

# The Disulfide Bond in S $L_{2,3}$ X-ray Emission Spectroscopy of Transition Metal Sulfides

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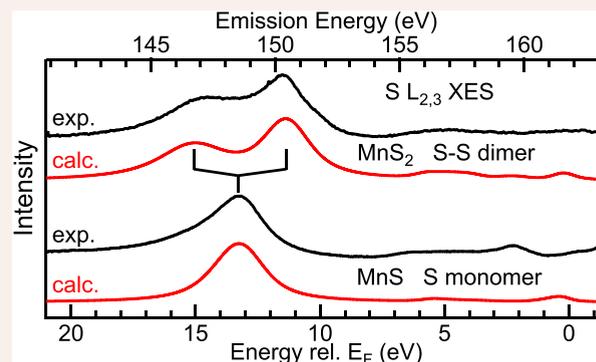
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**ABSTRACT:** S  $L_{2,3}$  X-ray emission spectroscopy (XES) and spectra calculations based on density functional theory are used to study the electronic valence structure of Mn, Fe, Co, and Cu transition metal sulfides, without (MnS, FeS, CoS, and Cu<sub>2</sub>S) and with (MnS<sub>2</sub>, FeS<sub>2</sub>, CoS<sub>2</sub>, and CuS) sulfur dimers in the crystal structure. We find the upper valence band region in the experimental spectra to be very sensitive to the transition metal, while a splitting in the transitions from the S 3s derived bands indicates the presence of sulfur dimers. This makes S  $L_{2,3}$  XES particularly well suited to the chemical speciation of these compounds. The spectral changes are reproduced by our spectra calculations, which allows for a detailed understanding of the electronic structure based on the band structure and the projected density of states. The splitting of the S 3s derived bands in the dimer compounds is found to increase with decreasing sulfur–sulfur distance, indicating an increase of covalency. The splitting in our experimental XES spectra is found to be smaller than in photoemission data and our spectra calculations, which we attribute to nuclear dynamics on the time scale of the X-ray emission process. In particular, the high sensitivity of S  $L_{2,3}$  XES to the presence of sulfur dimers, together with its suitability for operando studies, makes it a key technique for future studies on a wide range of relevant material systems.

**KEYWORDS:** X-ray emission spectroscopy (XES), transition metal sulfides, valence bond, sulfur dimers



## INTRODUCTION

Transition metal sulfides play an important role in many applications. These include the use of NiS<sub>2</sub> as a low-cost electrode material for high-performance supercapacitors,<sup>1</sup> NiS<sub>2</sub>, NiS, and Ni<sub>3</sub>S<sub>4</sub> electrocatalysts for hydrogen production in seawater,<sup>2</sup> or FeS<sub>2</sub>-based materials for batteries.<sup>3</sup> Further examples are CoS<sub>2</sub>, Co<sub>9</sub>S<sub>8</sub>, CoS, FeS<sub>2</sub>, and MoS<sub>2</sub> as anodes in Li-ion batteries<sup>4</sup> and M<sub>x</sub>S<sub>y</sub> (M = Fe, Co, Ni, etc.) as bifunctional electrocatalysts for water splitting.<sup>5</sup> In particular, earth-abundant metal pyrites (FeS<sub>2</sub>, CoS<sub>2</sub>, NiS<sub>2</sub>, and their alloys) are highly efficient for electrocatalytic hydrogen evolution and polysulfide reduction.<sup>6</sup> It is suggested that, in these materials, the transition metal as well as the disulfide anions (S<sub>2</sub><sup>2-</sup>) play an active role in high catalytic activity.<sup>6</sup>

Due to their importance, the geometric, electronic, and chemical structures of transition metal sulfides have been characterized extensively. For chemical speciation and electronic structure evaluation, electron and X-ray spectroscopy techniques have been employed. Much of that work has focused on the metal states, including 2p X-ray photoelectron spectroscopy (XPS)<sup>7–12</sup> and K and L<sub>2,3</sub> edge X-ray absorption spectroscopy (XAS).<sup>12–16</sup> But also S 2p XPS,<sup>7–9,12</sup> S K,<sup>13,15,17–19</sup> and L<sub>2,3</sub><sup>14,20,21</sup> XAS have been employed, giving detailed insights into these compounds.

