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**Forschungszentrum Karlsruhe**

Technik und Umwelt

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Institut für Reaktorsicherheit

# **Volume averaged conservation equations for volume-of-fluid interface tracking**

M. Wörner, W. Sabisch, G. Grötzbach, D.G. Cacuci

*Research Center Karlsruhe, Institute for Reactor Safety, Germany*

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## Problem and motivation

- **VOF computations in literature:**  
finite difference discretization of *local* equations

$$\frac{\partial f}{\partial t} + \nabla \cdot f \mathbf{v} = 0 \quad (0 \leq f \leq 1), \rho_m = f\rho_1 + (1-f)\rho_2, \mu_m = f\mu_1 + (1-f)\mu_2$$

$$\frac{\partial \rho_m \mathbf{v}}{\partial t} + \nabla \cdot \rho_m \mathbf{v} \mathbf{v} = -\nabla p + \nabla \cdot \mu_m (\nabla \mathbf{v} + \nabla \mathbf{v}^T) + \rho_m \mathbf{g} + \sigma \kappa \mathbf{n} \delta_S \quad \nabla \cdot \mathbf{v} = 0$$

- **Obvious inconsistency:**  
some quantities are averaged  $f, \rho_m, \mu_m$   
and some are not  $\mathbf{v}, p$

## (Problem and motivation)

- Both phases are assumed to move with same velocity, namely the center-of-mass velocity  
⇒ comput. grid must resolve boundary layer at interface
- Which are the additional terms if the boundary layer is not fully resolved, i.e. the phase velocities differ?
- **Goal:** Derive consistent volume-averaged equations for volume-of-fluid computations  
(VA-VOF eqs. for phases of constant density )

## Procedure to derive VA-VOF equations

1. local mass & momentum cons.-eqs. for phase  $k=1,2$  valid in  $\Omega_k(t)$
2. multiply eqs. by respective phase indicator function  $X_k(\mathbf{x},t)$
3. apply to each term VA operator  $\overline{\psi}_k^k \equiv \frac{1}{V_k} \int_V \psi_k(\mathbf{x} + \boldsymbol{\eta}, t) X_k(\mathbf{x} + \boldsymbol{\eta}, t) d\mathbf{x}$
4. apply Gauss & Leibnitz rules to obtain eqs. of two-fluid model:

$$\frac{\partial \alpha_k \rho_k}{\partial t} + \nabla \cdot \alpha_k \rho_k \overline{\mathbf{v}}_k^k = 0 \quad k = 1,2$$

$$\frac{\partial \alpha_k \rho_k \overline{\mathbf{v}}_k^k}{\partial t} + \nabla \cdot \alpha_k \rho_k \left( \overline{\mathbf{v}}_k^k \overline{\mathbf{v}}_k^k + \overline{\mathbf{v}'_k \mathbf{v}'_k}^k \right) = -\nabla \alpha_k \overline{p}_k^k + \alpha_k \rho_k \mathbf{g} + \nabla \cdot \alpha_k \overline{\boldsymbol{\tau}}_k^k + \mathbf{M}_k$$

## (Procedure to derive VA-VOF equations)

5. for simplified representation of VA-VOF eqs. we introduce

- liquid volumetric fraction  $f \equiv \alpha_1$

- mixture density  $\rho_m \equiv \sum_{k=1}^2 \alpha_k \rho_k = f\rho_1 + (1-f)\rho_2$

- center-of-mass velocity  $\mathbf{v}_m \equiv \frac{1}{\rho_m} \sum_{k=1}^2 \alpha_k \rho_k \overline{\mathbf{v}}_k = \frac{f\rho_1 \overline{\mathbf{v}}_1 + (1-f)\rho_2 \overline{\mathbf{v}}_2}{\rho_m}$

- relative velocity  $\mathbf{v}_r = \mathbf{v}_{21} \equiv \overline{\mathbf{v}}_2 - \overline{\mathbf{v}}_1$

$$\Rightarrow \overline{\mathbf{v}}_1 = \mathbf{v}_m - \frac{\alpha_2 \rho_2}{\rho_m} \mathbf{v}_r, \quad \overline{\mathbf{v}}_2 = \mathbf{v}_m + \frac{\alpha_1 \rho_1}{\rho_m} \mathbf{v}_r$$

