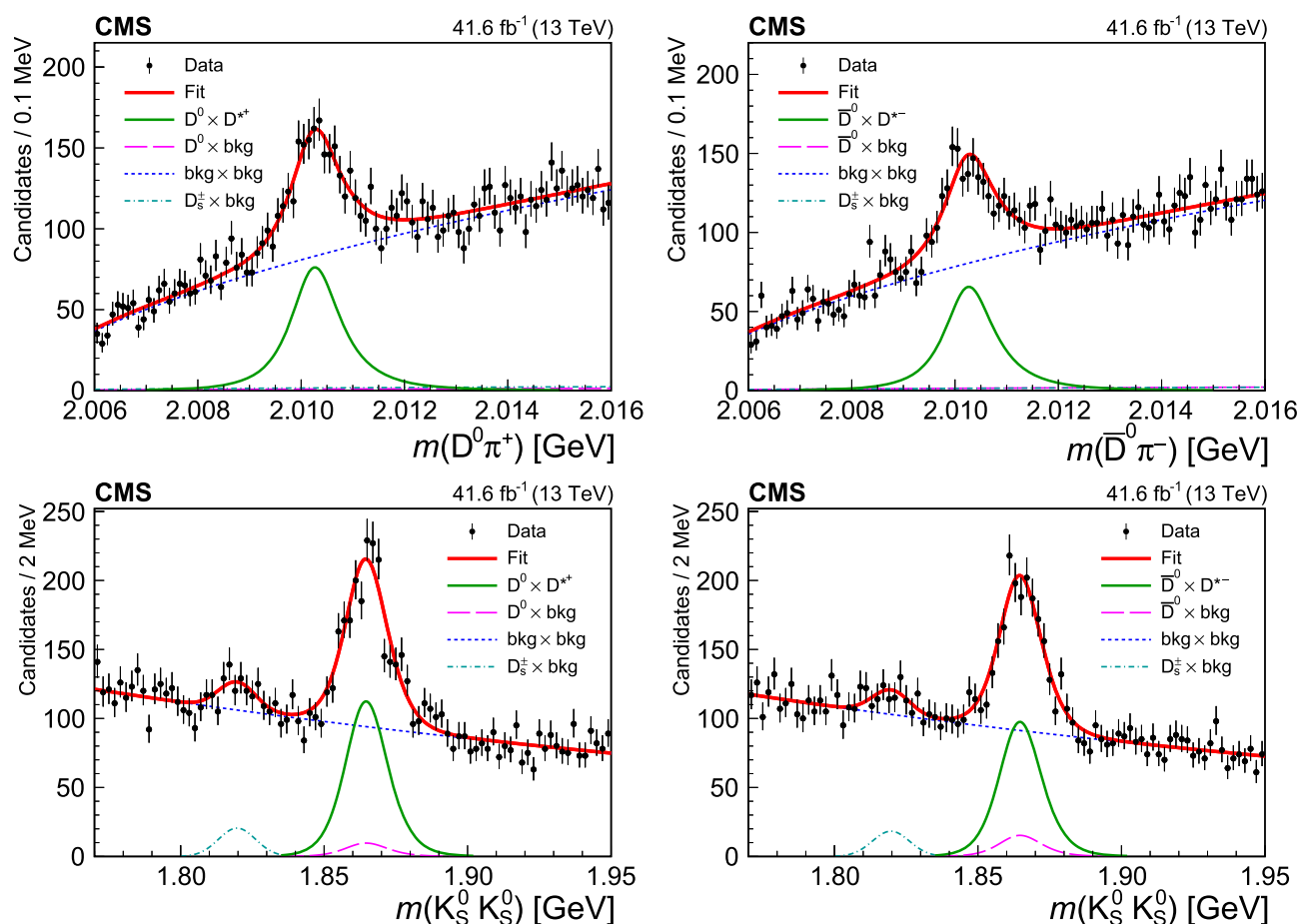


Fig. 3 The invariant mass distributions for D^{*+} candidates (left) and D^{*-} candidates (right), with the $m(D\pi^{\pm})$ distributions in the upper row and the $m(K_S^0 K_S^0)$ distributions in the lower row. Projections of the simultaneous 2D fit are also shown

