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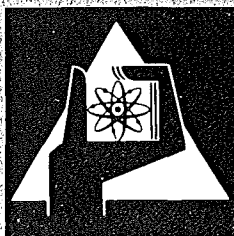
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Institut für Neutronenphysik und Reaktortechnik
Projekt Schneller Brüter

**MIGRØS - 2: A Program Written in FORTRAN for the Calculation
of Microscopic Group Constants from Nuclear Data**

Compiled by
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**GESELLSCHAFT
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Abstract

MIGRØS-2 allows the production of microscopic group constants of the ABN - type. The Karlsruhe nuclear data file KEDAK is used as the nuclear data basis. All group constants necessary for diffusion-, consistent P_1 -, and S_n -calculations, taking the anisotropy into account up to P_5 , can be calculated.

A description of the code and the underlying theory is given. An input and an output description, a sample problem and the program lists are provided.

Zusammenfassung:

MIGRØS-2 erlaubt die Bereitstellung von mikroskopischen Gruppenkonstanten vom ABN-Typ. Als Datenbasis wird die Karlsruher Kerndatenbibliothek KEDAK benutzt. Es können alle Gruppenkonstanten berechnet werden, die für Diffusions-, konsistente P_1 - und anisotrope S_n -Rechnungen benötigt werden.

Der Code und die zugrundeliegende Theorie werden beschrieben.

Es werden ein Eingabe - und Ausgabebeschreibung, ein Rechenbeispiel und die Programmlisten bereitgestellt.

Table of Contents:

Introduction

1. The organization of MIGRØS - 2
H. Huschke, B. Krieg
2. Description of the input and the output
H. Huschke, B. Krieg
3. The calculation of average group cross sections
H. Huschke, B. Krieg
4. The calculation of average group cross sections for infinite dilution and of energy resonance self shielding factors from resolved resonance parameters.
H. Huschke, B. Krieg
5. The calculation of average group cross sections for infinite dilution and of energy resonance self shielding factors from statistical resonance parameters.
H. Huschke, B. Krieg
6. The calculation of the matrices of the average zero'th moments of inelastic scattering.
B. Schatz, I. Siep
7. Microscopic multigroup scattering matrices for elastic neutron scattering in a P_1 - approximation.
H. Wiese
8. The calculation for the fine interval elastic scattering matrices up to P_5 approximation and other quantities required for the "Remo" - correction.
P. Vertes
9. The calculation of average group cross sections for infinite dilution and of energy resonance self shielding factors from energy point wise data in the resonance region.
H. Huschke, B. Krieg

10. The calculation of group constants in the thermal group.

H. Huschke, B. Krieg

11. The calculation of $1/v$ - average group values.

H. Huschke, B. Krieg

12. The calculation of fission spectra.

H. Huschke, B. Krieg

Appendix I. A modified version of the formalism of R. Froelich for the calculation of energy resonance self shielding factors in the range of unresolved resonances.

H. Huschke

Appendix II. An approximate calculation of current weighted resonance self shielding factors from statistical resonance parameters.

H. Huschke

Appendix III. The method of integration by a Romberg procedure.

P. Vertes

Appendix IV. The subintervals in the angular integration.

P. Vertes

Appendix V. Sample problem

Appendix VI. Program lists

INTRODUCTION

The code MIGROS-2 has been established to produce microscopic group constants for arbitrary isotopes. MIGROS-2 allows the production of group cross section sets of the ABN-type [1]. The Karlsruhe nuclear data file KEDAK [2] is used as the nuclear data basis. The code has been further developed from the code MIGROS [3], which has been used so far for the production of the Karlsruhe 26-group cross section sets. The microscopic group constants, which are calculated by MIGROS-2, can be used for the production of macroscopic group cross sections for diffusion, consistent P_1 - and for S_n calculations, taking the anisotropy into account up to P_5 . They are :

- average group cross sections for infinite dilution (without resonance self shielding) for all types of neutron reactions, from energy dependent cross sections tabulated point by point;
- energy resonance self shielding factors and average group cross sections for infinite dilution for capture, fission, elastic scattering and total neutron reaction from resolved resonance parameters, from statistical resonance parameters and from energy dependent cross sections, tabulated point by point, in the resonance region;
- matrices of the elastic scattering up to the P_5 -approximation from angular distributions;
- matrices of the zero'th moment of inelastic scattering from inelastic excitation cross sections or by an evaporation model;
- $1/v$ -average group values;
- fission spectra;
- cross sections for the thermal energy region (one energy group).

The code MIGROS-2 has a modular structure, which is very flexible in its need of storage.

A standard input is provided internally for the production of a 26-group cross section set with the same group structure as the ABN-set. The output of the results is given in two forms; on paper and on an external unit.

The input of MIGROS-2 is very versatile. If the user does not use the standard

input, he has to prove whether or not the input is physically meaningful and the methods in the used modules are adequate. In the following a short description of the task of MIGROS-2 is given.

The calculation of average group cross sections from pointwise data on KEDAK.

Group averages for all types of cross sections that are available on KEDAK are calculated. A trapezoidal rule for the integration is used. As integration points the energy points on KEDAK are used. Besides the average group cross-sections group values for the average cosine of the elastic scattering, for the number of secondaries per fission, for α and η are calculated.

The calculation of energy resonance self shielding factors and average group cross sections from resolved resonance parameters and from statistical parameters.

For capture, fission and elastic scattering flux weighted ($w(E) \sim \frac{1}{\sigma_t(E)+\sigma_0}$) resonance self shielding factors and for elastic scattering and the total neutron reaction current weighted ($w(E) \sim \left\{ \frac{1}{\sigma_t(E)+\sigma_0} \right\}^2$) self shielding factors dependent on temperature and on the background cross section σ_0 and the average group cross sections are calculated from resonance parameters. In the resolved resonance region Breit-Wigner single level parameters and Doppler-broadened line shape functions are used. Interference between potential and resonance scattering and the overlapping of resonances is taken into account exactly.

In the unresolved resonance region an analytical model, based on average resonance parameters and χ^2 -distributions for the parameters, is used. Interference between potential and resonance scattering is partly taken into account. Overlapping of resonances is considered approximately.

The calculation of energy resonance self shielding factors from pointwise data.

The same types of self shielding factors that are calculated from resonance parameters can also be calculated from pointwise data by a trapezoidal rule. They can only be calculated for the temperature for which the cross sections are tabulated.

The calculation of the zero'th and higher moments of the elastic scattering transfer elements.

There exist two approaches. In the first one average transfer elements from one group into another or into the same group are calculated. The average over the outscattering group may be weighted.

In the second approach, the outscattering group is subdivided into several intervals. To describe the transfer from one group to another, for each outscattering group the average values over the intervals are calculated. In both cases an integration is done over the in-scattering groups.

Both approaches expect tabulated angular distributions. The angles must be the same for all energies, at which the distributions are tabulated. The two approaches use different interpolation and quadrature procedures. In the first one the angular distributions are transformed from the centre of mass system to the laboratory system. The interpolations in the energy and in the angle are done by polynomials of the order four. The integration over the in-scattering group is performed as integration over the corresponding angle interval in the laboratory system by a Simpson rule, the integration over the outscattering group is performed by a trapezoidal rule. In the second approach the interpolation for the angle in the centre of mass system is done by a linear interpolation, the interpolation in energy is performed logarithmic. The integration over the in-scattering groups is transformed to an integration over the corresponding angle intervals in the centre of mass system. Both integration, over the angle intervals and over the outscattering groups is done by a Romberg procedure.

The calculation of the inelastic scattering matrix.

For the calculation of the inelastic scattering matrix in the range of discrete levels excitation cross sections are used. The total inelastic scattering cross section and the excitation cross sections are interpolated linearly. All energy integrations are done by a trapezoidal rule. At higher energies the Weißkopf evaporation model is used.

The calculation of the fission spectra.

Fission spectra are calculated by a Watt-formula. The fragment kinetic energy per nucleon and the Watt fragment nuclear "temperature" are used.

Thermal cross sections.

One thermal group with a Maxwellian spectrum is assumed.

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KFK-tr-144
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1. The organization of MIGRØS-2

1.1 The logical structure of MIGRØS-2

The code MIGRØS-2 has a modular structure. It consists of a control module and of several computational modules. In the control module all the input information, except the nuclear data, is checked, the order of processing is fixed and the information for the single computational modules is prepared. A computational module is defined as a logically complete, executable computer procedure, consisting of several subroutines and functions. In a computational module a specific type of group constants is calculated. The necessary nuclear data are read from the Karlsruhe nuclear data library, KEDAK /1/, in the computational modules. Also in these the group constants are output on paper and on an external storage unit. The available computational modules are given in table I.

The control is transferred from the control module to a computational module by a CALL-statement which calls the entry subroutine of the module. It is not possible to transfer the control directly from one module to another. A computational module can be called by the control module several times.

All computational modules are written in FØRTRAN-IV. All arrays are variably dimensioned, when they are influenced by input information. The control module is written in FORTRAN-IV, except for the routine DATAMAIN, which is written in ASSEMBLER.

Since the input of nuclear data and the output of the results is performed in the computational modules, the only information that must be transferred between the control module and the computational modules is input information other than nuclear data and information necessary for the control of the program. The transfer of data is partly done by an unlabeled CØMMØN-array. However, most of the data transfer is accomplished by the parameter lists in the entry subroutines of the computational modules.

MIGROS-2 is organized in an overlay structure, given in table II.

1.2 The control module

The control module has three tasks:

- to read all input information except the nuclear data,
- to provide a central data array and the unlabeled COMMON-array,
- to control the order of computation of the computational modules.

In the routines DATAMAIN und MAIN a central data array is provided, whose length can be defined by the user in the IBM/OS control language (see chapter 2).

In the subroutine INPUT, all input information except the nuclear data is read and checked. The central data array is prepared and the computational modules are called.

The input information is stored and the working arrays for the computational modules are provided in the central data array. It consists of two sections. One is permanent and the same for all computational modules. The other section depends on the particular module being considered and is allocated only when the control is actually transferred to the module. All information of the central data array is transferred by parameter lists in the entry sub-routines of the computational modules. A description is given in table III.

Also there is an unlabeled COMMON-array, described in table IV.

1.3 The organization of the central data array XL

The ASSEMBLER-routine DATAMAIN provides a central array named XL for the subroutine MAIN. The length of XL is defined by the user in the PARM.G-parameter of the EXEC-card in the IBM/OS control language. The subroutine MAIN calls the subroutine EING, which first prints an input description of the MIGRØS-system. Then all the input for the first material is read and that necessary for the variable dimensioning of the input arrays in the XL-array is retained, e.g., the number of group boundaries, the number of points of the weighting function, the number of computational modules that shall be used etc. The input unit is "backspaced" to the beginning and the subroutine INPUT is called. The starting addresses of the input arrays in the XL-array are arguments of INPUT. The index NFR indicates the beginning of the non-permanent section of the XL-array and is also an argument of the subroutine INPUT. All input arrays are filled with the information from the input cards or the standard built-in data. The input is also checked for consistency. Then the computational modules are called in the order fixed by the input.

Before every call of a computational module the length of the non-permanent section of the XL-array is checked to determine, if it is long enough for the working arrays of the module. The length of these working arrays are dependent on the input information and in some cases, can finally be defined only in the computational module itself. In this case an iterative procedure is used to determine the proper length. As a first step, a minimum length of the working array is assumed. If it turns out to be too small, control is given back to the control module and the size of the array is increased. When the required size of the working array is so large that the XL-array is not long enough the computational module in question is passed and a message to the user is given.

After all the required computational modules for one material have been called, a RETURN-statement to the subroutine MAIN is given. If there are more materials, a new arrangement of the XL-array is made etc.

1.4 Auxiliary subroutines

1.4.1 Reading program for the Karlsruhe nuclear data library KEDAK

For reading KEDAK, the subroutine NDF is necessary. A description is given in /3/.

1.4.2 Weighting functions

In the case, that there are specified no weighting functions for the group cross sections and matrices by input, the following two functions are used:

```
FUNCTION PHI(E)
  PHI = 1./E
  RETURN
END
```

E is the energy in [eV].

This weighting function is used for all group constants except the first and higher Legendre moments for elastic scattering.

```
FUNCTION PHI1 (L,E)
  PHI1 = 1./E * * L
  RETURN
END
```

E is the energy in [eV]

L = l + 1, l = Legendre moment.

This weighting function is used for the first and higher Legendre moments in the module 6 (not in module 9).

Table I

number of the module	described in	content	names of the subroutines and functions	entry subroutine
1	chapter 4	energy resonance self shielding factors from resolved Breit-Wigner resonance parameters	DOPW, EXPPX, FGEM, NDF, PHI, PSIXI, STOSS, WIRQ, WIRQU	FGEM
2	chapter 5	energy resonance self shielding factors from statistical Breit-Wigner parameters	DELTA, DMIT, DØPW, EPSI, EZZ, FSTAT, GAFM, GAMG, GAMN, NDF, PHASE, PHI, PØL, QUER, SIGC, SUCH, TAB	FSTAT
3	chapter 9	energy resonance self shielding factors from neutron cross sections given point by point in energy	DØPW, FSTRUK, NDF, PHI	FSTRUK
4	chapter 3	average group cross sections, for infinite dilution (without resonance self shielding from point by point data in the energy and average group values $\bar{\mu}$, η , α , ν)	DØPW, NDF, PHI, SUND	SUND

number of the module	described in	content	names of the subroutines and functions	entry subroutine
5	chapter 6	zero'th moment of the inelastic scattering from discrete levels and from an evaporation model	AKED, DNFAK, DØUB, EXD, NDF PHI, SCAT, TRA, XKØN	SCAT
6	chapter 7	zero'th and higher moments of the elastic scattering from angular distributions	ADD, FLUMMI, GRUPIN, G, IN, FØRM, IPØLA, IPØLIN, LECAL, LEGANS, LEGINT, LEGIST, LEGPØL, LØØKO, LØØK1, LØØK2, LØØK3, MAKRØ, MIXSGT, MUKØN, NDF, PHI, PHIL, PRINT, PUNK, PUSUM, SPRAL, SUM, TRAFØ, TRAPEZ	FLUMMI
7	chapter 12	fission spectra	DØPW, NDF, SPALT	SPALT
8	chapter 11	1/V-average group values	DØPW, EDV, PHI	EDV
9	chapter 8	zero'th and higher moments of elastic scattering for the "REMO"-correction	AKØR, AMESH, ANINT, BCM, EGRENZ, FXINT, HIDR, ICSØP, INTEN, ISØFAL, IWIN, IWØ KEDDAT, LMI, MASSIN, NDF, NØRM, PHI, PTL, REMØ, REMP SEARCH, SINT, SMØRN, SUCHM, WAHRS, WINK, ZWIN	REMØ

number of the module	described in	content	names of the subroutines and functions	entry subroutine
10	chapter 10	group cross sections in the thermal group	DØPW, NDF, THERM	THERM
control module	chapter 2		DATAMAIN, DØPW, EING, FREEFØ, FSTAE, INPUT, MAIN, NDF	

Table II The overlay structure of MIGRØS-2

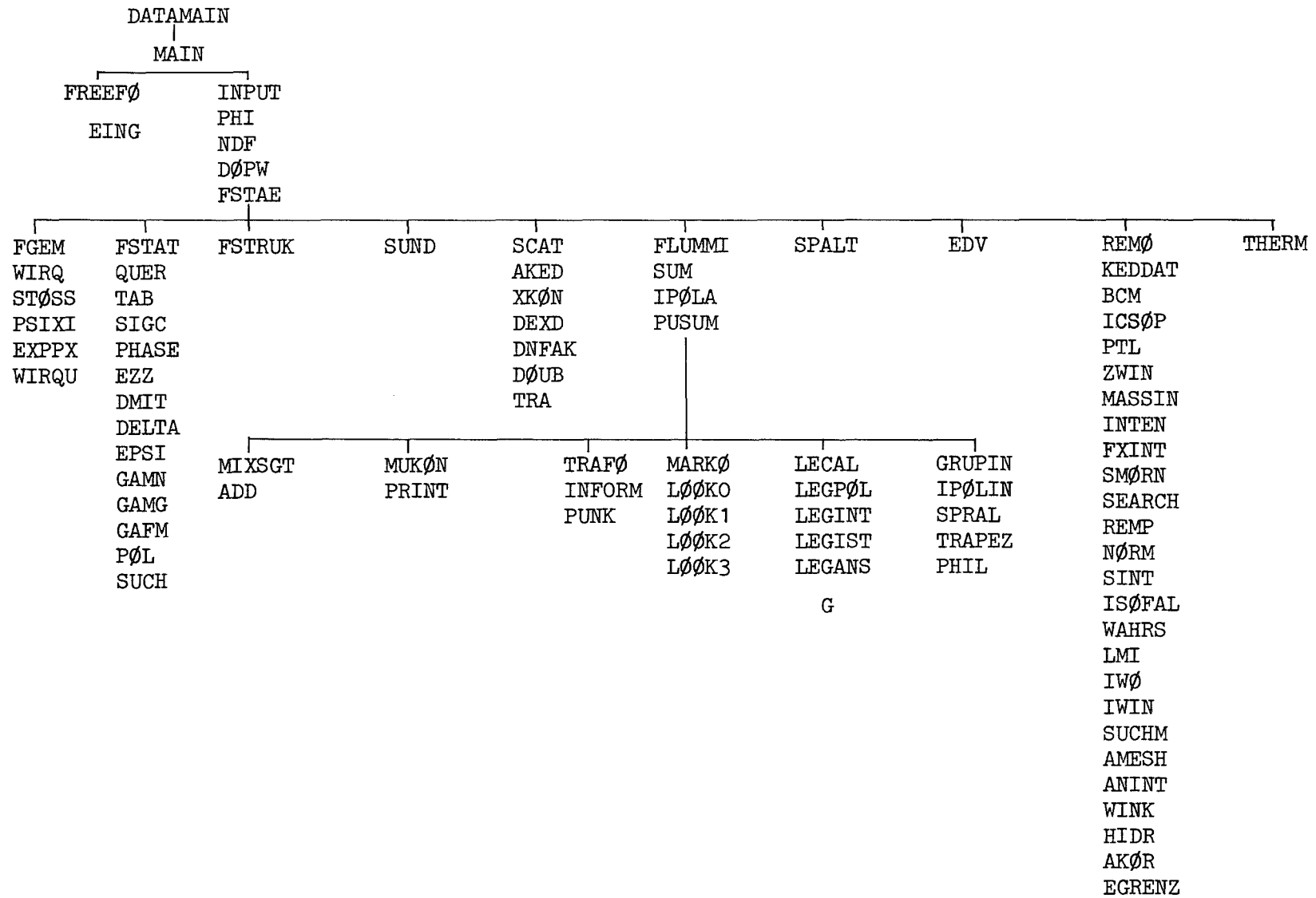


Table III

Structure of the central data array XL in the subroutine MAIN.

a) The permanent section

The symbols of the input description in chapter 2 are used.

length of the subarray	type	content of the subarray	see card n of the input description
NTYP	Real 8	ITYP(I), I=1 NTYP	15
NMAT	Real 8	NAME(I), I=1, NMAT	25
NMAT	Real 4	TZ(I), I=1, NMAT	25
NE	Real 4	ENG(I), I=1, NE	5
NFE	Real 4	EF(I), I=1, NFE	7
NFE x NSPEC	Real 4	(F(I,J), I=1, NFE),	}
		J=1, NSPEC	
3 x NA	Integer 4	(NR(I,J),I=1,3)J=1, NA	9
NT	Real 4	TEMP(I), I=1, NT	11
MI+1	Real 4	SIGO(I),I=1, MI	13
NA	Integer 4	NGRE(K), K=1, NCALL	9, 34, 33
NA x NE	Integer 4	(NI (K,I),I=1, NGRE) K=1, NCALL	9, 5, 34, 33
NA x NE	Integer 4	(N2)K,I), I=1, NGRE K=1, NCALL	9, 5, 34, 33
NA x NE	Integer 4	(NFG(K,I=1, NGRE) K=1, NCALL	9, 5, 34, 33
NA x NE	Integer 4	(NFI(K,I), I=1, NGRE K=1, NCALL	9, 5, 34, 33
2 x NE	Integer 4	this subarray is de- fined, but information is actually only stored if module 9 is called	5

Table III, cont.

b) The non-permanent section

Besides the permanent section a non-permanent section is defined, when a computational module is actually called. This non-permanent part exists only during the execution of a computational module. The non-permanent section starts with the index NFR, defined in the subroutine MAIN. The names of the auxiliary arrays are the same as used in the formal parameters of the entry subroutines of the computational modules. All other symbols are the same as in the input description of chapter 2.

b1) Module 1

length of the subarray	type	content of the subarray	see card n of the input description
MI x 7	Real 4	SUM, auxiliary array	13
NE x 3	Real 4	SUØ, " "	5
max (300,r)	Real 4	ER, " "	
max (300,r)	Integer 4	L, " "	
max (300,r)	Real 4	GJ, " "	
max (300,r)	Real 4	GAT, " "	
max (300,r)	Real 4	GAN, " "	
max (300,r)	Real 4	GAG, " "	
max (300,r)	Real 4	GAF, " "	
5 x max (400,i)	Real 4	STE, " "	

r = number of resonances on the KEDAK-library for the isotope in question

i = number of integration points for the integration of the effective group cross section in one energy group.