X-ray emission spectroscopy (XES) has also been used to study transition metal sulfides, both for metal<sup>15,22–25</sup> and sulfur<sup>15,17,26–28</sup> K and L<sub>2,3</sub> edges. The local character of XES, probing the chemical environment around a selected core level, makes it very powerful for chemical speciation, particularly in mixed compounds with many constituents. This, together with its suitability for operando studies, makes XES a very powerful technique for applied materials studies. To probe the local environment around the sulfur atoms, most commonly S K XES has been employed. In this study, we focus on S L<sub>2,3</sub> XES, with a particular emphasis on the S 3s derived valence bands, which are not visible in S K XES due to the dipole selection rules.

The S L<sub>2,3</sub> XES spectra are very sensitive to the chemical environment of sulfur, probing valence states of s and d symmetry, which was already recognized in an early pioneering study by O'Bryan and Skinner.<sup>28</sup> Besides their sensitivity to the

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S oxidation state,<sup>27</sup> sulfur compounds with the same oxidation state, which would not show a discernible core level shift in XPS, like sulfides<sup>27,29,30</sup> or sulfates,<sup>31</sup> can be easily distinguished. Of the transition metal sulfides addressed in the present study, S  $L_{2,3}$  spectra for MnS,<sup>28</sup> Cu<sub>2</sub>S,<sup>14,19,20,27,28,32,33</sup> CuS,<sup>14,19,27,34</sup> and FeS<sub>2</sub><sup>34</sup> are found in the literature.

In this study, we investigate a systematic series of transition metal sulfides *without* (MnS, FeS, CoS, and Cu<sub>2</sub>S) and with (MnS<sub>2</sub>, FeS<sub>2</sub>, CoS<sub>2</sub>, and CuS) sulfur dimers in the crystal structure to derive and understand the signature of the disulfide bond in S  $L_{2,3}$  XES.

## EXPERIMENTAL SECTION

Experiments were performed in two different experimental setups. The spectra for MnS, FeS, FeS<sub>2</sub>, Cu<sub>2</sub>S, and CuS were measured using the high-transmission spectrometer<sup>35</sup> of the Solid and Liquid Spectroscopic Analysis (SALSA) roll-up experimental station<sup>36</sup> at Beamline 8.0.1 of the Advanced Light Source (ALS), Lawrence Berkeley National Laboratory (Berkeley, USA). The energy scale for the SALSA data was calibrated as described in ref 29, giving an absolute (relative) uncertainty of 0.1 eV (0.03 eV). The MnS<sub>2</sub>, CoS, and CoS<sub>2</sub> data were measured with the high-transmission spectrometer<sup>37</sup> of the X-SPEC beamline<sup>38</sup> at the KIT Light Source. The energy scale for these measurements was calibrated using a CaSO<sub>4</sub> reference measurement for an absolute uncertainty of 0.1 eV. For both instruments, the experiments were run with a combined energy resolution (beamline and spectrometer) of 0.25 eV, as determined from the full-width at half-maximum of the elastic scattered peaks.

For the experiments at the ALS, MnS, FeS, FeS<sub>2</sub>, Cu<sub>2</sub>S, and CuS powders were purchased and shipped under argon from Alfa Aesar. The powders were fixed to the sample holder by pressing them onto double-sided and ultrahigh vacuum-compatible tape in a nitrogen-filled glove box. The samples were then transported in a nitrogen-filled container to the experimental station and loaded into vacuum without air exposure using a nitrogen-filled glove bag. For the experiments at the KIT Light Source, CoS and CoS<sub>2</sub> powders were purchased and shipped under argon from Fisher Scientific and pressed into thin pellets in an argon-filled glove box. For MnS<sub>2</sub>, a hauerite crystal from the Jeziórko mine in Tarnobrzeg (Poland) was cleaved in an argon-filled glove box. The samples were then transported to the beamline in an argon-filled sample container, which was then directly attached to the experimental station and transferred without any air exposure.

