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KERNFORSCHUNGSZENTRUM KARLSRUHE

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Externer Bericht

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Institut für Experimentelle Kernphysik

FORTTRAN-Program for the computation of the
Surface Impedance of Superconductors

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KARLSRUHE

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Surface Impedance of Superconductors

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Externer Bericht

Abstract

The FORTRAN-program presented in this report computes the surface impedance Z of superconductors in the Meissner state for frequencies lower than the energy gap. The formulae used are exact within the frame work of the BCS-theory for weak-coupling superconductors. Strong coupling effects can be introduced by using measured Δ/kT_c values. In the program the temperature dependence of the gap parameter $\Delta(T)$ is approximated by $(\Delta(T)/\Delta(0))^2 = \cos^2 \frac{\pi}{2} (T/T_c)^2$.

I. SUMMARY

The program (see tables) computes the surface impedance Z for the actual parameters of measurements:

[FO] : frequency; [TE(..)] : temperature,

the symbols in the brackets [...] are used in the FORTRAN-program (see tables). The number [IT] of temperature values is limited to ≤ 20 by the program.

The material parameters for a superconductor in the weak coupling limit are

[TC] : critical temperature;
[FAK] : observed $\Delta(0)/kT_c$ value;
[DLON] : London penetration depth at zero temperature and infinite mean free path λ :

$$\delta_L = \sqrt{m/(\mu_0 e^2 N)}$$

with m the mass, e the charge and N the density of (transport) electrons;

[XKOH] : dimension of a cooper pair (coherence length)
at $T = 0$ and $l = \infty$.

$$\xi_F = \hbar v_F / 2\Delta(0) \quad (2)$$

with the Fermi velocity v_F ;

[FREI] : mean free path of the electrons in the normal
state.

In the program the unit for temperature is degree Kelvin
and the unit for length Angstrom.

The accuracy of computation can be increased by lowering
the factor [REST]. The lower limit of this factor in our
program (see tables) is given by about $5 \cdot 10^{-7}$ because
the number of abscissas [IQ] is limited by 10^3 . For
 $REST \sim 10^{-4}$ the absolut accuracy of Z will be about 1 %,
the relative accuracy (for normal materials like lead
or niobium) will be about 1 %. For a further increase of
accuracy one should use adapted abscissas.

With the described input data the program computes the
surface impedance $Z = R + iX$ for specular "S" and diffuse
"D" reflection of the charge carriers at the surface. In
the output the surface resistance R [RS] or [RD] is prin-
ted in Ohm, whereas the surface reactance $X = \omega \mu_0 \lambda$ is
given through the penetration depth $\lambda = \int_0^{\infty} \text{Re}\{H(z,t)/H(0,t)\} dz$
[XS] or [XD].

The program (see tables) can be used also for a computation of the penetration for $\omega = 0$ (or $\hbar\omega \ll \Delta$) if one omits the subroutine NONAN in the program-deck.

II. INTRODUCTION

In plane geometry the surface impedance $Z = R + iX$ is given by the penetration of the magnetic field $H(z,t)$ into the wall:

$$Z = i\omega\mu_0 \int_0^{\infty} dz H(z,t)/H(0,t) = i\omega\mu_0 A(0,t)/B(0,t) \quad (3)$$

with z the distance from the surface and A the vector potential: $\vec{H} = \nabla \times \vec{A}$. The fields $H(z,t)$ or $A(z,t)$ are the solutions of the following equations and boundary conditions (1,3-5).

Maxwell's equation ($\partial^2/\partial t^2$ can be neglected for frequencies lower than 10^{13} Hz):

$$-\partial^2 A(z,t)/\partial z^2 = \mu_0 j(z,t);$$

material equation:

$$j(z,t) = - \int_{-\infty}^0 dt' \int dz' Q(z-z', t+t') A(z',t');$$

boundary conditions for the fields:

$$\partial A(z,t)/\partial z \Big|_{z=0} = B(0,t), \quad A(\infty,t) = 0;$$

boundary conditions for the charge carriers at the surface:

specular or diffuse reflection.

