

# Diffraktionstomographie innerhalb der Bornschen Näherung

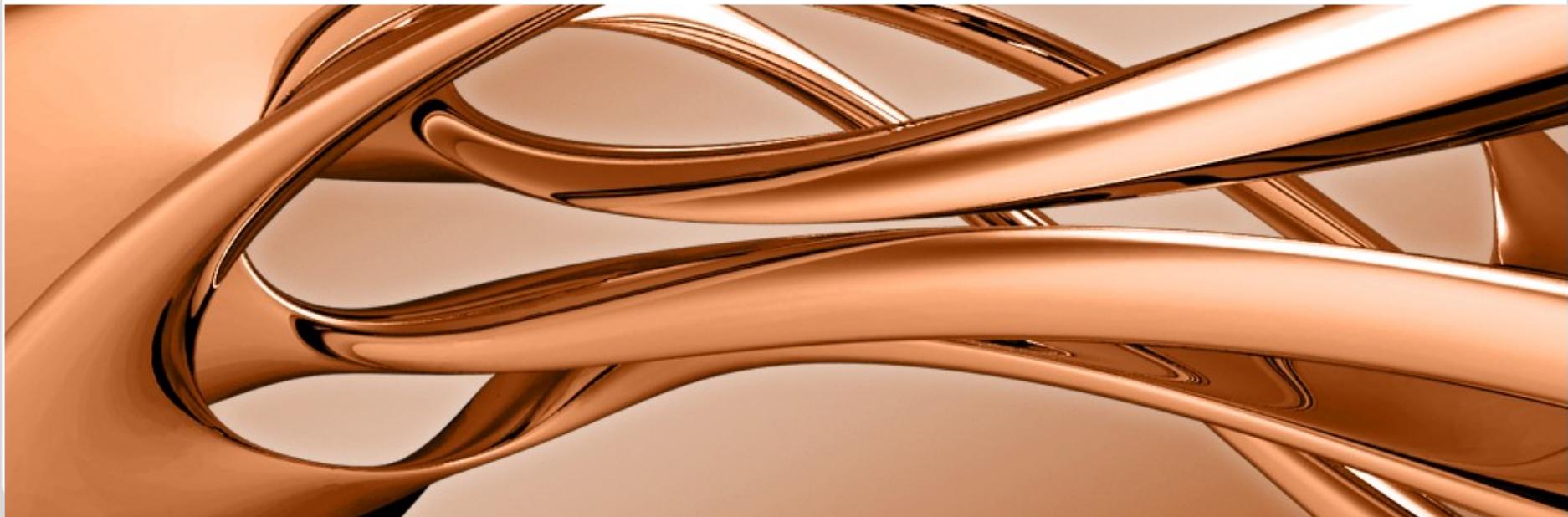
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Forschungszentrum Karlsruhe  
in der Helmholtz-Gemeinschaft



Universität Karlsruhe (TH)  
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# Inhalt

- I: Infrastruktur für die Parallelverarbeitung
- II: Physik: Lösung der fundamentalen Gleichungen
- III: Implementierung der Lösung
- IV: 1001 Parameter
- V: Mehrfachstreuung

# Part I: Infrastruktur für die Parallelverarbeitung

## ■ Ziele:

1. aus Matlab heraus
2. parallel
3. möglichst viele CPUs

zur Lösung von komplexen Aufgaben nutzbar machen.

# Infrastruktur

## Glite grid Middleware

- Viele Cluster, weltweit verteilt

- Struktur:

- Pro Grid

- Monitoring, Resource Broker
- Datenkatalog

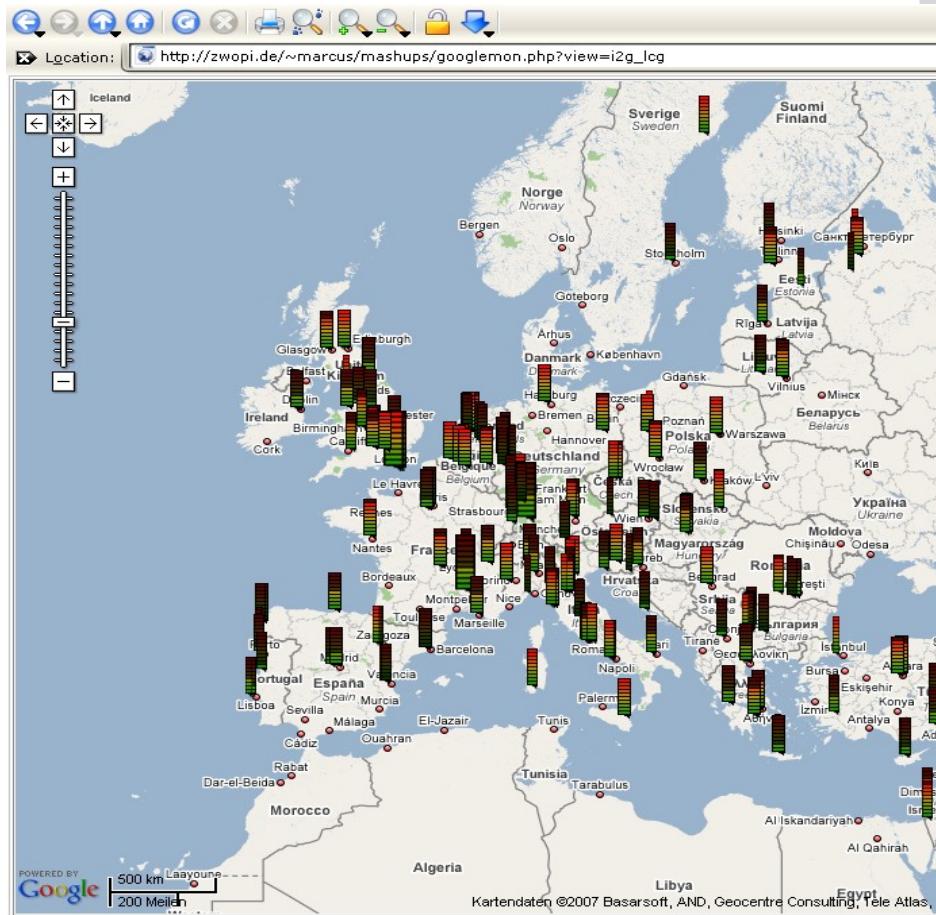
- Pro Rechenzentrum:

- 1 Compute Element (CE)  
=> N WorkerNodes (WN)

- 1 Storage Element (SE)  
=> Disk / Tape

## Eigenschaften:

- + Viele Ressourcen
- Komplexer Zugriff
- Job basiertes Paradigma
- Kein Zugriff aus Matlab



- Sites: 243 (in 49 countries)
- CPUS: 42798 (176 per site)
- RAM: 19TB
- RAM/CPU: 468MB
- DISK [Tot / Avail]: [8042TB / 5408TB] ([33892GB / 22792GB] per site)

# Verbesserung des Zugriffs

## ■ GridSolve

- Tool für (Remote Procedure Calls) RPC in global verteilten Umgebungen => Passt gut auf gLite.
- "Local" side
  - Schnittstellen für C, C++, ..., **Matlab**, Octave
- "Remote" side
  - Schnittstelle für **C** und Fortran
  - Via **Matlab Compiler** Einbindung von Matlab Code möglich

# Source code in Matlab



- Easy to use: (Example in Matlab code)

```
y=problem(x) <=> y=gs_call('problem', x)
```

- Transport input parameters to remote side
- Execute “problem”
- Transport result back
- Server executes C and Fortran libraries
- Can be extended by the C-function `system`

=> Reduce complexity of the grid to one function call

# Integration von GridSolve und gLite

## ■ Integration

- gLite jobs starten GridSolve Komponenten auf den gLite WorkerNodes (WN)

## ■ Neue Probleme:

- Infrastruktur:
  - Installation von GridSolve und Matlab Komponenten
  - Inkompatibilitäten der fundamentalen Bibliotheken (glibc)
- Schnittstelle: Muss für jede remote funktion neu definiert werden
  - Lösung: Verallgemeinerte Schnittstelle  
=> Bachelor / Masterarbeit, pending

# Part II:

## The Basic equations

- Based on fundamental equations of physics:
  - Newton's law  $F=ma$
  - Hooke's law  $F=Ds$
- We can derive a differential equation for pressure and particle velocity (wave equation):

$$\partial_k p(x) + \rho(x) \partial_t v_k(x) = f_k$$

$$\partial_k v_k(x) + \kappa(x) \partial_t p(x) = q$$

- Exact solution:

$$\hat{p}^{sct}(x) = \int G(x - x') s^2 \rho_o (\kappa_0 - \kappa_s) \hat{p}^{tot}(x') dV(x') + \int G(x - x') \partial_k \left[ \frac{\rho_0 - \rho_s}{\rho_s} \partial_k \hat{p}^{tot} \right] dV(x')$$

$$\hat{v}_k^{sct}(x) = \int G(x - x') s^2 \kappa_0 (\rho_0 - \rho_s) \hat{v}_k^{tot}(x') dV(x') + \int G(x - x') \partial_k \left[ \frac{\kappa_0 - \kappa_s}{\kappa_s} \partial_k \hat{v}_k^{tot} \right] dV(x')$$

$$\hat{p}^{tot} = \hat{p}^{sct} + \hat{p}^{inc}$$

$$\hat{v}^{tot} = - \frac{\partial_k \hat{p}^{tot}}{s \rho_s}$$



# Approximations

## ■ Exact solution:

$$\hat{p}^{sct}(x) = \int G(x - x') s^2 \rho_o (\kappa_0 - \kappa_s) \hat{p}^{tot}(x') dV(x') + \int G(x - x') \partial_k \left[ \frac{\rho_0 - \rho_s}{\rho_s} \partial_k \hat{p}^{tot} \right] dV(x')$$

$$\hat{v}_k^{sct}(x) = \int G(x - x') s^2 \kappa_0 (\rho_0 - \rho_s) \hat{v}_k^{tot}(x') dV(x') + \int G(x - x') \partial_k \left[ \frac{\kappa_0 - \kappa_s}{\kappa_s} \partial_k \hat{v}_k^{tot} \right] dV(x')$$

## ■ Born Approximation

$p^{tot} = p^{inc}$  (on the right side)

## ■ Kompressibility-only Approximation:

$$\rho_s(x) = \rho_0$$

■ Not equivalent to neglecting soundspeed Variations:

$$c = \frac{1}{\sqrt{\rho \kappa}}$$

## ■ Results in the "**forward solution**" within the Born Approximation

$$\hat{p}_{approx}^{sct} = \int G(x - x') s^2 \rho_0 (\kappa_0 - \kappa(x')) \hat{p}^{inc}(x') dV(x')$$