Table 3 Results of the 2D fit to the selected $D^{*+} \rightarrow D^0 \pi^+$ and $D^{*-} \rightarrow \bar{D}^0 \pi^-$ candidates, where $D^0 (\bar{D}^0) \rightarrow K_S^0 K_S^0$. The $D^{*\pm}$ signal yields N given in the second column are used in the evaluation of A_{CP}^{raw} . The χ^2 corresponds to the fit projection with 100 bins in the $x = m(D\pi^{\pm})$ axis and 90 bins in the $y = m(K_S^0 K_S^0)$ axis, as shown in Fig. 3. The uncertainties are statistical only

Decay	N	χ^2 (x axis)	χ^2 (y axis)
$D^{*+} \rightarrow D^0 \pi^+$	1095 ± 46	77	90
$D^{*-} \rightarrow \bar{D}^0 \pi^-$	951 ± 44	93	62

These variations are taken as independent systematic uncertainties.

In the signal channel fit, there is a contribution from the $D_s^{\pm} \rightarrow K_S^0 K_S^0 \pi^{\pm}$ decay, which is modeled by a Gaussian with free parameters. As an alternative, we remove this reflection by restricting the fit range to be $m(K_S^0 K_S^0) > 1.835$ GeV, and the deviation from the baseline is included as a systematic uncertainty.

To assess the systematic uncertainty related to the p_T reweighting, we vary the parameters of the reweighting func-

tion within their uncertainties. As an alternative, we consider the weights depending on the p_T of the low-momentum pion that is used for the flavor tagging instead of $p_T(D^{*\pm})$. The largest change is taken as a systematic uncertainty related to the reweighting.

Differences in A_{CP}^{pro} and A_{CP}^{det} between the two channels are expected to be reproduced by the simulation of the processes and the detector. Checking the reweighted reference channel and signal channel in simulation show that the p_T -, η -, and ϕ -dependent asymmetries are consistent with zero as is the integrated value of $(-0.13 \pm 0.34)\%$. Therefore, no systematic uncertainty is assessed.

If multiple candidates in the same event are removed by keeping only the one with the highest $D^{*\pm}$ vertex fit probability, the resulting ΔA_{CP} changes negligibly and no corresponding systematic uncertainty is assigned. Pion charge misidentification was shown to have a negligible effect as well.

All systematic uncertainties described above are uncorrelated and summarized in Table 4 together with the total

Table 4 Absolute systematic uncertainties in the measurement of ΔA_{CP}

Source	Uncertainty (%)
$m(D\pi^\pm)$ signal model	0.10
$m(D\pi^\pm)$ background model	0.02
$m(K_S^0 K_S^0)$ signal model	0.04
$m(K_S^0 K_S^0)$ background model	0.02
$m(K_S^0 K_S^0)$ fit range	0.06
Reweighting	0.09
Total	0.16

systematic uncertainty, calculated as the sum in quadrature of the different contributions.

9 Summary

A measurement of CP violation in D^0 decays is reported, using proton–proton collision data collected at $\sqrt{s} = 13$ TeV with a novel high-rate data stream (B parking). These data correspond to an integrated luminosity of 41.6 fb^{-1} and include about 10 billion events containing beauty hadron decays. The difference in the CP asymmetries between $D^0 \rightarrow K_S^0 K_S^0$ and $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ is measured to be:

$$\begin{aligned} \Delta A_{CP} &\equiv A_{CP}(K_S^0 K_S^0) - A_{CP}(K_S^0 \pi^+ \pi^-) \\ &= (6.3 \pm 3.0 (\text{stat}) \pm 0.2 (\text{syst})) \%. \end{aligned} \quad (4)$$