According to ³⁾ with the Fourier transformed integral kernel $Q(k,\omega)$ the surface impedance can then be written:

specular reflection:
$$Z_S = i\mu_0\omega \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{dk}{k^2 + \mu_0 Q(k, \omega)}$$
 (4)

diffuse reflection:
$$Z_D = i\mu_0\omega \frac{\pi}{\int_0^{\infty} dk \ln\{1 + \mu_0 Q(k, \omega)/k^2\}}$$

For Q evaluation of (4) the k-integration is transformed in a q-integration by: $q = k \epsilon_F \Delta(O)/\Delta(T)$ [DX(...)], i.e. the momenta k will be measured in units of a coherence length. With the abbreviations

$$\delta_L(T, \ell, \omega) = 1/\sqrt{\mu_0 |Q(O, \omega)|} \quad [\text{ENP}]$$

$$\kappa = \delta_L(T, \ell, \omega) / \{\epsilon_F \Delta(O)/\Delta(T)\} \quad [\text{AK}]$$

(4) can be written:

$$Z_S = i\mu_0\omega \delta_L(T, \ell, \omega) \frac{2\kappa}{\pi} \int_0^{\infty} \frac{dq}{\kappa^2 q^2 + Q(q, \omega) / |Q(O, \omega)|}$$
 (5)

$$Z_D = i\mu_0\omega \frac{\pi \delta_L(T, \ell, \omega)}{\int_0^{\infty} dq \ln(1 + Q(q, \omega) / (|Q(O, \omega)| \kappa^2 q^2))}$$

Before we can start with the integration dq in subroutine HAUPT and compute the integrand in subroutine INTQ (...) the integralkernel Q(k, ω) should be evaluated. According to ¹⁾ we split Q in an in ω analytical term Q_A and a non-analytical term Q_D:

$$\mu_0 \delta_L^2 Q(k, \omega) = \bar{Q}_A(k, \omega) + \bar{Q}_D(k, \omega) \quad (6)$$

The analytical term Q_A is defined by:

$$\bar{Q}_A(k, \omega) = a \sum_{\substack{m=(2n+1)a \\ n \geq 1}} \left\{ \frac{1 - m(m+i\tilde{\omega}) + \sqrt{m^2+1} \sqrt{(m+i\tilde{\omega})^2+1}}{\sqrt{m^2+1} \sqrt{(m+i\tilde{\omega})^2+1}} \right. \\ \left. 4 J \left[\frac{\sqrt{m^2+1} + \sqrt{(m+i\tilde{\omega})^2+1}}{2} + \frac{\hbar v_F}{2\Delta(T)\ell} \right] + c.c. \right\} \quad (7)$$

($a = \pi kT/\Delta(T)$; $\tilde{\omega} = \hbar\omega/\Delta(T)$: [0]). \bar{Q}_A is computed in subroutine ISUM, \bar{Q}_A is purely real and therefore gives for $\tilde{\omega} \ll 1$ the penetration depth, i.e. \bar{Q}_A describes the current of "Cooper" pairs. \bar{Q}_D which disappears for $\tilde{\omega} \rightarrow 0$ and $T \rightarrow 0$ consists of a real term \bar{Q}_P and a term \bar{Q}_{SCH} , describing the rf absorption:

$$\bar{Q}_D(k, \omega) = \bar{Q}_P(k, \omega) + \bar{Q}_{SCH}(k, \omega) \quad (8)$$

$$\bar{Q}_P(k, \omega) = \int_{1-\tilde{\omega}}^1 d\bar{\omega} \left\{ \text{th} \frac{\Delta(T)}{2k_B T} \bar{\omega} - \text{th} \frac{\Delta(T)}{2k_B T} (\bar{\omega} + \tilde{\omega}) \right\} \cdot \\ \left\{ \frac{1 + \bar{\omega}(\bar{\omega} - \tilde{\omega}) + i\sqrt{-(\bar{\omega} - \tilde{\omega})^2 + 1} \sqrt{\bar{\omega}^2 - 1}}{\sqrt{\bar{\omega}^2 - 1} \sqrt{-(\bar{\omega} - \tilde{\omega})^2 + 1}} J \left[\frac{\sqrt{-(\bar{\omega} - \tilde{\omega})^2 + 1} + i\sqrt{\bar{\omega}^2 - 1}}{2} + \frac{\hbar v_F}{2\Delta(T)\ell} \right] + c.c. \right\} \quad (9)$$

$$\bar{Q}_{SCH}(k, \omega) = 21 \int_1^{\infty} d\bar{\omega} \left\{ \text{th} \frac{\Delta(T)}{2k_B T} \bar{\omega} - \text{th} \frac{\Delta(T)}{2k_B T} (\bar{\omega} + \tilde{\omega}) \right\} \cdot$$

