

Numerical simulation of a RF – RF hybrid plasma torch with argon at atmospheric pressure

Evaluation of minimum sustaining current (MSI) in dual frequency (ICP) torch

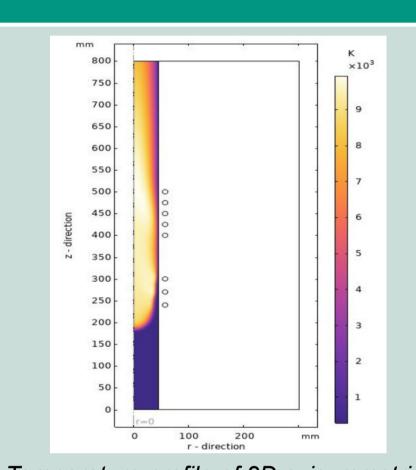
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Overview

- Inductively coupled plasma (ICP) at atmospheric pressure is an alternative to the classic gasbased combustion in the context of industrial heating applications
- Working with high-frequency electronics is feasible for higher plasma power coupling, but not at a desirable cost
- We investigate a hybrid approach, in which the high frequency (HF) coil ignites the plasma at a relatively low power, while the medium frequency (MF) coil supplies the remaining power in a cost-effective manner
- We present how the MSI of the MF coil changes as we vary the distance between the coils and the amount of power supplied to the HF coil

Implementation and Simplification

- Pure Argon gas at 1 atmosphere
- Local thermodynamics equilibrium (LTE) plasma assumed
- No kinetic reaction implemented
- Electrically neutral plasma
- 2D-axisymmetric simulation
- Radiation is not considered for the heat transfer
- The burner wall is modeled as a fixed temperature boundary condition at 293.15 *K*

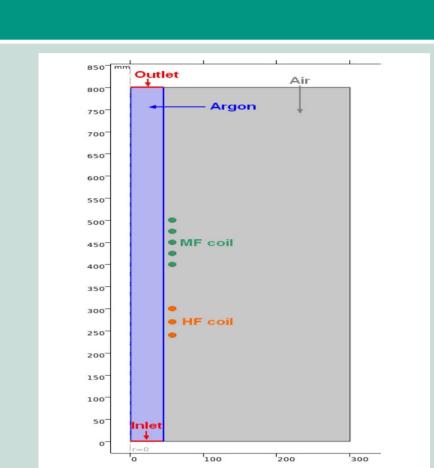


Temperature profile of 2D axisymmetric model of the dual frequency torch

Geometry

Integration of High frequency (HF) and Medium frequency (MF) coils

- Diameter of the tube $D_{burner} = 90 \ mm$
- Length of the tube $L_{burner} = 800 \ mm$
- Number of HF windings $N_{HF} = 3$
- Length of the HF coil $L_{HF} = 70 \ mm$
- Number of MF windings $N_{MF} = 5$
- Radius of copper coil $R_{coil} = 5 mm$
- Length of the MF coil $I = -110 \, mm$
- Length of the MF coil $L_{MF} = 110 \ mm$



2D axisymmetric model of the plasma torch

Equations

The equilibrium discharge model from COMSOL Multiphysics couples the following interfaces

- Magnetic fields
- Heat transfer in fluids R
- Laminar flow

$$\nabla \times H = J$$

$$B = \nabla \times A$$

$$J = \sigma E + j\omega D + \sigma u \times B + J_e$$

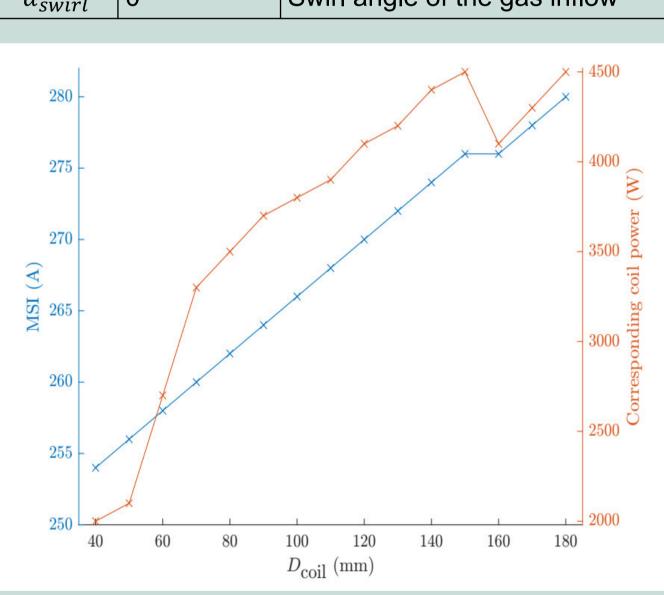
 $E = -j\omega A$

 $\rho C_p \frac{\partial T}{\partial t} + \rho C_p u. \nabla T + \nabla \cdot q = Q + Q_p + Q_{vp}$ $q = -\kappa \nabla T$

$$\rho \frac{\partial u}{\partial t} + \rho(u.\nabla)u = \nabla \cdot [-pI + K] + F$$
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho u) = 0$$

