

# **Contributions of the direct numerical simulation method to the validation of statistical turbulence models for natural convection**

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# Introduction

- Passive decay heat removal in metal and water cooled reactors
  - experiments (scaled reactor models, water as model fluid)
  - engineering computer codes (e.g. FLUTAN)
- Validation of statistical turbulence models for natural convection in liquid metals
  - basic experiments in simple geometries
  - direct numerical simulation

## Method of direct numerical simulation

- Conservation equations of mass, momentum and energy

$$\nabla \hat{\vec{u}} = 0$$

$$\partial \hat{\vec{u}} / \partial t + (\hat{\vec{u}} \nabla) \hat{\vec{u}} = - \nabla p + \frac{1}{\sqrt{Gr}} \nabla^2 \hat{\vec{u}} + \left( T_{ref} - T \right) \frac{\hat{\vec{g}}}{|\hat{\vec{g}}|}$$

$$\partial T / \partial t + (\hat{\vec{u}} \nabla) T = \frac{1}{Pr \sqrt{Gr}} \nabla^2 T$$

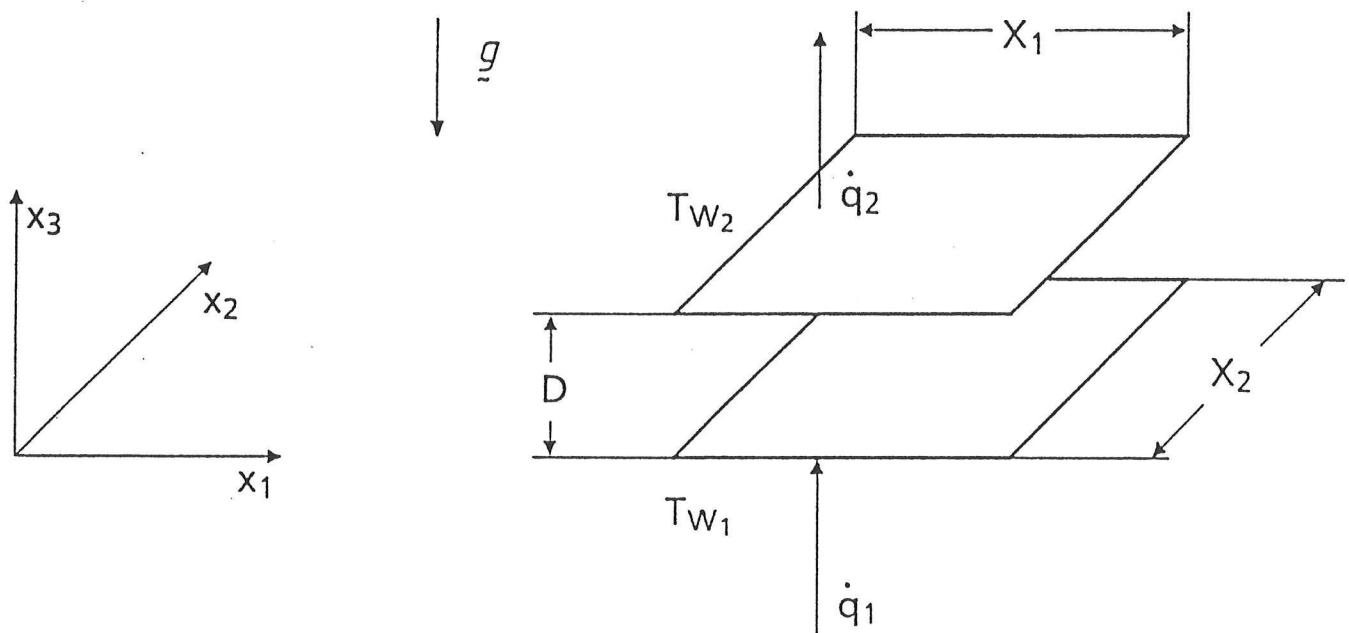
- 3d, time dependent
- resolve all spatial scales of turbulence
  - + no model assumptions, no parameters
  - limited to small turbulence levels and simple geometries