# The underlying Model

- Green's function contains wave response of the system:

$$G(x, w) = e^{i2\pi f * \Delta x / c_{bg}}$$

- No geometrical damping
- Hygens scattering
- Green's function can be adapted to reality

# Part II:

## The Basic equations

- Based on fundamental equations of physics:

- Newton's law  $F=ma$

- Deformation Equation  $F=Ds$

- We can derive a wave equation for pressure and particle velocity:

$$\partial_k p(x) + \rho(x) \partial_t v_k(x) = f_k \quad (\nabla^2 - k^2) \hat{u}_k(x) = \hat{f}_k \hat{u}_k(x)$$

$$\partial_k v_k(x) + \kappa(x) \partial_t p(x) = q \quad (\text{Inhomogeneous Helmholtz Equation})$$

- Exact solution:

$$\begin{aligned} \hat{p}^{sct}(x) &= \int G(x - x') s^2 \rho_o (\kappa_0 - \kappa_s) \hat{p}^{tot}(x') dV(x') + \int G(x - x') \partial_k \left[ \frac{\rho_0 - \rho_s}{\rho_s} \partial_k \hat{p}^{tot} \right] dV(x') \\ \hat{v}_k^{sct}(x) &= \int G(x - x') s^2 \kappa_0 (\rho_0 - \rho_s) \hat{v}_k^{tot}(x') dV(x') + \int G(x - x') \partial_k \left[ \frac{\kappa_0 - \kappa_s}{\kappa_s} \partial_k \hat{v}_k^{tot} \right] dV(x') \end{aligned}$$

### 3.3 Single Step Inversion

The error made within this Approximation can be written as:

$$ERR = \sum_S \sum_R \sum_\omega \left| \hat{p}_{real}^{sct} - \underbrace{\int G(x-x') s^2 \rho_0 \chi(x') \hat{p}^{inc}(x') V(x')}_{=\hat{p}_{approx}^{sct}} \right|^2 \quad \text{with: } \chi(x) = \kappa_0 - \kappa_s(x)$$

$$= \Delta\chi(x) \cdot \alpha$$

This Error is positive definite, hence minimal, when the derivative against the variation parameter  $\frac{\partial ERR}{\partial \alpha} = 0$ . Solving this for  $\alpha$  gives:

$$\alpha = \frac{Re \left\{ \sum_{S,R,\omega} \hat{p}^{sct} \int (s^2 \rho_o \Delta\chi G \hat{p}^{inc})^* dV \right\}}{\sum_{S,R,\omega} \left| \int \Delta\chi G \hat{p}^{inc} s^2 \rho_0 dV \right|^2}$$

$\alpha$  is maximal, if the numerator is maximal. This is the case when  $\Delta\chi$  is parallel to the rest of the expression in the numerator. I.e.

$$\Delta\chi(x) = \sum_{S,R,\omega} (G(x-x') \hat{p}^{inc}(x'))^* \hat{p}^{sct}(x') \quad (11)$$

# Inversion

## ■ Inversion via "Backpropagation" (BP)

$$\Delta\chi(x) = \sum_{S,R,\omega} (G(x - x') \hat{p}^{inc}(x'))^* \hat{p}^{sct}(x')$$

- Equivalent to first step of gradient method
- Similar result as for SAFT **but in frequency domain**
  - Advantage:
    - Frequency dependend corrections can be used
  - Disadvantage:
    - Each frequency needs to be computed individually  
=> Factor  $\sim 500$  slower
- Exploiting similarity of codes:
  - Re-use assembler parts of SAFT
  - Can benefit from preprocessing codes of SAFT

# Part III: Implementation

- Data access mode
  - Data parallel
    - => Send **subset of data** to a remote computer
    - => **Return all voxels** of image
    - => **Add images** locally
      - + Re-use Assembler optimisations
      - Large data transfers for large images
  - Volume parallel
    - => Send **all data** to each remote computer
    - => **Return subset** of voxels
    - => Combine **subimages** locally
    - One big data transfer at experiment start
      - + Re-use input data for several experiments
      - Rewrite Assembler code
- Code Snippets
  - Forward Solution Kernel
  - Backward Solution Kernel

# Implementation

## ■ Forward Solution Kernel (data parallel)

```

%% Forward Solution:
for emit = emitterStart:emitterEnd
  for rec = 1:n_rec
    for freq = freq_start:freq_end
      pSctTemp = 0;
      i_zwopi_f_by_cbg = i * 2 * pi * frequency / c_bg;
      for x=1:x_max
        for y = 1:y_max
          dist_src = sqrt((x - x_src(emit))^2 + (y - y_src(emit))^2);
          dist_rec = sqrt((x - x_rec(rec)) ^2 + (y - y_rec(rec)) ^2);

          greenValue = exp(-i_zwopi_f_by_cbg * (dist_rec + dist_src));

          pSctTemp = pSctTemp + (chi(x,y) * greenValue) ...
            * (p_inc_ft(emit, rec, freq) + p_sct_ft(emit, rec, freq));
        end
      end
      pSct(emit,rec,freq) = pSctTemp;
    end
  end
end

```

# Implementation

## ■ Backpropagation (BP) Solution Kernel (Volume parallel)

```

%% Backward Solution
for x=x_start:x_end;
    for y=y_start:y_end
        chi_tmp = 0;
        for freq = freq_start:freq_end
            i_zwopi_f_by_cbg = i * 2 * pi * frequency / c_bg;
            for emit=1:n_src
                for receiver=1:n_rec
                    dist_src = sqrt((x - x_src(emit))^2 + (y - y_src(emit))^2);
                    dist_rec = sqrt((x - x_rec(receiver))^2 + (y - y_rec(receiver))^2);

                    green = exp(i_zwopi_f_by_cbg*(dist_rec+dist_src));

                    chi_tmp = chi_tmp + green ...
                        * data_sct_ft(emitter, receiver, omega);
                end
            end
        end
        chi (x-x_start+1, y-y_start+1) = chi_tmp;
    end
end

```

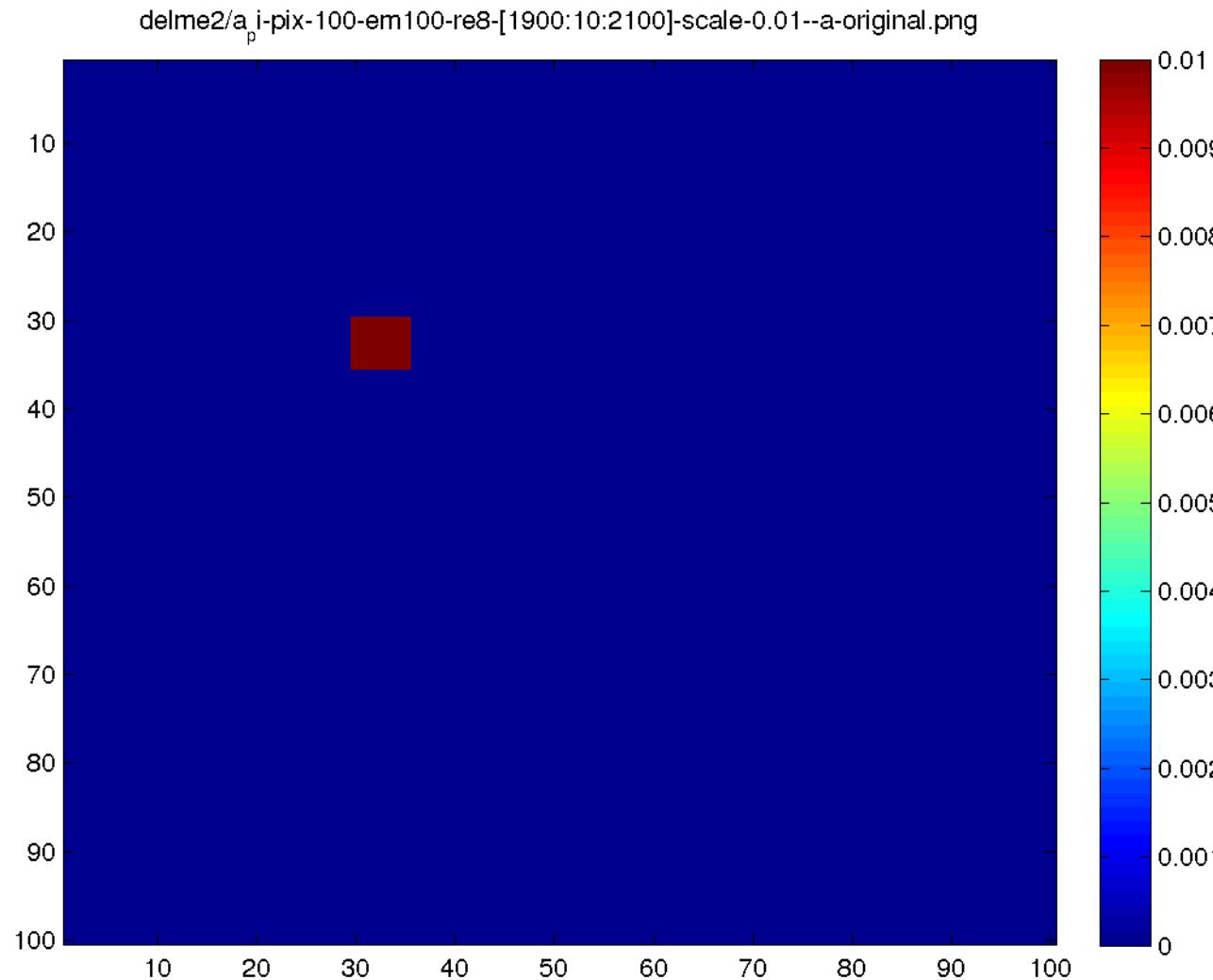
# Implementation

- Data parallel mode was accelerated in Assembler
  - Speedup ~ factor 50 (thanks to Michael Zapf)
  
- We can now
  - Run simulations to create A-Scans (for given potential)
  - Run reconstructions to find original potential (from A-Scans)