## THEORY

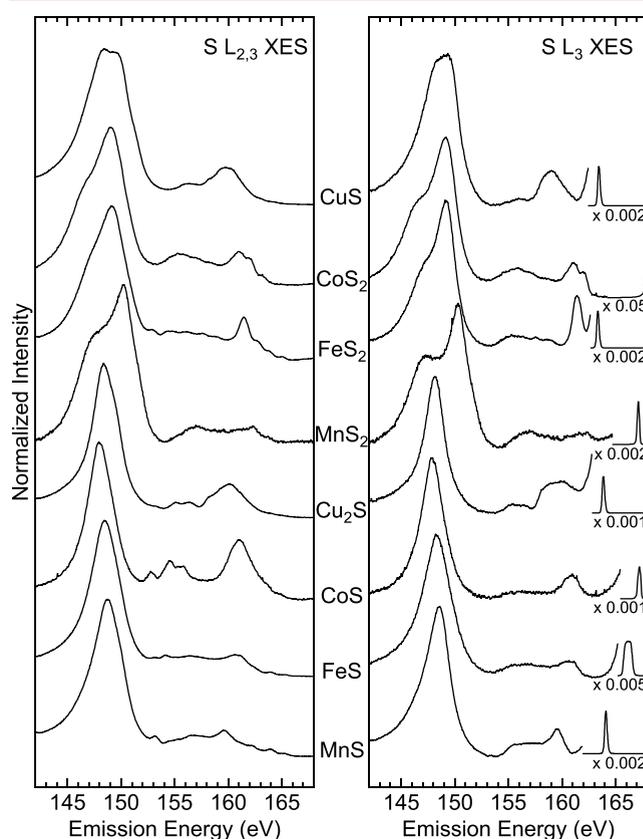
To calculate band structures, projected densities of states (PDOS), and XES spectra, we have used Wien2k, which uses the full-potential augmented plane wave plus local orbitals (APW+lo) method to solve the Kohn–Sham density functional theory (DFT) equations.<sup>39</sup> Exchange–correlation effects were described by the generalized gradient approximation (GGA) as parameterized by Perdew, Burke, and Ernzerhof (PBE).<sup>40</sup> XES spectra were calculated from the PDOS using the dipole approximation and Fermi's Golden Rule in the formalism described by Schwarz et al.<sup>41–43</sup> 1000  $k$  points were used for the self-consistent field (SCF) cycle, except for Cu<sub>2</sub>S with 100  $k$  points. For the PDOS and XES spectra calculations, 10,000  $k$  points were used (1000  $k$  points for Cu<sub>2</sub>S). The

crystal structures (see Figures S1–S8 in the Supporting Information) were obtained from the Materials Project Database<sup>44</sup> with the materials IDs mp-504 (CuS), mp-2070 (CoS<sub>2</sub>), mp-226 (FeS<sub>2</sub>), mp-581212 (Cu<sub>2</sub>S), mp-1274 (CoS), mp-2099 (FeS), and mp-2065 (MnS) and the Crystallography Open Database, ID 1011237 (MnS<sub>2</sub>).<sup>45</sup>

In addition to the calculations of the transition metal sulfide crystals, we have used the StoBe-DeMon 3.3 DFT code<sup>46</sup> to calculate the orbital energies of an isolated S<sub>2</sub><sup>2-</sup> ion as a function of sulfur–sulfur distance. The calculations used the Becke–Perdew gradient-corrected exchange and correlation functionals,<sup>47–49</sup> triple  $\zeta$  bases<sup>50</sup> to describe the valence electrons, and effective core potentials (ECPs)<sup>51</sup> with 73111/6111/1 basis sets. The auxiliary basis sets were composed of five  $s$  and four  $spd$  functions for sulfur to fit the Coulomb and exchange–correlation potentials.