# Computer code TURBIT

- finite volume method
  - central finite differences
  - staggered grid
- time integration
  - momentum equation: explicit
  - energy equation: explicit or semi-implicit
  - Poisson equation for pressure
- geometry
  - plane channel
  - annulus

# Rayleigh-Bénard convection

- geometry



- dimensionless numbers

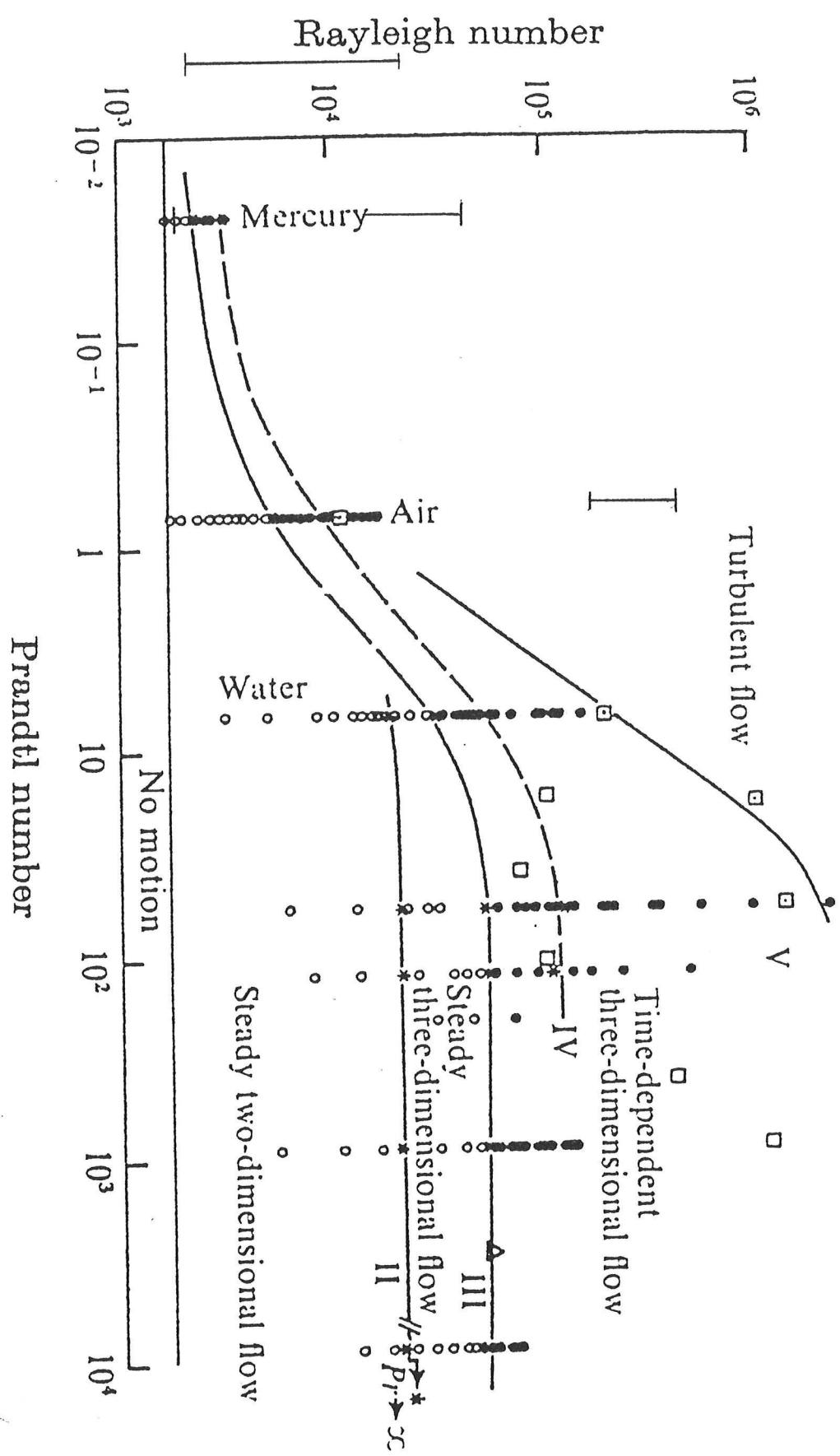
- Rayleigh number  $Ra \sim \Delta T_w$
- Prandtl number  $Pr = \nu/\kappa$