b2) Module 2

length of the subarray	type	content of the subarray	see card n of the input description
3 x 5 x MI	Real 4	SE auxiliary array	13
3 x 5 x MI	Real 4	SM " "	13
5 x MI	Real 4	XEUGZ " "	13

b3) Module 3

length of the subarray	type	content of the subarray	see card n of the input description
NE x 4	Real 4	XINTE, auxiliary array	5
MI x NE x 5	Real 4	ZA, " "	13, 5
MI x NE x 4	Real 4	XN, " "	13, 5
MI x NE x 6	Real 4	SE, " "	13, 5
NE	Real 4	XII, " "	5
MI x NE	Real 4	ZB, " "	13, 5
NE	Real 4	XI, " "	5
max (1500,i)	Real 4	EMU, " "	-
max (1500,i)	Real 4	XMU, " "	-
3 x max (1500,j)	Real 4	SN, " "	-
3 x max (1500,j)	Real 4	EN, " "	-

i = number of $\bar{\mu}_e$ -values on the KEDAK library for the isotope and energy range in question.

j = maximum number of σ_a^- , σ_e^- , or σ_f^- -values on the KEDAK library for an isotope and energy range in question.

b4) Module 4

length of the subarray	type	content of the subarray	see card n of the input description
NE	Real 4	SGC, auxiliary array	5
NE	Real 4	DUE, " "	5
NE	Real 4	XINTE, " "	5
NE	Real 4	ZINT, " "	5
NE	Real 4	XNEN, " "	5
NE	Real 4	STREU, " "	5
max (1500,i)	Real 4	SE, " "	-
max (1500,i)	Real 4	FSE, " "	-

i = number of $\bar{\mu}_e$ -values or ν_f -values for one isotope on the KEDAK library.

b5) Module 5

length of the subarray	type	content of the subarray	see card n of the input description
NE	Real 8	WAHR, auxiliary array	5
NE	Real 8	PRØB, " "	5
NE	Real 8	VW, " "	5
NE	Real 8	QUER, " "	5
NE + 1	Real 8	E, " "	5
max (700,i)	Real 8	ET, " "	-
max (25,j)	Real 8	AE, " "	-
NE	Real 4	SU, " "	5
NE	Real 4	QUØT, " "	5
NE	Integer 4	LBA, " "	5
NE	Real 4	WEIN, " "	5
NE + 1	Real 4	AG, " "	5

length of the subarray	type	content of the subarray	see card n of the input description
max (700,i)	Real 4	SGIT, auxiliary array	-
max (25,j)	Integer 4	KMAX, " "	-
max (400,j,k)	Real 4	SGIP, " "	-
max (400,i)	Real 4	WERT, " "	-

i = number of σ_{in} -values for one isotope on the KEDAK-library.

j = number of excitation levels for one isotope on the KEDAK-library.

k = number of excitation cross sections for one isotope on KEDAK-library.

b6) Module 6

length of the subarray	type	content of the subarray	see card n of the input description
max (i,400)	Real 4	A, auxiliary array	-
max (i,400)	Real 4	B, " "	-
max (i,400)	Real 4	EA, " "	-
max (i,400)	Real 4	EB, " "	-
max (i,400)	Real 4	E, " "	-
max (j,1000)	Real 4	EN, " "	-
max (j,1000)	Real 4	SGN, " "	-
max (j,1000)	Real 4	EC \emptyset , " "	-
max (j,1000)	Real 4	SC \emptyset , " "	-
max (j,1000)	Real 4	V, " "	-
max (j,1000)	Real 4	W, " "	-
max (j,1000)	Real 4	F, " "	-
k	Real 4	AR, " "	-
k	Real 4	FU, " "	-
NE + 1	Real 4	ABN, " "	5

length of the subarray	type	content of the subarray	see card n of the input description
NE + 1	Integer 4	INT, auxiliary array	5
NE + 1	Real 4	GR, " "	5
NE + 1	Real 4	FEKØE, " "	5
NE	Real 4	R, " "	5
NE	Real 4	RSP, " "	5
k x i	Real 4	SGNC, " "	-
l x i	Real 4	FEKØ, " "	-
6 x l x NE	Real 4	ELSIG, " "	5
2 x NE	Real 4	ELTØT, " "	5
max (m,3000)	Real 4	ET, " "	-
max (m,3000)	Real 4	ST, " "	-

i = 300 + maximum number of energy points for the reactor type SGNC on the KEDAK-library for four neighbouring energy groups.

j = i + maximum number of energy points for the reaction types SGN or MUEL or SGT on KEDAK in one energy group.

k = number of angle-mesh-points for the reaction type SGNC on KEDAK.

l = maximum number of elements of the elastic scattering matrix for one outscattering group.

m = maximum number of energy points for the reaction type SGT on KEDAK in four neighbouring groups for a mixture of isotopes.

b7) Module 7

length of the subarray	type	content of the subarray	see card n of the input description
NE	Real 4	X, auxiliary array	5

b8) Module 8

length of the subarray	type	content of the subarray	see card n of the input description
NE	Real 4	E, auxiliary array	5
NE	Real 4	V, " "	5
NE	Real 4	Y, " "	5
max (NFE, 10 x NE)	Real 4	ES, " "	5
max (NFE, 10 x NE)	Real 4	F, " "	5

b9) Module 9

length of the subarray	type	content of the subarray	see card n of the input description
XL-array minus the permanent section	Real 4	WORK, auxiliary array	-

b10) Module 10

This module uses only the permanent section of the XL-array.

Table IV Structure of the unlabeled ~~COMMON~~-array

length of the subarray	type	content of the subarray	see card n of the input description
1	Real 8	MAT	1
1	Integer 4	ISTRUK	1
1	Integer 4	ISPA	3
1	Integer 4	unit number for print output	
1	Integer 4	external storage number for the output of the results	
2	Integer 4	NR(2,J), NR(3,J) J is the index of the module actually called	9
1	Integer 4	J, at the end of each module J must be enlarged by one	9

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/3/ B. Krieg

Handling and Service Programs for the Karlsruhe Nuclear Data
File KEDAK

Part I: Management and Retrieval Programs

KFK 1725, 1972

2. Input and output of MIGRØS-2

2.1 Description of the input

In the following input description the expression "card" is used for the information that is read by one READ-statement. The input is unformatted. The following rules must be observed: The information of one "card" may actually be punched on several cards. Column one of the first physical card must not be blank, for the following physical cards belonging to the same READ-statement it must be blank. The information of two words must be separated by at least one blank. Alphanumeric information may consist of eight symbols or less. It is represented in the form 'TEXT bbbb'. The blanks and apostrophes are obligatory. For alphanumeric information consisting of five or less symbols the following representation is also allowed: α TEXT α . In this case no blanks are necessary. They are only necessary, if several words are linked: α TEX1 α α TEX2 α or linked α TEX1bTEX2b α .

2.1.1 Input description

Card 1

MAT	name of the material in the KEDAK-nomenclature /1/
ISTRUK	= 0: light or medium weight material, no temperature dependent resonance cross- sections = 1: heavy isotopes, temperature dependent resonance cross sections

must cover the whole energy range, in which weighted group constants shall be calculated.

If card 6 and card 7 are not provided, the following standard is used:

F(I,1) is given by the FUNCTION PHI (E) and (F(I,L), L=2, NLE + 1) by FUNCTION PHI1 (L,E) (NLE is defined in card 23).

FUNCTION PHI(E)

PHI = 1./E

RETURN

END

FUNCTION PHI1 (L,E)

PHI1 = 1./E x x L)

RETURN

END

E: energy

See also card 27

Card 8

α BL ϕ C α

4 constant

Card 9

NA

number of calls of any computational modules

((NR(I,J), I=1,3), J=1,NA)

NR(1,J) identification number of the module to be called (see table I. in chapter 1)

NR(2,J) number of the first
energy group,

NR(3,J) number of the last energy
group, for which the
module NR(1,J) is called.

Comment:

The numbers of the energy groups
are fixed by the chosen group
boundaries in card 4 and card 5.
the groups are numerated from 1
to NE-1, starting with the group
of the highest energy.

NR(2,J), NR(3,J) must be ordered
with increasing energy, e.g.
 $NR(2,J) \geq NR(3,J)$. The module
NR(1,J) is executed for all
groups from NR(2,J) to NR(3,J).
A module can be specified several
times. There is no limitation for
NA. In the case of scattering
matrices, NR(2,I) and NR(3,J)
are the outscattering groups.
The transfer elements for all
inscattering groups belonging
to an outscattering group are
calculated.

In the case when one of the modules 1 or 2 for the calculation of
selfshielding factors is called, the following two cards are
necessary:

Card 1a

$\alpha B L \phi C \alpha$

5 constant

Card 11

NT number of temperatures
(TEMP(I), I=1, NT) temperatures in $^{\circ}\text{K}$

Comment: If card 10 and card 11 are not specified, the following standards are used:

300 $^{\circ}\text{K}$, 900 $^{\circ}\text{K}$, 2100 $^{\circ}\text{K}$ for heavy isotopes (ISTRUK = 1)
0 $^{\circ}\text{K}$ for light and medium weight isotopes, (ISTRUK = 0).

There is no limitations for the number of temperatures.

For the case when one of the modules 1, 2 or 3 for the calculation of self shielding factors is called, the following two cards are necessary:

Card 12

$\alpha\text{BL}\phi\text{C}\alpha$ 6 constant

Card 13

MI number of background cross sections σ_0

(SIGO(I), I=1, MI) background cross sections σ_0 in (barn)

Comment: If card 12 and card 13 are not specified, the following standards are used: MI = 7,
 $\sigma_0 = 0, 10, 10^2, 10^3, 10^3, 10^4, 10^5, 10^6$ barns
There is no limitation for the number of σ_0 -values.

If the module 6 is called, the following eight cards are necessary:

Card 22

$\alpha_{BL} \phi_{Ca}$ 11 constant

Card 23

ISEL parameter, which defines the averaging procedure for the transferelements. ISEL = 2: The transferelements are the group averaged transfer cross sections normalized to the group averaged total elastic cross section.

ISEL = 1: The transferelements are the group averaged, to the total elastic cross sections normalized elastic transfer cross sections.

NLA lowest Legendre moment

NLE highest Legendre moment

Comment: NLA ist always set to zero by the program.
NLE must be smaller or equal to five. If $NLE < 1$, NLE is set to one by the program.
If card 22 and card 23 are not specified, the following standards are used:

ISEL = 2 for ISTRUK = 0
ISEL = 1 for ISTRUK = 1
NLA = 0
NLE = 5

Card 24

$\alpha_{BL\emptyset C\alpha}$ 12 constant

Card 25

NMAT number of materials in a composition

(NAME(I), TZ(I), I=1, NMAT)

NAME: names of the materials
TZ: atomic densities x 10^{-24}

Comment: card 24 and card 25 are only necessary if the fine structure of the weighting function for a defined composition is to be taken into account (see chapter 7)

Card 26

$\alpha_{BL\emptyset C\alpha}$ 13 constant

Card 27

MAZ(1) = 0: all Legendre moments are weighted with the macroscopic weighting function of the zero'th moment (see weighting function in card 7, F(I,1)).
= 1: each Legendre moment is weighted with it's own macroscopic weighting function (see card 7, F(I,J)).

MAZ(2) = 0: all Legendre moments are weighted with the microscopic weighting function of the zero'th moment $\sim 1/\Sigma(E)$, $\Sigma(E)$!

macroscopic total neutron
cross section for a composition.

= 1: the microscopic weighting
function for the l -th Legendre
moment is $\sim \{ 1/\Sigma \}^{l+1}$

Comment:

If card 26 and card 27 are not
specified, the following
standards are used:

MAZ(1) = 0, MAZ(2) = 0

Card 28

$\alpha B L \phi C \alpha$

14 constant

Card 29

NS

maximum number of groups in
one "macro" group.

NK

number of basic energypoints
used in that energy range in
a group, for which scattering
into neighbouring groups is
possible.

NR

Number of basic energy points
used in that energy range in
a group, for which no out-
scattering into neighbouring
groups is possible.

Comment:

If card 28 and card 29 are
not specified, the following
standards are used:

NS = 4, NK = 70, NR = 16

If module 9 is called, the following five cards are necessary:

Card 30

$\alpha_{BL} \phi C \alpha$ 15 constant

Card 31

ERR allowed integration error.

NJM $2^{NJM} + 1$ is the number angles for the integration.

NUJM $2^{NUJM} + 1$ is the maximum number of energy points for the integration.

Comment: If card 30 and card 31 are not specified, the following standards are used:

ERR = 0,05

NJM = 6

NUJM = 10

ERR, NJM, NUJM correspond to ϵ , NJM, NUJM in appendix IV

Card 32

$\alpha_{BL} \phi C \alpha$ 16 constant

Card 33

ISELR = 1: The transferelements from subintervals of the out-scattering group into other groups are interval averaged transfer cross sections, normalized to the interval averaged total elastic cross section.

ISELR = 0: The transfer cross sections are first normalized to the total elastic cross sections. The average of the outscattering interval is performed afterwards.

NLRA lowest Legendre moment

NLRE highest Legendre moment

NCALL = 0: The way of subdividing the groups, given by card 5, shall be the same for all groups. It is assumed, that the groups are subdivided into 14 fine groups and the fine groups are subdivided into five intervals. The subdivision is always performed equidistant in lethargy. The result of this procedure are 70 intervals per group.

NCALL > 0 : If the intervals shall be constructed in a different way as for NCALL = 0, NCALL is the number of calls of module 9. This number must be the same as specified in card 9.

If NCALL > 0, the next card must be specified NCALL-times.

Card 34

NGRE number of subdividing procedures in one call of module 9.

(N1(I),N2(I),NFG(I),NFI(I),I=1,NGRE)

N1(I) number of the first energy
 group subdivided by the I-th
 procedure

N2(I) number of the last energy
 group subdivided by the I-th
 procedure

NFG(I) number of fine groups per
 group in the I-th procedure

NFI(I) number of intervals per
 fine group in the I-th
 procedure.

Comment: If card 32 and card 33 are not
 specified, the following
 standards are used:

ISELR = 1 for ISTRUK = 0

ISELR = 0 for ISTRUK = 1

NLRA = 0

NLRE = 5

NCALL = 0

NFG = 14

NFI = 5

In this case, card 34 must be
omitted. In the case that
card 32, card 33 and card 34
are specified, it must be
noted, that $N1(I) \geq N2(I)$ and
 $N1(I) > N1(I+1)$, $N2(I) > N2(I+1)$.
(See also comment to card 9).

Card 35

α ENDE α 17 constant

Comment: This is the last card of the input for one material. It must not be omitted.

The cards 1 to 35 must be repeated for each material. The last card of the input must be

α ENDE α 18 constant

2.1.2 Description of the necessary external units and memory capacity

MIGRØS-2 needs the following external units:

unit number

8	for the input with FREEFØ.
1	Karlsruhe nuclear data library KEDAK.
3	unit for the unformatted output of the results.
10	is only necessary in the module 6 as working unit.

The necessary memory capacity must be specified by the user by the PARM.G-parameter in the EXEC-card of the IBM-job control language. The length of the PARM.G-array depends on the input information. The memory capacity, necessary for the single modules, is given in the following:

(The symbols of the input description are used).

Control module

$2 \times \text{NTYP} + 3 \times \text{NMAT} + \text{NFE} \times (\text{NSPEC} + 1) + \text{NT} + \text{MI} + 1 + 3 \times \text{NE} + 4 \times \text{NA} \times (1 + \text{NE})$

words. In the case, that one of these values are not specified by the input, the standard must be used: NTYP = 6, NFE = 1, NSPEC = 1, NT = 3, MI = 7, NMAT = 1.

The memory demand of that called computational module with the largest memory demand must be added to the memory demand of the control module.

Number of the
module

Memory demand of the module

1

$7 \times MI + 3 \times NE + 7 \times \max(300, \text{number of resonances on KEDAK})$
 $+ 5 \times \max(400, \text{maximum number of energy points, necessary for the integration over one energy group})$ -
words.

Comment: The number of integration points can not be estimated at the beginning of a calculation. If the assumed 400 words are not sufficient, the dimension of the corresponding array is enlarged automatically by the control module, provided that the PARM.G-array is large enough. If the PARM.G-array is not large enough, the module is passed. In this case it is recommended to duplicate the number of words for the integration array in the PARM.G-array.

2

$35 \times MI$ - words

3

$6 \times NE + 16 \times MI \times NE + 2 \times \max(1500, \text{number of MUEL-values on KEDAK in the wanted energy range})$
 $+ 6 \times \max(1500, \text{number of SGA-values, number of SGN-values, number of SGF-values in KEDAK in the wanted energy range})$ -
words.

- 4 $6 \times NE + 2 \times \max(1500, \text{number of SGN-values, number of SGF-values on KEDAK in the wanted energy range})$ -words
- 5 $12 \times NE + 3 \times (NE + 1) + 3 \times \max(700, \text{number of SGI-values on KEDAK above the threshold for inelastic scattering}) + 3 \times \max(25, \text{number of inelastic excitation levels}) + \max(25 \times 400, \text{number of inelastic excitation levels times the number of the energy points for the excitation cross section belonging to the first level}) + \max(400, \text{number of SGI-values above the threshold for inelastic scattering})$ -words.
- 6 $5 \times \max(400, ISM) + 7 \times \max(1000, ISD) + 4 \times (NE + 1) + 4 \times NE + \max(400 \times 21, ISM \times IC\emptyset S) + \max(400 \times 2, ISM \times NECU) + 6 \times \max(2, NECU) \times NE + 2 \times IC\emptyset S + 2 \times \max(4000, NTT)$ -words
- ISM: 300 + maximum number of energy points for the SGNC on KEDAK within four neighbouring energy groups.
- ISD: ISM + maximum number of SGN- or MUEL-values on KEDAK within one energy group.
- IC \emptyset S: number of angle points for the SGNC on KEDAK.

NECU: maximum energy loss by elastic scattering measured in energy groups (= 2 for scattering into the neighbouring group).

NTT: number of SGT-values for a given composition in four neighbouring energy groups.

7

NE-words

8

$3 \times \text{NE} + 2 \times \max(10 \times \text{NE}, \text{NFE})$ -words

9

$39 + \text{NM1} \times \text{IMAX} + (3 + \text{NUJM}) + \text{ICØS} \times (2 + \text{NIV}) + \text{NIV} + \text{IMAX} + 2 \times \text{NJM} + 3 \times \text{NDAT} + \text{NE} + \text{NFIN} \times \text{IZV} + \text{BUF}$ - words.

NM1: NLRE +1

NFIN: number of energy points in one energy group

NDAT: maximum number of SGT- or SGN-values on KEDAK in one energy group.

NIV: number of energy points for the SGNC on KEDAK

BUF: $\max(\text{NDAT}, 2 + \text{IMAX} \times \text{NFIN})$

IZV: $\max(4, \text{IMAX} \times \text{NM1})$

IMAX: maximum energy loss by elastic scattering measured in energy groups (identical with NECU in module 6).

ICØS: identical to module 6

10

-

The REGION-parameter on the JOB-card is given by the length of the PARM.G-array in K-bytes + 170 K-bytes for the MIGROS-programs and the buffer.

2.2 Description of the output

2.2.1 Printed output

Each computational module has an output of the results on paper. In table V an explanation of the symbols is given.

2.2.2 Unformatted output

Besides the printed output an unformatted output on unit 3 is provided. The arrangement of this output is described in table VI.

2.2.3 Output of error messages

In table VII a list of error messages and warnings is given. In this list only those messages are given which are not self explaining.

Table V

Module 1

Symbol	content	defined by formula	required KEDAK-types
SIGMA G	average group cross section for infinite dilution, radiative capture	(4.6)	} RES
SIGMA N	average group cross section for infinite dilution, elastic scattering	(4.6)	
SIGMA F	average group cross section for infinite dilution, fission	(4.6)	
SIGMAN1	identical with SIGMA N	-	
SIGMAT1	SIGM G + SIGMA N + SIGMA F	-	
SIGMA o	background cross section σ_o	(4.1)	
FG	flux weighted resonance self shielding factors } radiative capture	(4.1)	
FN			elastic scattering
FF			
FN1	current weighted resonance self shielding factors } elastic scattering	(4.3)	
FT1			total
Temperatures are given in $^{\circ}\text{K}$, the group boundaries in eV, cross sections in barns.			

