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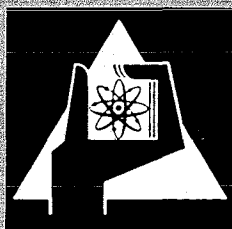
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Institut für Neutronenphysik und Reaktortechnik

**The Numerical Solution of the Three Dimensional
Neutron Transport Equation**

W. Kinnebrock



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On the Numerical Solution of the Three Dimensional
Neutron Transport Equation

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Abstract

M.R. Wagner, D.A. Sargis and S.C. Cohen describe a method of solving the three-dimensional neutron transport equation by using time-saving difference equations. The discrete angular flux is computed in 14 directions. This paper shows how to use a set of 24 angular directions identical with those proposed by B.G. Carlson and K.D. Lathrop for S_4 -calculations.

Zusammenfassung

M.R. Wagner, D.A. Sargis und S.C. Cohen beschreiben eine Methode zur Lösung der dreidimensionalen Neutronen-Transportgleichung mit 14 diskreten Winkelrichtungen. Wie in der vorliegenden Arbeit gezeigt wird, ist es möglich, diese Methode zu kombinieren mit einer diskreten Winkelverteilung, wie sie von B.G. Carlson und K.D. Lathrop für eine S_4 -Rechnung (24 Richtungen) vorgeschlagen wurde.

1. The first part of the document discusses the importance of maintaining accurate records of all transactions and activities. It emphasizes that proper record-keeping is essential for ensuring transparency and accountability in financial reporting. This section also highlights the role of internal controls in preventing errors and fraud, and the need for regular audits to verify the accuracy of the data.

2. The second part of the document focuses on the importance of clear communication and collaboration between all stakeholders involved in the process. It stresses that effective communication is key to ensuring that everyone is on the same page and that all necessary information is shared in a timely and accurate manner. This section also discusses the importance of documenting all decisions and actions taken, and the need for regular updates and reports to keep everyone informed of the progress and any issues that arise.

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Introduction

The long computing time of two-dimensional SN-programs leads to the assumption that three-dimensional SN-programs working with conventional difference equations are ineffective in time for present day computers.

M.R. Wagner, D.A. Sargis and S.C. Cohen, therefore, proposed a method based on time-saving difference equations (2, 3). In connection with this method two problems arise:

- 1) The existence of only 14 discrete angular directions allows the appearance of the well known Ray-Effect.
- 2) Six of the fourteen directions are parallel to the boundaries of zones. In most cases the angular flux is not continuous (at least in one derivative) for these directions; consequently these directions are not favorable for a discrete numerical integration.

Subsequently a method is described using the difference equations of 2, 3. The set of angular directions is identical with that of a conventional S_4 -calculation with 24 directions, proposed by B.G. Carlson and K.D. Lathrop in 1.

It can be assumed that by using these directions we have a method being effective in computing time and accuracy of angular quadrature.

1. Difference Equations and Spatial Mesh

The transport equation

$$(1) \quad \Omega \cdot \nabla \cdot \phi + \Sigma_a \cdot \phi = S \quad (S=\text{source})$$

can be written as

$$(2) \quad \frac{dg(t)}{dt} + \Sigma_a g(t) = S$$

$$\text{if } g(t) = \phi(w+t \cdot \Omega) = \phi(w + t \Omega, \Omega, E)$$

$$\text{For constant } S \text{ and } \Sigma_a \text{ we get } g(t) = g(0) \exp(-\Sigma_a \cdot t) + S/\Sigma_a$$

It follows

$$(3) \quad \phi(w+t \cdot \Omega) = \exp(-\Sigma_a \cdot t) \cdot \phi(w) + S/\Sigma_a$$

and (3) is for piece-wise constant S and Σ_a a suitable difference equation.

We now consider a rectangular cubic mesh with the constant spacing a . S and Σ_a may be constant within a cube. The flux may be defined in the following points: the central point P_0 of the cube, the central points of the six cube sides, the eight corner points. The corner points may be defined through P_1, P_2, \dots, P_8 corresponding to the arrangement of figure 1, and the central points of the sides may be defined through P_{ijk1} , if P_i, P_j, P_k, P_l are the corner points of the sides.