Using the world-average value of $A_{CP}(K_S^0 \pi^+ \pi^-) = (-0.1 \pm 0.8)\%$ [4, 18, 35], we report the measurement

$$A_{CP}(K_S^0 K_S^0) = (6.2 \pm 3.0 \pm 0.2 \pm 0.8)\%, \quad (5)$$

where the three uncertainties represent the statistical uncertainty, the systematic uncertainty, and the uncertainty in the measurement of the CP asymmetry in the $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ decay. The measured value is consistent with no CP violation within 2.0 standard deviations. Likewise, it is consistent with the LHCb [16] and the Belle measurements [17] at the level of 2.7 and 1.8 standard deviations, respectively. Tabulated results are provided in the HEPData record for this analysis [36]. This is the first CMS search for CP violation in the charm sector, paving the way for future measurements with more data, using new techniques, and in other channels.

Acknowledgements We congratulate our colleagues in the CERN accelerator departments for the excellent performance of the LHC and thank the technical and administrative staffs at CERN and at other CMS institutes for their contributions to the success of the CMS effort. In addition, we gratefully acknowledge the computing centers

and personnel of the Worldwide LHC Computing Grid and other centers for delivering so effectively the computing infrastructure essential to our analyses. Finally, we acknowledge the enduring support for the construction and operation of the LHC, the CMS detector, and the supporting computing infrastructure provided by the following funding agencies: SC (Armenia), BMBWF and FWF (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ, FAPERGS, and FAPESP (Brazil); MES and BNSF (Bulgaria); CERN; CAS, MoST, and NSFC (China); Minciencias (Colombia); MSES and CSF (Croatia); RIF (Cyprus); SENESCYT (Ecuador); ERC PRG, RVTT3 and MoER TK202 (Estonia); Academy of Finland, MEC, and HIP (Finland); CEA and CNRS/IN2P3 (France); SRNSF (Georgia); BMBF, DFG, and HGF (Germany); GSRI (Greece); NKFIH (Hungary); DAE and DST (India); IPM (Iran); SFI (Ireland); INFN (Italy); MSIP and NRF (Republic of Korea); MES (Latvia); LMTLT (Lithuania); MOE and UM (Malaysia); BUAP, CINVESTAV, CONACYT, LNS, SEP, and UASLP-FAI (Mexico); MOS (Montenegro); MBIE (New Zealand); PAEC (Pakistan); MES and NSC (Poland); FCT (Portugal); MESTD (Serbia); MCIN/AEI and PCTI (Spain); MOSTR (Sri Lanka); Swiss Funding Agencies (Switzerland); MST (Taipei); MHEI and NSTDA (Thailand); TUBITAK and TENMAK (Turkey); NASU (Ukraine); STFC (United Kingdom); DOE and NSF (USA). Individuals have received support from the Marie-Curie program and the European Research Council and Horizon 2020 