## VA-VOF mass conservation equations

- two-fluid model mass conservation equations for  $k=1,2$ :

$$\frac{\partial f \rho_1}{\partial t} + \nabla \cdot f \rho_1 \mathbf{v}_m = \nabla \cdot f(1-f) \frac{\rho_1 \rho_2}{\rho_m} \mathbf{v}_r \quad (1)$$

$$\frac{\partial (1-f) \rho_2}{\partial t} + \nabla \cdot (1-f) \rho_2 \mathbf{v}_m = -\nabla \cdot f(1-f) \frac{\rho_1 \rho_2}{\rho_m} \mathbf{v}_r \quad (2)$$

- rearrange to yield:

$$\frac{\partial f}{\partial t} + \nabla \cdot f \mathbf{v}_m = \nabla \cdot f(1-f) \frac{\rho_2}{\rho_m} \mathbf{v}_r \quad (1)/\rho_1$$

$$\nabla \cdot \mathbf{v}_m = -\nabla \cdot f(1-f) \frac{\rho_1 - \rho_2}{\rho_m} \mathbf{v}_r \quad (1)/\rho_1 + (2)/\rho_2$$

## VA single-field momentum equation

- sum up the momentum eqs. of the two-fluid model for phase 1 and 2

$$\frac{\partial}{\partial t} \sum_{k=1}^2 \alpha_k \rho_k \overline{\mathbf{v}_k} + \nabla \cdot \sum_{k=1}^2 \alpha_k \rho_k \overline{\mathbf{v}_k \mathbf{v}_k} =$$
$$-\nabla \sum_{k=1}^2 \alpha_k \overline{p_k} + \sum_{k=1}^2 \alpha_k \rho_k \mathbf{g} + \nabla \cdot \sum_{k=1}^2 \alpha_k (\overline{\boldsymbol{\tau}_k} + \boldsymbol{\tau}_{\text{sgs}}^k) + \mathbf{M}_1 + \mathbf{M}_2$$

where  $\underline{\boldsymbol{\tau}}_{\text{sgs}} = -\sum_{k=1}^2 \alpha_k \rho_k \overline{\mathbf{v}'_k \mathbf{v}'_k}$

## (VA single-field momentum equation)

- sum of momentum transfer terms (jump condition)

$$\mathbf{M}_1 + \mathbf{M}_2 = \frac{\sigma}{V} \int_{S_i(\mathbf{x},t)} \kappa \mathbf{n} dS \cong \sigma \bar{\kappa}^V \bar{\mathbf{n}}^V \frac{1}{V} \int_{S_i(\mathbf{x},t)} dS = \sigma \bar{\kappa}^V \bar{\mathbf{n}}^V a_{int}$$

- sum of convective terms

$$\sum_{k=1}^2 \alpha_k \rho_k \overline{\mathbf{v}_k \mathbf{v}_k} = \rho_m \mathbf{v}_m \mathbf{v}_m + \underline{\mathbf{D}}_{int},$$

interfacial area concentration

where  $\underline{\mathbf{D}}_{int} \equiv \alpha_1 \alpha_2 \frac{\rho_1 \rho_2}{\rho_m} \mathbf{v}_r \mathbf{v}_r =$  "momentum drift flux term"