Module 2

Symbol	content	defined by formula	required KEDAK-types
SIGMA G	average group cross sections for infinite dilution	radiative capture	}
SIGMA N		elastic scattering	
SIGMA F		fission	
SIGMAN1	identical with SIGMA N	-	ISØT1
SIGMAT1	SIGMA G + SIGMA N + SIGMA F	-	ISØT2
SIGMA o	background cross section σ_o	(5.1)	ST
FG	flux weighted resonance self shielding factors	radiative capture	STGF
FN		elastic scattering	
FF		fission	
FN1	current weighted resonance self shielding factors	elastic scattering	}
FT1		total	
Temperatures are given in $^{\circ}\text{K}$, group boundaries in eV, cross sections in barns			

Module 3

Symbol	content	defined by formula	required KEDAK-types
SIGMA A	average group cross sections for infinite dilution,	(9.10)	SGA
SIGMA C			
SIGMA N			
SIGMA F			
SIGMA No1	$\sigma_e \cdot \bar{\mu}_e$	(9.11)	SGN
SIGMA N1	identical to SIGMA No1	-	SGF
SIGMA T1	SIGMA C + SIGMA N + SIGMA F	-	SGT
FA	flux weighted resonance self shielding factors,	(9.1)	MUEL
FC			
FN			
FF			
FNo1	flux weighted resonance self shielding factor for $\sigma_e \cdot \bar{\mu}_e$	(9.3)	
FN1	current weighted resonance self shielding factor for $\sigma_e \cdot \bar{\mu}_e$	(9.8)	
FT1	current weighted total resonance self shielding factor	(9.5)	
SIGO	background cross section σ_0		
Group boundaries are given in eV, cross sections in barns			

Module 4

Symbol	content	defined by formula	required KEDAK-types
SGN	average group cross section for	(3.1)	SGN
SGA			SGA
SGF			SGF
SGI			SGI
SGC			SGA, SGF
SG2N			SG2N
MUEL	average cosine of elastic scattering	(3.2)	SGN, MUEL
NUE	average number of secondaries per fission	(3.3)	SGF, NUE
ALPHA	average α -value	(3.4)	SGG, SGF
ETA	average η -value	(3.5)	NUE,SGG,SGF

If group averages of other types of cross sections available on KEDAK are calculated, averages in the sense of formula (3.1) are calculated. The names are the same as the names of the underlying KEDAK-types. Cross sections are given in barns.

If SGN for the material Hbbbb01 is calculated, also a type STR is calculated, which corresponds to formula (3.6).

Module 5

Symbol	content	defined by formula	underlying KEDAK-type
PRØBSGI	inelastic scattering probabilities, in the discrete region	(6.6), (6.7)	SGI, ISØT2 SGI,SGIZ,ISOT2
	inelastic scattering probabilities, in the continuous range	(6.15)	
In the first column the numbers of the outscattering groups are printed. The numbers above the values of PRØBSGI are the inscattering groups.			

Module 6

Symbol	content	defined by formula	underlying KEDAK-type
SGNC0	zero'th moment		SGN, SGNC, MUEL, SGT
SGNC1	1 st moment		
SGNC2	2 nd moment		
SGNC3	3 th moment		
SGNC4	4 th moment		
SGNC5	5 th moment		
		of elastic scattering matrix, normalized to the total elastic cross section	

Module 7

Symbol	content	defined by formula	required KEDAK-types
CHI	fission spectrum	(12.2),(12.1)	CHICR

Module 8

Symbol	content	defined by formula	required KEDAK-types
1/V	1/v-group averages in [sec/cm]	(11.1)	-

Module 9

Symbol	content	defined by formula	required KEDAK-types
SGT	average total cross section	(8.1)	SGT
SGN	average elastic scattering cross section	(8.1)	SGN
MUEL	average cosine of elastic scattering	(8.2)	MUEL
FLUX	average flux (weighting function)	-	-
SGNCn	the n-th legendre moment for elastic scattering	(8.3)	SGN, SGNC

Cross sections are given in barns, the legendre moments are normalized to the total elastic cross section. The interval averages within one group are arranged with increasing energy.

Module 10

In this module the group constants for the thermal group are calculated. The same types as in module 4 are possible. The group constants are defined by the formulas (10.1) to (10.5).

Table VI (The symbols of table V are used)

Module 1

- 1. record: o , 'MIGR----'
- 2. record: n , material name, temperature in [°K], number of the energy group, lower group boundary in [eV], upper group boundary in [eV]
- 3. record: n , SIGMA G, SIGMA N, SIGMA F, SIGMA N1, SIGMAT1
- 4. record: n , SIGMA O, FG, FN, FF, FN1, FT1
- . . .
- . . .
- . . .
- . . . for all SIGMA O-values

n is an Integer $I \times 4$ and gives the number of 4-bytes words in the record. The material names are 8-bytes alphanumerical words. All other words are $R \times 4$. The 8-bytes alphanumerical words are counted as two words in n. These conventions are the same for all other modules. Exceptions are mentioned. "-" in the label names must be interpreted as blank.

Module 2

- 1. record: o , 'FSTAT---'
- 2. record: n , material name, temperature in [°K], number of the energy group, lower group boundary in [eV], upper group boundary in [eV].
- 3. record: n , SIGMA G, SIGMA N, SIGMA F, SIGMAN1, SIGMAT1
- 4. record: n , SIGMA O, FG, FN, FF, FN1, FT1
- . . .
- . . .
- . . . for all SIGMA O-values

Module 3

1. record: o , 'FSTRK---'
2. record: n , material name, number of the energy group, lower group boundary in [eV], upper group boundary in [eV]
3. record: n , SIGMA A, SIGMA N, SIGMA NO1, SIGMA N1, SIGMA T1
for non heavy isotopes (ISTRUK = 0)
n , SIGMA C, SIGMA N, SIGMA F, SIGMA NO1, SIGMA N1,
SIGMA T1
for heavy isotopes (ISTRUK = 1)
4. record: n , SIGO, FA, FN, FNO1, FN1, FT1
. for non heavy isotopes (ISTRUK = 0)
. n , SIGO, FC, FN, FF, FNO1, FN1, FT1
. for heavy isotopes (ISTRUK = 1)
repeated for all SIGO-values

Module 4

1. record: o , 'SGKE-----'
2. record: n , highest energy group (lowest number), lowest energy group
3. record: n , material name, name of reaction type (8-byte, alphanumeric)
4. record: n , group constant of the type defined by the 3rd word in the 3.rd record, for all energy groups, specified by the 2nd record. The values are ordered with increasing energy and decreasing group numbers.

The records 2 to 4 are repeated for all reaction types, specified by the input.

Module 5

1. record: o , 'SMTØT---'
2. record: n , material name, total number of outscattering groups
3. record: n , number of the outscattering group, elements of
. the matrix PRØBSGI in the sense that the first
. element describes scattering within the group,
. then scattering into the neighbouring group etc.

repeated for all outscattering groups.

Module 6

1. record: o , 'FLUM----'
2. record: n , material name, number of outscattering groups,
number of Legendre moments.
3. record: n , number of the outscattering groups, group
averaged total elastic cross section, group
averaged cosine for elastic scattering.

repeated for all outscattering groups.

4. record: n , number of the Legendre moment, number of the
outscattering group, matrix elements SGNCi
(i = number of the moment) arranged in the
sense that the first word describes scattering
within the group, the next one scattering
into the neighbouring group etc.

repeated for all outscattering groups (for a
certain Legendre moment)

repeated for all Legendre moments.

Module 7

1. record: o , 'SPALT---'
2. record: n , material name, energy of the fission inducing neutron in [eV], number of the lowest energy group, number of the highest energy group.
3. record: n , CHI for all groups specified in the 2nd record, arranged with increasing energy.

Module 8

1. record: o , 'S1/V----'
2. record: n , 1/V-values for all energy groups defined by the group boundaries in the input, arranged with increasing energy and decreasing group numbers.

Module 9

1. record: o , 'REMØ----'
2. record: n , material name, total number of outscattering groups, total number of Legendre moments.
3. record: n , number of the outscattering group, number of subgroups in the outscattering group, number of intervals in one subgroup (the same for all subgroups in one outscattering group), SGT for all intervals in one group (product of the number of subgroups and the number of intervals in one subgroup) arranged with increasing energy, SGN for all intervals in one group, MUEL for all intervals in one group, FLUX for all intervals in one group.
repeated for all outscattering groups.
4. record: n , number of the Legendre moment, number of the outscattering group, SGNCi (i = Legendre moment) for all intervals in one group arranged in the sense

that first the elements describing scattering within the group, then the elements describing scattering into the neighbouring group etc. are stored.

repeated for all outscattering groups (for a fixed Legendre moment)

repeated for all Legendre moments.

Module 10

1. record: o , 'THERM---'
2. record: n , material name, name of the reaction type, number of the thermal group.
3. record: n , value of the group constant specified in the 2nd record

Table VII Error messages and warnings

*** WARNING 2.04

Comment: To get a correct overlapping correction, Γ/Δ should be much smaller than one. Normally it is assumed, that $\Gamma/\Delta < 0.5$ is small enough. See chapter 5, formula (5.24).

*** WARNING 2.05

Comment: For the calculation of the current weighted resonance self shielding factors an approximation is used, which in some cases is not valid for small σ_0 -values. In this cases, the self shielding factors are not calculated and negative figures are printed. See chapter 5, and Appendix III, formula (25).

*** WARNING 2.06

Comment: The reason for this warning might be the same as in *** WARNING 2.04. If not, the formalisme for approximating the overlapping correction is insufficient.

The warning is given, when the denominator in formula (5.14) becomes negative. If the resulting $^s\sigma_{r,g}$ in (5.14) is small against the contribution of all other resonance series, this warning may be ignored.

References

/1/ B. Krieg

Handling and Service Program for the Karlsruhe Nuclear Data File
KEDAK.

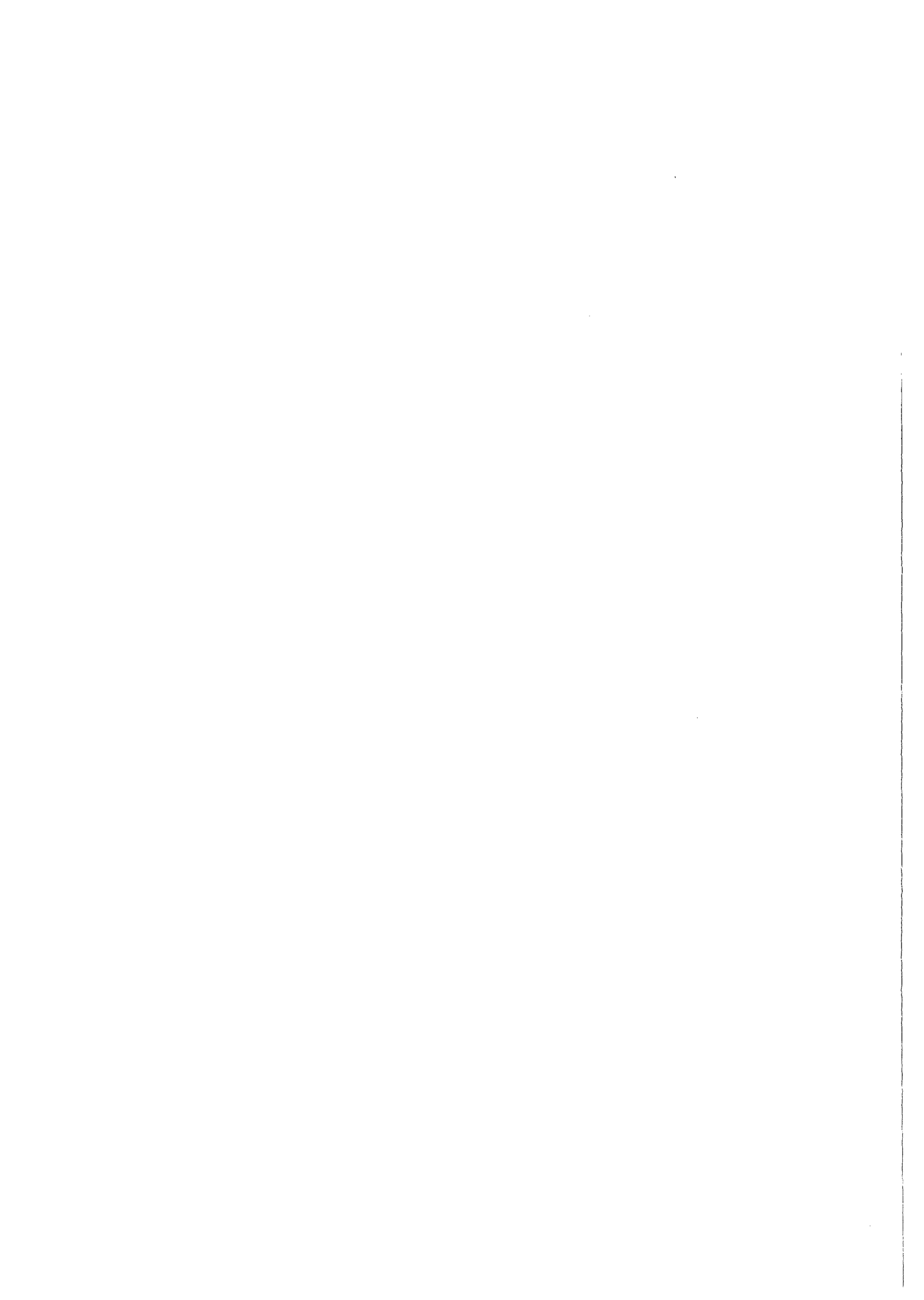
Part I: Management and Retrieval Programs

KFK 1725, 1972

/2/ L.P. Abayjan, N.O. Bazazjanc, I.I. Bondarenko, M.N. Nikolaev

Gruppenkonstanten schneller und intermediärer Neutronen für
die Berechnung von Kernreaktoren

KFK-tr-144



3. The calculation of average group cross sections from tabulated neutron cross sections. Module 4.

3.1 The average group cross sections are defined by

$${}^k\tilde{\sigma}_{x,g} = \frac{\int_{(g)} {}^k\sigma_x(E) \phi(E) dE}{\int_{(g)} \phi(E) dE} \quad (3.1)$$

k = isotope

x = neutron reaction (n,x)

g = energy group

$\phi(E)$ is the neutron flux density. In the resonance region, normally a slowly varying collision density is used, so that formula (3.1) defines the average group cross section for infinite dilution. The average group cross sections for all types of neutron reactions are calculated by formula (3.1), except:

- the average cosine of the elastic scattering

$${}^k\bar{\mu}_{e,g} = \frac{\int_{(g)} {}^k\bar{\mu}_e(E) {}^k\tilde{\sigma}_e(E) \phi(E) dE}{\int_{(g)} {}^k\tilde{\sigma}_e(E) \phi(E) dE} \quad (3.2)$$

- the average number of secondaries per fission

$${}^k\nu_g = \frac{\int_{(g)} {}^k\nu(E) {}^k\tilde{\sigma}_f(E) \phi(E) dE}{\int_{(g)} {}^k\tilde{\sigma}_f(E) \phi(E) dE} \quad (3.3)$$

f = fission

- the average α -value

$$\alpha_g = \frac{\int_{(g)} k \tilde{\sigma}_f(E) \phi(E) dE}{\int_{(g)} k \tilde{\sigma}_f(E) \phi(E) dE} \quad (3.4)$$

γ = radiative capture

- the average η -value

$$\eta_g = \frac{\nu'_g}{1 + \alpha_g} \quad (3.5)$$

For which type of neutron reaction an average cross-section shall be calculated can be chosen by input. There must be information available on KEDAK-library.

Besides that, the following group average

$$\tilde{\sigma}_{beH} = \frac{\int_{(g)} \frac{\sigma_e(E)}{E} F(E) dE}{\int_{(g)} F(E) dE} \quad (3.6)$$

is calculated for hydrogen. This group average can be used for the production of the elastic scattering matrix.

3.2 The integration in the averages (3.1) to (3.4) are done by a trapezoidal rule. As integration points the energy points at which the nuclear data on the KEDAK-library are tabulated and the group boundaries are used. The values of the cross sections at the boundaries are calculated by linear interpolation or extrapolation from the tables on KEDAK. The weighting function is either interpolated linearly from tables or is given by a function.

3.3 The following subroutines is necessary:

```
SUBROUTINE SUND (MM, ENG, NFE, REFE, EFE, ITYP, ITNAM, SGC, DUE,  
                XINTE, ZINT, XNEN, STREU, LDIM, LDIMP, SE, FSE)
```

The following parameters are defined by the control program :

MM : Number of energy group boundaries.

ENG : one-dimensional field containing the energy group boundaries in [eV].

NFE : number of points of the neutron flux density.
= 0, if the function PHI(E) is used .

REFE : one-dimensional field containing the energy points of the neutron flux density in [eV].

EFE : one-dimensional field containing the neutron flux density.

ITYP : number of cross section types .

ITNAM : one-dimensional field containing the names of these reaction types for which the average group cross sections shall be calculated .

LDIM : dimension of the following fields SE and FSE which is $\max (1500, \text{number of } \sigma_e(E)\text{-values on KEDAK, number of } \sigma_f(E)\text{-values on KEDAK in all regarded energy groups})$.

The following parameters are work fields used by the subroutine SUND :

SGC, DUE, XINTE, ZINT, SNEN, STREU one-dimensional fields of the
length MM

SE, FSE one-dimensional fields of the
length LDIM .

The following parameter is calculated in the subroutine

LDIMP: 0 , if the length of the working fields dimensioned with
LDIM is sufficient.

: n , number, by which LDIM should be increased.

4. The calculation of average group cross sections for infinite dilution and of energy resonance self shielding factors from resolved resonance parameters. Module 1.

4.1 Flux-weighted energy resonance self shielding factors are calculated for capture, fission and elastic scattering. They are defined as

$${}^k f_{x,g}(\tilde{\sigma}_o, T) = \frac{{}^k \tilde{\sigma}_{x,g}(\tilde{\sigma}_o, T)}{{}^k \tilde{\sigma}_{x,g}(\tilde{\sigma}_o \rightarrow \infty, T)} \quad (4.1)$$

- k = isotope
- x = neutron reaction (n,x)
- g = energy group
- T = temperature in °K
- σ_o = background cross section in barns

where

$${}^k \tilde{\sigma}_{x,g}(\tilde{\sigma}_o, T) = \frac{\left\langle \frac{{}^k \tilde{\sigma}_x(T, E) F(E)}{{}^k \tilde{\sigma}_t(T, E) + \tilde{\sigma}_o} \right\rangle}{\left\langle \frac{F(E)}{{}^k \tilde{\sigma}_t(T, E) + \tilde{\sigma}_o} \right\rangle} \quad (4.2)$$

is the microscopic effective cross section.

The weighting function $F(E)$ is the collision density and is assumed to have no resonance structure in the NR-approximation.

The average is defined as

$$\langle y(E) \rangle = \frac{1}{\Delta E_g} \int_{\Delta E_g} y(E) dE$$

where ΔE_g is the group width of group g .

The current-weighted resonance self shielding factors are calculated for elastic scattering and for the total neutron reaction

$${}^k f_{x,g}(\sigma_o, T) = \frac{{}^k \tilde{\sigma}_{x,g}(\sigma_o, T)}{{}^k \tilde{\sigma}_{x,g}(\sigma_o \rightarrow \infty, T)} \quad (4.3)$$

where

$${}^k \tilde{\sigma}_{x,g}(\sigma_o, T) = \frac{\left\langle \frac{{}^k \tilde{\sigma}_x(T, E) F(E)}{\{ {}^k \tilde{\sigma}_t(T, E) + \sigma_o \}^2} \right\rangle}{\left\langle \frac{F(E)}{\{ {}^k \tilde{\sigma}_t(T, E) + \sigma_o \}^2} \right\rangle} \quad (4.4)$$

is the current weighted microscopic effective cross section. For infinite dilution

$$\lim_{\sigma_o \rightarrow \infty} {}^k \tilde{\sigma}_{x,g}(\sigma_o, T) = \lim_{\sigma_o \rightarrow \infty} {}^k \tilde{\sigma}_{x,g}(\sigma_o, T) = {}^k \tilde{\sigma}_{x,g}^\infty \quad (4.5)$$

and

$$k_{x,g}^{\infty}(T) = \frac{\langle k_{\tilde{\sigma}_x}(E, T) F(E) \rangle}{\langle F(E) \rangle} \quad (4.6)$$

Therefore the denominators of (4.1) and (4.3) may be replaced by (4.6).

Normally $k_{\sigma_{x,g}}^{\infty}$ is independent of temperature, but in cases when only few resonances are within an energy group, and the group boundary cuts the wings of an important resonance, $k_{\sigma_{x,g}}^{\infty}$ may be temperature dependent.