Grant, contract Nos. 675440, 724704, 752730, 758316, 765710, 824093, 101115353, 101002207, and COST Action CA16108 (European Union); the Leventis Foundation; the Alfred P. Sloan Foundation; the Alexander von Humboldt Foundation; the Science Committee, project no. 22r1-037 (Armenia); the Belgian Federal Science Policy Office; the Fonds pour la Formation à la Recherche dans l'Industrie et dans l'Agriculture (FRIA-Belgium); the Agentschap voor Innovatie door Wetenschap en Technologie (IWT-Belgium); the F.R.S.-FNRS and FWO (Belgium) under the “Excellence of Science – EOS” – be.h project n. 30820817; the Beijing Municipal Science & Technology Commission, No. Z191100007219010 and Fundamental Research Funds for the Central Universities (China); the Ministry of Education, Youth and Sports (MEYS) of the Czech Republic; the Shota Rustaveli National Science Foundation, grant FR-22-985 (Georgia); the Deutsche Forschungsgemeinschaft (DFG), under Germany's Excellence Strategy – EXC 2121 “Quantum Universe” – 390833306, and under project number 400140256 – GRK2497; the Hellenic Foundation for Research and Innovation (HFRI), Project Number 2288 (Greece); the Hungarian Academy of Sciences, the New National Excellence Program – UNKP, the NKFIH research grants K 131991, K 133046, K 138136, K 143460, K 143477, K 146913, K 146914, K 147048, 2020-2.2.1-ED-2021-00181, and TKP2021-NKTA-64 (Hungary); the Council of Science and Industrial Research, India; ICSC – National Research Center for High Performance Computing, Big Data and Quantum Computing and FAIR – Future Artificial Intelligence Research, funded by the NextGenerationEU program (Italy); the Latvian Council of Science; the Ministry of Education and Science, project no. 2022/WK/14, and the National Science Center, contracts Opus 2021/41/B/ST2/01369 and 2021/43/B/ST2/01552 (Poland); the Fundação para a Ciência e a Tecnologia, grant CEECIND/01334/2018 (Portugal); the National Priorities Research Program by Qatar National Research Fund; MCIN/AEI/10.13039/501100011033, ERDF “a way of making Europe”, and the Programa Estatal de Fomento de la Investigación Científica y Técnica de Excelencia María de Maeztu, grant MDM-2017-0765 and Programa Severo Ochoa del Principado de Asturias (Spain); the Chulalongkorn Academic into Its 2nd Century Project Advancement Project, and the National Science, Research and Innovation Fund via the Program Management Unit for Human Resources & Institutional Development, Research and Innovation, grant B37G660013 (Thailand); the Kavli Foundation; the Nvidia Corporation; the SuperMicro Corporation; the Welch Foundation, contract C-1845; and the Weston Havens Foundation (USA).