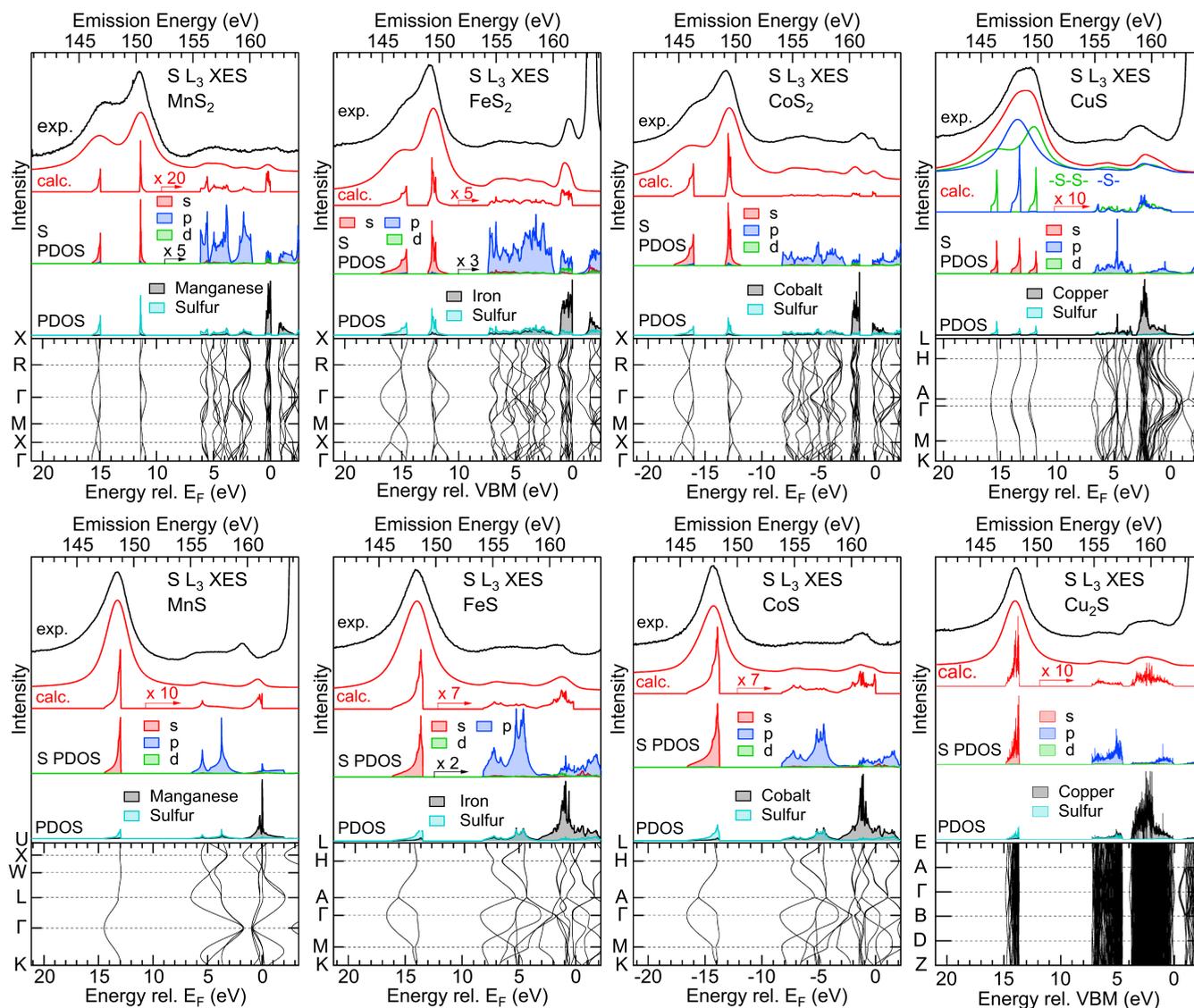
## RESULTS AND DISCUSSION

Figure 1 shows the experimental XES spectra of the investigated samples, excited with photon energies well above



**Figure 1.** Experimental S  $L_{2,3}$  (left) and S  $L_3$  (right) XES spectra of (from top to bottom) CuS, CoS<sub>2</sub>, FeS<sub>2</sub>, MnS<sub>2</sub>, Cu<sub>2</sub>S, CoS, FeS, and MnS. For S  $L_{2,3}$ , excitation energies of 184.6 eV (CuS, FeS<sub>2</sub>, Cu<sub>2</sub>S, FeS, and MnS) and 200.3 eV (CoS<sub>2</sub>, MnS<sub>2</sub>, and CoS) were used. The S  $L_3$  spectra were excited just below the S  $L_2$  onset, i.e., with 163.4 eV (CuS and FeS<sub>2</sub>), 168.2 eV (CoS<sub>2</sub>), 167.1 eV (MnS<sub>2</sub>), 163.9 eV (Cu<sub>2</sub>S), 167.2 eV (CoS), 166.2 eV (FeS), and 164.1 eV (MnS).