The averages of (4.3), (4.4) and (4.6) are calculated by numerical integration. The microscopic cross sections $\sigma_x(E)^*$ are calculated from resonance parameters by a Breit-Wigner single level formula.

$$\tilde{\sigma}_x(E, T) = \sum_{r=1}^R \tilde{\sigma}_x(E, E_r, T) \quad (4.7)$$

r = resonance

R = number of resonances taken into account

E_r = resonance energy

$$\tilde{\sigma}_x(E, E_r, T) = \tilde{\sigma}_{o_x} \cdot \Upsilon_r(\theta, X) \quad (4.8a)$$

for capture and fission,

* (the index for the isotope will be neglected)

$$\sigma_t(E, E_r, T) = \sigma_p + \sigma_{oc} \cdot \{ \Psi_r \cdot \cos 2\delta_L + \chi_r \cdot \sin 2\delta_L \} \quad (4.8b)$$

and for elastic scattering

$$\sigma_e(E, E_r, T) = \sigma_t(E, E_r, T) - \sum_x \sigma_x(E, E_r, T) \quad (4.8c)$$

x = capture, fission, where

$$\sigma_{ox} = 4\pi \lambda^2(E_r) g \cdot \frac{\Gamma_n(E_r) \Gamma_x(E_r)}{\Gamma^2(E_r)} \cdot \left(\frac{E}{E_r} \right)^{l-1/2} \quad (4.9)$$

$$\sigma_{oc} = 4\pi \lambda^2(E_r) \cdot g \cdot \frac{\Gamma_n(E_r)}{\Gamma(E_r)} \cdot \left(\frac{E}{E_r} \right)^{l-1/2} \quad (4.10)$$

$$\Psi(\theta, x) = \frac{\theta}{2\sqrt{\pi}} \int_{-\infty}^{+\infty} \frac{\exp\left\{-\frac{\theta^2}{4} \cdot (x-y)^2\right\}}{1+y^2} \quad (4.11)$$

$$\chi(\theta, x) = \frac{\theta}{2\sqrt{\pi}} \int_{-\infty}^{+\infty} \frac{y \cdot \exp\left\{-\frac{\theta^2}{4} \cdot (x-y)^2\right\}}{1+y^2} dy \quad (4.12)$$

$$\theta = \frac{\Gamma}{\Delta}, \quad x = \frac{2(E-E_r)}{\Gamma}, \quad y = \frac{2(E'-E_r)}{\Gamma} \quad (4.13)$$

$$\Delta = \frac{4kT E_r}{A} = \text{Doppler width}$$

Γ = total width

$$g = \frac{2I+1}{2(2i+1)}, \quad I = \text{total spin, } i = \text{spin of the nucleus}$$

$$\delta_1 = \frac{R'}{\lambda} - \arctg \frac{R'}{\lambda} \cdot 1 \quad \text{for } l = 0, 1$$

R' = effective radius of the nucleus

λ = reduced neutron length

Γ_n = neutron width

Γ_x = fission or capture width

σ_p = potential cross section

For $T \rightarrow 0$, that means $\theta \rightarrow \infty$, and

$$\psi(\theta \rightarrow \infty, x) = \frac{1}{1+x^2} \quad (4.14)$$

$$\chi(\theta \rightarrow \infty, x) = \frac{x}{1+x^2} \quad (4.15)$$

For resonances at negative resonance energies the formulas (4.9) and (4.10) are modified to

$$r\delta_{ox} = 4\pi\Lambda_0^2 \cdot g \cdot \frac{\Gamma_n(E_r) \Gamma_x(E_r)}{\Gamma^2(E_r)} \cdot E^{l-1/2} \quad (4.16)$$

$$r\delta_{oc} = 4\pi\Lambda_0^2 \cdot g \cdot \frac{\Gamma_n(E_r)}{\Gamma(E_r)} \cdot E^{l-1/2} \quad (4.17)$$

where $\Lambda_0 = \lambda (1 \text{ eV})$.

4.2 The method of integration for the averages in (4.2), (4.4) and (4.6).

The expressions (4.2), (4.4) and (4.6) are calculated by numerical integration. A trapezoidal rule is used. To keep the number of integration points as small as possible, which is necessary, because the calculation of the function values is very time consuming, the integration points are not taken equidistant. To be sure that the resonances are not lost during the integration,

the resonance energies are chosen as integration points. Let X_i be the integration points and Y_i the corresponding function values. Then the integration algorithm can be described in the following way:

First step: $X^0 = (X_1^0, X_2^0, \dots, X_n^0) = (E_{g+1}, E_{r,i} (i=1, n-2), E_g)$

$$Y^0 = (Y_1^0, Y_2^0, \dots, Y_n^0)$$

E_{g+1} lower group limit

E_g upper group limit

$E_{r,i}$ resonance energy of the i -th resonance in energy group g .

Second step: $X^1 = (X_1^0, \frac{X_1^0 + X_2^0}{2}, X_2^0, \dots, X_n^0) = (X_1^1, X_2^1, \dots, X_{n+1}^1)$

$$Y^1 = (Y_1^1, \dots, Y_{n+1}^1)$$

$$a = \frac{1}{2} (X_2^1 - X_1^1) \cdot (Y_1^1 + Y_2^1)$$

$$b = \frac{1}{2} (X_2^1 - X_1^1) \cdot (Y_1^1 + \tilde{Y}_2^1)$$

$$\tilde{Y}_2^1 = \frac{1}{2} (Y_1^1 + Y_3^1)$$

The condition

$$\frac{b-a}{a} < \varepsilon \tag{4.18}$$

where ε is a given error limit, must be satisfied.

If condition (4.18) is not fulfilled, the second step is repeated until condition (4.18) is fulfilled, let's say, after repeating the second step $(l-1)$ times.

$$X^1 = (X_1^{1-1}, \frac{X_1^{1-1} + X_2^{1-1}}{2}, X_2^{1-1}, \dots, X_{n+1-1}^{1-1}) = (X_1^1, \dots, X_{n+1}^1)$$

$$Y^1 = (Y_1^1, \dots, Y_{n+1}^1)$$

Third step: $a = \frac{1}{2} (X_3^1 - X_2^1) \cdot (Y_2^1 + Y_3^1)$

$$b = \frac{1}{2} (X_3^1 - X_2^1) \cdot (\tilde{Y}_2^1 + Y_3^1)$$

$$\tilde{Y}_2^1 = \frac{1}{2} (Y_1^1 + Y_3^1)$$

The condition

$$\frac{b-a}{a} < \varepsilon \quad \text{must be satisfied,} \quad (4.19)$$

If the condition is fulfilled, the integral over the interval $[X_1^1, X_3^1]$ is given as

$$I_1 = \frac{1}{2} (X_3^1 - X_1^1) \cdot (Y_1^1 + Y_3^1)$$

and the algorithm is repeated in the interval $[X_3^1, X_{n+1}^1]$ starting with the integration points

$$X^0 = (X_3^1, \dots, X_{n+1}^1) = (X_1^0, X_2^0, \dots, X_{n+1-2}^0).$$

If condition (4.19) is not fulfilled, the integral in the interval $[X_1^1, X_2^1]$ is calculated as

$$I_1 = \frac{1}{2} (X_2^1 - X_1^1) \cdot (Y_1^1 + Y_2^1)$$

and the algorithm is repeated in the interval $[X_2^1, X_{n+1}^1]$ starting with the integration points

$$X^0 = (X_2^1, \dots, X_{n+1}^1) = (X_1^0, X_2^0, \dots, X_{n+1-1}^0)$$

The algorithm is repeated until the integration over the total energy width is performed. The integral over the energy group is given by the sum.

$$I = \sum_{i=1}^{N+1} I_i, \text{ where } N \text{ is the number of repetitions of the algorithm.}$$

4.3 For the calculation of energy resonance self shielding factors and average group constants from resolved resonance parameters the following subroutine are used:

FGEM (NS, SIGO, NE, ENG, NEF, ES, F, NT, TEMP, PR, NMR, NFST, SUM, SUO, IRE, IREP, ER, L, GJ, GAT, GAN, GAG, GAF, ISTE, ISTEP, STE)

In this subroutine the input of the nuclear data from the Karlsruhe nuclear data library KEDAK and the output of the f-factors and the average cross sections is organized. Also the integration algorithm is performed in this subroutine. The following parameters must be defined:

NS : number of σ_0 -values.
SIGO(NS) : one-dimensional field containing the σ_0 -values [barn].
NE : number of the group boundaries.
ENG(NE) : one-dimensional field containing the group boundaries [eV].
NEF : number of energy points of the weighting spectrum.
ES(NEF) : one-dimensional field containing the energy points of the weighting spectrum [eV].
F(NEF) : one-dimensional field containing the weighting function at the energy points of field ES.

NT : number of temperatures.
TEMP(NT) : one-dimensional field containing the temperatures
[°K].
PR : error limit ϵ , as defined in (4.18).
NMR : R/2 as defined in (4.7).
IRE : length of working fields, should be longer than
the number of resonances for one isotope.
ISTE : length of a working field for integration, should
be longer than the number of integration points
necessary for the integration of effective group
cross sections within one group.

Working fields:

SUM(NS,7), SUO (NE,3), ER (IRE), GJ (IRE), GAT (IRE), GAN (IRE),
GAG (IRE), GAF (IRE), STE (5, ISTE).

The following parameters are calculated in the subroutine:

NFST : number of the first energy group, for which f-factors
can be calculated only from statistical resonance
parameters. This parameter is important only, if the
f-factors are calculated for an energy range,
covering both, the range of resolved and the range
of unresolved resonances.
IREP : 0, if the length of the working fields dimensioned
with IRE is sufficient.
: n number, by which IRE should be increased.
ISTEP : 0, if the length of the working field dimensioned
with ISTE is sufficient.
: n number, by which ISTE should be increased.

Information is also transferred by the unlabeled COMMON. It is
described in chapter 1.

A description of the output is given in chapter 1.

WIRQ (INR, E, SF, SG, ST, IRE, ER, GJ, GAN, GAT, GAF, L, GAG, T, A, R,
RLA)

In this subroutine the contribution of one resonance to the energy dependent cross section at a given energy is calculated by a Breit-Wigner single level formula.

The following parameters must be defined:

INR : number of the resonance (all resonances of an isotope are numbered with increasing resonance energies).

E : energy, [eV].

ER(IRE) : one-dimensional field containing the resonance energies for one isotope, [eV].

IRE : must be greater or equal to the number of resonances for one isotope.

GJ(IRE) : one dimensional field containing the statistical parameters g_j for all resonances of one isotope.

GAN(IRE) : one dimensional field containing the neutron half widths for all resonances of one isotope, [eV].

GAT(IRE) : one dimensional field containing the total half widths for all resonances of one isotope, [eV].

GAF(IRE) : one dimensional field containing the fission widths for all resonances of one isotope, [eV].

L(IRE) : one dimensional field containing the neutron angular momentum for all resonances of one isotope.

GAG(IRE) : one dimensional field containing the capture widths for all resonances of one isotope, [eV].

T : temperature in $^{\circ}\text{K}$.

A : atomic weight.

R : radius of the nucleus, [$\sqrt{\text{barn}}$].

RLA : reduced neutron wave length, [$\sqrt{\text{barn} \cdot \text{eV}}$].

The following parameters are calculated in the subroutine:

SF : contribution to the fission cross section at the energy E from the resonance with number INR, [barn].
SG : contribution to the capture cross sections at the energy E from the resonance with number INR [barn].
ST : contribution to the total resonance cross section at the energy E from the resonance with number INR, [barn].

WIRQU (NR1, E, SIFG, SIGG, SIGT, NR, IRE, ER, GJ, GAN, GAT, GAF, L, GAG, T, A, R, RLA, NMIN)

In this subroutine the sum of the contribution to the cross section at the energy E from all resonances, that are taken into account, is performed.

The following parameters must be defined:

NR1 : the number of the resonance belonging to the first resonance energy that is greater than the higher energy limit of the actual energy group. All resonances of an isotope are numbered with increasing resonance energies).
E : energy [eV].
NR : maximum number of resonances for one isotope.
IRE : must be greater or equal to NR.
ER(IRE) : one-dimensional field containing the resonance energies for one isotope.
GJ(IRE) : one-dimensional field containing the statistical parameters g_j for all resonance of one isotope.
GAN(IRE) : one-dimensional field containing the neutron half widths for all resonance of one isotope [eV].
GAT(IRE) : one-dimensional field containing the total half widths for all resonances of one isotope. [eV].
GAF(IRE) : one-dimensional field containing the fission half widths for all resonances of one isotope. [eV].

L(IRE) : one-dimensional field containing the neutron angular momentum for resonances of one isotope.
GAG(IRE) : one dimensional field containing the capture widths for all resonances of one isotope. [eV].
T : temperature in °K.
A : atomic weight.
R : radius of the nucleus in [√barn]
RLA : reduced neutron wave length, [√barn · √eV].
NMIN : the number of resonances at higher and at lower resonance energies than the energy E, that are taken into account for the calculation of the cross sections at the energy E.

The following parameters are calculated in the subroutine:

SIFG : fission cross section at the energy E, [barn].
SIGG : capture cross section at the energy E, [barn].
SIGT : total cross section at the energy E, [barn].

STOSS (E1, E2, E3, F1, F2, F3, NFE, ES, F)

In this subroutine the macroscopic weighting function (collision density) at the energies E1, E2, E3 is calculated either from an energy point wise given weighting function by interpolation or from a function PHI(E). As standard PHI(E) = 1/E is used.

The following parameters must be defined:

E1, E2, E3 : energies in [eV].
NFE : number of energy points of the weighting spectrum.
If NFE = 0,1, the standard PHI(E) = 1/E is used.
ES(NFE) : one-dimensional field containing the energy points of the weighting spectrum [eV].
F(NFE) : one-dimensional field containing the weighting spectrum.

The following parameters are calculated in this subroutine:

F1, F2, F3 : weighting spectrum at the energies E1, E2, E3.

PSIXI (X, T, U, V) /1/, /2/.

In this subroutine the functions (4.11) and (4.12) are calculated.

The following parameters must be defined:

X : defined as (4.13)

T : is defined as $\left\{\frac{1}{\theta}\right\}^2$, where θ is given in (4.13).

The following parameters are calculated in the subroutine:

U : defined by (4.11)

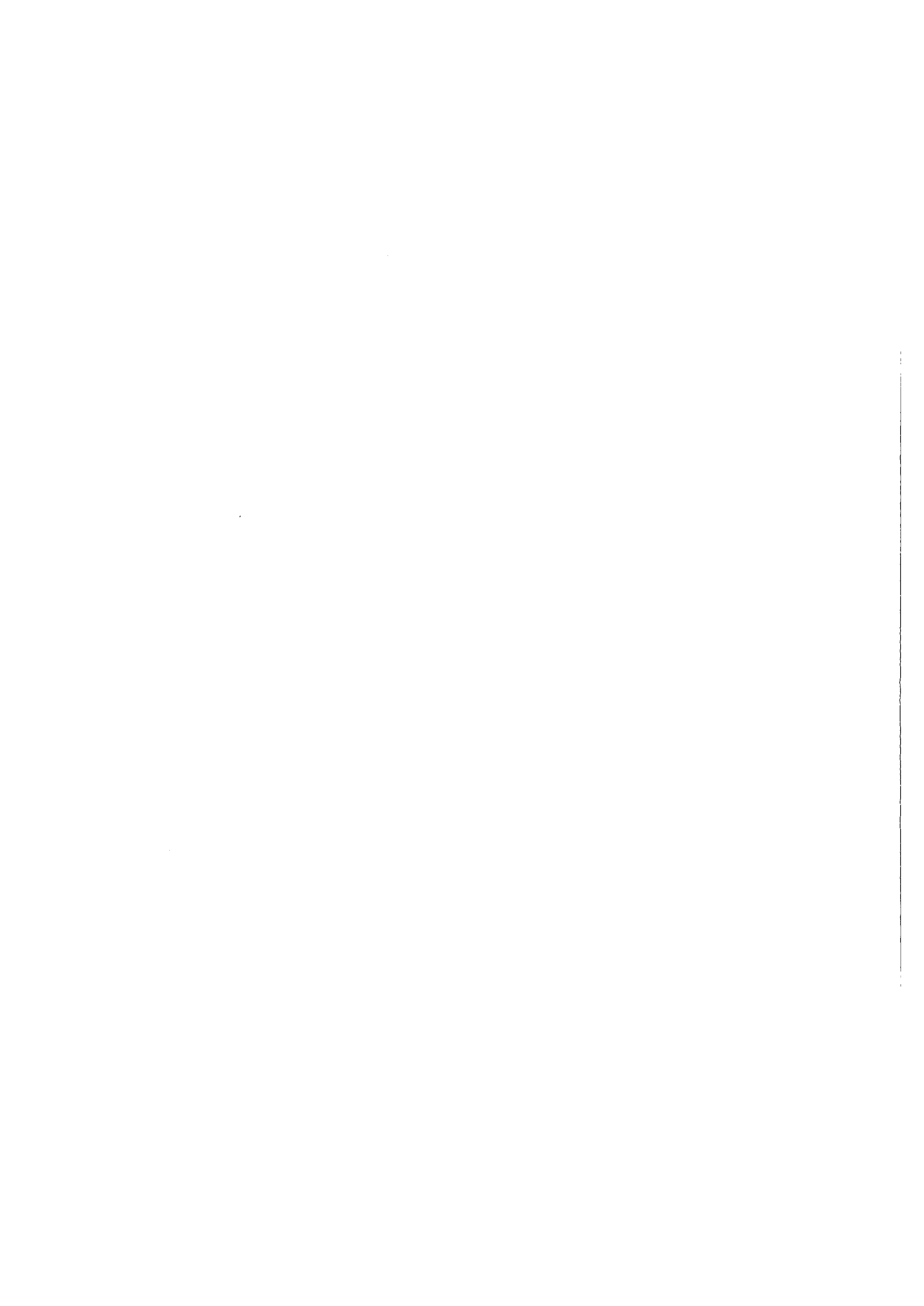
V : defined by (4.12)

References

/1/ H. Späth, INR-Arbeitsbericht, private Communication.

/2/ C. Chiarella, A. Reichel

On the Evaluation of Integrals Related to the Error Function
Math. of Computation 22, 1968, p. 137 - 143.



5. The calculation of average group cross sections for infinite dilution and of energy resonance self shielding factors from statistical resonance parameters. Module 2.

5.1 The flux weighted resonance self shielding factors are calculated for capture, fission and elastic scattering. They are defined as

$$k_{f,x,g}^f(\tilde{\sigma}_0, T) = \frac{k_{\tilde{\sigma}_{x,g}}^k(\tilde{\sigma}_0, T)}{k_{\tilde{\sigma}_{x,g}}^k(\tilde{\sigma}_0 \rightarrow \infty, T)} \quad (5.1)$$

$k_{\sigma_{x,g}}^k$ = microscopic effective group cross section

k = isotope

x = neutron reaction (n,x)

g = energy group

T = temperatur in °K

σ_0 = background cross section in barns.

The following approximation is used:

$$k_{\tilde{\sigma}_{x,g}}^k(\tilde{\sigma}_0 \rightarrow \infty, T) \cong k_{\tilde{\sigma}_{x,g}}^k(\tilde{\sigma}_{0,max}, T) \quad (5.2)$$

with

$$\tilde{\sigma}_{0,max} = \max(10^6, \text{highest } \sigma_0 \text{ value given by the input}) \quad (5.3)$$

The microscopic effective group cross section may be represented by

$$\bar{\sigma}_{x,g}(\bar{\sigma}_o, T) = \frac{\sum_j \phi_j \cdot \bar{\sigma}_{x,g}(\bar{\sigma}_o, T, E_j) \cdot \Delta E_j}{\sum_j \phi_j \cdot \Delta E_j} \quad (5.4)$$

ϕ_j = is a slowly varying, resonance free flux density at the energy E_j , normally approximated by a collision density

ΔE_j = energy interval around E_j

$\sigma_{x,g}(E_j)$ = is the effective cross section in the energy interval ΔE_j around E_j for the temperature T and the background cross section σ_o .

The $\sigma_{x,g}(E_j)$ are slowly varying with energy when calculated from statistical resonance parameters. The group average of (5.4) is therefore performed in the following approximate way

$$\bar{\sigma}_{x,g}(\bar{\sigma}_o, T) \simeq \frac{\sum_j \phi_j \cdot \bar{\sigma}_{x,g}(E_j)}{\sum_j \phi_j} \quad (5.5)$$

where for the E_j are chosen the upper and the lower group boundary of group g , E_g and E_{g+1} , and $(E_g + E_{g+1})/2$.

$\sigma_{x,g}(E_j)$ is composed of the contributions of all resonance series s .

$$\tilde{\sigma}_{x,g}(E_j) = \sum_s^s \tilde{\sigma}_{x,g}^s(E_j) \quad (5.6)$$

for capture and fission, and

$$\tilde{\sigma}_{e,g}(E_j) = \sum_s^s \left\{ \tilde{\sigma}_{r,g}^s(E_j) - \tilde{\sigma}_{c,g}^s(E_j) - \tilde{\sigma}_{f,g}^s(E_j) \right\} + \tilde{\sigma}_p \quad (5.7)$$

$\sum_s^s \sigma_{r,g}(E_j)$ is the total resonance cross section, (c = capture, f = fission).

The current weighted resonance self shielding factors are calculated for the total neutron reaction and for elastic scattering. They are defined as :

$$k_{rx}^k(\tilde{\sigma}_o, T) = \frac{k_{rx,g}^k(\tilde{\sigma}_o, T)}{k_{rx,g}^k(\tilde{\sigma}_o \rightarrow \infty, T)} \quad (5.8)$$

where

$$k_{rx,g}^k(\tilde{\sigma}_o \rightarrow \infty, T) \cong k_{rx,g}^k(\tilde{\sigma}_{o,max}, T) \quad (5.9)$$

$\sigma_{o,max}$ is given by (5.3).

The microscopic effective cross section for the neutron reaction (n,x) in the energy group g may be written as :

$${}_1\tilde{\sigma}_{x,g} = \frac{\sum_j {}_1\phi_j \cdot {}_1\tilde{\sigma}_{x,g}(E_j) \cdot \Delta E_j}{\sum_j {}_1\phi_j \cdot \Delta E_j} \quad (5.10)$$

${}_1\phi_j$ is a slowly varying, resonance free current density, normally approximated by a collision density.

For (5.10) the same approximation for averaging as for (5.4) is used.