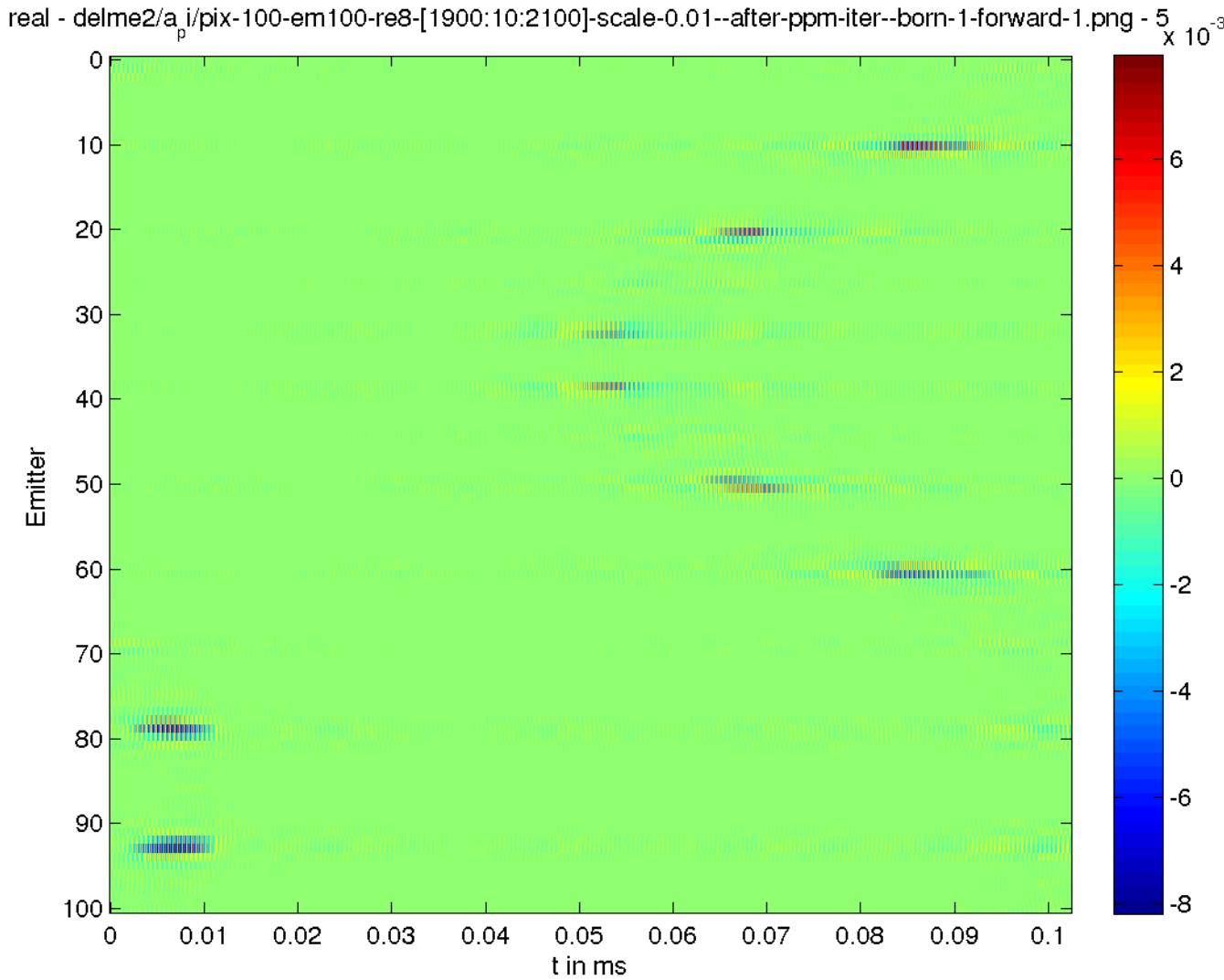
# Implementation

- Data parallel mode was accelerated in Assembler
  - Speedup ~ factor 50 (thanks to Michael Zapf)
- We can now
  - Run simulations to create A-Scans (for given potential)
  - Run reconstructions to find original potential (from A-Scans)
- First images looked like these

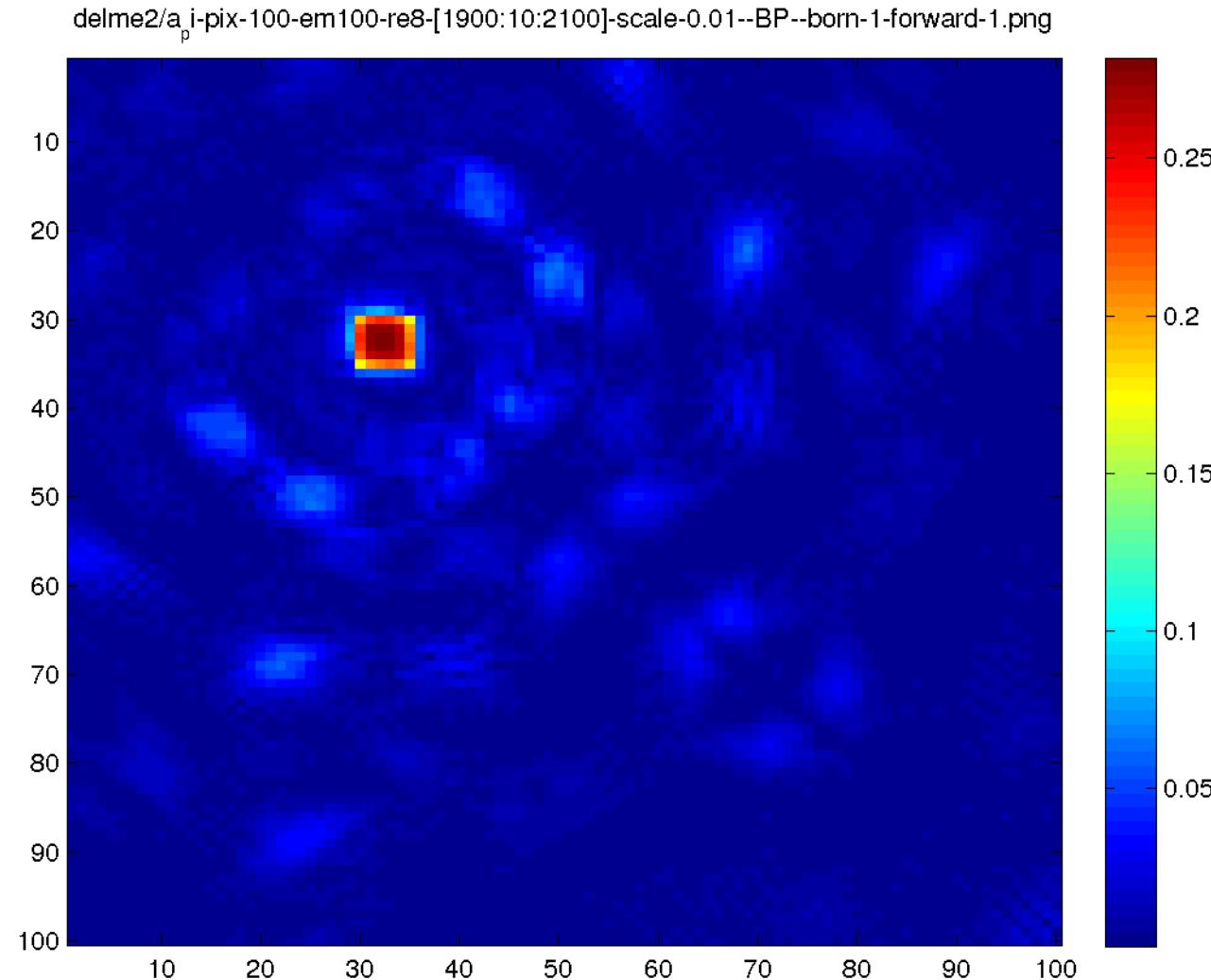
# Original Potential



# B-Scan

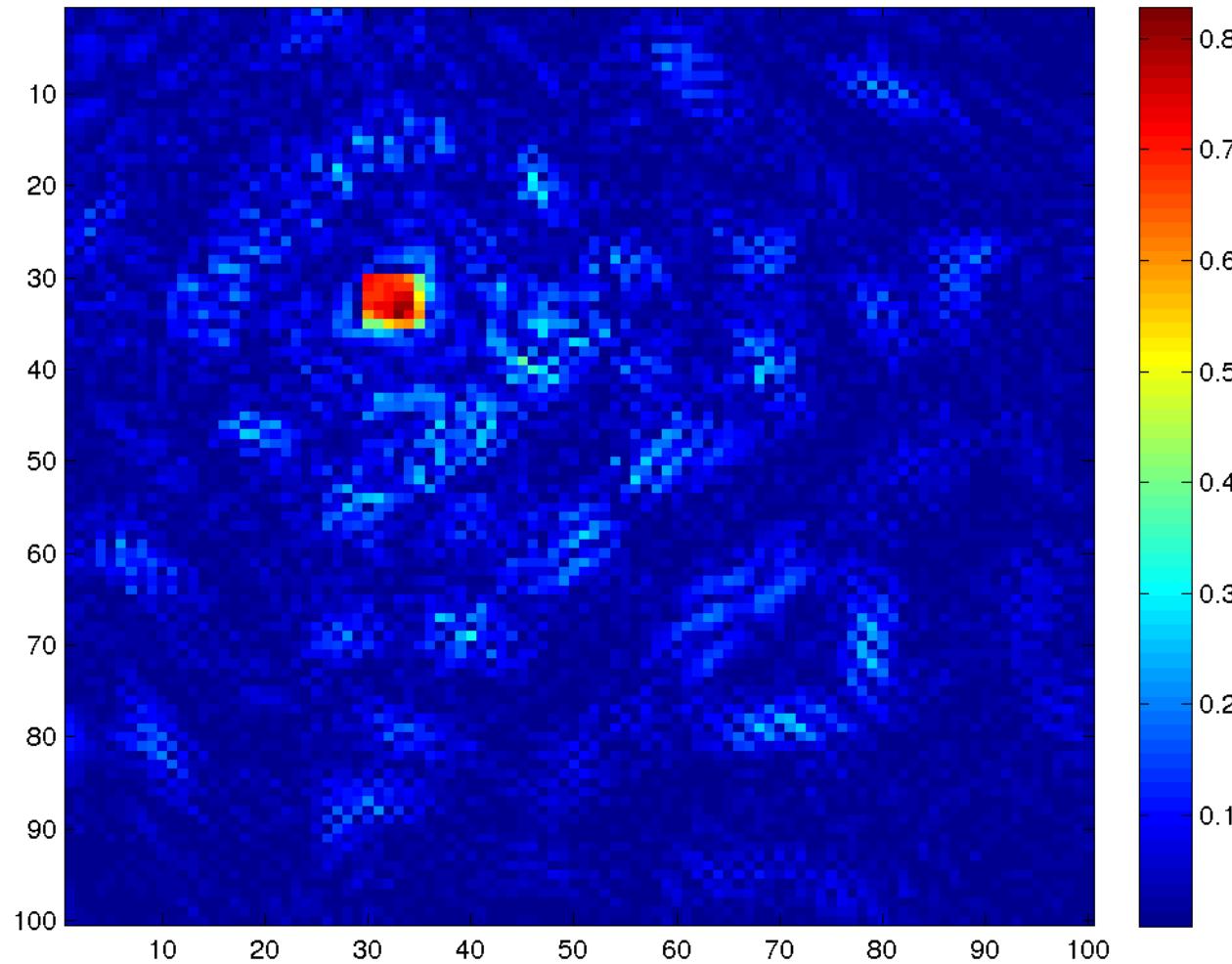


# Reconstruction via backprojection



# Reconstruction via SAFT

delme2/a\_i-pix-100-em100-re8-[1900:10:2100]-scale-0.01--SAFT-imag--born-1-forward-1.png