Data Availability Statement This manuscript has no associated data. [Authors' comment: Release and preservation of data used by the CMS Collaboration as the basis for publications is guided by the [CMS data preservation, re-use, and open access policy](#).]

Code Availability Statement The CMS core software is publicly available on GitHub (<https://github.com/cms-sw/cmssw>).

Declarations

Conflict of interest The authors declare that they have no Conflict of interest.

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Funded by SCOAP³.

Appendix A: Additional projections of the 2D fit

Figures 4 and 5 show the projections of the 2D fit in the signal channel in subranges of the mass variables. The top and middle rows show 1D projections of the 2D fit on

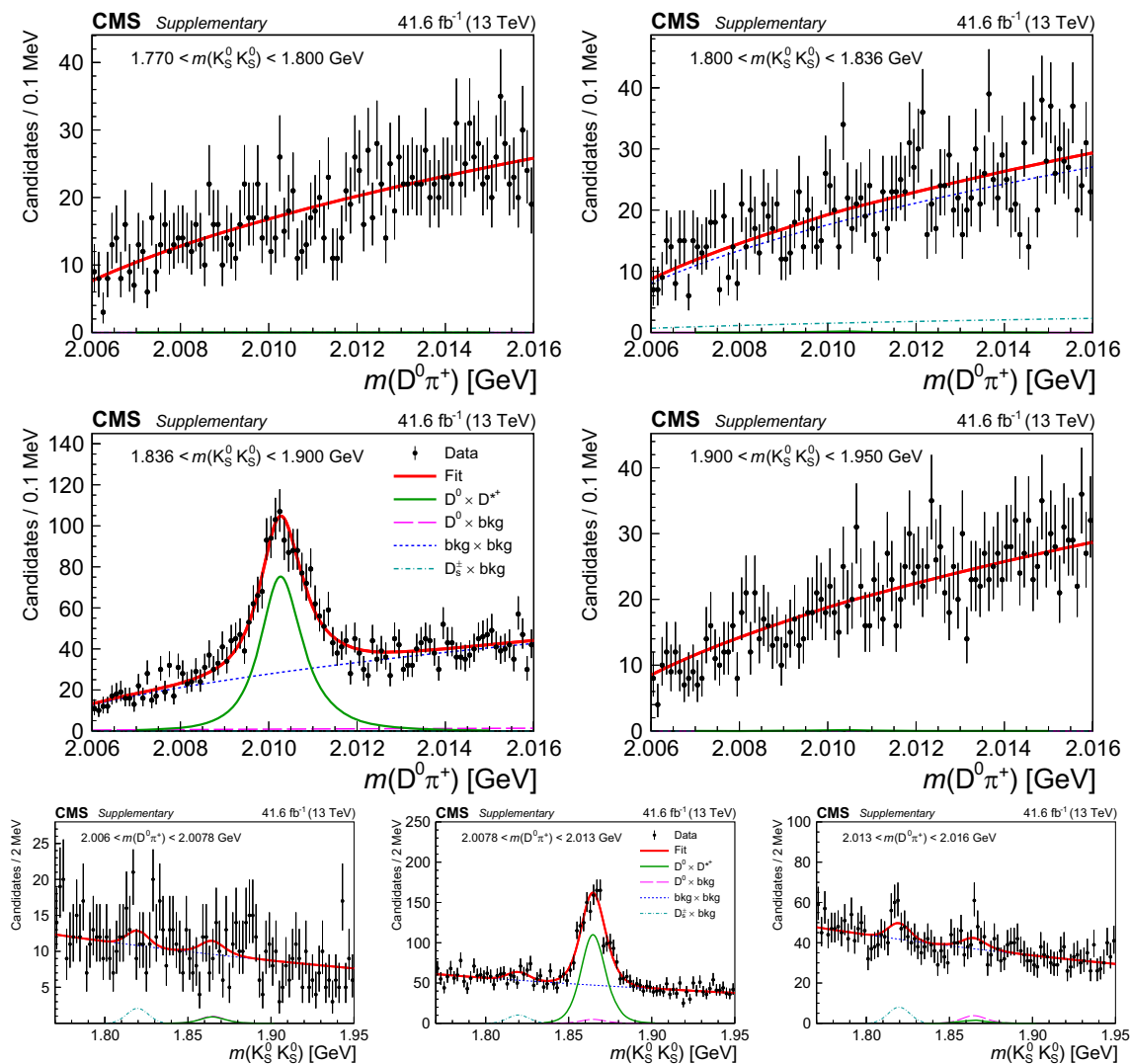


Fig. 4 Results of the 2D fit to the $m(D\pi^\pm) \times m(K_S^0 K_S^0)$ for the signal channel, D^{*+} candidates. Upper and middle rows show 1D projections of the 2D fit on $m(D^0 \pi^+)$ in ranges of $m(K_S^0 K_S^0)$: left sideband (upper left), region of $D_S^\pm \rightarrow K_S^0 K_S^0 \pi^\pm$ contamination (upper right),

signal region of $K_S^0 K_S^0$ (middle left), and right sideband (middle right). Lower row shows 1D projections of the 2D fit on $m(K_S^0 K_S^0)$ in ranges of $m(D^0 \pi^+)$: left sideband (left), signal region of $D^0 \pi^+$ (center), and right sideband (right)