the absorption onset (S  $L_{2,3}$ , left) and in-between the  $L_3$  and  $L_2$  edges ( $L_3$ , right). The spectra contain transitions from valence states that exhibit wave function overlap with the S 2p core levels and obey the dipole selection rule  $\Delta l = \pm 1$ , i.e., probing valence states of  $s$  and  $d$  symmetry. The valence bands of the



**Figure 2.** Comparison of the experimental S  $L_3$  XES spectra of the investigated compounds (“exp.”, top spectrum in each panel) with calculations based on DFT (“calc.”). The bottom and top rows show compounds with S monomers and dimers in the crystal structure, respectively. Each panel shows (from bottom to top): calculated band structure, density of states projected onto atomic species (PDOS), symmetries for the sulfur atoms (S PDOS), calculated S  $L_3$  XES spectrum without and with experimental and lifetime broadening (both red), and the experimental spectrum (black). The theory energy scales on the bottom and the experimental emission energy scales on the top were shifted to align the most prominent spectral features in theory and experiment. In some cases, the ordinates of the S PDOS and the unbroadened S  $L_3$  calculations were multiplied by the factors given (upper valence band region only).

transition metal sulfides investigated are derived from the S 3s and 3p as well as the metal 3d and 4s states. Due to suitable symmetry and large overlap with the S 2p states, all spectra are dominated by transitions from S 3s derived bands around 148.5 eV. The upper valence bands are dominated by S 3p derived bands around the sulfur atoms, and we thus only find weak emission above 154 eV from sulfur s and d admixtures to these bands. As can be seen in Figure 1, this latter spectral region is very sensitive to the specific metal bonding partner and can serve as a “fingerprint” tool. While compounds with the same metal exhibit similarities in this region, we still find distinct differences that allow for a speciation of the respective compound.

Different from the upper valence band region, the peak assigned to transitions from S 3s derived bands at  $\sim 148.5$  eV (“S 3s” in the following) shows much less dependence on the

metal of the compound. In fact, only small changes in line widths are observed for this feature between the spectra of  $\text{Cu}_2\text{S}$ ,  $\text{CoS}$ ,  $\text{FeS}$ , and  $\text{MnS}$  due to differences in lifetime broadening. In contrast, a splitting of the “S 3s” peak is observed for  $\text{CuS}$  (also observed in refs 19,27,34),  $\text{CoS}_2$ ,  $\text{FeS}_2$  (also in ref 34), and  $\text{MnS}_2$ . This is a consequence of the splitting of the S 3s derived bands by the covalent disulfide bond in these compounds and will be discussed in more detail below.

We note that the S  $L_{2,3}$  spectra mainly differ from the  $L_3$  spectra by containing all features twice, with an  $L_3/L_2$  intensity ratio of 2:1 and separated by the spin–orbit splitting of the S 2p core levels (1.20 eV<sup>52,53</sup>). The one exception is found for our  $\text{CoS}$  data; this sample was partially oxidized, leading to sulfate and sulfite contributing to the S  $L_{2,3}$  XES spectrum at 152.8, 154.5, 155.7, and  $\sim 161$  eV. In contrast, the  $L_3$  spectra

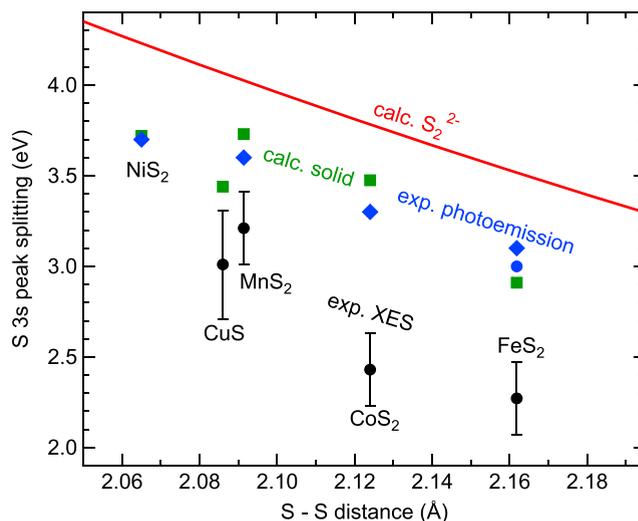
are free of these artifacts since they are excited below the absorption edge of sulfates and sulfites. In the following, we will thus discuss the “simpler”  $L_3$  spectra in more detail to avoid any sulfate, sulfite, and  $L_2$  contributions.