Simulation results

Default values			
	Variable	Values	Description
	f_{HF}	13.56 MHz	HF coil frequency
	P_{HF}	3 kW	HF coil excitation power
	f_{MF}	200 kHz	MF coil frequency
	I_{MF}	5000 A	MF coil excitation current
	D_{coil}	100 mm	Distance between the two coils
	Q_{inlet}	100 L/min	Argon flow at the inlet
	acwirl	0°	Swirl angle of the gas inflow



MSI of the MF coil for different coil distance with $P_{HF}=3\ kW$

Varying the power of HF coil

 $341 A \text{ at } P_{HF} = 1.2 kW$

1.2 *kW*

Conclusion

The coil distance fixed at 100 mm

The MSI increases from 266 A at $P_{HF} = 3 kW$ to

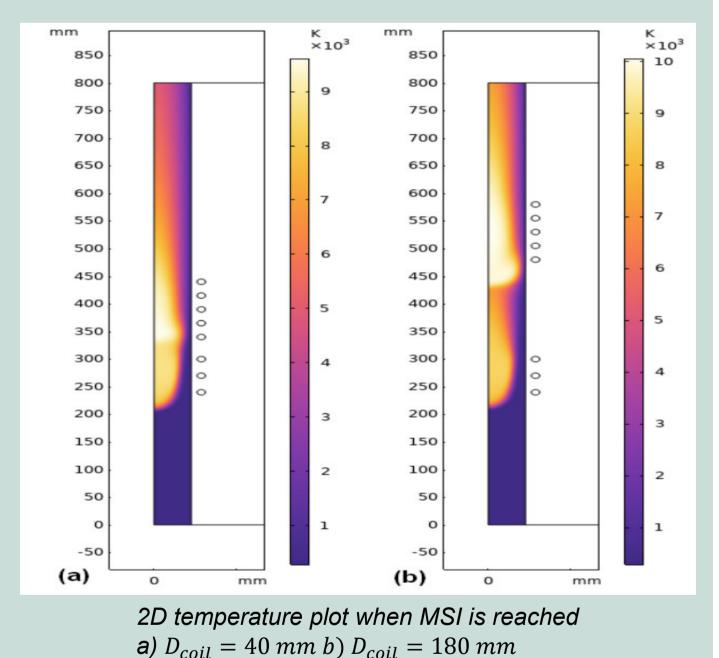
It is more reasonable to choose the HF power

complications form the engineering point of view

around 2.5 kW or 3 kW rather than around

Varying the distance between the coils

- MSI values for MF coil at varied distance between $40 \ mm$ and $180 \ mm$
- 3 kW power excitation used for HF coil
- The MSI increases almost linearly with the coil distance, the shorter the distance, the hotter the gas at the entrance of the MF coil