$$\cdot \left\{ \frac{1 + \bar{\omega}(\bar{\omega} + \tilde{\omega}) + \sqrt{\bar{\omega}^2 - 1} \sqrt{(\bar{\omega} + \tilde{\omega})^2 - 1}}{\sqrt{\bar{\omega}^2 - 1} \sqrt{(\bar{\omega} + \tilde{\omega})^2 - 1}} J \left[-1 \frac{\sqrt{(\bar{\omega} + \tilde{\omega})^2 - 1} + \sqrt{\bar{\omega}^2 - 1}}{2} + \frac{\hbar v_F}{2\Delta(T)\ell} \right] \right. \quad (10)$$

$$\left. \frac{1 + \bar{\omega}(\bar{\omega} + \tilde{\omega}) - \sqrt{\bar{\omega}^2 - 1} \sqrt{(\bar{\omega} + \tilde{\omega})^2 - 1}}{\sqrt{\bar{\omega}^2 - 1} \sqrt{(\bar{\omega} + \tilde{\omega})^2 - 1}} \text{Re} J \left[-1 \frac{\sqrt{(\bar{\omega} + \tilde{\omega})^2 - 1} + \sqrt{\bar{\omega}^2 - 1}}{2} + \frac{\hbar v_F}{2\Delta(T)\ell} \right] \right\}$$

\bar{Q}_P and \bar{Q}_{SCH} will be computed in subroutine NONAN.

The integrals (7), (9) and (10) depend on the function J which results from an angle integration:

$$J(u) = 6 \frac{1}{\hbar v_F k} \left\{ - \frac{u 2\Delta(T)}{\hbar v_F k} + \left[\frac{u 2\Delta(T)}{\hbar v_F k} + 1 \right] \text{arctg} \frac{\hbar v_F k}{u 2\Delta(T)} \right\} \quad (11)$$

J(u) will be evaluated in subroutine WINK (...).

In the following sections we will discuss in more detail the various summations or integrations in the subroutines.

III. MOMENTUM INTEGRATION

The q-integration is done by Simpson's rule (formula 25.4.6 in ⁶). The q-range $\{0, \infty\}$ is subdivided into the intervals $[I] = 1, \dots, 5$ ($[G(1)]$, $[G(2)]$..) and the 6th inter-

val $\{ [G(6)] / [G(1)] \}$, in which q is substituted by $\frac{1}{q}$. The abscissas $[DX ()]$ within these intervals are evaluated in the main program, their number is given by $[ID(I)]$ (with $[ID(I)] = u$ and $[DELT] = h$ in 25.4.6 in ⁶⁾). The subdivision into 6 intervals before using Simpson's rule has been introduced to improve the accuracy by giving certain q -regions larger weight. For materials with, e.g. $\delta_L / \xi_F \ll 1$ small momenta $k \approx (\delta_L \xi_F)^{-1/2}$ are most important. After evaluation of Q in subroutines NONAN and ISUM the integrand is calculated in subroutine INTQ (,) and summed in subroutine HAUPT. The result $- Z / (\omega \mu_0 \delta_L(T, \ell, \omega))$ is transferred to the main program, where after some multiplications the surface resistance in Ohm and the penetration depth $\lambda(T, \ell, \omega)$ in Angstrom are printed out.

IV. EVALUATION OF $Q(k, \omega)$

We start with the smallest fraction of Q , namely the \bar{Q}_D in subroutine NONAN to reduce rounding errors, The larger \bar{Q}_A is computed afterwards in subroutine ISUM. A numerical computation of \bar{Q}_D with Simpson's rule is not very precise because of the square root singularity of the integrand. More adequate for numerical integration of square root singularities is the integration by Chebyshev Polynomials of the First Kind, i.e. the use of formula 25.4.39 of ⁶⁾:

$$\int_a^b \frac{f(y)dy}{\sqrt{(y-a)(b-y)}} = \sum_{i=1}^n \omega_i f(y) + R_n \quad (12)$$

$$y_i = \frac{b+a}{2} + \frac{b-a}{2} x_i, \quad x_i = \cos \frac{(2i-1)\pi}{2n}, \quad \omega_i = \frac{\pi}{n}$$

The evaluation of the residue R_n for (9) and (10) is too difficult and it is omitted, therefore. A sufficient n is chosen according to the requirement that there must be enough abscissas (> 5) in the region where $f(y)$ has a sharp maximum.