${}_1\tilde{\sigma}_{x,g}(E_j)$ is composed of the contributions of the single resonance series s.

$${}_1\tilde{\sigma}_{e,g}(E) = \sum_s \left\{ {}_1^s\tilde{\sigma}_{r,g}(E_j) - {}_1^s\tilde{\sigma}_{c,g}(E_j) - {}_1^s\tilde{\sigma}_{f,g}(E_j) + \tilde{\sigma}_p \right\} \quad (5.11)$$

$${}_1\tilde{\sigma}_{t,g}(E) = \sum_s {}_1^s\tilde{\sigma}_{r,g} + \tilde{\sigma}_p \quad (5.12)$$

The effective resonance cross section at the energy E_j and the series s is calculated by a modified theory*, first developed by R. Froelich /1/, /2/.

$${}_1^s\tilde{\sigma}_{x,g}(E) = \tilde{\sigma}_{p,eff}^s \frac{\left[1 + \frac{\langle \tilde{\sigma}_r \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma_x \cdot J(\beta, \theta)}{{}^sD \cdot \cos 2\delta_L} - \frac{{}^sD \langle \tilde{\sigma}_x \rangle \langle \tilde{\sigma}_r \rangle}{\Delta \sqrt{2\pi} \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon}{1 - \left[1 + \frac{\langle \tilde{\sigma}_r \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma_x \cdot J(\beta, \theta)}{{}^sD} + \frac{{}^sD \langle \tilde{\sigma}_r \rangle^2}{\Delta \sqrt{2\pi} \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon} \quad (5.13)$$

for fission and capture

* see appendix I

and

$${}^s\tilde{\sigma}_{rg}(E) = {}^s\tilde{\sigma}_{p,eff} \frac{\left[1 + \frac{\langle {}^s\tilde{\sigma}_r \rangle}{\langle {}^s\tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma \cdot J({}^s\beta, {}^s\theta)}{{}^sD} - \frac{{}^sD \langle {}^s\tilde{\sigma}_r \rangle^2}{\Delta \sqrt{2\pi} \langle {}^s\tilde{\sigma}_t \rangle^2} \cdot \varepsilon}{1 - \left[1 + \frac{\langle {}^s\tilde{\sigma}_r \rangle}{\langle {}^s\tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma \cdot J({}^s\beta, {}^s\theta)}{{}^sD} + \frac{{}^sD \langle {}^s\tilde{\sigma}_r \rangle^2}{\Delta \sqrt{2\pi} \langle {}^s\tilde{\sigma}_t \rangle^2} \cdot \varepsilon} \quad (5.14)$$

for the total resonance cross section.

${}^s\Gamma_x$ width for the neutron reaction (n,x) and the resonance series s

${}^s\Gamma$ total width for the resonance series s.

$$J(\beta, \theta) = \int_0^{\infty} \frac{\psi(\theta, x)}{\psi(\theta, x) + \beta} dx \quad (5.15)$$

$$\psi(\theta, x) = \frac{\theta}{2\sqrt{\pi}} \int_{-\infty}^{+\infty} \frac{\exp\left\{-\frac{\theta^2}{4} \cdot (x-y)^2\right\}}{1+y^2} dy \quad (5.16)$$

$$\theta = \frac{\Gamma}{\Delta} \quad , \quad \Delta = \sqrt{\frac{4kTE}{A}} \quad (5.17)$$

k = Boltzmann constant

E = energy

A = mass number

T = temperatur

$\delta_l = \frac{R}{\lambda} - \text{arctg} \frac{R}{\lambda} \cdot l$ for $l = 0, 1$

R = effective radius of the nucleus

λ = reduced neutron wave length

$\overline{s_D}$ = average level distance for the resonance series s.

$$\overline{\Gamma_f} \equiv \int_0^{\infty} \Gamma_n \cdot \Gamma_f \cdot {}^s F_n(\Gamma_n) \cdot {}^s F_f(\Gamma_f) d\Gamma_n d\Gamma_f \quad (5.18)$$

where ${}^s F_n(\Gamma_n)$ is the distribution of the neutron width of the resonance series s and ${}^s F_f(\Gamma_f)$ is the distribution of the fission width of the resonance series s.

$\langle {}^s \tilde{\sigma}_x \rangle$ average cross section of the neutron reaction (n,x) for the series s

$\langle {}^s \tilde{\sigma}_r \rangle$ average resonance cross section of the series s.

$$\langle \tilde{\sigma}_t \rangle \equiv \tilde{\sigma}_0 + \tilde{\sigma}_p + \sum_s \langle {}^s \tilde{\sigma}_r \rangle \quad , \quad (5.19)$$

with

$$\tilde{\sigma}_0 = \frac{1}{N} \left\{ \Sigma_t - N \cdot (\tilde{\sigma}_p + \sum_s \langle {}^s \tilde{\sigma}_r \rangle) \right\} \quad (5.20)$$

the average background cross section.

Σ_t is the total cross section of the mixture.

$${}^s\tilde{\sigma}_{p,eff} = \langle \tilde{\sigma}_t \rangle - \langle {}^s\tilde{\sigma}_r \rangle \quad (5.21)$$

$${}^s\beta = \frac{\langle \tilde{\sigma}_t \rangle}{{}^s\tilde{\sigma}_{oc}} \quad (5.22)$$

$${}^s\tilde{\sigma}_{oc} = 4\pi \lambda^2 \cdot g \cdot \frac{{}^s\Gamma_n}{{}^s\Gamma} \cdot \cos 2\delta_l \quad (5.23)$$

$$\epsilon = 2 \cdot \int_0^{\infty} \exp\left\{-\frac{D^2}{2\Delta^2}\right\} \cdot \Omega(D) dD \quad (5.24)$$

There must be noted that the expression (5.24) for the correction is only correct for Doppler broadened resonances, so that $\Gamma/\Delta \ll 1$. The formalism cannot be used for the natural line shape of the resonances.

$$\Omega(D) = \frac{1}{D} \int_{\mathcal{R}=1}^{\nu/2} \cos\left\{\frac{4\pi}{\nu} \mathcal{R} + \frac{\nu \cdot D}{2D} \sin \frac{4\pi}{\nu} \mathcal{R}\right\} \exp\left\{\frac{\nu D}{2D} \left(\cos \frac{4\pi}{\nu} \mathcal{R} - 1\right)\right\} \quad (5.25)$$

with $\nu = 10$.

To calculate the current weighted effective cross sections for the resonance series s , an approximation of H. Huschke* is used. That means, that the current weighted effective cross section can be calculated by the formulas (5.13) to (5.25), when $\langle \sigma_t \rangle$ is replaced by $\langle \sigma_t \rangle / 2$. This also means, that ${}^s\sigma_{p,eff}$ is replaced by $\{ {}^s\sigma_{p,eff} - \langle \sigma_r \rangle \} / 2$.

The average level distance $\overline{{}^sD}$ for the resonance series s at the energy E is given by

$$\overline{{}^sD} = \overline{{}^osD} \cdot \frac{\{ {}^sE_B + E \}^2}{\{ {}^sE_B \}^2} \cdot \exp \left\{ -\sqrt{89,72 ({}^sE_B - E)} + \sqrt{89,72 {}^sE_B} \right\} \quad (5.26)$$

$\overline{{}^osD}$ average level distance of resonance series s for low energies, $E \ll E_B$.

E_B binding energy of the last neutron in the compound nucleus in MeV.

E energy in MeV.

The statistical mean values of $\overline{{}^s\Gamma_x}$ (${}^s\beta$, ${}^s\theta$) are normally calculated by numerical integration. For the distribution of the neutron half widths $F_n(\Gamma_n)$ and the fission half widths $F_f(\Gamma_f)$ χ^2 -distribution are used.

$$F(\Gamma) d\Gamma = \frac{\nu}{2\Gamma \cdot G(\frac{\nu}{2})} \left(\frac{\nu \cdot \Gamma}{2\Gamma} \right)^{-1/2} \cdot \exp \left\{ -\frac{\nu}{2} \frac{\Gamma}{\Gamma} \right\} d\Gamma \quad (5.27)$$

ν degree of freedom

G here is the Γ -function

* see appendix II

In the case, where

$$\frac{\overline{sD}}{\sqrt{2\pi} \cdot \Delta} \cdot \frac{\langle s\tilde{\sigma}_r \rangle}{s\tilde{\sigma}_{p,eff}} \cdot sE_r < 0.1 \quad (5.28)$$

the following approximations are used:

$$\overline{s\Gamma \cdot J(s\beta, s\theta)} \cong \frac{\overline{sD} \cdot \langle s\tilde{\sigma}_r \rangle}{s\tilde{\sigma}_{p,eff}} \left[1 - \frac{\overline{sD}}{\sqrt{2\pi} \Delta} \cdot \frac{\langle s\tilde{\sigma}_r \rangle}{s\tilde{\sigma}_{p,eff}} \cdot sE_r \right] \quad (5.29)$$

$$\overline{s\Gamma_x \cdot J(s\beta, s\theta)} \cong \frac{\overline{sD} \cdot \langle s\tilde{\sigma}_x \rangle \cos 2\delta_L}{s\tilde{\sigma}_{p,eff}} \left[1 - \frac{\overline{sD}}{\sqrt{2\pi} \Delta} \cdot \frac{\langle s\tilde{\sigma}_r \rangle}{s\tilde{\sigma}_{p,eff}} \cdot sE_x \right] \quad (5.30)$$

for $x = \text{capture, fission}$.

The sE_r , sE_x are defined as

$$sE_r = \frac{\overline{\{s\Gamma_n\}^2}}{\{s\Gamma_n\}^2} = 1 + \frac{2}{\nu_n} \quad (5.31)$$

ν_n degree of freedom of the χ^2 -distribution for the neutron half width.

$${}^s E_x = \frac{\left[\frac{\{s\Gamma_n\}^2 \cdot s\Gamma_x}{s\Gamma} \right]}{s\Gamma_n \left[\frac{s\Gamma_n \cdot s\Gamma_x}{s\Gamma} \right]} \quad (5.32)$$

The ${}^s E_x$ are calculated by numerical integration, with the χ^2 -distributions defined in (5.21).

The average cross sections of the neutron reaction (n,x) for the resonance series s are given by

$$\langle {}^s \tilde{\sigma}_x \rangle = \frac{2\pi^2}{sD} \cdot \lambda^2 \cdot g \cdot \overline{s\Gamma_n} \cdot {}^s S_x \quad (5.33)$$

with

$${}^s S_x = \frac{\left[\frac{s\Gamma_n \cdot s\Gamma_x}{s\Gamma} \right]}{\overline{s\Gamma_n}} \quad (5.34)$$

and

$$\langle {}^s \tilde{\sigma}_r \rangle = \frac{2\pi^2}{sD} \cdot \lambda^2 \cdot g \cdot \overline{s\Gamma_n} \cdot \cos 2\delta_l \quad (5.35)$$

The sS_x are calculated by numerical integration with the χ^2 -distributions defined in (5.21)

The average neutron half widths are calculated from the reduced average neutron half widths

$$\overline{\sigma}_n = \overline{\sigma}_n^0 \cdot \sqrt{E} \cdot \frac{R^2 + (1+l)\lambda^2}{R^2 + \lambda^2} \quad (5.36)$$

for $l = 0, 1$.

The average group cross sections for infinite dilution are calculated in the following way:

$$\tilde{\sigma}_{x,g}^\infty = \frac{\sum_j \phi_j \left\{ \sum_s \langle {}^s\sigma_x \rangle \right\}}{\sum_j \phi_j} \quad (5.37)$$

for capture and fission,

$$\tilde{\sigma}_{e,g}^\infty = \frac{\sum_j \phi_j \left\{ \sum_s (\langle {}^s\sigma_r \rangle - \langle {}^s\sigma_c \rangle - \langle {}^s\sigma_f \rangle) + \sigma_p \right\}}{\sum_j \phi_j} \quad (5.38)$$

for elastic scattering, where ϕ_j is the weighting function at the group boundaries and in the middle of the group.

5.2 The numerical procedure in calculating the statistical averages.

Statistical averages of the following type have to be calculated:

$$\overline{f(\Gamma)} = \int_0^{\infty} f(\Gamma) F(\Gamma) d\Gamma \quad (5.39)$$

where $f(\Gamma)$ is a function, which depends on Γ , $F(\Gamma)$ is a probability distribution. $F(\Gamma)$ is a χ^2 -distribution of the degree of freedom ν .

The procedure of integration is the following. The whole range of integration is divided into n intervals so that

$$\int_{\Gamma_i}^{\Gamma_{i+1}} F(\Gamma) d\Gamma = \frac{1}{n} \quad \text{for all intervals } i=1, n \quad (5.40)$$

In the interval i the average $\overline{\Gamma}_i$ is calculated

$$\overline{\Gamma}_i = \frac{\int_{\Gamma_i}^{\Gamma_{i+1}} \Gamma \cdot F(\Gamma) d\Gamma}{\int_{\Gamma_i}^{\Gamma_{i+1}} F(\Gamma) d\Gamma} = n \cdot \int_{\Gamma_i}^{\Gamma_{i+1}} \Gamma \cdot F(\Gamma) d\Gamma \quad (5.41)$$

For χ^2 -distributions for several degrees of freedom

$$\chi_i = \frac{\overline{\Gamma}_i}{\overline{\Gamma}} \quad ; \quad i = 1, n \quad (5.42)$$

are tabulated.

$$\bar{\Gamma} \equiv \int_0^{\infty} \Gamma \cdot F(\Gamma) d\Gamma \quad (5.43)$$

The integral (5.39) then is approximated by

$$\overline{f(\Gamma)} = \sum_{i=1}^n f(\bar{\Gamma} \cdot \chi_i) \cdot \int_{\bar{\Gamma}_i}^{\bar{\Gamma}_{i+1}} F(\Gamma) d\Gamma = \frac{1}{n} \sum_{i=1}^n f(\bar{\Gamma} \cdot \chi_i) \quad (5.44)$$

This approximation reproduces $\bar{\Gamma}$ exactly

$$\bar{\Gamma} = \frac{1}{n} \sum_{i=1}^n \bar{\Gamma} \cdot \chi_i = \frac{1}{n} \sum_{i=1}^n \bar{\Gamma}_i$$

Using equation (5.41) one gets

$$\bar{\Gamma} = \sum_{i=1}^n \int_{\bar{\Gamma}_i}^{\bar{\Gamma}_{i+1}} \Gamma \cdot F(\Gamma) d\Gamma = \int_0^{\infty} \Gamma \cdot F(\Gamma) d\Gamma.$$

In the same way the statistical averages of the type

$$\overline{f(\Gamma_1, \Gamma_2)} = \int_0^{\infty} d\Gamma_1 \int_0^{\infty} d\Gamma_2 f(\Gamma_1, \Gamma_2) \cdot F_1(\Gamma_1) \cdot F_2(\Gamma_2) \quad (5.45)$$

are approximated by

$$\overline{f(\Gamma_1, \Gamma_2)} = \frac{1}{n_1 \cdot n_2} \sum_{i=1}^{n_1} \sum_{j=1}^{n_2} f(\bar{\Gamma}_1 \cdot \chi_{1,i}, \bar{\Gamma}_2 \cdot \chi_{2,j}) \quad (5.46)$$

with

$$\chi_{1,i} = \frac{\overline{\Gamma_{1,i}}}{\overline{\Gamma_1}} \quad , \quad \chi_{2,j} = \frac{\overline{\Gamma_{2,j}}}{\overline{\Gamma_2}} \quad . \quad (5.47)$$

5.3 For the calculation of energy resonance self shielding factors and average group constants from statistical resonance parameters the following subroutines are used:

FSTAT (MI, SIGO, J, ENG, NFE, REFE, EFE, MT, TEMP, SE, SM, XEUGZ)

In this subroutine the average group cross sections and the energy resonance self shielding factors are calculated from statistical parameters residing on the KEDAK-file. After the subroutine QUER has calculated the microscopic average cross sections and the effective microscopic cross sections at three energy points E numerical integration is performed and the average group cross sections and the self shielding factors are printed on paper and are stored on an external unit.

MI : number of values of the background cross section σ_0 .
SIGO(MI) : one-dimensional field containing the background cross sections [barns] .
J : number of energy group limits .
ENG(J) : one-dimensional field containing the energy group limits in [eV] .
NFE : number of points of the collision density,
0 or 1, if the function $\text{PHI}(E) = 1/E$ is used .
REFE(NFE) : one-dimensional field containing the energy points of the collision density in [eV] .
EFE(NFE) : one-dimensional field containing the flux density for all energy points in REFE .
MT : number of temperatures .

TEMP(MT) : one-dimensional field containing the temperatures in $[^{\circ}\text{K}]$.

SE(3,5,MI) : three-dimensional field which is calculated in the subroutine QUER .

SM(3,5,MI) : three-dimensional field which is calculated in the subroutine QUER .

XEUG(5,MI) : two-dimensional field containing the resonance self shielding factors for all σ_0 -values .
The first index indicates the reaction type

- 1: capture, flux-weighted
- 2: elastic scattering, flux-weighted
- 3: fission, flux-weighted
- 4: elastic scattering, current-weighted
- 5: total, current-weighted

A description of the COMMON and of the output is given in chapters 1 and 2.

QUER (E, TT, N, SE, SM, XNYN, FXNY, DEL, RQU, CHI, DN, EBI, XA, XL, GN, GG, NEY, EY, GFA, GI, IT, MI, SIGO)

In this subroutine the effective microscopic cross sections for the resonance series s at the energy E as defined by equation (5.13) and (5.14) are calculated. Also the effective microscopic cross sections at the energy E including the contributions of all resonance series s as defined by the equations (5.6), (5.11) and (5.12) are calculated in this subroutine. The following parameters must be defined:

E : energy in $[eV]$, at which the effective cross section should be calculated .

TT : temperature in $[^{\circ}\text{K}]$.

N : n-th energy point within an energy group at which the microscopic effective cross sections are calculated for the evaluation of the effective group cross sections.

- XNYN(10) : one-dimensional field containing the ν_n , the degree of freedom of the χ^2 -distributions for the neutron width, for all resonance series s ($s \leq 10$).
- FXNY(10) : one-dimensional field containing the ν_f , the degree of freedom of the χ^2 -distribution for the fission width, for all resonance series.
- DEL : $(4k/A) \cdot 10^6$, k = Boltzmann constant in $[eV/^\circ K]$, A = mass number.
- RQU : R^2 [barn], R = radius of the nucleus.
- CHI(25,4) : 2-dimensional field, containing the χ_i as defined in equation (5.42) for second degrees of freedom for the χ^2 -distribution ($n = 25$, $\nu = 1, 2, 3, 4$).
- DN(10) : one-dimensional field containing the average level distance at low energies \overline{os}_D for all resonance series s ($s \leq 10$) [meV].
- EBI : binding energy of the last neutron in the compound nucleus [eV].
- XA : $\Lambda_0^2 = E \cdot \lambda^2$ [eV barn].
- XL(10) : one-dimensional field containing the quantum numbers l for the angular momentum in all resonance series.
- GN(10) : one-dimensional field containing the average reduced neutron widths $\overline{\Gamma}_n^0 \cdot 10^3$ for all resonance series, $\overline{\Gamma}_n^0$ in $[eV^{1/2}]$.
- GG(10) : one-dimensional field containing the average capture widths for all resonance series in [meV].
- NEY : number of energy points, for which the fission widths are tabulated on the KEDAK-library.
- EY(100) : one-dimensional field containing the energy points, for which the fission widths are read from the nuclear data library [eV]; maximum number of points is 100.
- GFA(10, 100) : two dimensional field containing the average fission widths at the energies EY for all resonance series in [meV],
first parameter for the series,
second parameter for the energy points.

GI(10) : one-dimensional field containing the statistical parameter $\frac{2J+1}{2(2I+1)}$ for all resonance series.
IT : number of resonance series.
MI : number of values of the background cross section σ_0 .
SIGO(MI) : one-dimensional field containing the values of the background cross sections.

The following parameters are calculated in the subroutine

SE(3,5,MI) : three-dimensional field containing the microscopic effective cross section in [barn],
first index for energy,
second index for the reaction type.
1: capture, see formula (5.6)
2: elastic scattering, see formula (5.7)
3: fission, see formula (5.6)
4: elastic scattering, see formula (5.11)
5: total, see formula (5.12)
third index for the background cross section σ_0 .
SM(3,5,MI) : three-dimensional field containing the microscopic average cross sections in [barn],
first index for energy,
second index for the reaction type as for SE,
third index for the background cross section σ_0 ,
(not significant, because the average cross section does not depend on σ_0).

EZZ (E, NS, EZG, EZC, EZF, DRSG, DRFS, CHI, XNYN, FXNYN, GN, RQU,
XL, XA, GG, NEY, EY, GF)

In this subroutine the statistical mean values, defined by equation (5.32) and (5.34), are calculated.

The following parameters must be defined:

E : energy [eV].

NS : number of the resonance series to be calculated.

CHI(25,4) : two-dimensional field containing the χ_i as defined in equation (5.42) for several degrees of freedom for the χ^2 -distribution ($n = 25$, $\nu = 1, 2, 3, 4$).

XNYN(10) : one-dimensional field containing the ν_n , the degree of freedom of the χ^2 -distribution for the neutron half width, for all resonance series s ($s \leq 10$).

FXNYN(10) : one-dimensional field containing the ν_f , the degree of freedom of the χ^2 -distribution for the fission width, for all resonances.