# Results from first images

- SAFT produces similar images like BP
- Physics and Signalprocessing Theory not respected so far
  - Artefacts
  - No correlation between simulation parameters and reality
  - Impossible to compare results
- Consequence
  - Major rewrite of code (to properly respect reality)
  - Improved understanding of parameters and artefacts
  - Limits for the frequencies used

# Limits for the frequency range ( $c = \lambda f$ )

- Diameter of USCT

$$\lambda_{max} < D$$

- Pixelsize

$$\lambda_{min} > 2 \cdot dr$$

- Samplingrate (Nyquist)

$$f_{max} < \frac{1}{2} f_{sample}$$

- Spacial Nyquist

$$f_{max} < \frac{c \cdot ROI}{2\sqrt{2} \cdot dr \cdot D}$$

- Grating Lobes

$$\lambda_{min} > 2 * d_{Sensor-Sensor} = \frac{2\pi \cdot D}{N_{Sensors}}$$

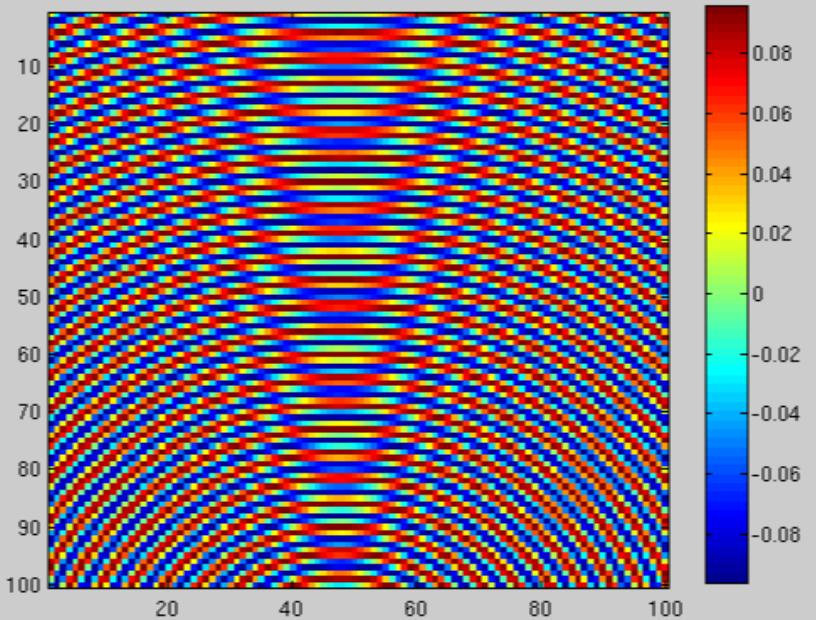
- Center frequency

- In range of simulation frequencies

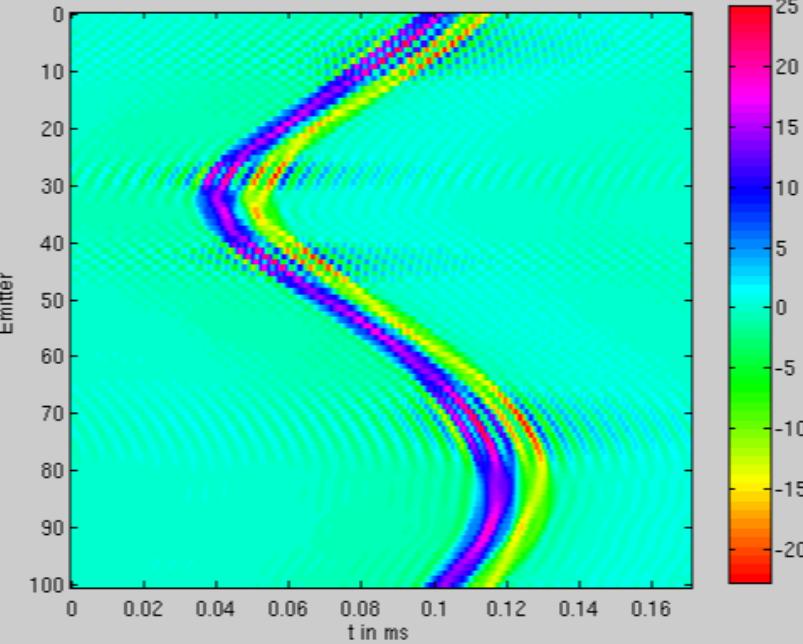
- Bandwidth

- Small bandwidth => Long signals => Bad Images

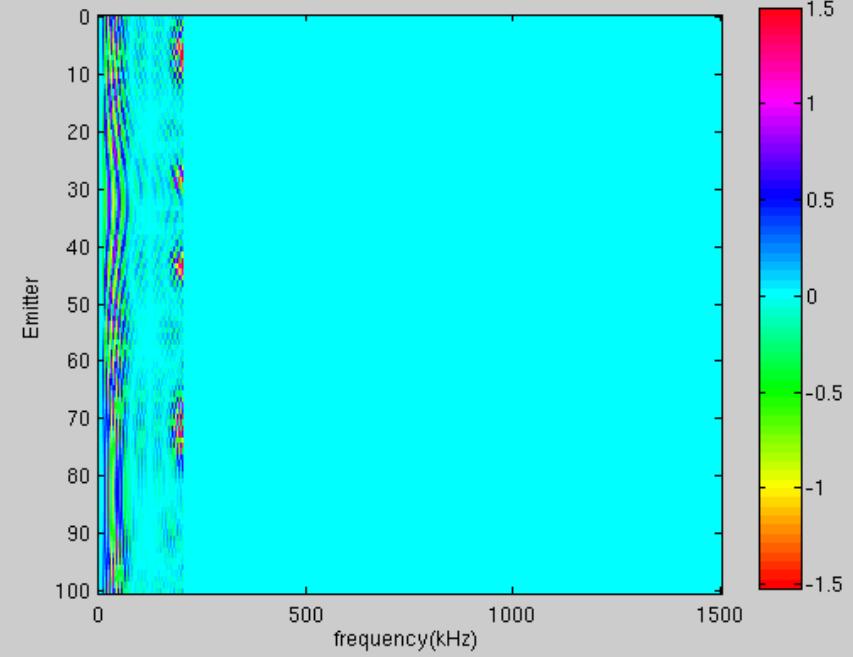
freq[66]: 208.33 KHz.png



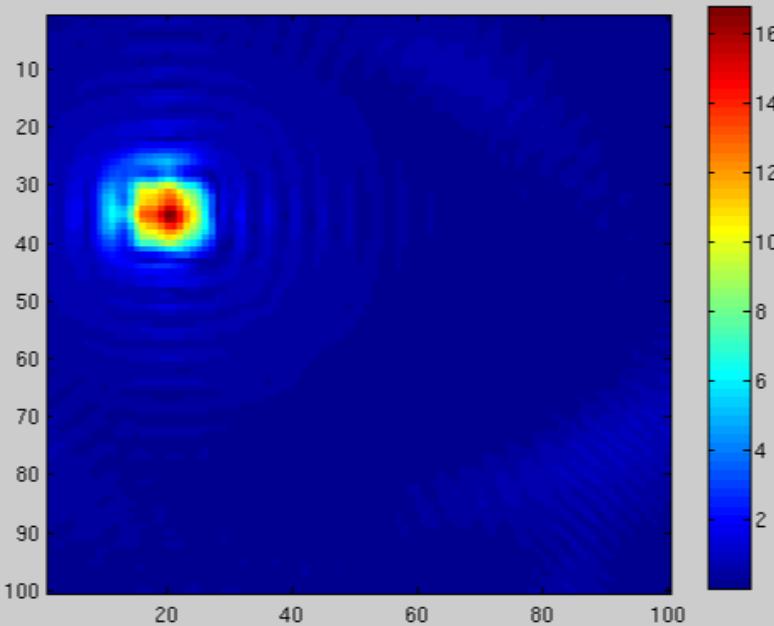
text3/a-100-e100-r1-[12.5:208.3]-sr3.00MHz-cf500000.00-bw500000.00--ba1-fw1.png-re