IV.1 $\bar{Q}_{SCH}(k, \omega)$

This term describes the rf absorption, i.e. \bar{Q}_{SCH} includes the imaginary part of Q . For computation integrand of (10) is transformed by substituting $u = 1/\bar{\omega}$; i.e. the interval (a, b) in (12) is: $(0, 1)$. For the integral the maximum of $f(y) = 1/\sqrt{1+u\bar{\omega}-u}$ becomes more important for lower $\bar{\omega}$, i.e. the number of abscissas [IS] will increase with $1/\sqrt{\bar{\omega}}$.

IV.2 $\bar{Q}_P(k, \omega)$

For $\bar{\omega} < 1$ $\bar{Q}_P(k, \omega)$ is a slowly varying function with $\bar{\omega}$ the number of abscissas [IP] will therefore be small. Moreover there is no need of a more exact computation of \bar{Q}_P for $\bar{\omega} \ll 1$, because \bar{Q}_P is real and small compared to \bar{Q}_A .

IV.3 $\bar{Q}_A(k, \omega)$

\bar{Q}_A (7) is computed in subroutine ISUM. For $\bar{\omega} = 0$ and $T = 0$ ($\bar{\omega} \leq 2$) Q_A is the only non-vanishing part. For the numerical computation the temperature sum in ISUM yields easily a good accuracy. The number of terms in the sum [MS] must be increased for $T \rightarrow 0$ because the terms are $\propto 1/(T/T_c)^2$. To reduce rounding errors in computing the sum ISUM starts with the smallest terms.

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ment Printing Office, Washington, C.D., 1964)