GN(10) : one-dimensional field containing the average reduced neutron widths $\bar{\sigma}_n^0 \cdot 10^3$ for all resonance series. $\bar{\sigma}_n^0$ in [\sqrt{eV}],

RQU : R^2 [barn], R = radius of the nucleus.

XL(10) : one-dimensional field containing the quantum numbers l for the angular momentum for all resonance series.

XA : $\Lambda_0^2 = E \cdot \lambda^2$ [eV barn].

GG(10) : one-dimensional field containing the average capture widths for all resonance series [meV].

EY(100) : one-dimensional field containing the energy points, for which the fission widths are read from the nuclear data library [eV]; maximum number of points is 100.

GF(10, 100) : two-dimensional field containing the average fission widths at the energies EY for all resonance series in [meV],
first index for the series,
second index for the energy points.

The following parameters are calculated in the subroutine:

EZG : defined by equation (5.32), for capture.

EZC : defined by equation (5.31).

EZF : defined by equation (5.32), for fission.
DRSG : defined by equation (5.34), for capture.
DRFS : defined by equation (5.34), for fission.

TAB(XT, XK, LG, DSJ, XABCJ, DJK)

In this subroutine the $J(\beta, \theta)$ -function and the derivatives $\frac{\partial J}{\partial \beta}$ and $\frac{\partial J}{\partial \theta}$ are calculated.

The following parameters must be defined:

XT : $\theta = \frac{\Gamma}{\Delta}$, Γ and Δ in [eV], (5.17).
XK : k, where k is defined by $\beta = 2^k \cdot 10^{-5}$, β is defined by (5.22).
LG : 1, if XABCJ is calculated
2, if DJK is calculated
0, if DSJ is calculated
negative, if XABCJ, DJK and DSJ are calculated

The following parameters are calculated by the subroutine:

XABCJ : $J(\beta, \theta)$ -function as defined by (5.15).
DSJ : $\frac{\partial J(\beta, \theta)}{\partial \theta}$
DJK : $\frac{\partial J(\beta, \theta)}{\partial k}$, with $\beta = 2^k \cdot 10^{-5}$.

SUCH (E, NSU, EY)

This is an auxiliary subroutine

Besides the subroutines the following functions are necessary:

DMIT (E, NS, DN, EBI)

In this function the average level distance at the energy E is calculated by formula (5.26).

E : energy in [eV].
NS : number of the resonance series .
DN(10) : one-dimensional-field containing the $\overline{\sigma_D^s}$ for all resonance series s [meV].
EBI : binding energy E_B of the last neutron in the compound nucleus [eV].
DMIT : $\overline{\sigma_D}$ at the energy E for the resonance series NS in [meV].

GAMN (E, NS, GN, RQU, XL, XA)

In this function the average neutron width at the energy E is calculated by the formula (5.36).

E : energy in [eV].
NS : number of the resonance series.
GN(10) : one-dimensional field containing the reduced average neutron half widths $\overline{\Gamma_n^0}$ for all resonance series in [meV].
RQU : R^2 in [barn], R = radius of the nucleus
XL(10) : one-dimensional field containing the quantum numbers for the angular momentum for all resonance series.
XA : $\lambda^2 \cdot E$ in [barn], λ is the neutron wave length.
GAMN : $\overline{\Gamma_n}$, the average neutron half width at the energy E for the resonance series NS [meV].

GAMG (E, NS, GG)

In this subroutine the average capture half width at the energy E is provided.

E : energy in [eV].
NS : number of the resonance series.
GG(10) : one-dimensional field containing the average capture half widths for all resonance series in [meV].
GAMG : average capture half widths at the energy E for the resonance series NS in [meV].

GAFM (E, NS, NEY, EY, GF)

In this subroutine the average fission half width at the energy E is interpolated from the tabulated values.

E : energy [eV].
NS : number of the resonance series.
EY(100) : one-dimensional field, containing the energies, at which the average fission widths are tabulated [eV].
GF(10, 100) : one-dimensional field, containing the tabulated average fission widths in [meV].
GAFM : average fission width at the energy E for the resonance series NS in [meV]

POL (X, X1, FX1)

Auxiliary function for a linear interpolation .

X : argument, for which an arbitrary function should be interpolated .
X1(2) : one-dimensional field, containing two arguments of the function .
FX1(2) : one-dimensional field, containing the values of the function at the arguments X1.
POL : value of the function at X.

EPSI (E, T, NS, DN, EBI, DEL)

In this function the overlapping correction ϵ , defined by formula (5.24) is calculated.

E : energy in [eV],
T : temperature in [°K],
NS : number of resonance series.
DN(10) : one-dimensional field containing the $\overline{\sigma}_D$ for all resonance series s [meV],
EBI : binding energy E_B of the last neutron in the compound nucleus [eV],
DEL : $(4k/A) \cdot 10^6$, k = Boltzmann constant in [eV/°K],
A = mass number.

DELTA (E, T, NS, DEL)

In this function the Doppler width, formula (5.17), is calculated.

E : energy in [eV],
T : temperature in [°K],
NS : number of resonance series.
DEL : $\frac{4k}{A} \cdot 10^6$; k = Boltzmann constant in [eV/°K],
A = mass number.

PHASE (E, NS, XA, RQU, XL)

In this function $\cos 2\delta_1$ is calculated, where $\delta_1 = R/\lambda - \arctg \frac{R}{\lambda} \cdot 1$.
R is the effective radius of the nucleus, λ is the reduced neutron wave length.

E : energy in [eV].
NS : number of resonance series.

XA : $\Lambda_0^2 = E \cdot \lambda^2$ [eV·barn].
XL(10) : one-dimensional field containing the quantum numbers l for the angular momentum in all resonance series .

SIGC(E, NS, XA, GI, DN, EBI, GN, RQU, XL)

In this function the average capture cross section at the energy E for a single resonance series is calculated

E : energy in [eV].
NS : number of resonance series .
XA : $\Lambda_0^2 = E \cdot \lambda^2$ [eV·barn].
GI(10) : one-dimensional field containing the statistical parameter $\frac{2J+1}{2(2I+1)}$ for all resonance series.
DN(10) : one-dimensional field containing the average level distance at low energies $\overline{\sigma}_D$ for all resonance series in [meV].
EBI : binding energy of the last neutron in the compound nucleus in [eV]
GN(10) : one-dimensional field containing the average reduced neutron widths $\overline{\Gamma}_n^0 \cdot 10^3$ for all resonance series, $\overline{\Gamma}_n^0$ in [eV^{1/2}].
RQU : R^2 in [barn], R = effective radius of the nucleus.
XL(10) : one-dimensional field containing the quantum number l of the angular momentum for all resonance series.

References

/1/ R. Froelich

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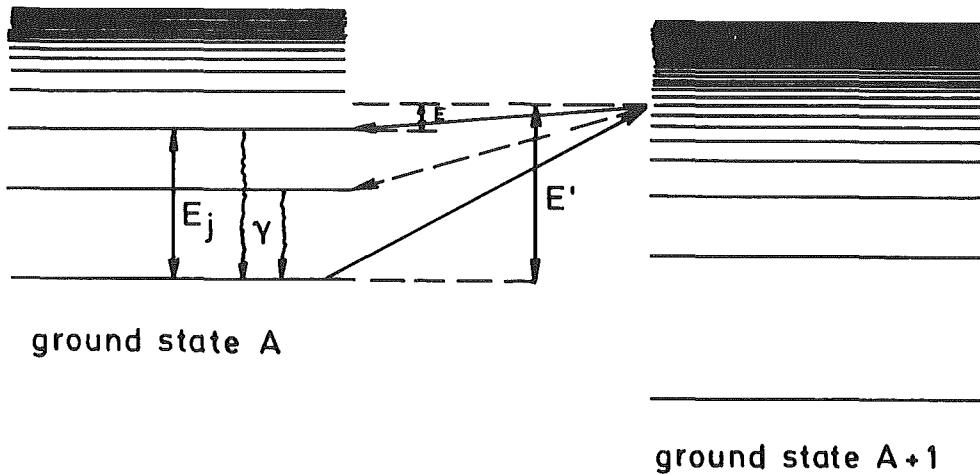
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6. The calculation of the matrices of the average zeroth moments of inelastic scattering. Module 5.

6.1 Description of the inelastic scattering process

Capture of a neutron with incident energy E' by a nucleus of atomic mass A leads to the formation of a compound nucleus $A + 1$ with excitation energy $AE = E_B + E'$ where E_B is the binding energy of the neutron in the compound nucleus. The compound nucleus decays then by emission of a neutron to an excited level of the residual nucleus, and the transition to the ground state occurs by γ -emission.



The excitation levels of the residual nucleus are either discrete ones or are belonging to the so-called continuum. The levels in the continuum cannot be resolved experimentally and concerning their position one can make statements only from statistical theory. In the discrete region the energies of the excitation levels are well-known from experiment.

The general expression for the probability of inelastic scattering out of any energy group g to any group h for an isotope k is given by Yiftah, Okrent, Moldauer /3/ in the following form:

$$k_{\sigma_i, g \rightarrow h}^{\infty} = \frac{\int_{E_{h+1}}^{E_h} dE \int_{E_{g+1}}^{E_g} dE' k_{\sigma_i}(E') \phi(E') P(E' \rightarrow E)}{\int_{E_{g+1}}^{E_g} \phi(E') dE'} \quad (6.1)$$

Here is E' the incident neutron energy and E the energy of the scattered neutron. $\phi(E')$ represents the energy dependent neutron flux. $P(E' \rightarrow E)$ indicates the transition probability of the inelastically scattered neutrons. For its determination one has to distinguish between the range of resolved excitation levels of the residual nucleus and the so-called continuum range in which the levels of the residual nucleus are undistinguishable.

6.1.1 Energy groups with discrete excitation levels for inelastic scattering

For given incident neutron energies E' in group g and outgoing neutron energies E in group h inelastic scattering can occur only to those excitation levels for which $E = E' - E_j$ where E_j is the energy of the j th level.

The transition probability therefore is a δ -function which has the value one, if the above condition is fulfilled and zero, if not. Thus we have

$$\begin{aligned} k_{\sigma_i}(E') P(E' \rightarrow E) &= \sum_{j=1}^N \sigma_i^{E_j}(E') P_j(E' \rightarrow E) \\ &= \sum_{j=1}^N \sigma_i^{E_j}(E') \delta(E' - E - E_j) \end{aligned} \quad (6.2)$$

where N is the number of excitation levels of the isotope considered,

P_j the probability of the transition to the j th level of the residual nucleus and $\sigma_i^{E_j}$ the inelastic excitation cross section of the j th level. With this special form of the transition - probability the double integration in equation (6.1) can be reduced to a single one over the energy group h to which the scattering occurs.

$$k_{i,g \rightarrow h}^{\infty} = \frac{\sum_j \int_{(h)} dE k_{\sigma_i}^{E_j}(E + E_j) \psi(E + E_j)}{\int_{(g)} \psi(E') dE'} \quad (6.3)$$

With the transformation $E + E_j \rightarrow E'$ it follows:

$$k_{i,g \rightarrow h}^{\infty} = \frac{\sum_j \int_{E_{h+1} + E_j}^{E_h + E_j} dE' k_{\sigma_i}^{E_j}(E') \psi(E') \chi(E_{g+1} \leq E' \leq E_g)}{\int_{E_{g+1}}^{E_g} \psi(E') dE'} \quad (6.4)$$

By the distribution function χ it was taken into account that the incident neutron energy E' should be contained in a 26-group g . On the other hand it must be of course

$$E_{h+1} + E_j \leq E' \leq E_h + E_j$$

because of the integration over this range. The integration therefore has to be extended only over the interval $[E_u, E_o]$ with the following definition

$$E_u = \begin{cases} E_{g+1}, & \text{if } E_{h+1} + E_j < E_{g+1} \\ E_{h+1} + E_j, & \text{if } E_{h+1} + E_j \geq E_{g+1} \end{cases} \quad (6.5)$$

$$E_o = \begin{cases} E_g, & \text{if } E_h + E_j > E_g \\ E_h + E_j, & \text{if } E_h + E_j \leq E_g \end{cases}$$

Instead of the neutron flux the collision density $F(E') = \Sigma_t(E') \phi(E')$ was taken as weighting spectrum since outside the resonance region $\Sigma_t(E')$ can be considered as constant and then $F(E') \sim \phi(E')$ (see also /5/). The formula (6.4) then becomes the form:

$$k \sigma_{i,g \rightarrow h}^{\infty} = \frac{\sum_j \int_{E_{h+1} + E_j}^{E_h + E_j} dE' k \sigma_i^{E_j}(E') F(E') \chi(E_{g+1} \leq E' \leq E_g)}{\int_{E_{g+1}}^{E_g} F(E') dE'} \quad (6.6)$$

The matrix elements $\sigma_{i,g \rightarrow h}$ are normalized in the computer program SCAT (described under section 6.2.1) to the total inelastic scattering cross section value over group g

i.e.

$$k \sigma_{i,g}^{\infty} = \frac{\int_{(g)} \sigma_i(E') F(E') dE'}{\int_{(g)} F(E') dE'} \quad (6.7)$$

The total inelastic scattering cross sections and the inelastic excitation cross sections are taken from the Karlsruhe Neutron Nuclear Data File KEDAK /6/.

6.1.2 Energy groups with unresolved excitation levels for inelastic scattering

If the incident neutron energy is high enough, to excite energy levels in the "continuum region" of the residual nucleus A, the level densities are calculated using the formula of Weisskopf.

Let U be the excitation energy of the nucleus of mass A and C and ν free parameters, then according to Weisskopf the level density $w(U)$ is

$$w(U) = C \exp(2\sqrt{\nu AU}) \quad (6.8)$$

The probability of inelastic scattering of a neutron from incident energy E' to final energy E is given by the statistical model as

$$P_i(E' \rightarrow E) = NE \exp(-E/\theta(E')) \quad (6.9)$$

where N is a normalization constant and θ is the "nuclear temperature", which is related to the level density of the nucleus by

$$\frac{1}{\theta(E')} = \frac{d}{dU} \ln w(U) \Big|_{U=E'} \quad (6.10)$$

For our special formula (6.8) this results in

$$\Theta(E') = \sqrt{\frac{E'}{\nu A}} \quad [\text{MeV}] \quad \begin{cases} E' \text{ in MeV} \\ \nu \text{ in MeV}^{-1} \end{cases} \quad (6.11)$$

The free parameter ν is an input quantity for the program. Szwarcbaum et al. /2/ recommend $\nu = 0.16 \text{ MeV}^{-1}$ for all nuclei, whereas in the YOM-cross-section-set /3/ $\nu = 0.0961 \text{ MeV}^{-1}$ had been used. The figure 1 shows the influence of the parameter ν on the spectrum of the inelastically scattered neutrons. As an example the inelastic scattering of neutrons with an incident energy $E' = 8.5 \text{ MeV}$ on U 238 was taken, using $\nu = 0.16$, $\nu = 0.0961$ and $\nu = 0.06$. Figure 2 shows the spectra for neutrons inelastically scattered by U 238 using $\nu = 0.16$ and for $E' = 10.5 \text{ MeV}$, $E' = 8.5 \text{ MeV}$ and $E' = 6.5 \text{ MeV}$. With (6.8) and (6.11) inserted into (6.1) the cross section for inelastic scattering of a neutron on material k from group g to group h is given by

$$k \sigma_{i,g \rightarrow h}^{\infty} = \frac{\int_{(g)} dE' k \sigma_i(E') F(E') N \Theta(E')^2 \exp\left(-\frac{E_{h+1}}{\Theta(E')}\right) \left(\frac{E_{h+1}}{\Theta(E')} + 1\right)}{\int_{(g)} F(E') dE'} \quad (6.12)$$

$$= \frac{\int_{(g)} dE' k \sigma_i(E') F(E') N \Theta(E')^2 \exp\left(-\frac{E_h}{\Theta(E')}\right) \left(\frac{E_h}{\Theta(E')} + 1\right)}{\int_{(g)} F(E') dE'}$$

In (6.12) the integration over the final energy group, which can be done analytically, has already been carried out.

The integration over E' can be carried out numerically only. As the original program was written for the IBM 7074 computer and we had only a limited region of computer storage for the inelastic scattering routine, the following approximation, the same as in /2/ and /3/ was used. For the nuclear temperature $\Theta(E')$ the value averaged by $F(E')$ over the energy-group g is taken. With this the integration over E' in (6.12) can be carried out easily with the result

$$\begin{aligned}
 k_{\sigma_{i,g \rightarrow h}}^{\infty} &= k_{\sigma_{i,g}}^{\infty} \cdot N \left[\Theta_g^2 \exp\left(-\frac{E_{h+1}}{\Theta_g}\right) \left(\frac{E_{h+1}}{\Theta_g} + 1\right) \right. \\
 &\quad \left. - \Theta_g^2 \exp\left(-\frac{E_h}{\Theta_g}\right) \left(\frac{E_h}{\Theta_g} + 1\right) \right] \quad (6.13) \\
 &= k_{\sigma_{i,g}}^{\infty} k_{\tilde{P}_{i,g \rightarrow h}} \cdot N
 \end{aligned}$$

the normalization constant N is obtained by the relation.

$$\sum_{h \geq g} k_{\sigma_{i,g \rightarrow h}}^{\infty} = k_{\sigma_{i,g}}^{\infty} \quad \text{that means} \quad (6.14)$$

$$N = \frac{1}{\sum_{h \geq g} k_{\tilde{P}_{i,g \rightarrow h}}}$$

so that finally the probability for inelastic scattering of a neutron by material k from group g to group h is

$${}^k P_{i,g \rightarrow h} = \frac{{}^k \tilde{P}_{i,g \rightarrow h}}{\sum_{h \geq g} {}^k \tilde{P}_{i,g \rightarrow h}} \quad (6.15)$$

Different from Szwarcbaum et al. /2/ and from the ABN-set /4/ we calculate the transition probabilities ${}^k P_{i,g \rightarrow h}$ for all groups h , which can be reached from group g until down to thermal energies. For groups of small energy however, the quantities E_{h+1}/θ_g and E_h/θ_g are very small compared to one (θ_g is around 0.5 MeV for U 238 and neutron incident energy $E' = 8.5$ MeV). When calculating the expression (6.13) on the computer two almost equal quantities have to be subtracted, which results in numerical inaccuracies. As a consequence i.e. the values of ${}^k P_{i,g \rightarrow h}$ were not monotonically descending for $h \geq 20$ in the energy-structure of the 26-group-ABN-set /4/ and sometimes even became negative. As a direct double precision calculation on the IBM 360/65 did not completely eliminate the inaccuracies, we programmed a double precision function, that calculates

$$\sum_{n=1}^{\infty} \frac{x^n}{n!} = \exp(x) - 1 = \tilde{F}(x)$$

By use of this function and using the following abbreviations

$$x_h = \frac{E_h}{\theta_g} \quad x_{h+1} = \frac{E_{h+1}}{\theta_g} \quad \Delta x_h = x_h - x_{h+1}$$

the expression, which has to be calculated to obtain ${}^k P_{i,g \rightarrow h}$ (6.13) is

$$\theta_g^2 \exp(-x_{h+1}) \left[-\Delta x_h - \tilde{F}(-\Delta x_h) - x_{h+1} \tilde{F}(-\Delta x_h) - \Delta x_h \tilde{F}(-\Delta x_h) \right] \quad (6.16)$$

With this numerical procedure values for ${}^k P_{i,g \rightarrow h}$ are calculated, that are monotonically descending for $h \geq 2o$.

6.1.3 Energy groups, where the excitation levels of the residual nucleus are partly resolved and partly unresolved.

For nearly all materials the last neutron incident energy, which can be treated with the discrete level method, lies within a 26 group, so that two different methods for calculation of the $P_{i,g \rightarrow h}$ have to be used in the same group. The results for the lower and upper part of the group are linked together as follows. Let ΔE_g^d be the energy range of group g , where the discrete level method is used, and ΔE_g^c the energy range, where the statistical model is used. Then for group g the transition probability for inelastic scattering is calculated according to the formula

$${}^k P_{i,g \rightarrow h}^{(d+c)} = d_g {}^k P_{i,g \rightarrow h}^d + c_g {}^k P_{i,g \rightarrow h}^c \quad (6.17)$$

In (6.17) $d_g {}^k P_{i,g \rightarrow h}^d$ is the transition probability to group h of neutrons, whose incident energies lie in the range of resolved excitation levels of the residual nucleus within group g , that means

$$d_g {}^k P_{i,g \rightarrow h}^d = \frac{\int dE' \int dE {}^k \sigma_i(E' \rightarrow E) F(E')}{\Delta E_g^d (h) \int_{(g)} dE' {}^k \sigma_i(E') F(E')}$$

and correspondingly

$$c_g {}^k P_{i,g \rightarrow h}^c = \frac{\int dE' \int dE {}^k \sigma_i(E' \rightarrow E) F(E')}{\Delta E_g^c (h) \int_{(g)} dE' {}^k \sigma_i(E') F(E')}$$

(6.18)

In the program ${}^k P_{i,g \rightarrow h}^d$ and ${}^k P_{i,g \rightarrow h}^c$ are both normalized in the following manner

$$\sum_{h \geq g} {}^k P_{i,g \rightarrow h}^d = 1 \quad \text{and} \quad \sum_{h \geq g} {}^k P_{i,g \rightarrow h}^c = 1 \quad (6.19)$$

so that

$$d_g = \frac{\int dE' {}^k \sigma_i(E') F(E')}{\int_{(g)} dE' {}^k \sigma_i(E') F(E')}, \quad c_g = \frac{\int dE' {}^k \sigma_i(E') F(E')}{\int_{(g)} dE' {}^k \sigma_i(E') F(E')} \quad (6.20)$$

A special subroutine for the calculation of (n, 2n) transition probabilities is being developed. So far the inelastic scattering transition probabilities are used for (n, 2n) reactions too.