sodium:  $Pr = 0.006$

mercury:  $Pr = 0.024$

air:  $Pr = 0.71$

# Rayleigh-Bénard convection



## Statistical Turbulence Models

- closure of  $\overline{u_i' u_j'}$ ,  $\overline{u_i' T'}$ 
  - solution of modelled transport equations for turbulent quantities  
e.g.:  $k, \varepsilon, \overline{T'^2}, \varepsilon_T, \overline{u_i' u_j'}, \overline{u_i' T'}$
- use DNS-data to
  - analyse exact transport equations for  $k, \varepsilon, \dots$
  - analyse model assumption for different terms

# Transport equation for turbulent heat flux $\overline{u_3' T'}$

$$0 = - \frac{\partial}{\partial x_3} \left( \frac{\overline{u_3'^2 T'}}{\overline{u_3'^2}} + p' T' - \frac{1}{Pr \sqrt{Gr}} \overline{u_3'} \frac{\partial T'}{\partial x_3} - \frac{1}{\sqrt{Gr}} \overline{T'} \frac{\partial u_3'}{\partial x_3} \right)$$

$D_q$

$$- \overline{u_3'^2} \frac{\partial \overline{T}}{\partial x_3} + \overline{T'^2} + p' \frac{\partial T'}{\partial x_3}$$

