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Volume-of-fluid method based numerical simulations of gas-liquid flow in confined geometries

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Merits of DNS of gas-liquid flow

- Allows to get deeper insight into flow mechanisms and thus fosters physical understanding

 Here: <u>bubble train flow in square mini-channel</u>
- Provides a complete database of the 3D velocity and pressure field and phase distribution with high spatial and temporal resolution
 - Here: <u>analysis of liquid phase turbulence kinetic</u> <u>energy equation for bubble swarm flow</u>

In-house code TURBIT-VOF

- Volume-of fluid method for interface tracking
 - Interface is locally approximated by plane (PLIC method)
- Governing equations for two incompressible fluids
 - Single field momentum equation with surface tension term
 - Zero divergence condition for center-of-mass velocity
 - Advection equation for liquid volumetric fraction f
- Solution strategy
 - Projection method resulting in pressure Poisson equation
 - Explicit third order Runge-Kutta time integration scheme
- Discretization in space
 - Finite volume formulation for regular staggered grid
 - Second order central difference approximations

Multi-phase micro process engineering

- Miniaturized devices offer certain advantages
 - High interfacial area per unit volume
 - \Rightarrow Enhanced heat and mass transfer
 - Defined interface geometry
 - \Rightarrow Numbering up instead of scaling up
 - Example: Micro bubble column*



- Flow pattern in single channel: bubble train flow
 - Bubbles of identical shape move with same velocity
 - Flow is fully characterized by a single "unit cell" of length L_{uc}



Numerical set up

- Square vertical channel with cross section 2 mm \times 2 mm
- Consideration of <u>one</u> flow unit cell only
- Account for influence of trailing/leading unit cells by <u>periodic boundary conditions</u> in axial direction
- Length of flow unit cell, L_{uc} , is input parameter
- Flow is driven in vertical direction (y) by specified axial pressure gradient and buoyancy

Physical parameters

• Fluid properties Factor 10 higher than ρ and μ of air

$ ho_{l}$	$ ho_{g}$	μ_l	μ_g	σ
957 kg/m ³	11.7 kg/m ³	0.048 Pa s	1.84×10 ⁻⁴ Pa s	0.022 N/m

• Initial bubble shapes (void fraction $\varepsilon = 33\%$)



Simulations are started from gas and liquid at rest

Computational parameters

Case	$L_{ m uc}$ / W	Domain	Grid	Time steps
A1	1	1 × 1 × 1	$48 \times 48 \times 48$	24,000
A2	1	1 × 1 × 1	$64 \times 64 \times 64$	60,000
В	1.25	1 × 1.25 × 1	$48 \times 60 \times 48$	24,000
С	1.5	1 × 1.5 × 1	$48\times72\times48$	26,000
D	1.75	1 × 1.75 × 1	$48 \times 84 \times 48$	26,000
Е	2	$1 \times 2 \times 1$	$48 \times 96 \times 48$	28,000

Results on both grids show only slight differences

Time history of mean velocities



Steady state values of bubble velocity U_B and mean liquid velocity U_l increase with increasing length of the flow unit cell

Bubble shape and velocity field

Velocity field in vertical mid-plane

Right half: frame of reference moving with bubble Left half: fixed frame of reference







Comparison with experiment

Non-dimensional bubble diameter			r Relativ	Relative velocity Non-dimensional U _B		
Case	L _{uc} / W	Ca _B	D_B/W	$(U_B - J_{total})/U_B$	U_{B}/J_{total}	
Α	1	0.204	0.81	1.80	0.445	
В	1.25	0.207	0.84	1.75	0.430	
С	1.5	0.215	0.85	1.75	0.430	
D	1.75	0.238	0.85	1.78	0.438	
Е	2	0.253	0.85	1.8	0.445	
Experimental data [*] correlated in terms of capillary number $Ca_B \equiv \mu_U U_B / \sigma$						
0.2 – 0.25		0.82 – 0.86	5 1.68 – 1.84	0.435 – 0.475		
			\checkmark	\checkmark	\checkmark	

D_B measured along channel diagonal^{*}



Bubble diameter in simulations



 D_B/W decreases with increase of Ca_B <u>only</u> if the bubble length L_B is larger than about 1.2 the channel width (this is the case in the experiments by Thulasidas et al.)