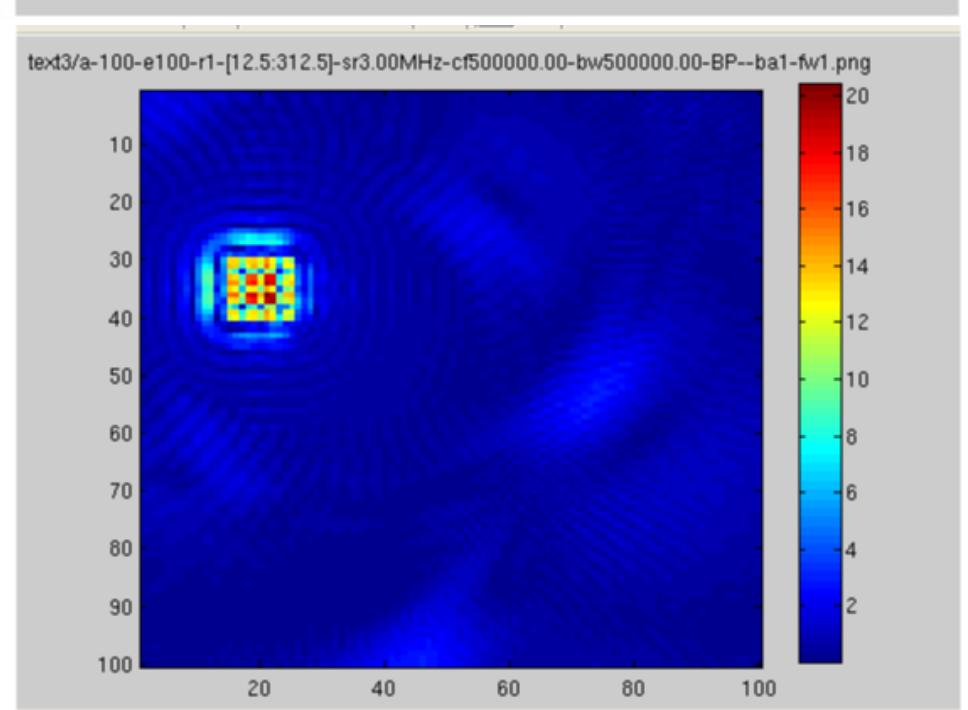
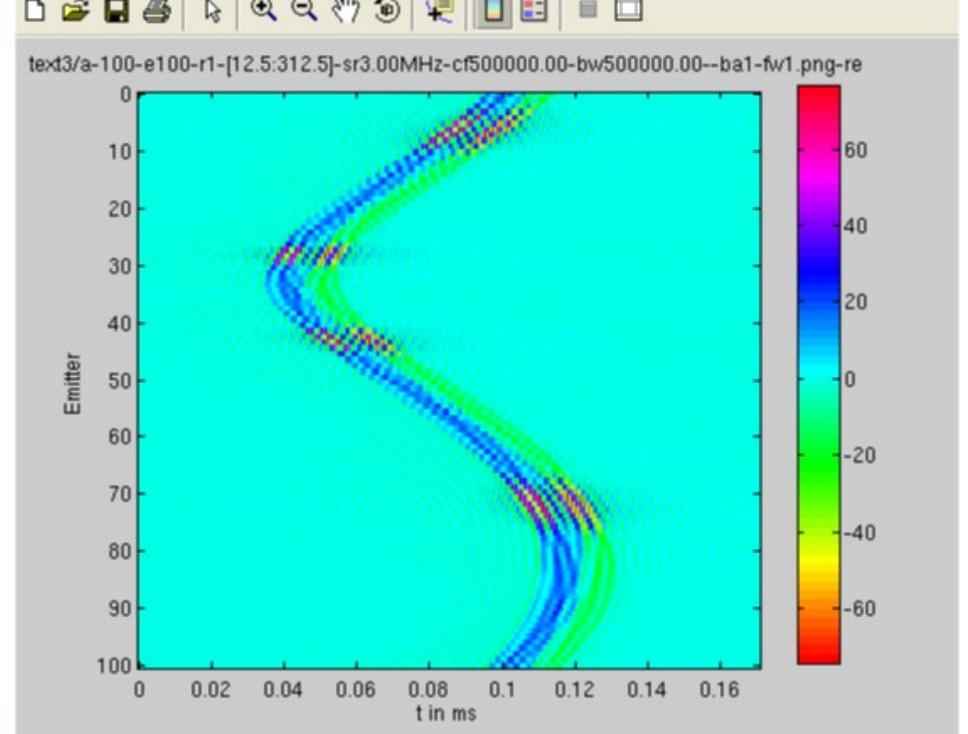
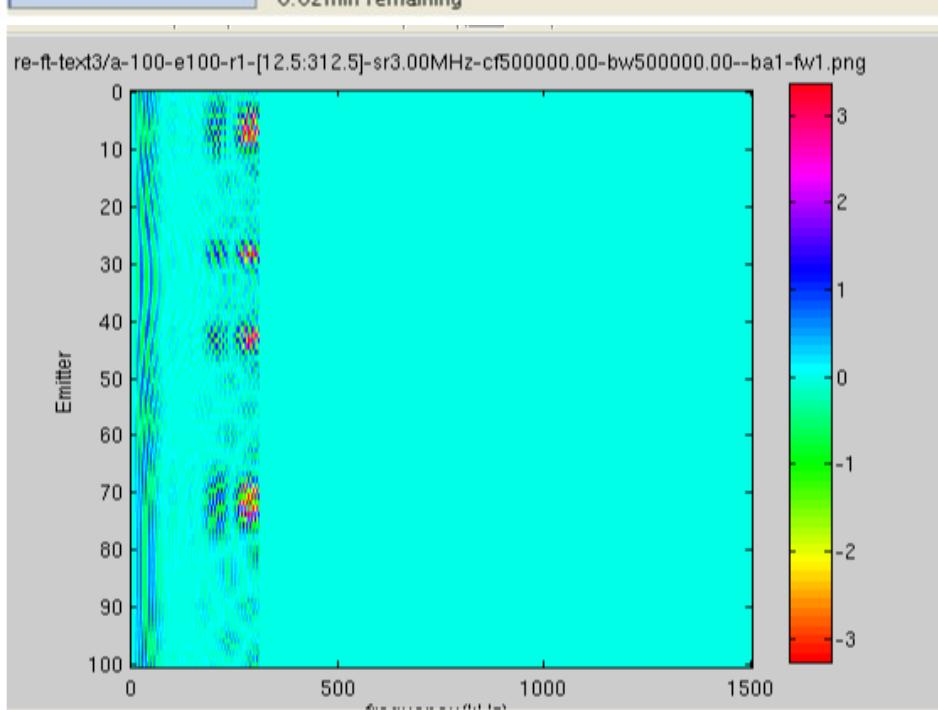
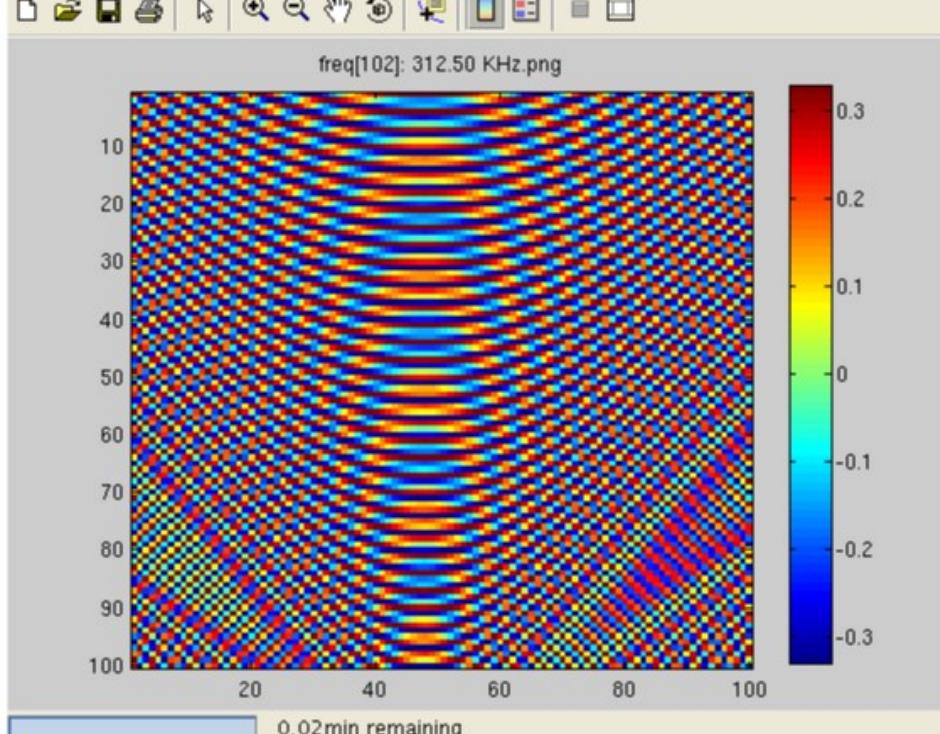


re-ft-text3/a-100-e100-r1-[12.5:208.3]-sr3.00MHz-cf500000.00-bw500000.00--ba1-fw1.png

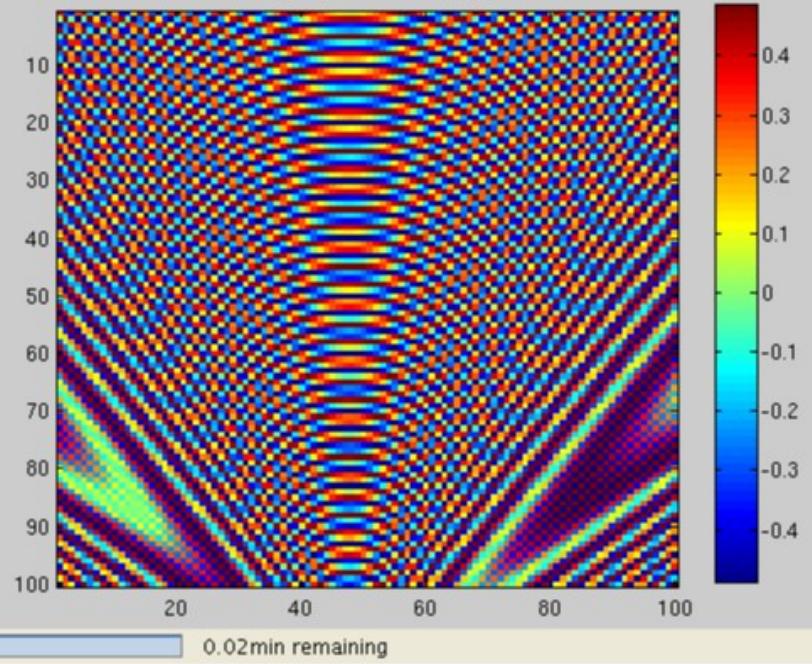


text3/a-100-e100-r1-[12.5:208.3]-sr3.00MHz-cf500000.00-bw500000.00-BP--ba1-fw1.png