$P_q$

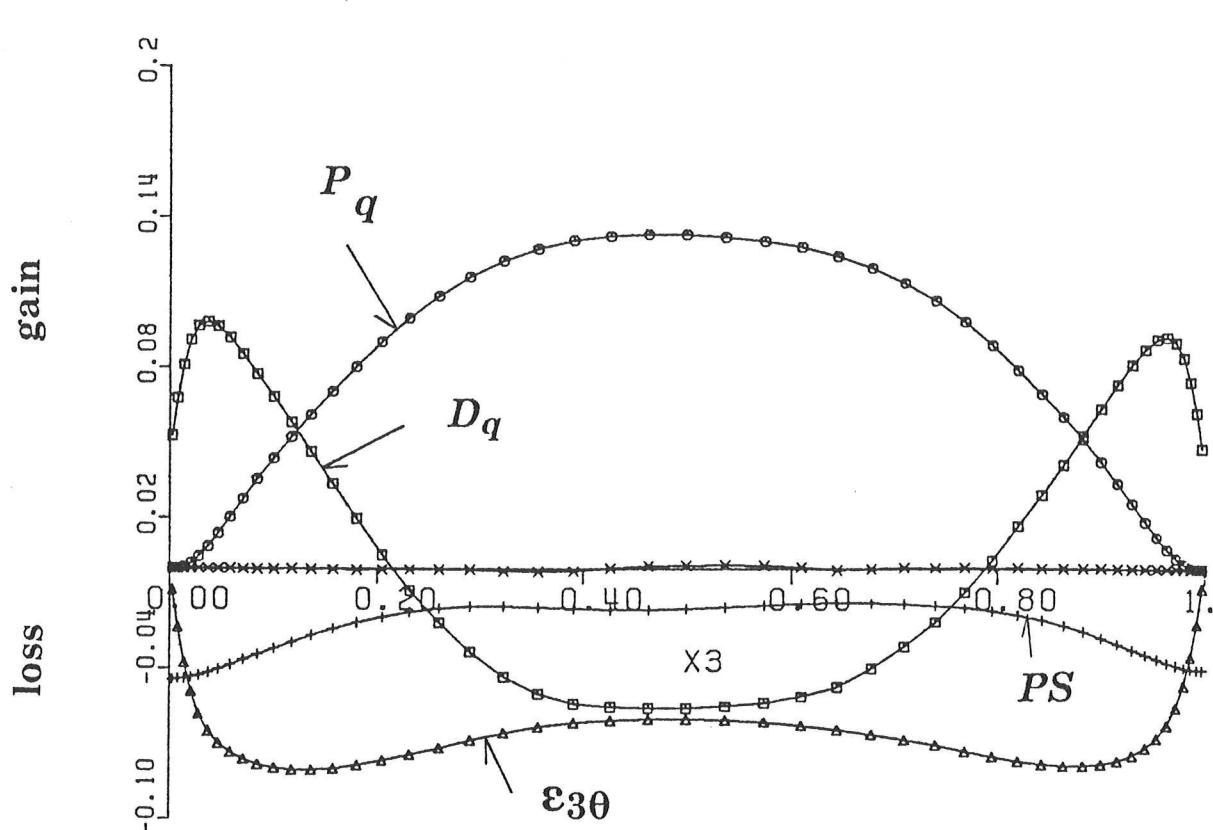
$PS$

$$- \frac{1}{\sqrt{Gr}} \left( 1 + \frac{1}{Pr} \right) \frac{\partial u_3'}{\partial x_i} \cdot \frac{\partial T'}{\partial x_i}$$