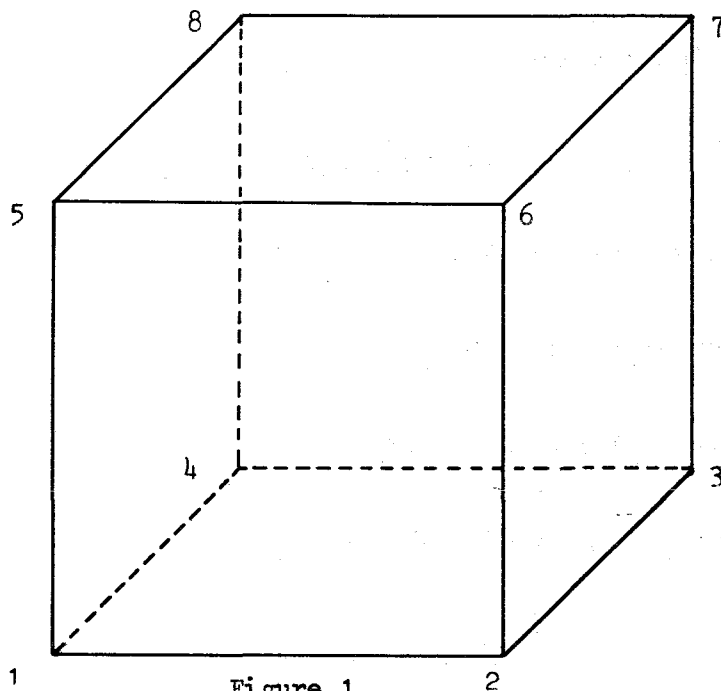


Figure 1

We now choose as discrete directions Ω_j the directions of the straight lines which connect the corner points with one of the central points of the three opposite sides. So we have, for instance,

$$\begin{aligned} \Omega_1 : P_1 &\rightarrow P_{2367} \\ \Omega_2 : P_1 &\rightarrow P_{3478} \\ \Omega_3 : P_1 &\rightarrow P_{5678} \end{aligned}$$

Let us now consider Ω_1 . Formula (3) enables us to compute the angular flux of this direction by the following equations:

$$\begin{aligned} (4) \quad \phi_{2367}^1 &= \exp(-\Sigma_a \cdot l) \cdot \phi_1^1 + S/\Sigma_a \\ \phi_7^1 &= \exp(-\Sigma_a \cdot l) \cdot \phi_{1458}^1 + S/\Sigma_a \end{aligned}$$

where ϕ_{ijke}^m denotes the flux with direction Ω_m at the point P_{ijke} , ϕ_q^m is the flux in P_q and l means the length of the line $\overline{P_1 P_{2367}}$. The angular

flux in P_0 is obtained from

$$\phi_0^1 = \frac{1}{4} \cdot \int \phi_{2367}^1 + \phi_1^1 + \phi_7^1 + \phi_{1458-7}^1$$

In the same way we have for Ω_2 :

$$\phi_{3478}^2 = \exp(-\Sigma_a \cdot 1) \cdot \phi_1^2 + S/\Sigma_a$$

$$\phi_7^2 = \exp(-\Sigma_a \cdot 1) \cdot \phi_{1256}^2 + S/\Sigma_a$$

$$\phi_0^2 = \frac{1}{4} \cdot \int \phi_{3478}^2 + \phi_1^2 + \phi_7^2 + \phi_{1256-7}^2$$

etc.

After computing all the angular dependent fluxes ϕ_0^j ($j=1,2,\dots,24$) we get the total flux $\hat{\phi}_0$ in P_0 by

$$\hat{\phi}_0 = \sum_{j=1}^{24} w_j \phi_0^j$$

with suitable weights w_j .

3. Weights and Directions

As a result of the distribution of the angular directions described in chapter 2 we have 24 directions or 24 points on the unit sphere. There are three points per octant and the set of points is invariant under all 90-degree rotations around one axis. Thus, the description of one octant is sufficient to describe the arrangement of all points.

The direction Ω_j may have the coordinates (μ_j, η_j, ζ_j) with

$$\mu_j^2 + \eta_j^2 + \zeta_j^2 = 1$$

Figure 2 shows the directions of one octant as a function of their coordinates ($A = \sqrt{1/6}$)