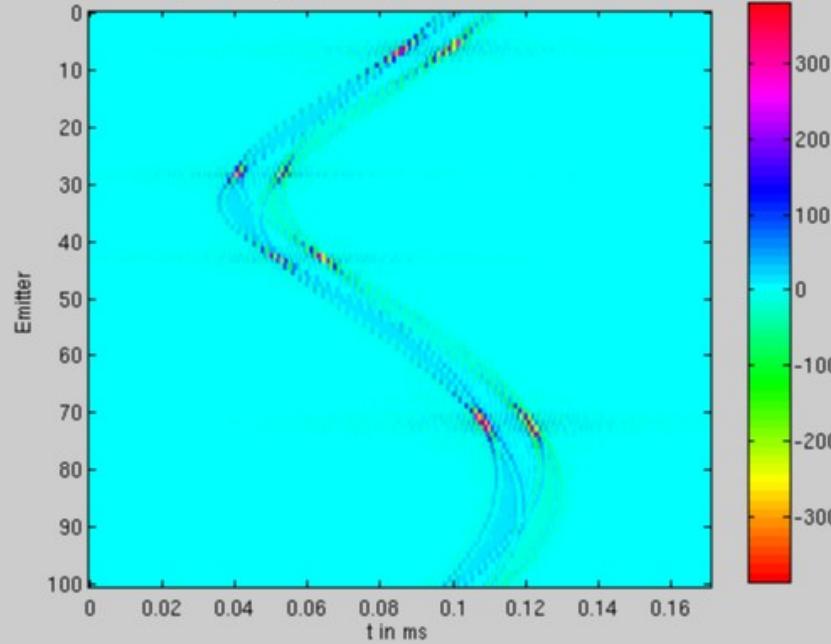


freq[208]: 625.00 KHz.png

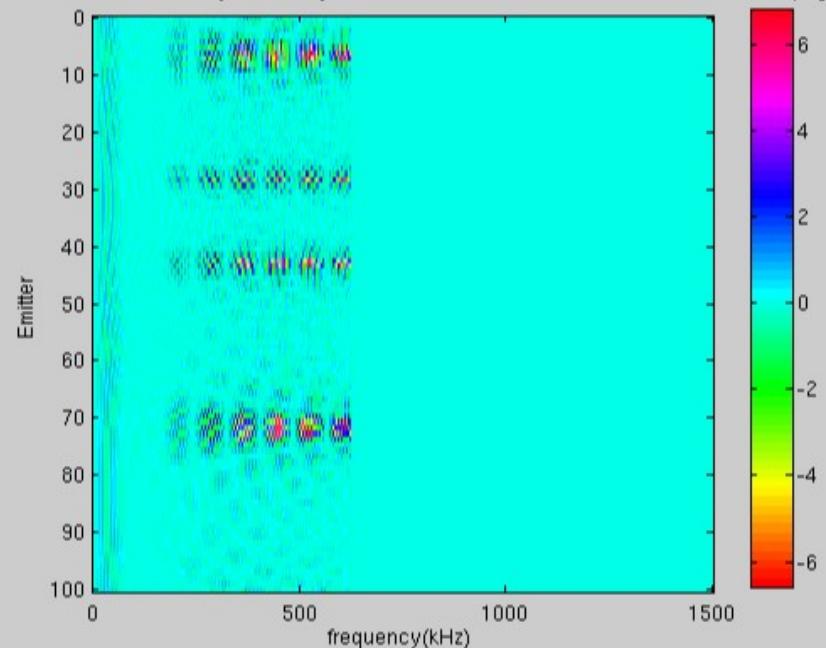


0.02min remaining

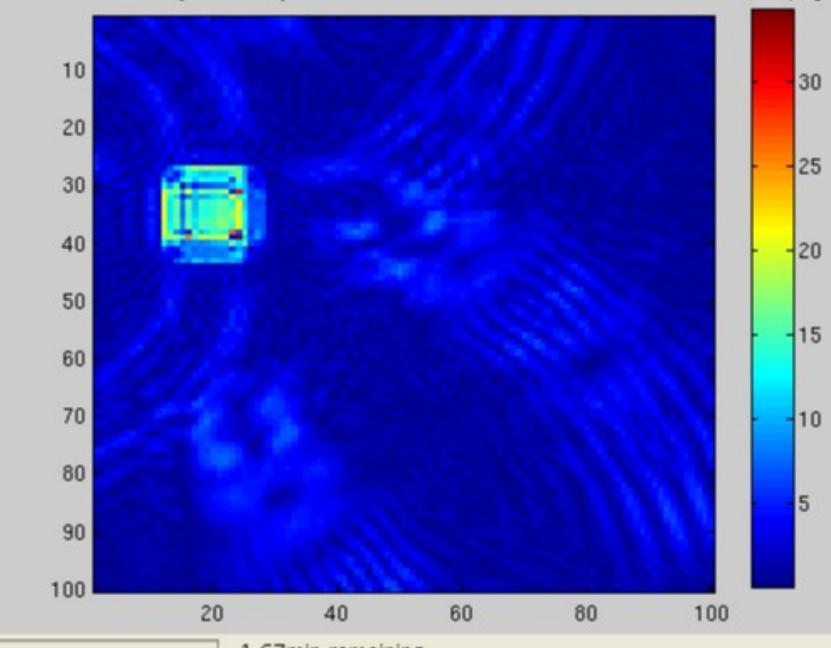
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re-ft-text3/a-100-e100-r1-[12.5:625.0]-sr3.00MHz-cf500000.00-bw500000.00--ba1-fw1.png

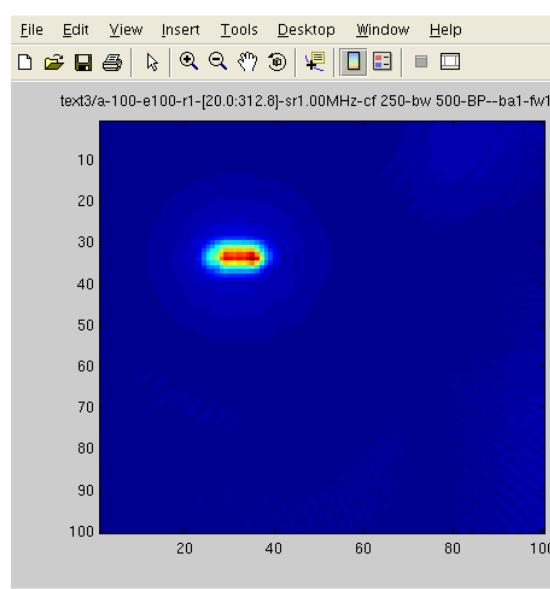
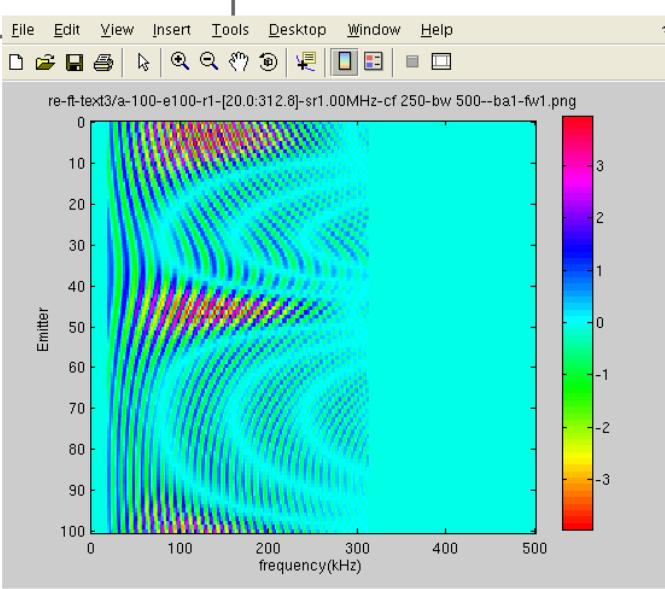
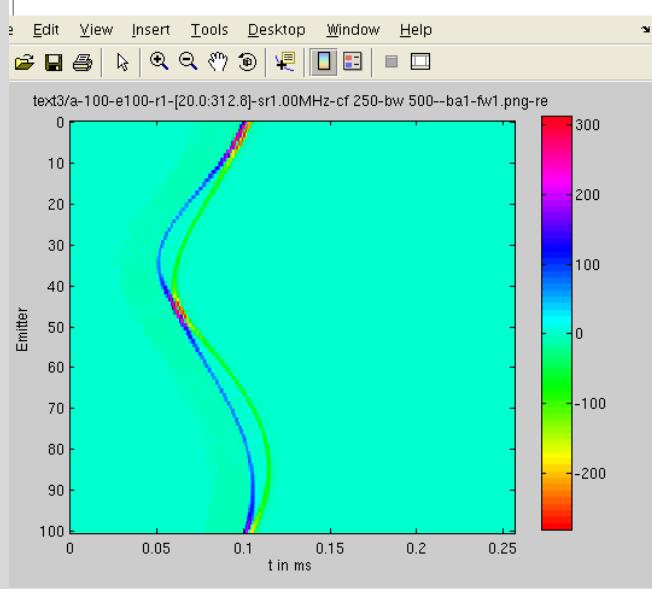
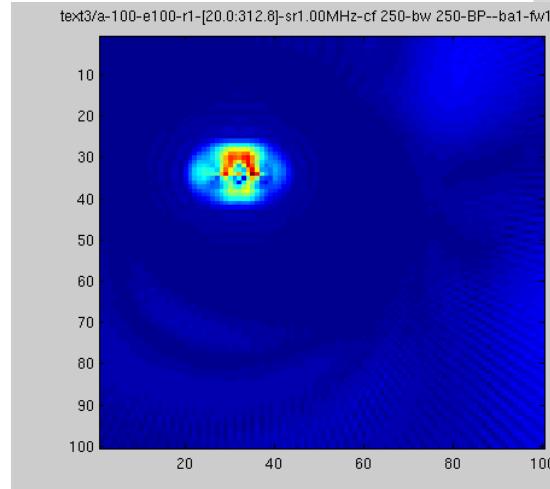
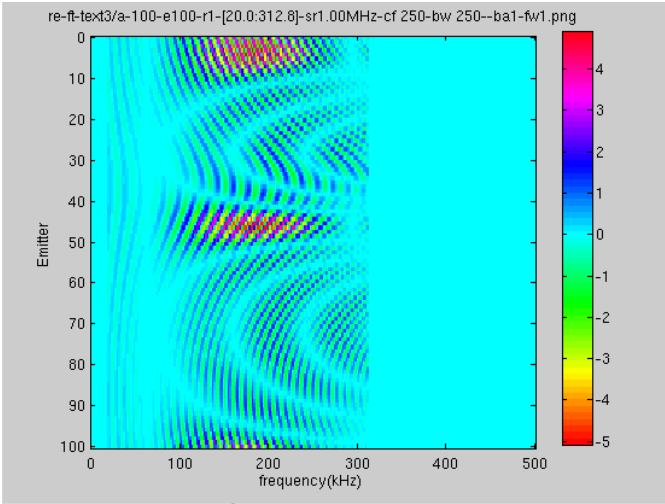
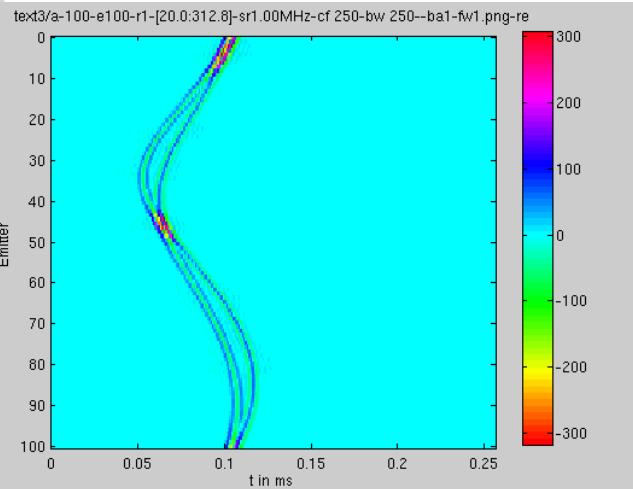