FORTRAN - PROGRAM

```
C SURFACE IMPEDANCE OF SUPERCONDUCTORS FOR FREQUENCIES SMALLER THAN
C THE ENERGY GAP.
  DIMENSION G(6),ID(6),TE(20)
  REAL*8 DELT,DET,DQ(6),DA,DE,DZ,DX(1000),PI      ,DQ1 ,DS(1000)
  COMPLEX*8 CHS,CHD,CI
  CCOMPLEX*16 CDS(1000),CT(1000),CS,CD
  COMMON/HL1/ID,IQ,MS,IS,IP
  COMMON/HL2/AK,A2,FL,O,GP
  COMMON/HL3/PI,DQ,DX,DQ1,DS
  COMMON/HL4/CHS,CHD,CI
  COMMON/HL5/CDS,CT,CS,CD

C IT=NUMBER OF TEMPERATUREVALUES.
  READ (5,1) IT
C FO=FREQUENCY (HERTZ),TE(I)=TEMPERATURES (KELVIN).
  READ (5,2) FO,(TE(I),I=1,IT)
C MATERIALPARAMETERS(UNITS:DEGREE KELVIN,ANGSTROM)
  READ (5,2) TC,FAK,DLON,XKOH,FREI
C REST GIVES THE ACCURACY OF THE COMPUTATION.
  READ (5,2) REST
  1 FORMAT(I5)
  2 FGMAT(8E10.4)
  WRITE(6,2) FO,(TE(I),I=1,IT),REST
C THE MOMENTUMINTEGRAL IS SPLITTED IN 6 PARTS OVER THE INTERVALLS (
C (I),G(I+1)).
  G(1)=1.E-6
  G(2)=1.E-2
  G(3)=0.2
  G(4)=2.
  G(5)=5.
  G(6)=50.
C THESE 6 INTEGRALS ARE EVALUATED BY SIMPSON'S RULE. 1/ID(I) GIVES
C THE SPACING OF THE ABCISSAS.
  ID(1)=5
  ID(2)=7
  ID(3)=15
  ID(4)=20
  ID(5)=20
  ID(6)=9
  IS=(6.E-4/REST)**0.25+1.
  DC 38 N=1,6
  38 ID(N)=ID(N)*IS
  PI=3.1415926535
  PH =0.479*10.**(-10)
  CI=(0.,1.)
  FLO=XKOH/FREI
  AL=DLON/XKOH
  DG=PH*FO/TC/FAK
  DX(1)=0.
  DX(2)=G(1)
  IQ=2
  DC 40 N=1,5
  LL=ID(N)-1
  D=ID(N)*2
  ID(N)=LL
  DA=G(N)
  DZ=G(N+1)
  DELT=(DZ-DA)/D
```

```
DET=DELT*2.
DQ(N)=DELT/3.
DO 39 K=1,LL
DA=DA+DET
DE=DA-DELT
IQ=IQ+1
DX(IQ)=DE
IQ=IQ+1
39 DX(IQ)=DA
DE=DZ-DELT
IQ=IQ+1
DX(IQ)=DE
IQ=IQ+1
40 DX(IQ)=DZ
LL=ID(6)-1
D=ID(6)*2
ID(6)=LL
DA=G(1)
DZ=1./G(6)
DELT=(DZ-DA)/D
DET=DELT*2.
DQ(6)=DELT/3.
IQ=IQ+1
DX(IQ)=1./DA
DO 41 K=1,LL
DA=DA*DET
DE=DA-DELT
IQ=IQ+1
DX(IQ)=1./DE
IQ=IQ+1
41 DX(IQ)=1./DA
DE=DZ-DELT
IQ=IQ+1
DX(IQ)=1./DE
IQ=IQ+1
DX(IQ)=1./DZ
WRITE(6,100)
WRITE(6,101)
WRITE(6,102)
WRITE(6,103)TC,FAK,DLON,XKOH,FREI
WRITE(6,107)FO
WRITE(6,108)
DO 22 K=1,IT
T=TE(K)/TC
B=PI/2.*(1.-T*T)
AG=SQRT(SIN(B))
FL=FLO/AG
O=CG/AG
IF(O-0.5)69,22,22
69 GP=AG/T*FAK
IS=1./SQRT(O*REST)/5.
IP=1./SQRT((2.-O)*REST)/5.
DO 19 I=1,IQ
DS(I)=0.
19 CDS(I)=0.
CALL NONAN
MS=GP/(SQRT(REST)*2.*PI)
CALL ISUM
```



```
DQ1=CDABS(CDS(1))
ENP=1./DSQRT(DQ1)
AK=AL*ENP*AG
A2=AK*AK
CALL HAUPT
CHS=CHS*ENP
CHD=CHD*ENP
XS=REAL(CHS)*DLON
XD=REAL(CHD)*DLON
B=(FO/10.**17)*8.*PI*PI*DLON
RS=AIMAG(CHS)*B
RD=AIMAG(CHD)*B
WRITE(6,109) TE(K),XS,RS,XD,RD,CHS,CHD
22 CONTINUE