6.2 Description of the computer program

The computer program for the calculation of the inelastic scattering matrices called SCAT is a subroutine of the managing program. It consists of the following subprograms:

A K E D, T R A, X K O N

Furthermore it uses the computer programs for retrieval of KEDAK-data NDFOPN, NDFLOC, NDFNXT. The function of the above computer-programs is outlined below.

6.2.1 S C A T

SCAT (NX, EG, XNUE, NFE, EF, FI, NE 27, WAHR, PROB, VW, QUER, E, NET, ET, NAE, AE, SU, QUOT, LBA, WEIN, AG, NETP, SGIT, NAEP, KMAX, ISG, ISGP, SGIP, IWE, IWP, WERT).

The parameters which have to be defined before calling SCAT are:

NX : number of energy group boundaries
NE 27 : NX + 1
EG(NX) : energy group boundaries in eV
NFE : = 0 in the case the weighting spectrum is given by
a function
= constant equal to the number of points at which
the weighting spectrum is given in this case
EF(NFE) : energy points of the weighting spectrum
FI(NFE) : values of the weighting spectrum
NET : dimension corresponding to the number of points
of the inelastic scattering cross section on
KEDAK for energies up to E (NAB + 1)
NAE : dimension corresponding to the number of
inelastic excitation levels
ISG : dimension corresponding to the maximum of the
number of KEDAK energy points for the different
inelastic excitation levels
IWE : dimension corresponding to the number of KEDAK
energy points in the "discrete" region (up to
E (NAB + 1)

Only working fields are the following quantities:

NETP, NAEP, ISGP, IWP, SGIT (NET), LBA(NX), VW(NX), QUER(NX), PROB(NX),
AG(NX), E(NX) and a group of parameters delivered by the subroutine
AKED:

ET(NET), SGIP (NAE, ISG), JMAT, KMAX (NAE), AE (NAE), SU (NX), QUOT (NX),
WERT (IWE)

Quantities calculated in SCAT are

WAHR(NX) = WEIN(NX) : normalized probabilities for inelastic scattering out of an energy group g into a group h.

The functions of the program SCAT are:

Calculation of the scattering matrices and organisation of the output for one material. This is done by the following successive steps:

1. Determination of the integration limits E_u, E_o for a given j, g, h.
2. linear interpolation of $k_{\sigma_i}^{E_j}(E')$ and $F(E')$ at E_u, E_o
3. calculation of the integral over E_u, E_o by trapezoidal rule
4. return to 1. and next j by keeping the g and h; summation over all j excitation levels for a given g, h, and division by $\int_{(g)} F(E') dE'$ (calculated in AKED)
5. return to 1. and next h by keeping g; calculation of $\sum_{h=1}^g k_{\sigma_{ig \rightarrow h}}^{\infty}$ and normalization of the probabilities

$$k P_{i,g \rightarrow h} = \frac{k \sigma_{i,g \rightarrow h}^{\infty}}{\sum_{h=1}^g k \sigma_{i,g \rightarrow h}^{\infty}}$$

6. return to 1. and next g and so on until g becomes the energy group, where the excitation levels are partly resolved and partly unresolved. Then instead of 6. the subroutine XKON is called and 5. is calculated only for the part of group g in which the excitation levels are resolved (see section 6.1.3).

6.2.2 AKED

The subroutine AKED is called in SCAT with the following arguments AKED (E, ET, SGIP, JMAT, KMAX, LMAX, SU, QUOT, WERT, AE, EMAX, EMIN, IA, NAB, NFE, EF, FI, NE, NX, NET, NETF, NETP, SGIT, NAE, NAEF, NAEP, ISG, ISGF, ISGP, IWE, IWF, IWP).

Parameters which have to be defined before calling AKED are:

NX : number of boundaries of the energy groups not necessarily including the thermal energy group

NE : number of energy group boundaries including the thermal energy group

= NX + 1 in the case the energy groups defined externally do not contain a thermal energy group (0.001 eV - 0.025 eV); NX in the case they do contain the thermal energy group

E(NE) : energy group boundaries

NFE : = 0 in the case the weighting spectrum is given by a function

= constant equal to the number of points at which the weighting spectrum is given in this case

EF(NFE) : energy points of the weighting spectrum

FI(NFE) : values of the weighting spectrum

NET : dimension corresponding to the number of points of the inelastic scattering cross section on KEDAK for energies up to E (NAB + 1)

NAE : dimension corresponding to the number of inelastic excitation levels

ISG : dimension corresponding to the maximum of the number of KEDAK energy points for the different inelastic excitation levels

IWE : dimension corresponding to the number of KEDAK energy points in the "discrete" region (up to E (NAB + 1))

working fields in AKED are the following:

SGIT (NET), FELD (6), KFEL (6), NFEL (6), NEFT, NETP, NAEF, NAEP, ISGF, ISGP, IWP, IWF.

Parameters which are determined in AKED and are transferred to SCAT are:

ET(NET) : KEDAK energies

SGIP(NAE,ISG) : two dimensional field filled by the inelastic excitation cross sections for a level NAE

JMAT : number of inelastic excitation levels for the material considered

KMAX(NAE) : number of energy points for the inelastic excitation cross sections of the different levels.

SU(NX) : one dimensional field giving the integral of the collision density $\int_{(g)} F(E) dE$ about a particular energy group g

QUOT(NX) : one dimensional field giving the value of the integral
$$\frac{\int_{(g)} {}^k\sigma_i(E') F(E') dE'}{SU}$$

WERT(IWE) : one dimensional field giving the value of the weighting function at the KEDAK energies

AE(NAE) : one dimensional field indicating the energies of the excitation levels

EMAX : upper energy limit of the "discrete" region

EMIN = ET(1) : last KEDAK energy at which ${}^k\sigma_i$ is still equal to zero

IA : number of the energy group in which EMIN is lying, i.e. lowest energy group out of which inelastic scattering is possible

NAB : number of the energy group in which EMAX is lying, i.e. where the excitation levels are to be considered to be partly resolved and partly unresolved.

AKED has the following functions:

1. reading of $\sigma_i(E')$ and the corresponding energy values E'
 reading of $\sigma_i^{Ej}(E')$ and the corresponding energy values E'
 above the treshold of the (n, n') process

2. linear interpolation of $F(E')$ at the KEDAK energy points
3. linear interpolation of $F(E')$ (by calling TRA) and $\sigma_i(E')$ at the energy boundaries of a given group structure
4. Integration of $\int_{(g)} F(E') dE'$ and $\int_{(g)} F(E') \sigma_i(E') dE'$ by using the trapezoidal rule
5. calculation of $\int_{(g)} F(E') \sigma_i(E') dE' / \int_{(g)} F(E') dE'$ and return to 3, and next g until the last g is reached i.e. the energy group NAB where the excitation levels are partly resolved and partly unresolved.

6.2.3 TRA

This subroutine has the following arguments:

TRA (EFI, FIS, NFE, EF, FI). The parameters which have to be defined before the subroutine call are:

NFE : = 0 in the case the weighting spectrum is given by a function
= constant equal to the number of points at which the weighting spectrum is given in this case.
EF(NFE) : energy points of the weighting spectrum
FI(NFE) : values of the weighting spectrum

The two other parameters are determined in TRA.

EFI : energy at which the weighting function is calculated
FIS : value of the weighting function at EFI.

The weighting spectrum is either given by a function ($NFE = 0$) and is then available over CALL FUNCTION PHI(E) or is given pointwise.

6.2.4 The subroutine XKON

XKON (NZG, NUGR, DI, EMAX, W, A, NE, ENG, NF, E, F, XNUE, NX)

Arguments submitted by the calling program:

- NZG : group number of that energy group, in which inelastic scattering occurs partially on unresolved excitation levels of the residual nucleus
- NUGR : number of energy group g , from which the neutron starts before the process of inelastic scattering.
- DI : normalization factor, that corresponds to c_g^k in formula (6.19).
DI is different from 0. only for NUGR = NZG.
- EMAX : highest energy of incident neutron, for which inelastic scattering occurs on resolved levels of the residual nucleus.
- A : mass number of the residual nucleus.
- ENG(NE) : one dimensional array, which contains the limits of the energy groups. (NE - 1) = number energy groups.
- NF : If NF is equal to 1, the weighting function for the averaging of the nuclear temperature is submitted by a function subprogram. Otherwise NF means the number of energy points, $E(NF)$, at which the value of the weighting function, $F(NF)$, is submitted by input.
- $\left. \begin{array}{l} E(NF) \\ F(NF) \end{array} \right\}$: energies and corresponding values of the weighting function, submitted by input. If NF is equal to 1, E and F are not used by XKON.
- NX : dimension of the vector field for the transition probabilities for inelastic scattering.

quantities, calculated by XKON

- W(NX) : double precision vector field, which by XKON is filled with $P_i, g \rightarrow h$ (see formula (6.15)) for a

group g corresponding to NUGR and all down-scattering groups h with equal or lower energy than g .

For

the numerical methods, using the function EXD(X) are applied as described above. The function EXD(X) called in XKON performs a double precision calculation of $\sum_{n=1}^{\infty} \frac{X^n}{n!} = \exp(X) - 1$. The function DNFAK(N) called in EXD(X) supplies the double precision values of $N!$

6.2.5 Output for one material

The program SCAT has an output on tape in the framework of the MIGROS - output and an output on listing. The output on listing has the following form:

1. line: SUBROUTINE SCAT
2. line:
2 Materialname Number of energy groups out of which inelastic scattering occurs
3. and eventually further lines

The scattering probabilities $P_i, g \rightarrow h$ in the form:

g	$g + 1$	$g + 2$	$g + 3$
$P_i, g \rightarrow h$	$P_i, g \rightarrow h + 1$	$P_i, g \rightarrow h + 2$	$P_i, g \rightarrow h + 3$	

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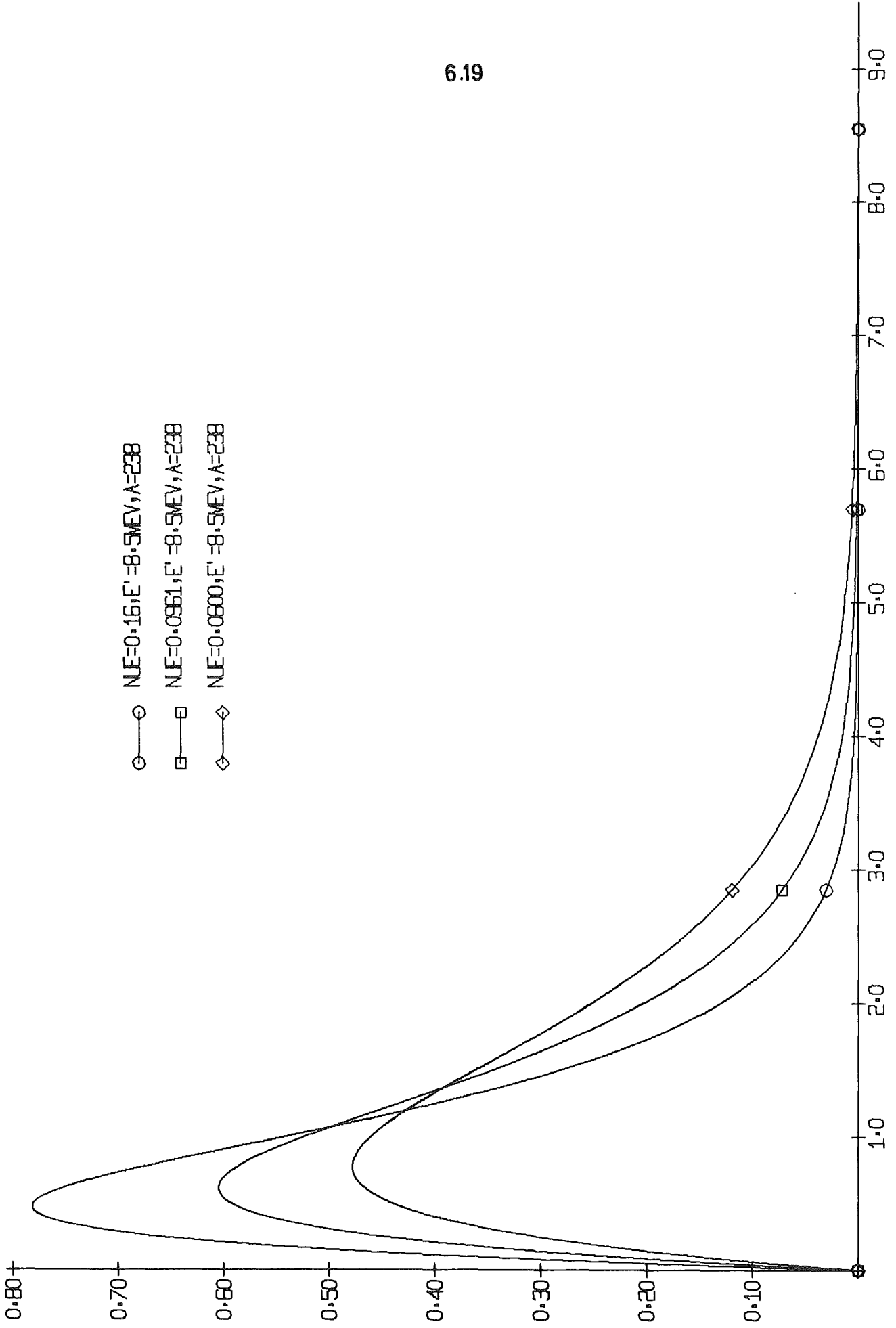


FIG.00001 P IN(E' -> E) = N * E * EXP(-E/THETA)

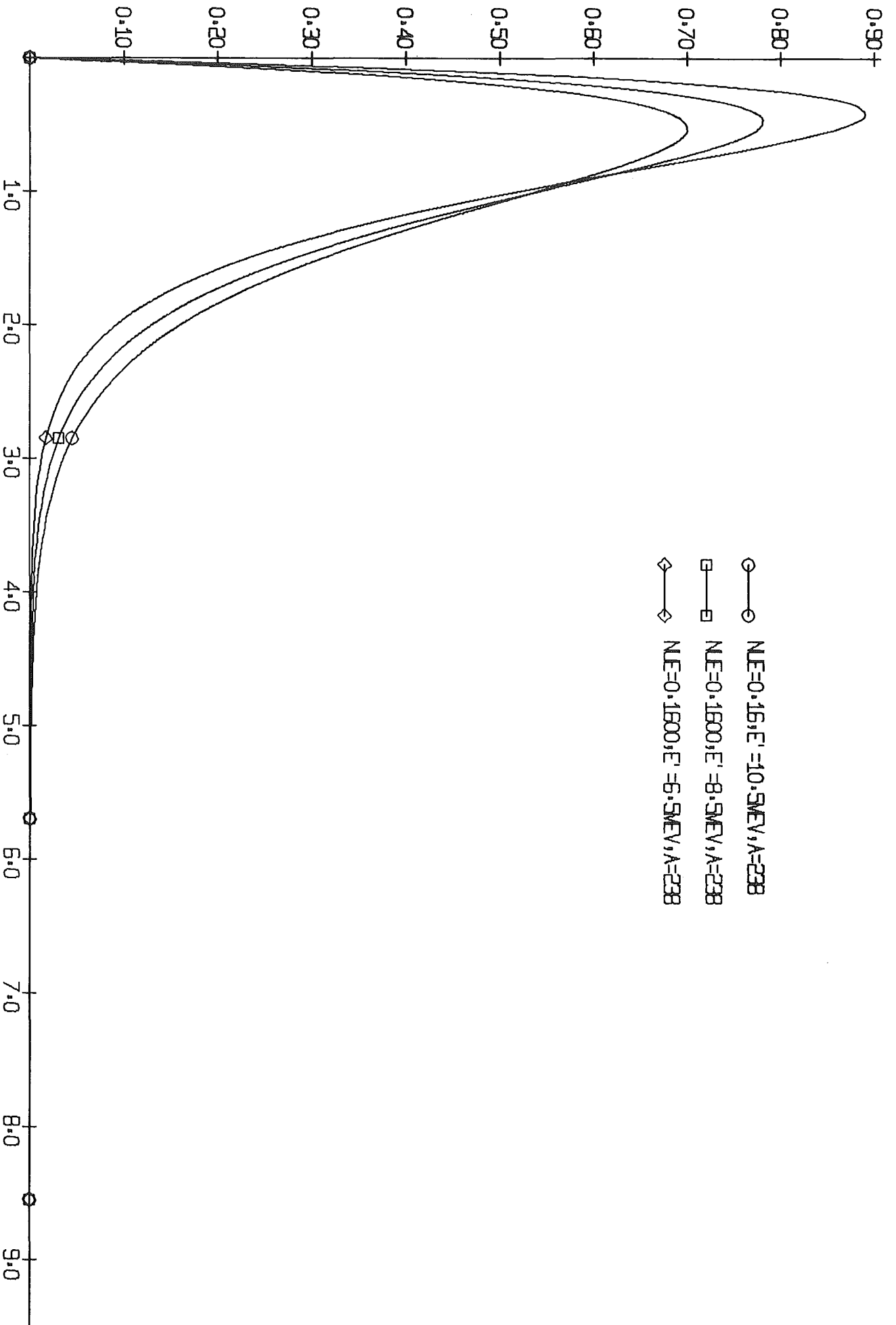


FIG.00002 P IN(E)→E=N*E*EXP(-E/THETA)

7. Microscopic Multigroup Scattering Matrices for Elastic Neutron Scattering in a P_1 -Approximation, $1 \leq 5$ *. Module 6.

7.1 Definition of the quantities calculated in this module

Given the KEDAK data

$\sigma(E)$ = microscopic cross section for elastic neutron scattering at energy E (KEDAK type SGN),

$\bar{w}(E, \bar{\mu}_0)$ = angular distribution in the c.m.-system for elastic neutron scattering with

$$2\pi \int_{-1}^{+1} d\bar{\mu}_0 \bar{w}(E, \bar{\mu}_0) = 1$$

E = energy of the incoming neutron,

$\bar{\mu}_0$ = scattering cosine in the c.m.-system, (KEDAK type SGNC),

$\bar{\mu}(E)$ = mean scattering cosine for elastic neutron scattering at energy E in the laboratory system (KEDAK type MUEL),

$\sigma_t(E)$ = total microscopic neutron reaction cross section at energy E (KEDAK type SGT),

and given a prescription for the approximation of the weighting functions $\Psi_1(E)$, $1 = 0, 1, \dots, 5$, this module calculates as final results three sets of quantities: the elastic cross section in the g -th group

$-\int_g$ means $\int_{E_g}^{E_{g+1}}$ with the group boundaries $E_1 > E_2 > \dots > E_G$ -

$$\sigma^g = \int_g dE \sigma(E) \Phi_0(E) / \int_g dE \Phi_0(E), \quad g = 1, 2, 3, \dots, G \quad (7.1)$$

* A detailed report (ref.1) of the code FLUMMI which is an extended version of this MIGRØS 2 module will be published as a KFK-report (in German). Further references will be found there.

the normalized matrixelements for scattering from group h into group g

$$S_e^{h \rightarrow g} = \frac{1}{\sigma^h} \int_h dE \sigma(E) w_\ell(E \rightarrow g) \Phi_\ell(E) / \int_h dE \Phi_\ell(E) \quad (7.2)$$

$$\ell = 0, 1, 2, 3, 4, 5, \quad g, h = 1, 2, \dots, G, \quad g \gg h$$

and the mean scattering cosine in the g-th group

$$\bar{\mu}^g = \sum_{h \gg g} S_1^{g \rightarrow h}, \quad g = 1, 2, 3, \dots, G \quad (7.3)$$

In (7.2) the quantity $w_1(E \rightarrow g)$ is defined as

$$w_\ell(E \rightarrow g) = \int_g dE' w_\ell(E, E'), \quad (7.4)$$

where

$$w_\ell(E, E') = \begin{cases} \mu_0(E, E') \\ \frac{d}{dE'} \int_{-1}^1 d\mu_0 \, 2\pi w(E, \mu_0) P_\ell(\mu_0) & \text{if } E' \leq E \leq E'/\alpha \\ 0 & \text{otherwise} \end{cases} \quad (7.5)$$

Here $\mu_0(E, E')$ is the scattering cosine for scattering from energy E to energy E' in the laboratory system:

$$\mu_0(E, E') = \frac{A+1}{2} \sqrt{\frac{E'}{E}} - \frac{A-1}{2} \sqrt{\frac{E}{E'}} \quad (7.6)$$

A = atomic weight of the scattering nucleus in neutron mass units,
 $P_l(\mu_0)$ is the l-th Legendre polynomial, $\alpha = \left(\frac{A-1}{A+1}\right)^2$

Further $w(E, \mu_0)$ is the angular distribution for elastic neutron scattering in the laboratory system, related to $\bar{w}(E, \bar{\mu}_0)$ through

$$w(E, \mu_0) = \bar{w}(E, \bar{\mu}_0(\mu_0)) \frac{d\bar{\mu}_0(\mu_0)}{d\mu_0} \quad (7.7)$$

where

$$\bar{\mu}_0(\mu_0) = \frac{1}{A} (\mu_0^2 - 1 + \mu_0 \sqrt{\mu_0^2 - 1 + A^2})$$

and therefore

$$\frac{d\bar{\mu}_0(\mu_0)}{d\mu_0} = \frac{1}{A} \frac{(\mu_0 + \sqrt{\mu_0^2 - 1 + A^2})^2}{\sqrt{\mu_0^2 - 1 + A^2}}$$

For computation the righthand side of the defining equation (7.4), which involves a double integration as it stands, can easily be simplified.