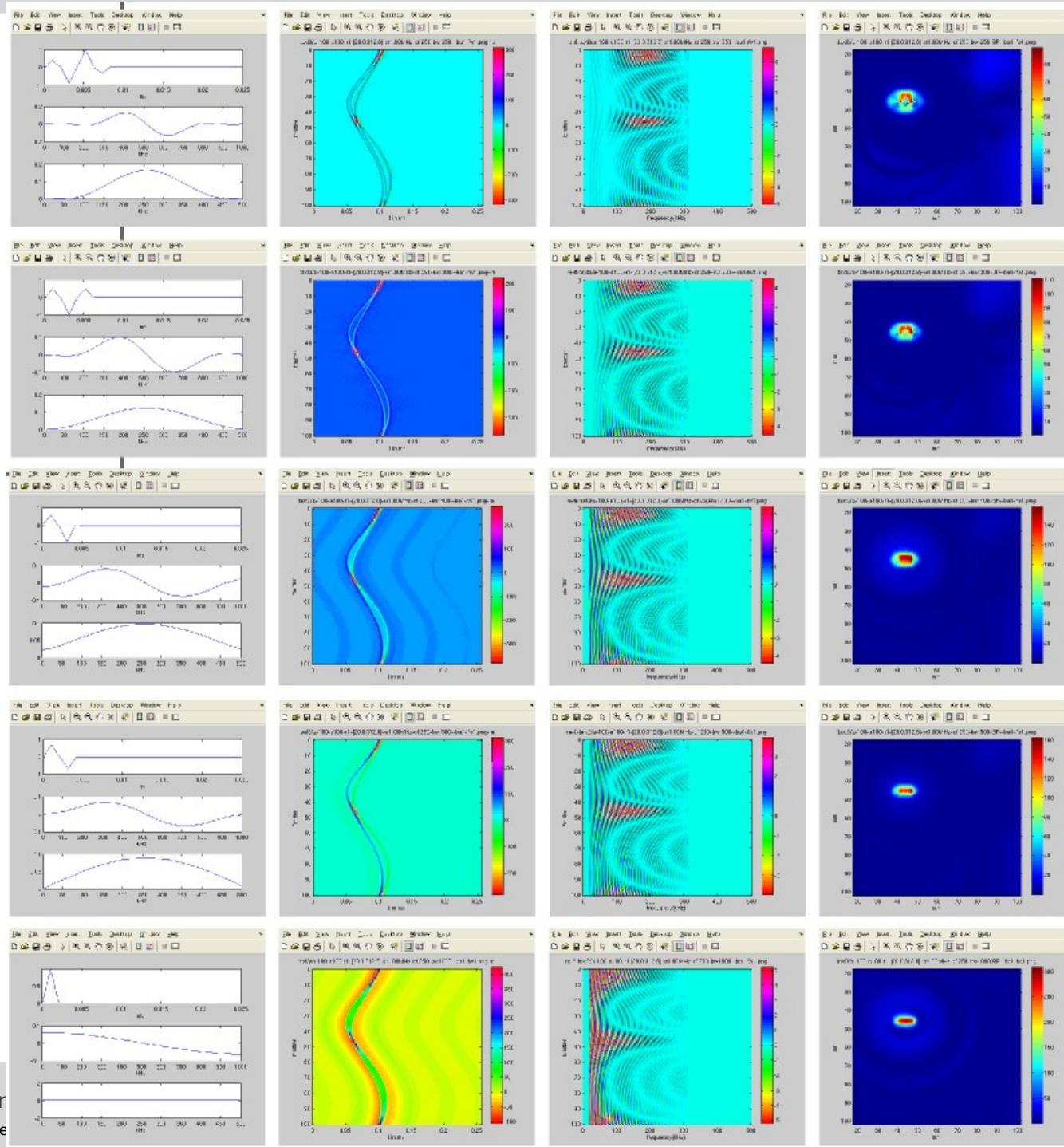
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1.67min remaining

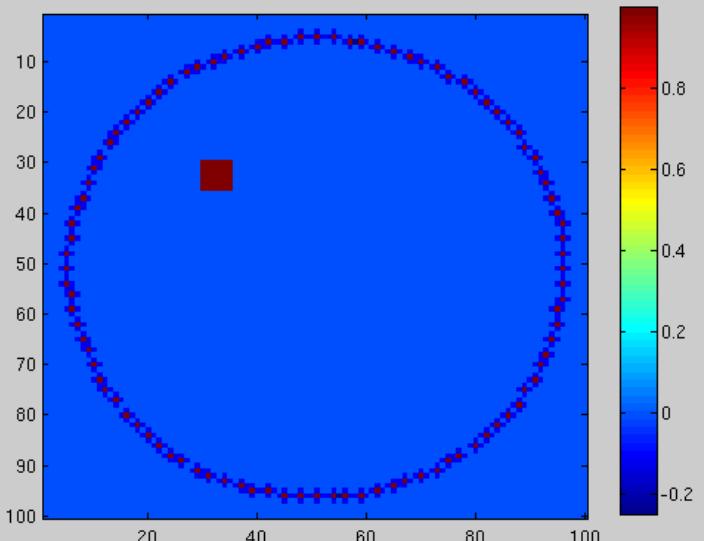
# Variation of the Bandwidth



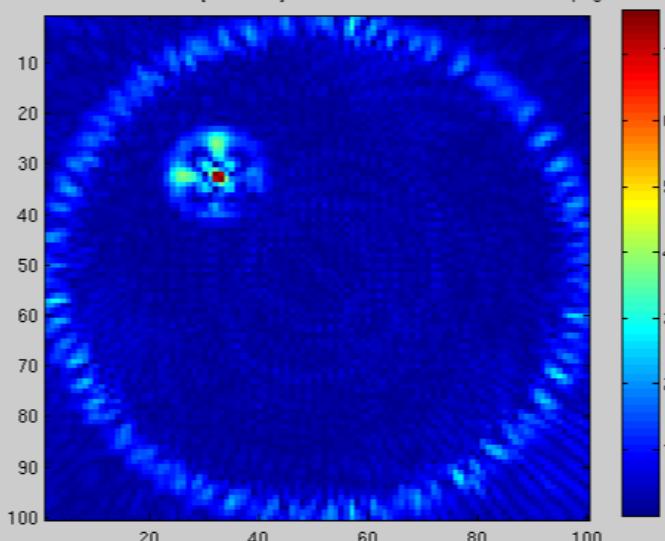


# Parameters defined for verification with Brno

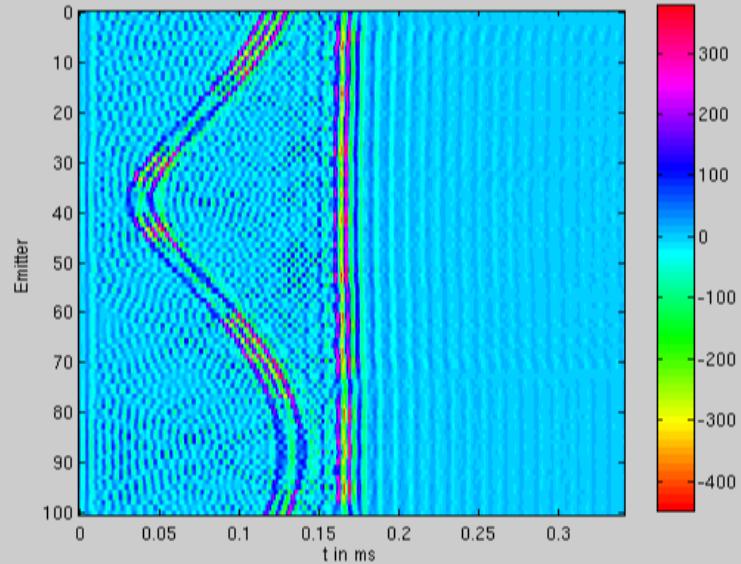
- Samplerate: 1 MHz
- Resolution: 1 x 100 x 100
- Centerfrequency: 250 kHz
- Bandwidth: 500 kHz
- Simulation range: 20 – 313 kHz (149 steps)



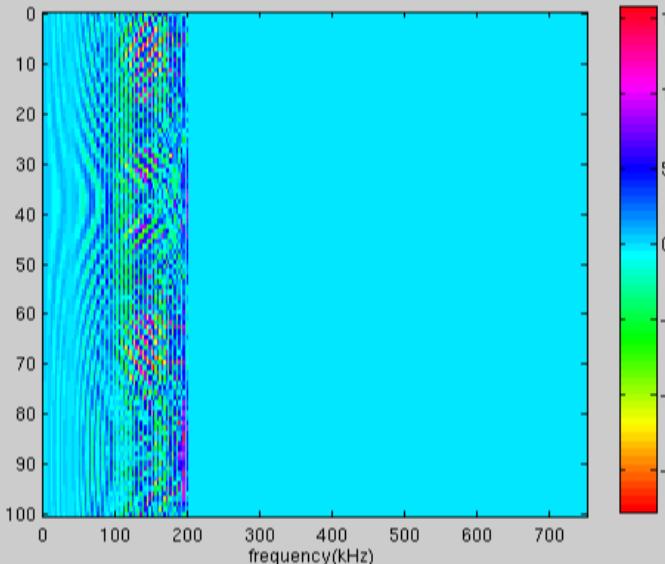
text3/a-100-e100-r1-[1.0:200.0]-sr1.50MHz-BP--born-1-forward-1.png



real - text3/a-100-e100-r1-[1.0:200.0]-sr1.50MHz-before-ppm-iter--born-1-forward-1.png - 1



real-ft - text3/a-100-e100-r1-[1.0:200.0]-sr1.50MHz-before-ppm-iter--born-1-forward-1.png - 1