	Ω_1	Ω_2	Ω_3
μ	2A	A	A
η	A	2A	A
ζ	A	A	2A

Figure 2

The geometrical arrangement on one octant is

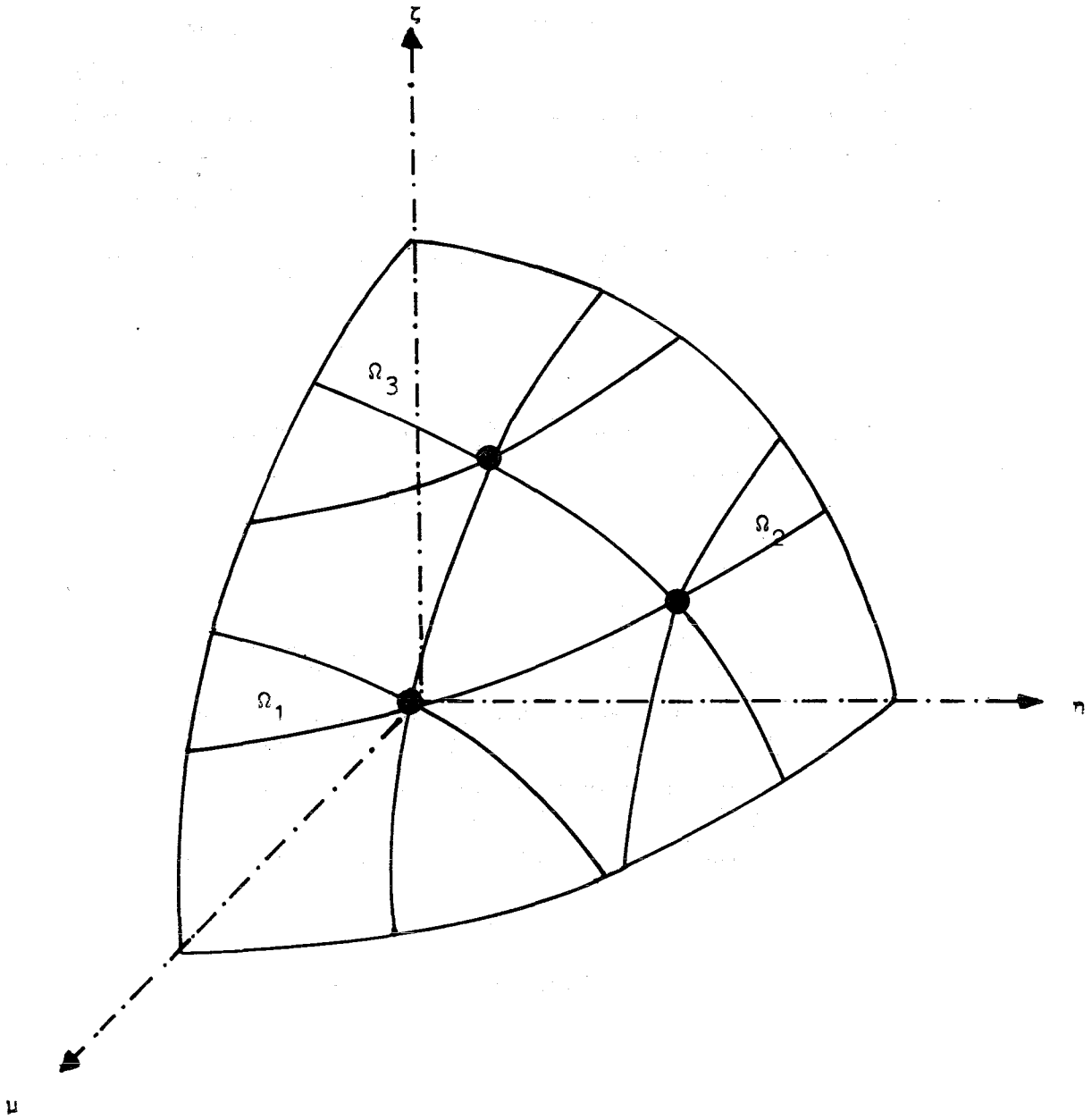


Figure 3

shown in figure 3. The directions satisfy the requirements given by B.G. Carlson and K.D. Lathrop for suitable S_4 -directions. They can be calculated by formula (4) in [3].

The weights w_j for numerical integration must be identical with those contained in [3]. We, therefore, choose $w_j = w = 1/24$ for all j . Then we have

$$(5) \quad \frac{1}{2} \int_{-1}^1 \mu^k d\mu = \sum_{j=1}^{24} \mu_j^k w = \begin{cases} 0 & (k \text{ odd}) \\ \frac{1}{k+1} & (k \text{ even}) \end{cases}$$

for $k=0,1,2,3$. For $k=2$ (5) is the well known "diffusion theory condition".

A code based on the theory described should be effective in computing-time and accuracy of the angular integration.

Literature

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