## (VA single-field momentum equation)

- sum of diffusive terms

$$\sum_{k=1}^2 \alpha_k \overline{\underline{\boldsymbol{\tau}}_k} = \sum_{k=1}^2 \alpha_k \mu_k \left( \nabla \overline{\mathbf{v}_k} + \nabla \overline{\mathbf{v}_k}^T \right) \equiv \underline{\boldsymbol{\tau}}_m + \underline{\boldsymbol{\tau}}_{\text{int}}$$

where  $\underline{\boldsymbol{\tau}}_m \equiv \mu_m (\nabla \mathbf{v}_m + \nabla \mathbf{v}_m^T)$ ,  $\mu_m \equiv f\mu_1 + (1-f)\mu_2$

$$\Rightarrow \underline{\boldsymbol{\tau}}_{\text{int}} = \alpha_2 \mu_2 \left( \nabla \frac{\alpha_1 \rho_1}{\rho_m} \mathbf{v}_r + \nabla \frac{\alpha_1 \rho_1}{\rho_m} \mathbf{v}_r^T \right) - \alpha_1 \mu_1 \left( \nabla \frac{\alpha_2 \rho_2}{\rho_m} \mathbf{v}_r + \nabla \frac{\alpha_2 \rho_2}{\rho_m} \mathbf{v}_r^T \right)$$

$\underline{\boldsymbol{\tau}}_{\text{int}}$  = "interfacial friction term"

## Set of VA-VOF equations

$$\frac{\partial f}{\partial t} + \nabla \cdot f \mathbf{v}_m = \nabla \cdot f(1-f) \frac{\rho_2}{\rho_m} \mathbf{v}_r$$

$$\nabla \cdot \mathbf{v}_m = -\nabla \cdot f(1-f) \frac{\rho_1 - \rho_2}{\rho_m} \mathbf{v}_r$$

$$\begin{aligned} \frac{\partial \rho_m \mathbf{v}_m}{\partial t} + \nabla \cdot (\rho_m \mathbf{v}_m \mathbf{v}_m + \underline{\mathbf{D}}_{\text{int}}) = & -\nabla \cdot \sum_{k=1}^2 \alpha_k \overline{p}_k + \rho_m \mathbf{g} + \sigma \overline{K}^V \overline{\mathbf{n}}^V a_{\text{int}} \\ & + \nabla \cdot \mu_m (\nabla \mathbf{v}_m + \nabla \mathbf{v}_m^T) + \nabla \cdot (\underline{\boldsymbol{\tau}}_{\text{int}} + \underline{\boldsymbol{\tau}}_{\text{sgs}}) \end{aligned}$$

## (Set of VA-VOF equations)

$$\frac{\partial f}{\partial t} + \nabla \cdot f \mathbf{v}_m = \nabla \cdot \cancel{f(1-f)} \frac{\rho_2}{\rho_m} \mathbf{v}_r = \mathbf{0}$$

limit  $V \rightarrow 0 \Rightarrow \mathbf{v}_r = \mathbf{0}$   
 $\Rightarrow$  local VOF equations

$$\nabla \cdot \mathbf{v}_m = -\nabla \cdot \cancel{f(1-f)} \frac{\rho_1 - \rho_2}{\rho_m} \mathbf{v}_r = \mathbf{0}$$

$$\begin{aligned} \frac{\partial \rho_m \mathbf{v}_m}{\partial t} + \nabla \cdot (\rho_m \mathbf{v}_m \mathbf{v}_m + \cancel{\mathbf{D}_{int}}) &= -\nabla \cdot \sum_{k=1}^2 \alpha_k \overline{p}_k + \rho_m \mathbf{g} + \sigma \overline{\kappa}^V \overline{\mathbf{n}}^V a_{int} \\ &+ \nabla \cdot \mu_m (\nabla \mathbf{v}_m + \nabla \mathbf{v}_m^T) + \nabla \cdot (\cancel{\boldsymbol{\tau}_{int}} + \boldsymbol{\tau}_{sgs}) \end{aligned}$$

## Closure of VA-VOF equations

- 3 equations:  $f$ -equation = liquid mass conservation eq.  
divergence condition = gas mass conserv. eq.  
mixture momentum equation
- 5 unknowns  $f, \mathbf{v}_m, \mathbf{v}_r, p_1, p_2$
- 1<sup>st</sup> assumption: both phases share same pressure field
- 2<sup>nd</sup> closure relation: constitutive equation for  $\mathbf{v}_r$ 
  - trivial closure: “homogenous model”  $\mathbf{v}_r = \mathbf{0}$