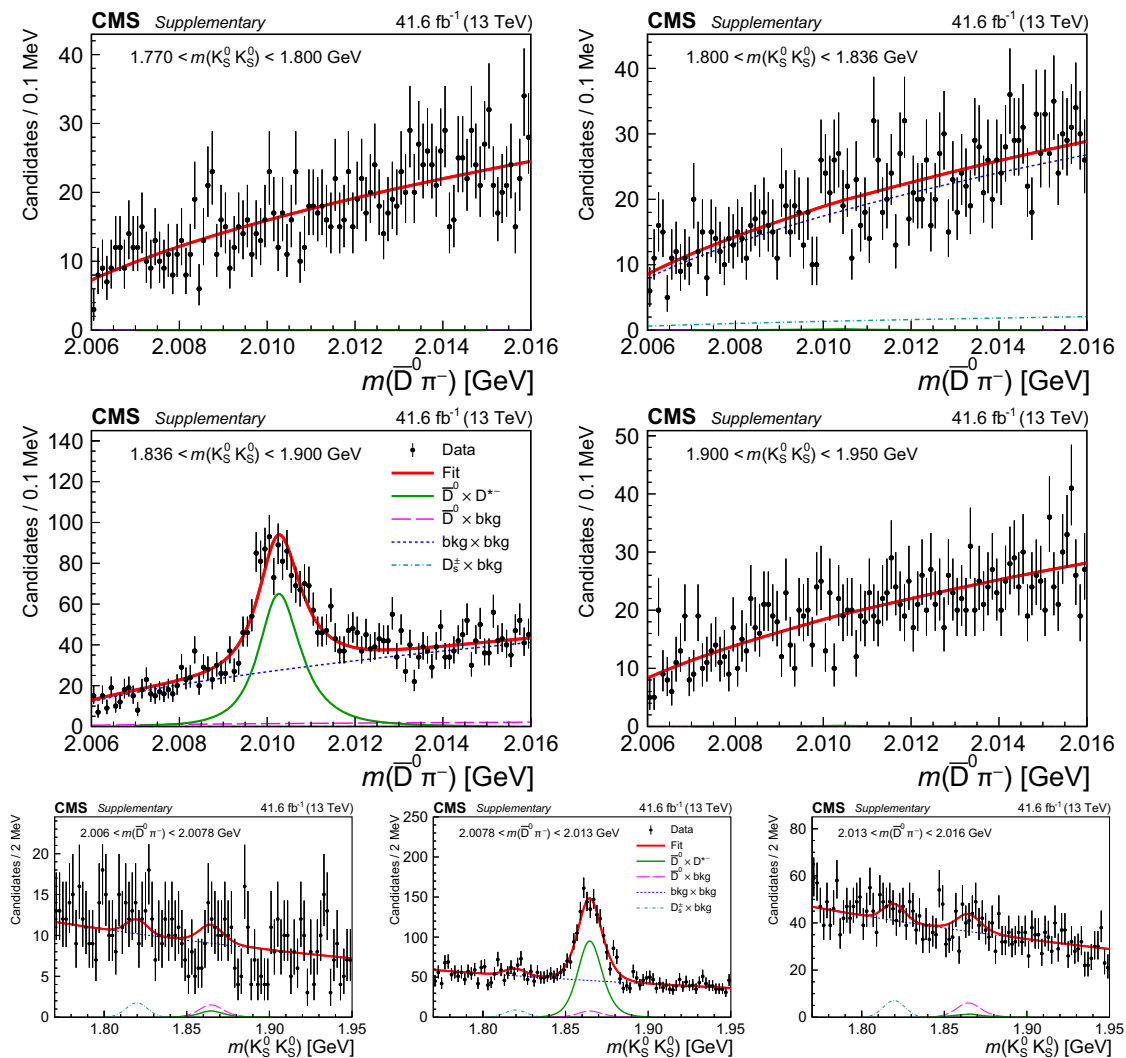


Fig. 5 Results of the 2D fit to the $m(D\pi^\pm) \times m(K_S^0 K_S^0)$ for the signal channel, $D^{*\pm}$ candidates. Upper and middle rows show 1D projections of the 2D fit on $m(\bar{D}^0 \pi^-)$ in ranges of $m(K_S^0 K_S^0)$: left sideband (upper left), region of $D_S^\pm \rightarrow K_S^0 K_S^0 \pi^\pm$ contamination (upper right),

signal region of $K_S^0 K_S^0$ (middle left), and right sideband (middle right). Lower row shows 1D projections of the 2D fit on $m(K_S^0 K_S^0)$ in ranges of $m(\bar{D}^0 \pi^-)$: left sideband (left), signal region of $\bar{D}^0 \pi^-$ (center), and right sideband (right)

$m(D\pi^\pm)$ in ranges of $m(K_S^0 K_S^0)$: left sideband, region of $D_S^\pm \rightarrow K_S^0 K_S^0 \pi^\pm$ contamination, signal region of $K_S^0 K_S^0$, and right sideband. The lower three plots show 1D projections of the 2D fit on $m(K_S^0 K_S^0)$ in ranges of $m(D\pi^\pm)$: left sideband, signal region of $D\pi^\pm$, and right sideband.

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