300 - 3750 - 3700 (M) - 3650 Jawood III Jawood II Jaw

MSI of the MF coil for different power excitation of the HF coil

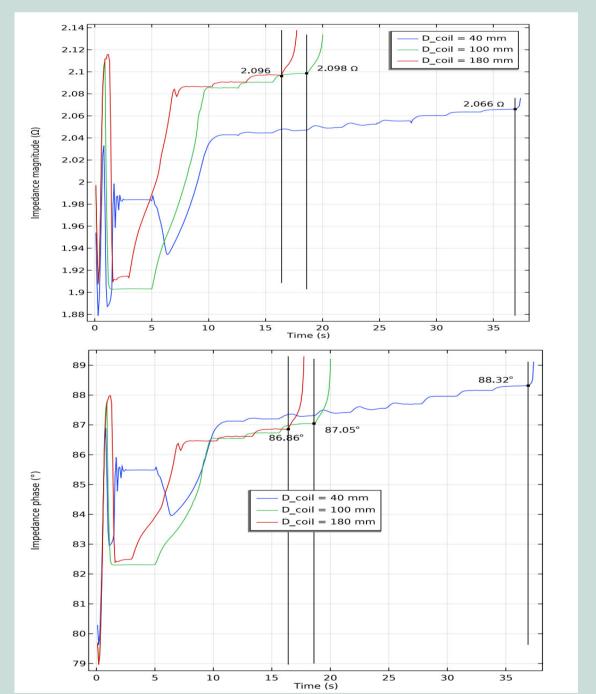
Additional quantities of interest

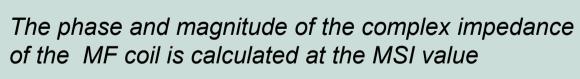
Impedance of the MF coil

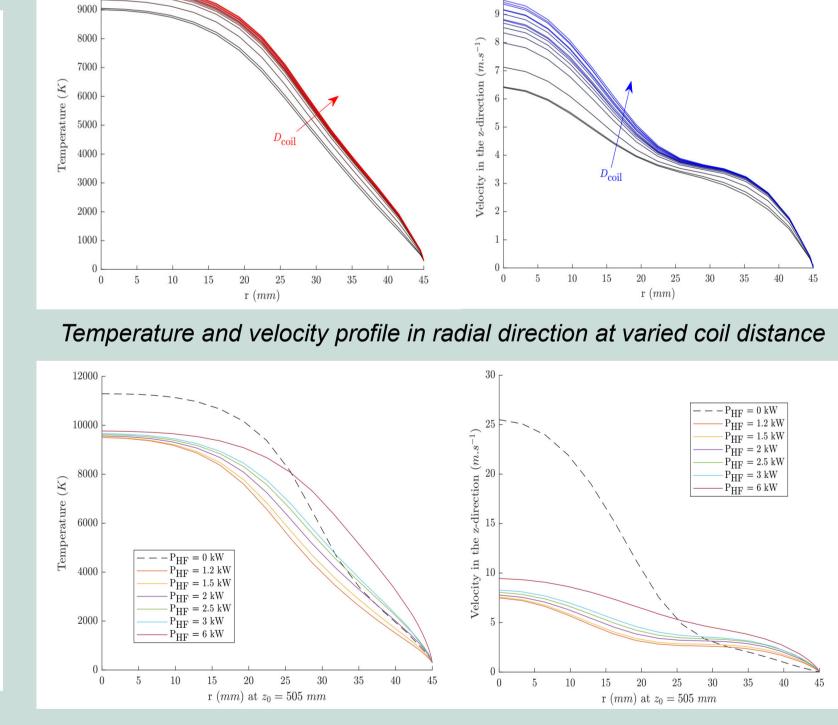
- Maximum and minimum impedance magnitude of $2.121~\Omega$ (at $P_{HF}=1.2~kW$) and $2.066~\Omega$ (at $D_{coil}=180~mm$) are respectively achieved
- Maximum and minimum impedance phase of 88.32° (at $D_{coil} = 40~mm$) and 86.86° (at $D_{coil} = 180~mm$ are also respectively obtained

Temperature and velocity profiles

- Temperature and velocity profiles along radial direction at the end of MF coil (z_0) are investigated when 1) at variable D_{coil} 2) at different P_{HF} level
- Both temperature and velocity increase as D_{coil} and P_{HF} increases



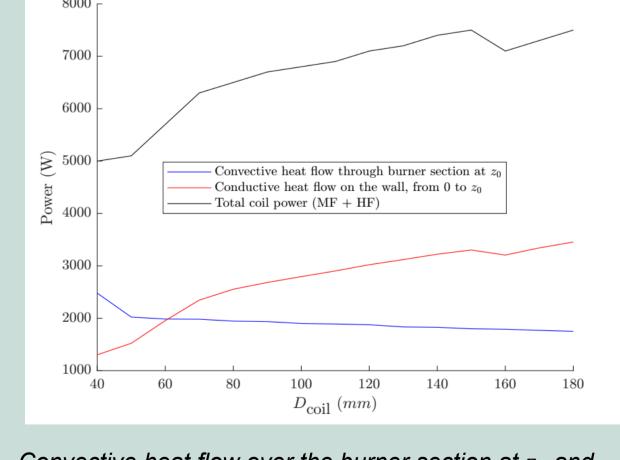




Temperature and velocity profile at different high frequency power level

Conductive and convective heat flows at the end of the MF coil

- $P_{HF} \text{ is fixed at } 3 \text{ } kW$
- The convective heat flux integrated over a disk of radius R_{burner} , located at z_0
- The conductive heat flux integrated over the burner wall on a cylindrical surface of length z_0



Convective heat flow over the burner section at z_0 and conductive heat flow over the burner wall up to z_0

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[3] D. Bernardi, V. Colombo, E. Ghedini, and A. Mentrelli. Numerical simulation for the characterization of operating conditions of rf –rf hybrid plasma torches. *The European Physical Journal D - Atomic, Molecular and Optical Physics*, 28:399–422, 2004.

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[4] V. Frolov, D. Ivanov, and V. Sosnin. Numerical simulation of high power rf–rf hybrid plasma torch. *IOP Conference Series: Materials Science and Engineering*, 643:012071, 2019.

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Igniting and sustaining an inductive plasma at a low/medium frequency may be impossible at

We quantified the minimum sustaining current for stable plasma operation while varying the

distance between the coils and the power from the HF support. The advantage of this design is

is expensive. The hybrid torch concept presented here is an attempt to solve this issue

Having a quasi-constant impendance means the design can be modified without further

lost if the coil distance increases too much (plasmas completely separated)

atmospheric pressure. On the other hand, coupling a lot of power with high-frequency electronics