## Relevance of modelling $v_r$

- estimate for thickness of boundary layer

$$\delta \approx \frac{d_b/2}{\sqrt{\text{Re}_b}}, \quad \text{where} \quad \text{Re}_b = \frac{d_b U_\infty}{\nu_l}$$

- resolve boundary layer by 3 mesh cells

$$\Rightarrow \frac{d_b}{\Delta x} = 6\sqrt{\text{Re}_b}$$

$\text{Re}_b$	10	100	1,000
$d_b / \Delta x$	19	60	189

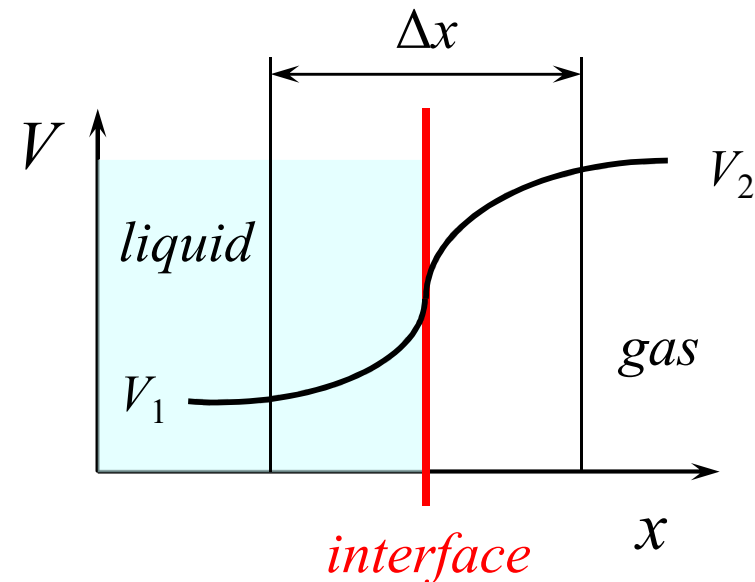
- current 3D VOF computations:  $\text{Re}_b \leq O(100)$ ,  $d_b/\Delta x \leq 50$

## Order of magnitude estimation

- momentum drift flux term\*

$$\frac{D_{int}}{\rho_m V_m V_m} = q(f, \rho_1, \rho_2, V_1, V_2)$$

$$\frac{f}{1-f} = \frac{\rho_1}{\rho_2} \Rightarrow q = q_{max} \cong \frac{1}{2} \frac{V_2 - V_1}{V_1 + V_2}$$



- interfacial friction term

$$\rho_1 = \rho_2, \mu_1 = \mu_2 \Rightarrow \frac{\tau_{int}}{\tau_m} \propto \frac{(V_2 - V_1) \frac{\partial f}{\partial x}}{\frac{\partial V_m}{\partial x}} \cong \frac{(V_2 - V_1) \frac{1-0}{\Delta x}}{\frac{V_2 - V_1}{\Delta x}} = 1$$

\* S. Mitran, FZKA Report 6357 (2000)

## A first model proposal for $v_r$

- **Local Uniform Relative Velocity model (LURV)**
  - algebraic model
  - conservative discretization of  $\underline{D}_{\text{int}}$  on staggered grid
- results for static 2D interface
  - “homogeneous model” in  $f$ -eq. and divergence equation
  - LURV model in momentum equation
  - *unphysical* spurious currents may be slightly reduced or increased
  - ⇒ LURV model will be tested for *physical* flow fields

## Conclusions and outlook

- Derivation of volume-averaged VOF equations
  - phase velocities may differ from center-of-mass velocity  
⇒ additional terms in interface cells as compared to local VOF eqs.
  - closure assumption for relative phase velocity  $\mathbf{v}_r$  is required  
⇒ suitable models for  $\mathbf{v}_r$  must be developed and tested
  - perspective of “low bubble resolution” VOF simulations for high  $Re_b$
  - VA momentum eq. applies also to level set and front-tracking method
- Two-phase “large-eddy simulation” with interface tracking
  - SGS model for unresolved velocity fluctuations in single phase cells
  - model for  $\mathbf{v}_r$  to account for unresolved boundary layer