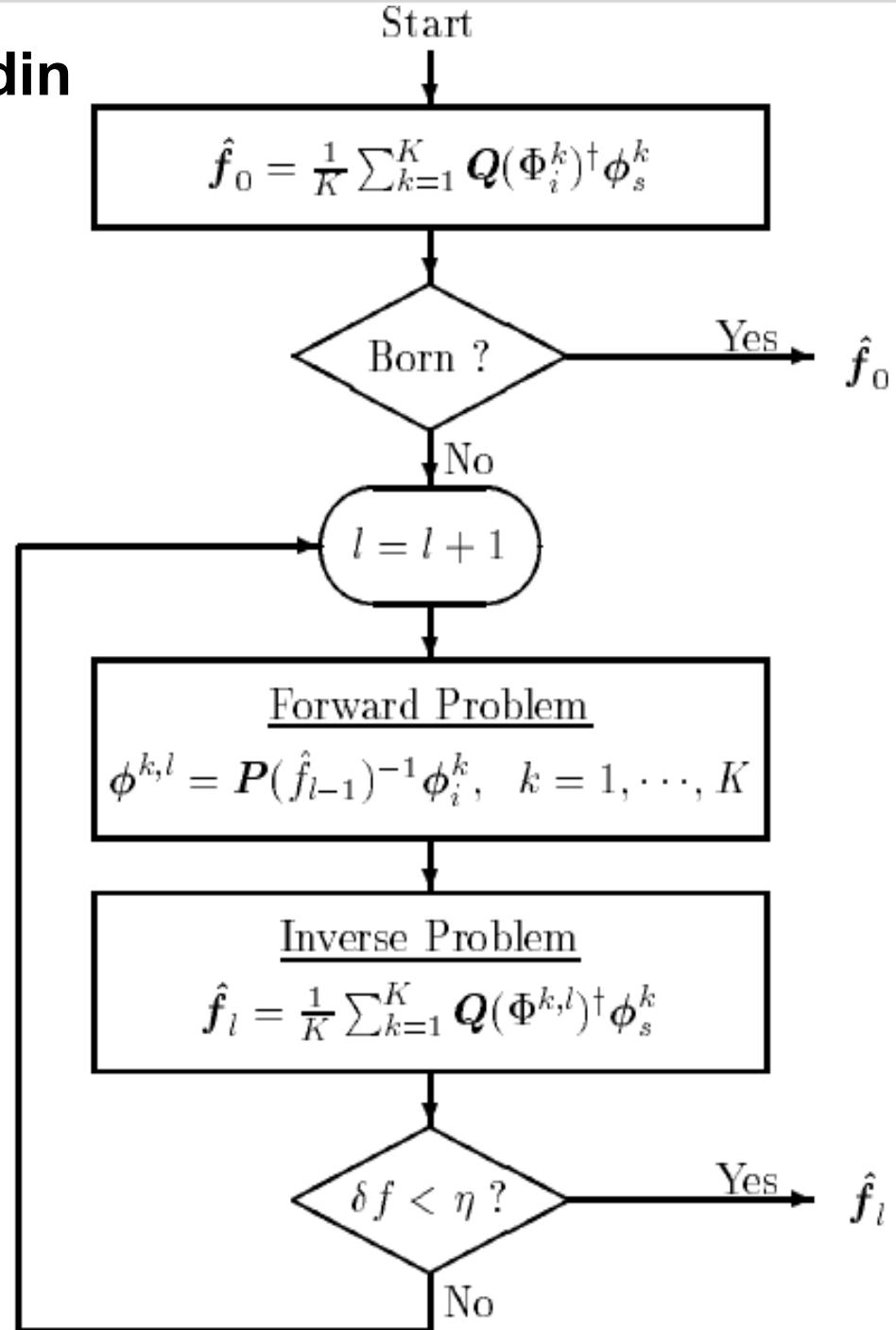
# Part IV: Multiple Scattering (work in progress)

- In principle (=literature) very simple:
  - Iterations over forward and backward solution
  - Changes are propagated via updates to the image and via the measured field

# Based on PhD of O.S.Haddadin

1. Compute inverse solution (image) from measured data
2. Compute forward solution to update the total pressure
3. Compute inverse solution with updated pressure

I know, this causes nightmares for experimentalists...



# My formulation of the iterative scheme

- Correction of the assumption of the total field by solving the forward problem
- Re-computing the inverse problem

$$\hat{p}_{\kappa, \text{Born}}^{sct, (l)} = \int G \chi_{\kappa}^{(l-1)} (\hat{p}^{\text{inc}} + \hat{p}_{\kappa, \text{Born}}^{sct, (l-1)}) dV$$

$$\Delta \chi_{\kappa}^{(l)} = \sum_{S, R, \omega} (G \hat{p}^{\text{inc}})^* \hat{p}_{\kappa, \text{Born}}^{sct, (l-1)}$$

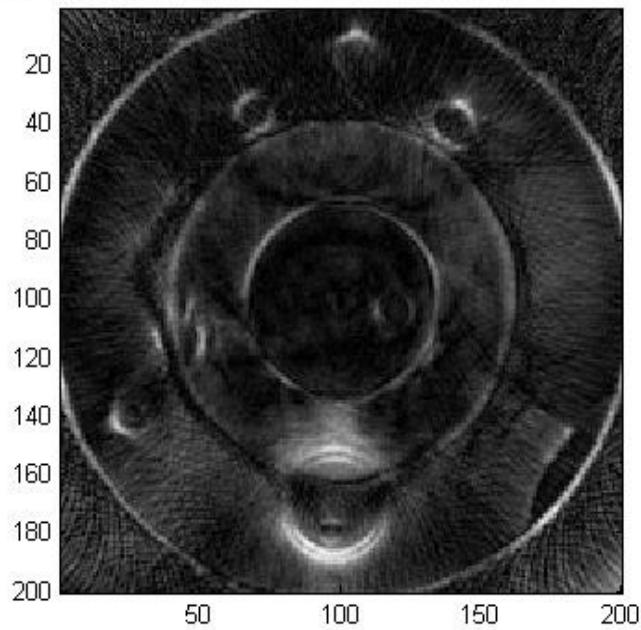
Iterations  $l=1, \dots, L$

- This is basically "averaging"

# First Results

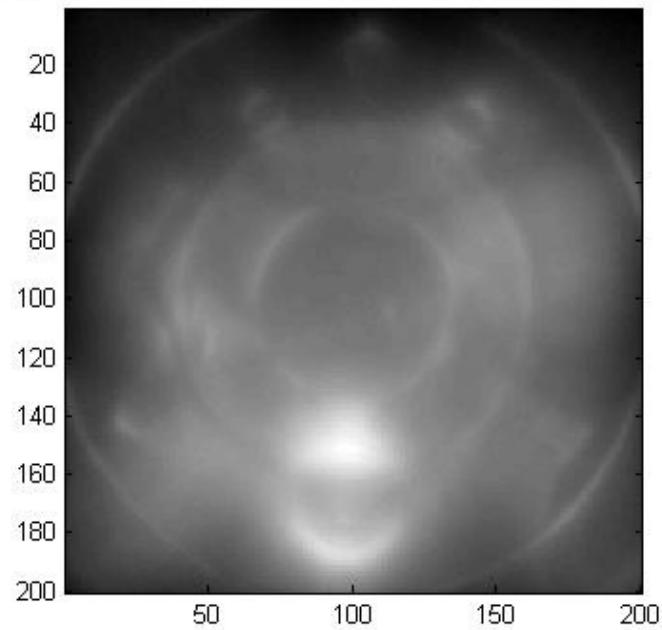
## „Milk-effect“ 😞

n\_rec:n\_src [5:100], backpropagation, omega [1:1:200] - string



Iteration 0

n\_rec:n\_src [5:100], backpropagation, omega [1:1:200] - string



Iteration 1

# Current work

- Modification of the **forward solution**:

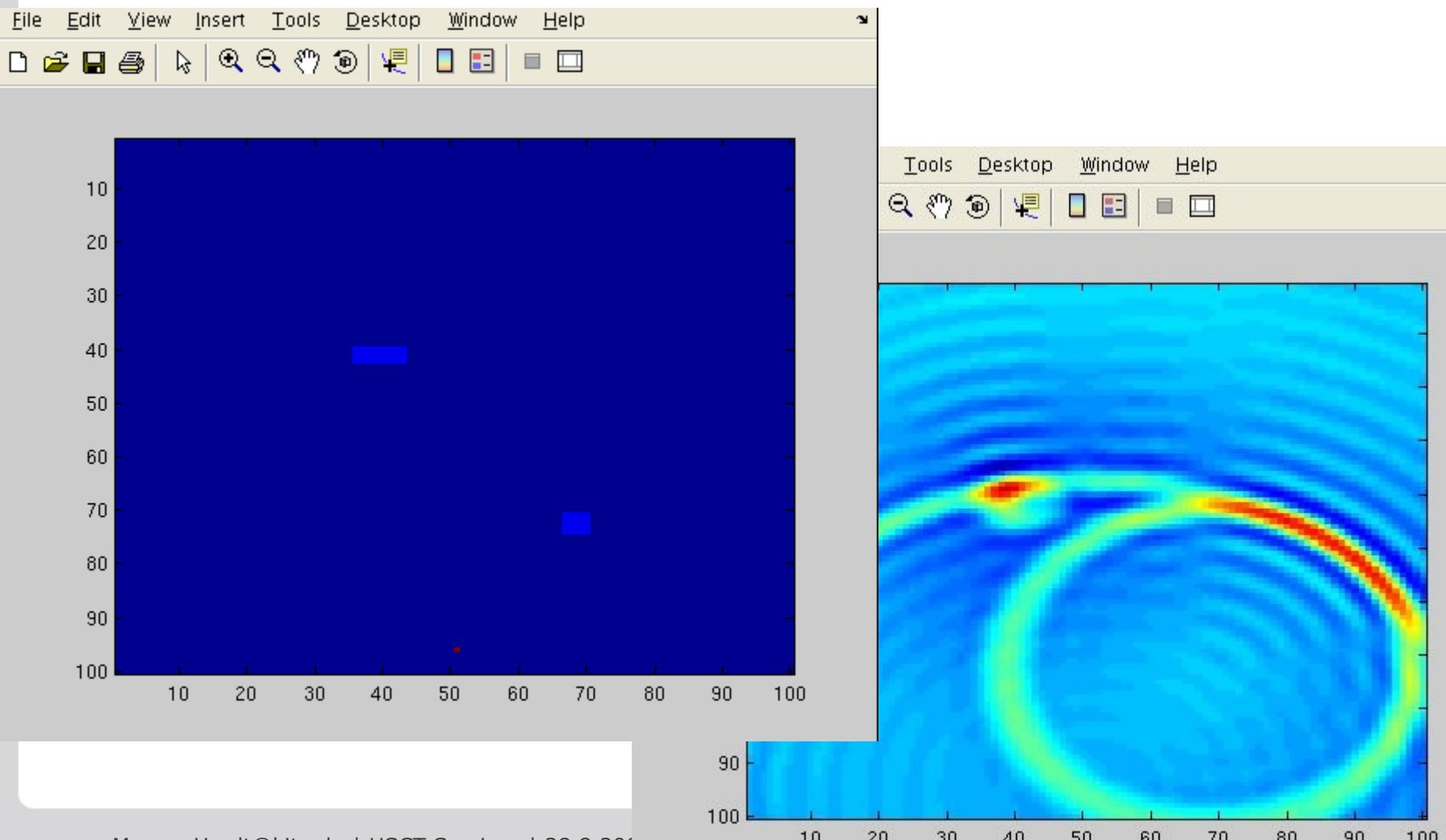
- Interpretation of:

$$\hat{p}_{\kappa, Born}^{sct, (l)} = \int G \chi_{\kappa}^{(l-1)} (\hat{p}^{inc} + \hat{p}_{\kappa, Born}^{sct, (l-1)}) dV$$

- Previously: Assumed surface integral over "dV"
  - Now: Volume Integral inside "dV"  
=> Simulate pressure at every pixel  
(instead of: at every receiver)
  - For 1 x 100 x 100 Voxels: Factor 10.000 slower
- => Able to "Send off" secondary waves in **forward iterations**

# Example

■ <file:///home/marcus/grid/talks/2009-08-28-IPE-Seminar>



# Conclusions

- Parallel computing facilities integrated into Matlab
- Simulations correlate to realistic parameters
- Reconstructions of BP and SAFT show similar images
- First order scattering can be simulated
  - Higher order scattering in progress
- No noise yet
- No Matrix operations required

# Questions?