D_B measured along channel diagonal^{*}



Influence of capillary number top row: Ca = 0.205, bottom row: Ca = 0.043



Turbulence modeling for bubbly flow

- No generally accepted model available in literature
- Analytical turbulence kinetic energy eq. for liquid phase*:

$$\frac{\partial}{\partial t}(\alpha_{L}k_{L}) + \nabla \cdot \left(\alpha_{L}k_{L}\overline{\mathbf{u}_{L}}\right) = \frac{1}{Re_{ref}} \nabla \cdot \left(\alpha_{L}\overline{\mathbb{T}_{L}^{'} \cdot \mathbf{u}_{L}^{'}}\right) - \nabla \cdot \left[\alpha_{L}\left(\overline{p_{L}^{'}\mathbf{u}_{L}^{'}} + \frac{1}{2}\overline{(\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'})\mathbf{u}_{L}^{'}}\right)\right]$$

$$\frac{\partial}{\partial t}(\alpha_{L}k_{L}) + \nabla \cdot \left[\alpha_{L}k_{L}\overline{\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'}}\right] + \frac{\partial}{\partial t}\left[\alpha_{L}\overline{\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'}\right] + \frac{\partial}{\partial t}\left[\alpha_{L}\overline{\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'} - \alpha_{L}\overline{\mathbf{u}_{L}^{'}}\right] + \frac{\partial}{\partial t}\left[\alpha_{L}\overline{\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'}}\right] + \frac{\partial}{\partial t}\left[\alpha_{L}\overline{\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'}}\right] + \frac{\partial}{\partial t}\left[\alpha_{L}\overline{\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'} - \alpha_{L}\overline{\mathbf{u}_{L}^{'} \cdot \mathbf{u}_{L}^{'}}\right] + \frac{\partial}{\partial t}\left[\alpha_{L}\overline{\mathbf{u}_{L}^{'} - \alpha_{L}\overline{\mathbf{u}_{L}^{'} -$$

Simulation of bubble swarm flow

X₁ ↑

- Simulation mimics section of a flat bubble column
 - periodic b.c. in vertical and span-wise directions
 - rigid lateral walls
- Domain: $1 \times 1 \times 1$, Grid: $64 \times 64 \times 64$
- Eight bubbles with $d_{\rm B}/W = 0.25$ ($\varepsilon = 6.5\%$)
- Phase density ratio: 0.5
- Phase viscosity ratio: 1
- Bubble Eötvös number: 3.065
- Morton number: $3.06 \cdot 10^{-6}$



Flow visualization





view from top

 X_2

Wall-normal profiles of mean quantities







Budget of k_L-equation



Models for interfacial term

Reference	Work of drag force, $W_{\rm D}^{*}$	Other contributions, $W_{\rm ND}^{*}$
Kataoka & Serizawa (1997) Model 1, KS	$0.075 f_{\rm w} \left[\frac{3}{4} \alpha_{\rm G} \frac{C_{\rm D}}{d_{\rm B}^*} U_{\rm T}^{*3} \right]$	$-\alpha_{\rm G}\frac{k_{\rm L}^{*3/2}}{d_{\rm B}^*}$
Hill et al. (1995) Model 2, HWGI	$\frac{3}{4} \frac{\alpha_{\rm G} C_{\rm D}}{d_{\rm B}^*} \left \overline{\mathbf{u}_{\rm R}^*} \right \left\{ \frac{\mu_{\rm L}^* \overline{\mathbf{u}_{\rm R}^*} \cdot \nabla^* \alpha_{\rm G}}{0.3 \rho_{\rm L}^* \alpha_{\rm L} \alpha_{\rm G}} + 2k_{\rm L}^* (C_{\rm t} - 1) \right\}$	None
Lahey & Drew (2000) Model 3, LD	$\frac{1}{4}\alpha_{\rm L}\left(1+C_{\rm D}^{4/3}\right)\alpha_{\rm G}\frac{\left \overline{\mathbf{u}_{\rm R}^*}\right ^3}{d_{\rm B}^*}$	None
Morel (1997) Model 4, M	$\frac{3}{4} \alpha_{\rm G} \frac{C_{\rm D}}{d_{\rm B}^*} \left \overline{\mathbf{u}_{\rm R}^*} \right ^3$	$\frac{1+2\alpha_{\rm G}}{2\alpha_{\rm L}}\alpha_{\rm G}\left\{\frac{{\rm D}_{\rm G}\overline{\mathbf{u}_{\rm G}^*}}{{\rm D}t^*}-\frac{{\rm D}_{\rm L}\overline{\mathbf{u}_{\rm L}^*}}{{\rm D}t^*}\right\}\cdot\overline{\mathbf{u}_{\rm R}^*}$
Pfleger & Becker (2001) Model 5, PB	$1.44\alpha_{\rm L} \left[\frac{3}{4} \alpha_{\rm G} \frac{C_{\rm D}}{d_{\rm B}^*} \left \overline{\mathbf{u}_{\rm R}^*} \right ^3 \right]$	None

Performance of models for interfacial term



Conclusions

- Numerical simulation of bubble train flow
 - Square vertical mini-channel (2 mm × 2 mm)
 - Good agreement with experimental data from literature
 - Investigation on influence of length of flow unit cell
 - Strong influence of capillary number
- Liquid turbulence kinetic energy equation for bubbly flow
 - production mainly by interfacial terms
 - no local equilibrium between interfacial production and dissipation
 \Rightarrow significant redistribution by diffusion
 - interfacial terms can well be modeled by work of drag forces