$\varepsilon_{30}$

## Budget of $\overline{u_3'T'}$

- Sodium,  $Ra = 24\,000$

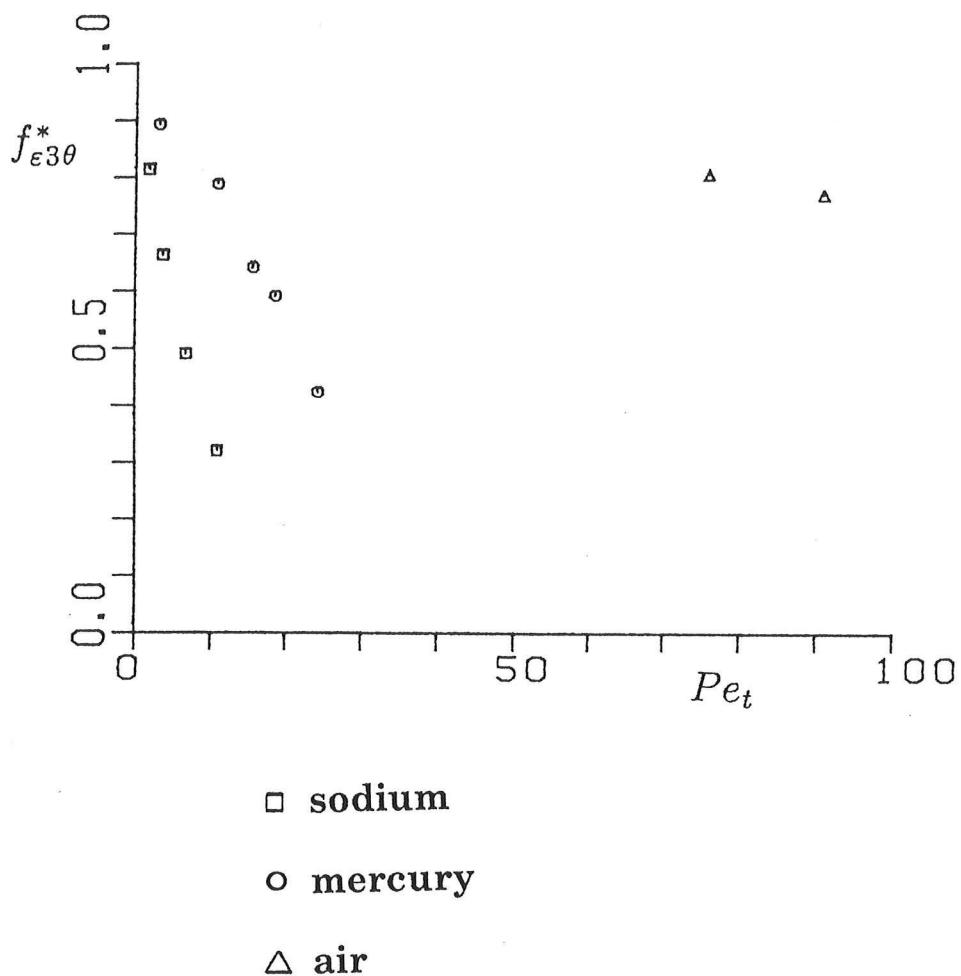


## Modelling of $\varepsilon_{i\theta}$

- Shikazono & Kasagi (1993):

$$\varepsilon_{i\theta} = f_{\varepsilon i\theta}^* \frac{1+Pr}{2\sqrt{Pr}} \frac{\varepsilon}{k} \overline{u_i' T}$$

- Analysis of  $f_{\varepsilon 3\theta}^*$  by DNS-data

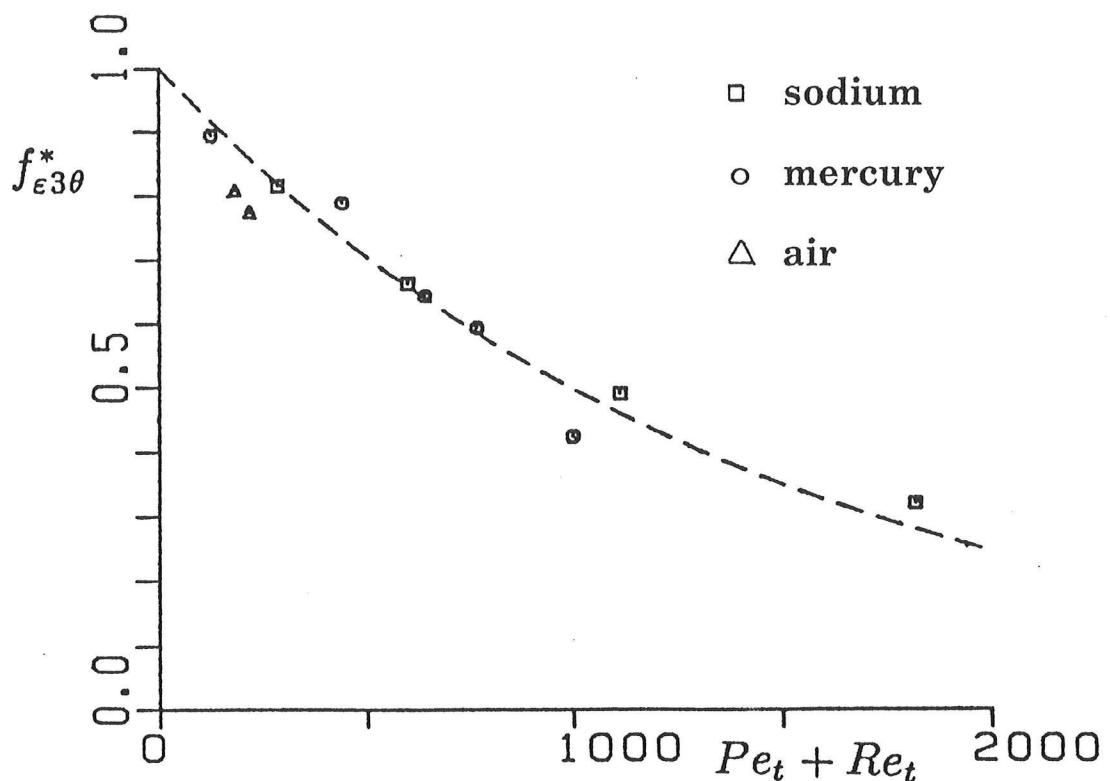


## Modelling of $\varepsilon_{i\theta}$

- Shikazono & Kasagi (1993):

$$\varepsilon_{i\theta} = f_{\varepsilon i\theta}^* \frac{1+Pr}{2\sqrt{Pr} \sqrt{R}} \frac{\varepsilon}{k} \overline{u_i' T}$$