For this discussion, Figure 2 compares the experimental  $S L_3$  spectra with calculations based on DFT. For each compound, Figure 2 shows (from bottom to top) the calculated band structure, density of states projected onto atomic species (PDOS) and symmetries for the sulfur atoms (S PDOS), the calculated  $S L_3$  XES spectrum without and with experimental and lifetime broadening (both red), and the experimental spectrum (black). Overall, the calculations give a very good description of the experimental spectra, and we observe the following trends: the upper valence bands are mainly comprised of states of p symmetry, located near the sulfur atoms (at higher binding energies), and of d symmetry, located near the metal atoms (at lower binding energies). Neither the states of pure p symmetry located at the sulfur atoms (being dipole-forbidden) nor the states of d symmetry localized at the metal atoms (due to missing wave function overlap with the  $S 2p$  core levels) give rise to spectral intensity in the  $S L_3$  XES spectra. However, we find some degree of hybridization between  $S 3s$  and  $S 3p$ , leading to s symmetric S PDOS in the region of the sulfur p bands. Also, hybridization between  $S 3d$  and the metal d states is observed, leading to d symmetric S PDOS in the region of the metal d bands. These states are reflected in the  $S L_3$  XES spectra as a broader structure between emission energies of approx. 154 to 158 eV and emission at or slightly above 158 eV, respectively. As a general trend, the increasing occupation of the metal d states with increasing atomic number is reflected in an increase of the occupied metal PDOS below the VBM (or  $E_F$ ) as well as an increase of hybridization with the  $S 3d$  states. This is most prominently visible for the copper sulfides, where the intensity increases above 158 eV.

As expected, we find predominantly s symmetric S PDOS in the region of the “ $S 3s$ ” peak(s). For the compounds with exclusively S monomers, we observe one band (MnS) or one group of bands (FeS, CoS, and  $Cu_2S$ ), while the compounds with S dimers exhibit two (MnS<sub>2</sub>, FeS<sub>2</sub>, and CoS<sub>2</sub>) clearly split groups of bands. For CuS, three groups of  $S 3s$  derived bands are found, two of which are due to sulfur dimers and the third due to sulfur monomers within the compound. In a simple molecular orbital picture, the splitting of the bands in the sulfur dimers is caused by hybridization of the  $S 3s$  orbitals, forming bonding and antibonding “molecular” orbitals (at higher and lower binding energies, respectively).

We find the magnitude of this splitting to vary among the different compounds, which is investigated in conjunction with Figure 3. There, values for the  $S 3s$  splitting are derived from our experimental XES spectra (black circles), from the corresponding DFT spectra calculations (green squares), from experimental photoemission spectra taken from literature<sup>10,54</sup> (blue diamonds), as well as from DFT calculations of an isolated  $S_2^{2-}$  ion (red line). Overall and as expected, the  $S 3s$  splitting decreases as a function of increasing S–S distance. The calculated values for the  $S_2^{2-}$  ion lie significantly above the values of the solids, which we attribute to the missing environment in the case of the isolated ion.

The splitting values calculated for the different solid compounds (i.e., those investigated here, plus NiS<sub>2</sub>) are in good agreement with the literature values<sup>10,54</sup> from photoemission experiments. In contrast, we find our XES experi-