100 FCRMAT(1H1,48H SURFACE IMPEDANCE IN THE SUPERCONDUCTING STATE )
101 FORMAT(1H0,48H=====)
102 FORMAT(1H0/2X,19HMATERIALPARAMETERS:)
103 FORMAT(1H0,6H TC = ,F4.2,24H K          GAP(T=0)/KTC = ,F4.2/32H LOND
100N PEN. DEPTH(T=0,1/L=0) = ,F7.1,9H ANGSTROM/32H COHERENCE LENGTH
2(T=0,1/L=0)F = ,F7.1,9H ANGSTROM/32H MEAN FREE PATH
3= ,F7.1,9H ANGSTROM)
107 FORMAT(1H0,9H FOR F = ,E10.4,29H HERTZ THE BCS-THEORY RESULTS/3H
1 /14X,9H SPECULAR-,10X,17H DIFFUSE REFLEXION)
108 FORMAT(1H0,3X,4H TEM-,6X,4H PEN.,4X,7H SURFACE,5X,4H PEN.,4X,7H SURFACE
1/2X,8H PERATURE,3X,5H DEPTH,3X,10H RESISTANCE,2X,5H DEPTH,3X,10H RESIST
ZANCE/4X,3H(K),7X,3H(A),6X,5H(OHM),6X,3H(A),6X,5H(OHM))
109 FORMAT(1H0,3X,F4.2,3X,F8.1,1X,E10.3,1X,F8.1,1X,E10.3,8X,4E11.4)
STOP
END
```

```
SUBROUTINE HAUPT
DIMENSION ID(6)
REAL*8 DQ(6),DX(1000),PI,          DQ1,DS(1000)
COMPLEX*8 CHS,CHD,CI
COMPLEX*16 CDS(1000),CS           ,CT(1000),CIS,CID,CES,CED,CAS,
1CAD,CD
COMMON/HL1/ID,IQ,MS,IS,IP
COMMON/HL2/AK,A2,FL,O,GP
COMMON/HL3/PI,DQ,DX,DQ1,DS
COMMON/HL4/CHS,CHD,CI
COMMON/HL5/CDS,CT,CS,CD
M=1
I=2
CALL INTQ(I,M)
CIS=DX(2)*CS
CID=DX(2)*CD+2.*DX(2)
DO 40 J=1,6
IF(J-6)11,10,10
10 M=2
CALL INTQ(I,M)
CIS=CIS+CS*DX(2)
CID=CID+CD*DX(2)
11 CES=0.
CED=0.
CAS=0.5*CS
CAD=0.5*CD
LL=ID(J)
DO 39 K=1,LL
CALL INTQ(I,M)
CES=CES+CS
CED=CED+CD
CALL INTQ(I,M)
CAS=CAS+CS
CAD=CAD+CD
39 CONTINUE
CALL INTQ(I,M)
CES=CES+CS
CED=CED+CD
CALL INTQ(I,M)
CIS=CIS+DQ(J)*(4.*CES+(2.*CAS+CS))
40 CID=CID+DQ(J)*(4.*CED+(2.*CAD+CD))
A=2./PI*AK
CHS=A*CIS
CHD=2./A/CID
RETURN
END
```

```
SUBROUTINE INTQ(L,M)
DIMENSION ID(6)
REAL*8 DQ(6),DX(1000),PI,          DX2,DA,  DQ1,DS(1000)
COMPLEX*8 CHS,CHD,CI
COMPLEX*16 CDS(1000),CT(1000),CS,CD,CDG
COMMON/HL1/ID,IQ,MS,IS,IP
COMMON/HL2/AK,A2,FL,O,GP
COMMON/HL3/PI,DQ,DX,DQ1,DS
COMMON/HL4/CHS,CHD,CI
COMMON/HL5/CDS,CT,CS,CD
CDG=CDS(L)/DQ1
DX2=DX(L)*DX(L)
DA=DX2*A2
GO TO (10,11),M
10 CS=1./(DA+CDG)
CD=CDLOG(1.+CDG/DA)
GO TO 9
11 CS=1./(A2+CDG/DX2)
CD=CDLOG(1.+CDG/DA)*DX2
9 L=L+1
RETURN
END
```

```
SUBROUTINE NONAN
DIMENSION ID(6)
REAL*8 D ,DX(1000),PI,          DQ1,DS(1000),DV,DZ,DU,DG,DA,DC,DP,DM,
IDW,DWO,DQ(6), OT,O2
COMPLEX*8 CHS,CHD,CI
COMPLEX*16 CDS(1000),CT(1000),CS,CD,CY
COMMON/HL1/ID,IQ,MS,IS,IP
COMMON/HL2/AK,A2,FL,O,GP
COMMON/HL3/PI,DQ,DX,DQ1,DS
COMMON/HL4/CHS,CHD,CI
COMMON/HL5/CDS,CT,CS,CD
OT=O/2.*GP
IF(CT-0.2)10,10,9
10 O2=CT*OT
DV=OT*(1.+O2/6.*(1.+O2/20.*(1.+O2/42.*(1.+O2/72.*(1.+O2/110.)))))*