- Analysis of  $f_{\varepsilon 3\theta}^*$  by DNS-data



$$f_{\varepsilon i\theta}^* = \exp(-C_{\varepsilon i\theta}(Re_t + Pe_t))$$

$$C_{\varepsilon 3\theta} \approx 0.0007$$

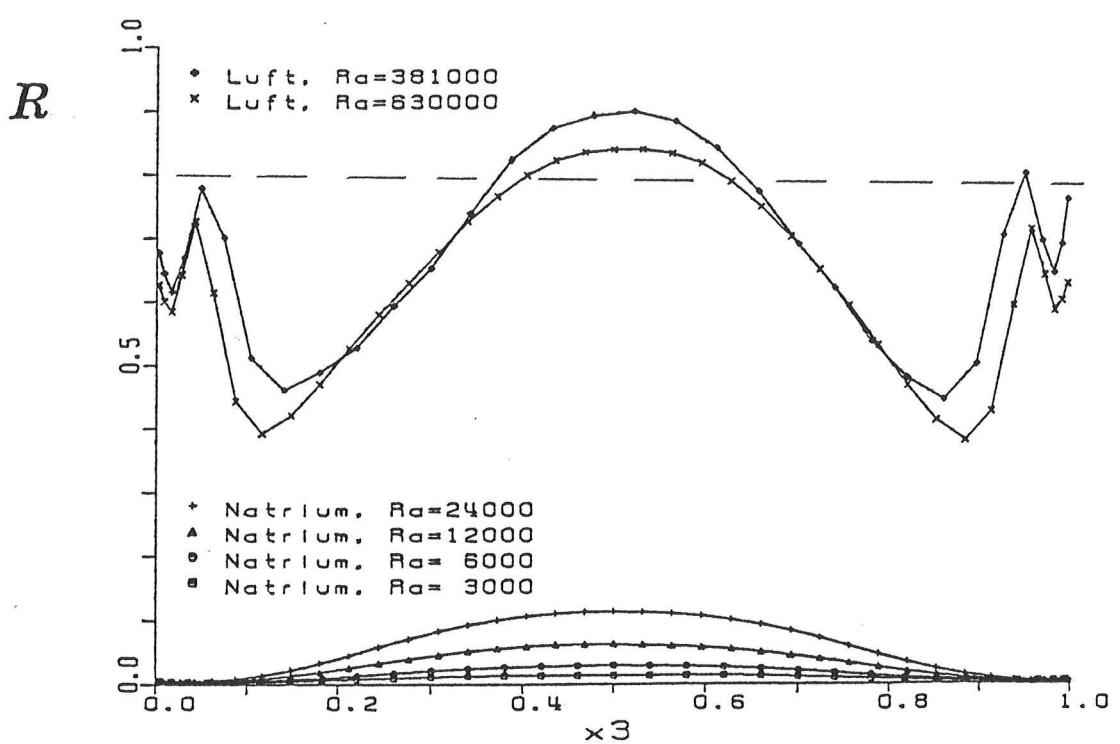
## Closure of sink term $\varepsilon_T$

- Launder (1975)

$$\varepsilon_T = \frac{1}{R} \frac{\overline{T^2}/2}{k} \varepsilon$$

$$R \approx 0.8$$

- evaluation of  $R$



## Conclusions

- Direct numerical simulation
  - no assumptions about turbulence
  - numerical experiment
  - simple geometries
  - small turbulence levels
- here: Rayleigh-Bénard convection,  
three fluids
- Validation of statistical turbulence models
  - evaluation of all correlations
  - analysis of budgets of turbulent quantities
  - analysis of model assumptions
  - calibration of coefficients and improvement of models