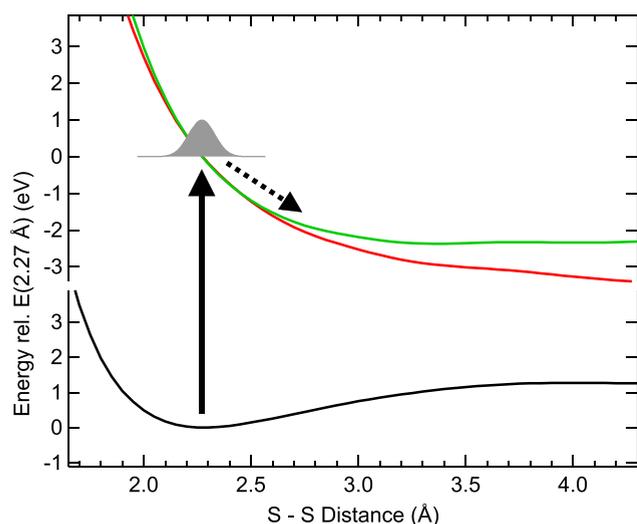
**Figure 3.**  $S 3s$  peak splitting for the studied compounds as a function of S–S distance (taken from the same crystal structures used in the calculations), determined experimentally by XES (black circles, this work) and photoelectron spectroscopy<sup>10,54</sup> (blue diamonds), as well as calculated based on DFT (green squares). For comparison, the splitting calculated for an isolated  $S_2^{2-}$  ion of varied S–S distance is shown as a red line.

ments to yield a significantly smaller splitting than in both calculations and photoemission experiments. While the former could potentially indicate shortcomings in the energy scale of the calculation (as is not unusual for DFT approaches), photoemission and XES nominally share the same final state and are thus expected to give the same energy difference based on the final-state rule.<sup>55</sup> We attribute the decreased splitting to the finite duration of the XES process, during which ultrafast nuclear dynamics occur. The time scale is governed by the core-hole lifetime ( $\sim 16$  fs for  $S 2p$ <sup>56</sup>), and thus symmetry breaking,<sup>57</sup> vibrational excitation,<sup>58–60</sup> as well as full dissociation of molecules<sup>61–65</sup> (among others) can occur on that time scale.

To understand these dynamics for the sulfur dimer, Figure 4 shows the potential energy curves of the  $S_2^{2-}$  ion in the ground (black), lowest core-excited (red), and core-ionized (green) states, as derived from our DFT calculations. The exciting X-ray will create a wave packet located at the equilibrium bond distance of the ground state, which will move along the potential energy surface of the excited state until emission takes place. In general, this will increase the S–S distance in the final state and thus lead to a smaller splitting (and to a peak broadening) in the XES spectra. If we take the slope of the red curve in Figure 3, we see a decrease in splitting of  $\sim 0.5$  eV, as roughly observed between exp. XES and PES/calc. solid, which would correspond to a S–S distance increase of  $\sim 0.07$  Å. This is a rather small change in distance, similar to the S–S distance difference between CuS and FeS<sub>2</sub>.

## CONCLUSIONS

The electronic valence band structure of a series of Mn, Fe, Co, and Cu transition metal sulfides has been investigated using  $S L_{2,3}$  soft X-ray emission spectroscopy (XES). We have studied compounds with sulfur dimers (CuS, CoS<sub>2</sub>, MnS<sub>2</sub>, and FeS<sub>2</sub>) and without sulfur dimers ( $Cu_2S$ , CoS, FeS, and MnS) in the structure. The XES spectra provide unique fingerprints for each compound that can be understood and used for chemical



**Figure 4.** Calculated potential energy curves of the  $S_2^{2-}$  ion in the ground (black), lowest core-excited (red), and core-ionized (green) states. The vertical arrow indicates the X-ray excitation that creates a wave packet in the excited state (gray Gaussian), which will move along the potential energy curve (dashed arrow).

speciation. We find that the upper valence band region in the spectra is very sensitive to the metal bonding partner of sulfur in the compounds, which can be understood in detail using spectra calculations based on DFT. Furthermore, we find a splitting of the S 3s band in the compounds containing S dimers, which decreases with increasing S–S distance. While this qualitative change agrees with our spectra calculations as well as with photoemission experiments in the literature, we find the splitting in XES to be less than the photoemission values and those predicted by our ground state calculations. This is assigned to ultrafast dynamics on the femtosecond time scale of the X-ray emission process that leads to a S–S bond elongation upon X-ray excitation. The demonstrated high sensitivity of S  $L_{2,3}$  XES to the S–S dimer bond is expected to give crucial insights for a wide range of applied material systems.

## ■ ASSOCIATED CONTENT

### SI Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/photonsci.6c00007>.

Crystal structures of the investigated compounds (PDF)

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## Notes

The authors declare no competing financial interest.

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