1DEXP(-OT)
CT=-CT*2.
GO TO 8
9 OT=-CT*2.
DV=(1.-DEXP(OT))/2.
8 B=4*IS
BS=PI/B
DO 48 K=1,IS
B=4*(IS-K)+2
D=B*BS
DZ=0.5*(1.+DCOS(D))
DU=GP/DZ
DA=DEXP(-DU)
DG=DEXP(-DU+OT)
DC=DA/((1.+DG)*(1.+DA))*DV
DP=DZ+1.
DM=1.-DZ
DU=DP*DM
DW=DSQRT(DU)
DA=DZ*O
DWO=DSQRT(DU+DA*(2.+DA))
DP=DZ*DSQRT(DZ*DP)
DU=DA/DU
IF(DU-0.1)7,7,6
7 DU=DU*(2.+DA)/2.
DM=DU*(1.-0.5*DU*(1.-DU*(1.-1.25*DU*(1.-1.4*DU*(1.-1.5*DU*(1.-11.*
1DU/7.*(1.-1.5*DU*(1.-13.*DU/9*(1.-1.4*DU)))))))*DW
GO TO 5
6 DM=DWO-DW
5 CY=-CI*DM/2./DZ+FL
CALL WINK(CY)
DU=DC/DP/DWO
DWU=DWO*DW
DM=1.+DZ*(DZ+O)
DPS=(DM+DWU)*DU
DMS=(DM-DWU)*DU
DO 20 I=1,IQ
20 CDS(I)=CDS(I)+DPS*CT(I)
DM=DWO+DW
CY=CI*DM/2./DZ+FL
CALL WINK(CY)
```

```
DC 21 I=1,IQ
21 DS(I)=DS(I)+DMS*CT(I)
48 CONTINUE
CY=-CI*BS*4.
DO 22 I=1,IQ
CDS(I)=(CDS(I)+DS(I))*CY
22 DS(I)=0.
B=4*IP
BS=PI/B
DO 47 K=1,IP
B=4*(IP-K)+2
D=B*BS
DZ=1.+0.5*Q*(1.+DCOS(D))
DU=GP*DZ
DA=DEXP(-DU)
DG=DEXP(-DU-OT)
DC=DG/((1.+DG)*(1.+DA))*DV
DP=DZ+1.
DM=DZ-1.
DW=DSQRT(DP*DM)
DA=G-DM
DU=DP-G
DWO=DSQRT(DA*DU)
CY=DWO+CI*DW+FL
CALL WINK(CY)
CY=(1.+DZ*(DZ-G)+CI*DW*DWO)/DSQRT(DP*DU)*DC
DO 23 I=1,IQ
23 DS(I)=DS(I)+CY*CT(I)
47 CONTINUE
B=+BS*4.
DO 24 I=1,IQ
24 CDS(I)=CDS(I)+DS(I)*B
RETURN
END
```

```
SUBROUTINE ISUM
DIMENSION ID(6)
REAL*8 DQ(6),DX(1000),PI, G,      DU,DW,DS(1000),DQ1
COMPLEX*8 CHS,CHD,CI
COMPLEX*16 CDS(1000),CS          ,CB,CU,CW,CL,CP,CA,CT(1000),CD
COMMON/HL1/ID,IQ,MS,IS,IP
COMMON/HL2/AK,A2,FL,O,GP
COMMON/HL3/PI,DQ,DX,DQ1,DS
COMMON/HL4/CHS,CHD,CI
COMMON/HL5/CDS,CT,CS,CD
G=PI/GP
DO 20 I=1,IQ
20 DS(I)=0.
DO 30 M=1,MS
B=2*(MS-M)+1
B=B*G
DU=B*B+1.
CB=B+O*CI
CU=CB*CB+1.
DW=DSQRT(DU)
CW=CDSQRT(CU)
CL=(DW+CW)/2.+FL
CP=DW*CW
CA=(1.-B*CB+CP)/CP
CALL WINK(CL)
DO 21 I=1,IQ
21 DS(I)=DS(I)+CA*CT(I)
30 CONTINUE
DO 22 I=1,IQ
22 CDS(I)=DS(I)*G +CDS(I)
RETURN
END
```

```
SUBROUTINE WINK(CL)
DIMENSION ID(6)
REAL*8 DQ(6),DX(1000),PI,      D,DQ1,DS(1000)
COMPLEX*8 CHS,CHD,CI
COMPLEX*16 CDS(1000),CS          ,CT(1000),CD,CC,C2,CL
COMMON/HL1/ID,IQ,MS,IS,IP
COMMON/HL2/AK,A2,FL,O,GP
COMMON/HL3/PI,DQ,DX,DQ1,DS
COMMON/HL4/CHS,CHD,CI
COMMON/HL5/CDS,CT,CS,CC
DO 20 I=1,IQ
D=DX(I)
CD=D/CL
C2=CD*CD
U=CDABS(CD)
IF(U-0.2)10,10,11
10 CT(I)=3./CL*(1./3.-C2*(1./15.-C2*(1./35.-C2*(1./63.-C2*(1./99.-
1C2*(1./143.-C2*(1./195.-C2/255.))))))
GO TO 20
11 CD=1./CD
CT(I)=0.75/D*(-2.*CD+(1.+CD*CD)/CI*(CDLOG(CD+CI)-CDLOG(CD-CI)))
20 CONTINUE
RETURN
END
```