In the case of isotropy in the c.m.-system we have

$$w_l(E, E') = \begin{cases} \frac{\mu_0(E, E')}{\frac{d}{dE'} \frac{1}{2A} \int_{-1}^1 d\mu_0 \frac{(\mu_0 + \sqrt{\mu_0^2 - 1 + A^2})^2}{\sqrt{\mu_0^2 - 1 + A^2}} P_l(\mu_0)} & \text{if } E' \leq E \leq E'/\alpha \\ 0 & \text{otherwise} \end{cases}$$

hence

$$w_{\ell}(E \rightarrow g) = \frac{\hat{\mu}_0(E, E_{g-1})}{2A} \int_{\hat{\mu}_0(E, E_g)} d\mu_0 \frac{(\mu_0 + \sqrt{\mu_0^2 - 1 + A^2})^2}{\sqrt{\mu_0^2 - 1 + A^2}} P_{\ell}(\mu_0) \quad (7.8)$$

where

$$\hat{\mu}_0(E, E_g) = \begin{cases} -1 & \text{if } E > E_g/d \text{ and } E_g > \hat{E} \\ +1 & \text{if } E \leq E_g \text{ and } E_g > \hat{E} \\ \mu_0(E, E_g) & \text{if } E_g < E < E_g/d \text{ and } E_g > \hat{E} \\ -1 & \text{if } E_g = \hat{E} \end{cases} \quad (7.9)$$

Here \hat{E} is the lower limit of the considered energy region (In LECAL $\hat{E} = ABN (IM + 1)$). With (7.4) scattering to energies below \hat{E} is avoided, in particular: there is not scattering out of the lowest group.

In the case of anisotropy in the c.m.-system we have

$$w_{\ell}(E \rightarrow g) = 2\pi \int_{\hat{\mu}_0(E, E_g)}^{\hat{\mu}_0(E, E_{g-1})} d\mu_0 w(E, \mu_0) P_{\ell}(\mu_0) \quad (7.10)$$

7.2 Integration, energy- and cosine-mesh, interpolation, and weighting

The multigroup constants to be determined according to (7.1) and (7.2) are integrals over weighting functions and cross sections given at discrete energies with a varying density of these energy points. In the energy integrals we have as weighting functions the Legendre moments $\Phi_1(E)$ of the angle- and energy-dependent flux density $\Psi(E,\mu)$, approximated by the product of a slowly varying macro-spectrum and a possibly rapidly varying micro-spectrum. In angular integrations the Legendre polynomials $P_1(\mu)$ appear as weighting functions.

The numerical technique applied for integration is as follows:

- integrals are approximated by finite sums;
angular integration is performed with Simpson's rule since the cosine mesh can easily be spaced uniformly;
- energy integration is performed with the trapezoidal rule because of the varying density of the discrete cross section energies and the need for a fine mesh in the outgroup scattering region (defined below);
- the number of discrete points is chosen so that the program calculates cases which can be treated analytically with a maximum absolute error of $4 \cdot 10^{-4}$ in the normalized scattering matrix elements.

Interpolation of the SGN, MUEL, and SGT is linear; energy and angle interpolation of the SGNC is performed using polynomials of degree four.

The angular integrals are those in equ. (7.8) and (7.10). Since the Legendre polynomial $P_1(\mu_0)$ has 1 zeros in $[-1, +1]$, the integrals change sign, and there is partial cancellation.

The effect of cancellation increases with l , but this can be overcome by a cosine mesh with a number of meshpoints increasing with l .

For the computation of the integral in (7.10) - unisotropy in the c.m.-system - the following table gives the number NST(L) of the cosine meshpoints in $[-1, +1]$ for the L-th moment:

L	0	1	2	3	4	5
NST(L)	81	81	161	161	321	321

The integration according to (7.10) is performed by LEGANS. The integral in (7.8) is analytically given and does not depend on energy. The integrals

$$B_\ell(\omega) = \frac{1}{2} \int_{-1}^{\omega} d\mu_0 \frac{d\bar{\mu}_0}{d\mu_0} P_\ell(\mu_0) \quad (7.11)$$

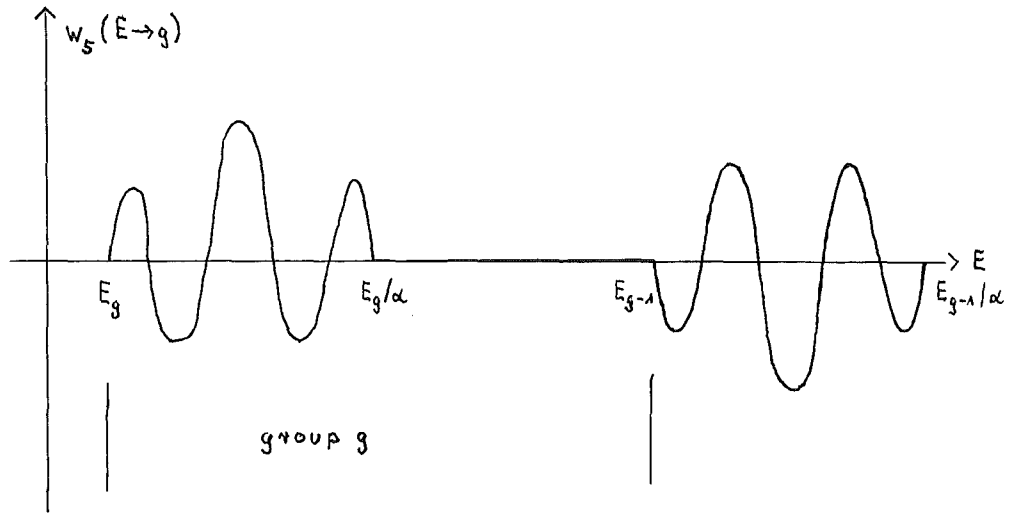
with $\bar{\mu}_0(\mu_0)$ according to (7.7) need be calculated only once for all groups in the region of isotropic scattering in the c.m.-system. So we can afford to determine $B_1(\omega)$ at NST(5) = 321 ω -points in $[-1, +1]$. The integration is done analytically, using different forms and approximations of the above integral for $A = 1$, $1 < A \leq 30$, $A > 90$, A = atomic weight of the scattering nucleus (see ref. 1).

The calculation of the $w_1(E \rightarrow g)$ according to (7.8) is performed by LEGIST, by interpolating in the table of the $B_1(\omega)$ determined by LEGINT.

The energy integral that requires the **finest** energy mesh is

$$\int_h dE \sigma(E) w_\ell(E \rightarrow g) \Phi_\ell(E) \quad (7.12)$$

The $w_1(E \rightarrow g)$ are of a form qualitatively shown in the following figure for $l = 5$:



The interval $[E_g, E_g/\alpha]$ from which scattering to energies $E < E_g$ is possible is called outgroup scattering region; if $E_g/\alpha < E_{g-1}$, the interval $[E_g/\alpha, E_{g-1}]$, from which scattering occurs only to energies $E \geq E_g$, is called ingroup scattering region.

In the outgroup scattering region, $w_1(E \rightarrow g)$ may change its sign several times. To describe this behavior properly we need a sufficiently fine energy mesh (independent of α). In the ingroup scattering region there is only a weak dependence on energy so that a coarse mesh will do.

Comparisons with analytical cases show that for $\ln \frac{E_{g-1}}{E_g} = \Delta u_g \lesssim 0.7$

NK = number of discrete energies in the outgroup scattering region, uniformly spaced in lethargy
= 70,

NR = number of discrete energies in the ingroup scattering region, uniformly spaced in lethargy
= 16

is an appropriate choice; other values can be specified by input.

NR should be large enough to guarantee a sufficiently accurate pointwise representation of the makro-weighting spectrum if there are only a few discrete SGN-energies.

The whole group is an outgroup scattering region if $E_g/\alpha \geq E_{g-1}$; in this case NK discrete energies are assigned to the group.

These NK + NR discrete energies per group are called basic points; they are determined in PUNK. Together with the discrete SGNC-energies found on KEDAK they form the coarse mesh corresponding to which the $w_1(E \rightarrow g)$ are calculated in LECAL.

In GRUPIN the coarse mesh and the discrete SGN-energies are joined by means of PUSUM to form the fine mesh. Linear interpolation of the $w_1(E \rightarrow g)$ and the $\sigma(E)$ to the discrete energy points of the fine mesh yields a pointwise representation of the $\sigma_1(E \rightarrow g) = \sigma(E) \cdot w_1(E \rightarrow g)$ which contains on one hand the information available from KEDAK (without smoothing by interpolation), and on the other hand the oscillatory behavior of the $\sigma_1(E \rightarrow g)$ in the outgroup scattering region.

The weighting function $\phi_1(E)$ is approximated by the product of a macro- and a micro-spectrum. In the following table the twelve possible forms of approximation of $\phi_1(E)$ are listed.

Macro-spectrum	Micro-spectrum	MAZ(1)	MAZ(2)	NSPEC
$1/E$	1	o	o	o
$1/E$	$1/\Sigma_t(E)$	o	o	o
$1/E$	$1/(\Sigma_t(E))^{1+1}$	o	1	o
$1/E^{1+1}$	1	1	o	o
$1/E^{1+1}$	$1/\Sigma_t(E)$	1	o	o
$1/E^{1+1}$	$1/(\Sigma_t(E))^{1+1}$	1	1	o
$F_o(E)$	1	o	o	1
$F_o(E)$	$1/\Sigma_t(E)$	o	o	1
$F_o(E)$	$1/(\Sigma_t(E))^{1+1}$	o	1	1
$F_1(E)$	1	1	o	1+1
$F_1(E)$	$1/\Sigma_t(E)$	1	o	1+1
$F_1(E)$	$1/(\Sigma_t(E))^{1+1}$	1	1	1+1

MAZ(1), MAZ(2) and NSPEC are input parameters controlling the weighting of the higher moments. The fine-structure mixture from which Σ_t is determined, and macro-pointspectra are input quantities.

Now the integral (7.12) is determined in the following way:

Let $E_1 < E_2 < \dots < E_{IKK}$ be the discrete energies of the fine mesh in group h, where E_1, E_{IKK} are the lower and upper boundary of group h. We put

$$f(E) = \sigma(E) \cdot w_e(E \rightarrow g) \cdot m(E)$$

- $m(E)$ denotes the macro-spectrum; then (7.12) reads

$$\int_h dE f(E) \cdot g(E)$$

where $g(E)$ denotes the micro-spectrum. After determination of $m(E_i)$, $i = 1, 2, \dots, IKK$, by interpolation or with FUNCTIØN PHI or FUNCTIØN PHIL, we have

$$\begin{aligned} \int_h dE f(E) g(E) &= \sum_{i=1}^{IKK-1} \int_{E_i}^{E_{i+1}} dE f(E) g(E) \\ &\approx \sum_{i=1}^{IKK-1} \bar{f}_i \cdot \Delta_i^g E \end{aligned} \tag{7.13}$$

where $\bar{f}_i = 0.5 (f(E_i) + f(E_{i+1}))$ and $\Delta_i^g E = \int_{E_i}^{E_{i+1}} dE g(E)$. The $\Delta_i^g E$ are called weighted intervall lengths (see GRUPIN, SPRAL, TRAPEZ). If fine-structure weighting is not taken into account, $\Delta_i^g E$ reduces to $\Delta E = E_{i+1} - E_i$ and (7.13) is the usual trapezoid sum. In the case of fine-structure weighting $\int dE g(E)$ is evaluated with the trapezoidal rule.

7.3 The routines of this module: purpose and a description of the
argument lists and the commons

Subroutine :	Purpose :
FLUMMI	main control program
SUM	summation
IPØLA	interpolation with polynomials of degree 4
PUSUM	joining of point sets
MAKRØ	construction of a macro-group structure
MIXSGT, ADD	calculation of the total macroscopic cross section for fine-structure weighting
TRAFØ	transformation of angular distributions
INFØRM	printing of additional information: groups, macro-groups, distributions of SGNC and SGN, group lethargy widths, group integrals of the zeroth moment of the weighting spectrum, number of discrete SGN- and SGT-energies per macro-group, number of fine-mesh energy points per macro-group
PUNK	generation of a basic energy mesh
(NDF)	reading of KEDAK
LØØKO	reading of ISØT1 (atomic weight) using NDF
LØØK1	reading of SGN, MUEL, SGT using NDF
LØØK2	reading of SGNC using NDF
LØØK3	reading of SGNC-energies using NDF
LECAL	control of the calculation of the $w_1(E \rightarrow g)$ according to (7.8) and (7.10)
LEGPØL	calculation of pointwise Legendre polynomials
LEGINT	calculation of Legendre integrale according to (7.11)
G(FUNCTIØN)	$G(N, X, Y) = (X^N - Y^N)/(X - Y)$
LEGIST	calculation of partial Legendre coefficients according to (7.8)
LEGANS	calculation of partial Legendre coefficients according to (7.10)
GRUPIN	weighted energy integration according to (7.2)

IPØLIN	linear interpolation
SPRAL	presentation of the moment-dependent total weighting spectrum as described in chapter 7.2
PHI(FUNCTIØN)	$PHI(E) = 1/E$
PHIL(FUNCTIØN)	$PHIL(L,E) = 1/E^L$
MUKØN	final determination of group constants according to (7.1), (7.2), (7.3)
PRINT	printing and reserving of final results

In the following a description of the argument lists and the commons of the subroutines is given. Under the heading "Arguments" those quantities are found which must be defined before computation will start. Under the heading "Results" those quantities are listed which are calculated by the routine. For a detailed description of the routines themselves see ref. 1.

SUBRØUTINE FLUMMI (A, B, EA, EB, E, EN, SGN, ECØ, SCØ, V, W, F, AR, FU, ABN, INT, GR, FEKØE, EG, R, RSP, ESP, SPEK, SGNC, FEKØ, ELSIG, ELTØT, ET, ST, MAT, DAT, NLA, NLE, ISEL, NMAX, NX, NE27, NSPEK, LSPEK, MAZ, NTK, KT, NTTTT, NTPP, ICØS, ICØSP, NECU, NECUP, ISM, ISMP, ISØ, ISØP, ISCØ, ISCØP, ISEC, ISECP, KIM, NS, NK, NR)

CØMMØN STØFF, ISTRUK, ISPA, NØUT, KPR, IM, IL, KL

Purpose:

- Construction of macro-groups by means of MAKRØ;
- calculation of unnormalized scattering matrices for each macro-group by means of LECAL and GRUPIN,
- calculation of normalized scattering matrices etc. by means of MUKØN,
- printing and reserving of results with PRINT

Arguments:

NX = number of energy group boundaries (input or standard)

(EG(I), I=1, NX) = energy group boundaries (increasing) (input or standard)

NE27 = NX + 1

Comment: the group boundaries will be modified and rearranged in the following manner:

ABN(I) = EG(NX - I + 1), I = 1, 2, ..., NX,

if ABN(NX) $\left\{ \begin{array}{l} > 10^{-3} & : \text{NEGR} = \text{NE27}, \text{ABN}(\text{NEGR}) = 10^{-3} \\ = 10^{-3} & : \text{NEGR} = \text{NX} \\ < 10^{-3} & : \text{ERRØR } 6.8 \end{array} \right.$

NGR = NEGR - 1

Now we have for further use in this module

NGR = number of energy groups

NEGR = number of group boundaries

(ABN(I), I=1, NEGR) = group boundaries (decreasing)

IL, IM = indices of the first and last energy group for which scattering matrices etc. shall be calculated (input)

NMAX = highest Legendre order which can be treated (NMAX = 5 in this version)

NLA = first Legendre moment to be calculated (input or standard, but finally NLA = 0 always)

NLE = last Legendre moment to be calculated (input or standard)

NSPEK = macro-weighting spectrum control parameter (NSPEK = 1: spectrum from FUNCTIØN PHIL; we have KSPE = KSPEK = 0 in GRUPIN; NSPEK > 1: NSPEK = number of discrete energies of the macro-weighting spectrum; we have KSPE = 1 and KSPEK = NSPEK in GRUPIN)

LSPEK = number of values of the macro-weighting spectrum if NSPEK > 1, LSPEK = 1 otherwise

(ESP(I),I=1,NSPEK) = discrete energies of the
macro-weighting spectrum if NSPEK > 1,
no meaning otherwise (input)

(SPEK(I),I=1,LSPEK) = values of the macro-
weighting spectrum (the first NSPEK
values for the 0. moment, the second
NSPEK values for the 1. moment etc.
if MAZ(1) = 1, the first NSPEK values
- there will be no more in this case -
for all moments if MAZ(1) = 0) (input)

NTK = micro-weighting spectrum pointer
(NTK = 1: the fine-structure mixture
from which the micro-weighting
spectrum is to be calculated is
given in the arrays MAT and DAT;
NTK = 0: no fine-structure weighting)

KT = number of different nuclei in the
fine-structure mixture (input)

(MAT(I),I=1,KT) = nuclei in KEDAK notation
in the fine-structure mixture (input)

(DAT(I),I=1,KT) = densities in particles per
ccm times 10^{-24} of the nuclei in the
fine-structure mixture (input)

NTTT = estimated number of discrete energies
of the micro-spectrum

(MAZ(I),I=1,2) = control parameters concerning
the weighting of the higher moments
(input)

ISM = estimated number of energy meshpoints
of the coarse mesh per macro-group

ISD = estimated number of energy meshpoints
of the fine mesh per macro-group

ICØS = number of angular meshpoints of the
angular distributions (SGNC) on KEDAK

NECU = scattering width

ISCØ = Max (NECU·ISD, ISM·ICØS)

ISEC = Max (ISD, ISM·NECU)
NS = maximum number of energy groups per macro-group (input or standard)
NK = number of basic points in the out-group scattering region per group (input or standard)
NR = number of basic points in the ingroup scattering region per group (input or standard)
STØFF = nucleus in KEDAK notation for which scattering matrices etc. shall be calculated
NØUT = printed-output unit (usually NØUT=6)
KPR = unit on which the final results are stored
ISEL = pointer for additional output (input or standard)
KL = MIGRØS counting index

Auxiliary arrays:

A(ISM), B(ISM), EA(ISM), EB(ISM), E(ISM),
EN(ISD), SGN(ISD), ECØ(ISD), SCØ(ISD),
V(ISD), W(ISD), F(ISD), AR(ICØS), FU(ICØS),
ABN(NE27), INT(NE27), GR(NE27), FEKØE(NE27),
R(NX), RSP(NX), SGNC(ICØS, ISM), FEKØ(ISM,
NECU), ELSIG(6, NECU, NX), ELTØT(2, NX),
ET(NTTT), ST(NTTT)

Results:

NTTP = increment of NTTT if NTTT is found to be for small during execution
ICØSP = increment of ICØS if ICØS is found to be too small during execution
NECUP }
ISMP } analogous
ISDP }
ISCØP }
ISECP }
KIM = downward extension in groups of the energy region to be calculated in order to treat outscattering in the low energy groups correctly

SUBROUTINE SUM(M, A, S)

Purpose: Summation of real numbers different in sign and magnitude with minimization of rounding errors

Arguments: M = number of summands
(A(I), I=1, M) = summands

Results:
$$S = \sum_{I=1}^M A(I)$$

SUBROUTINE IPOLA (M, A, B, N, X, Y, T)

Purpose: Continuous and differentiable interpolation with polynomials of degree four (for details see ref. 1)

Arguments: M = number of given points
(A(I), I=1, M) = given abscissae
(B(I), I=1, M) = given ordinates
N = number of the given new abscissae
(X(I), I=1, N) = given new abscissae
with $X(1) = A(1)$, $X(N) = A(M)$, $X(I-1) < X(I)$,
I = 2, N

Auxiliary array: T(M)

Results: (Y(I), I=1, N) = new ordinates corresponding with the given new abscissae

SUBROUTINE PUSUM (KA, A, KE, E, B)

Purpose: Joining of point sets

Arguments: KA = number of elements in the array A
(A(I), I=1, KA) = first point set (increasing)
KE = number of elements in the array E
(E(I), I=1, KE) = second point set (increasing)

Auxiliary array: B(KE)

Results: KE = number of different points in A and E
(E(I),I=1,KE) = different points in A and
the given E (increasing)

SUBROUTINE MAKRO (ALFA, ISØ, IL, IM, NGR, ABN, ISD, ISM, ISDP, NS,
NK, NECU, NUEB, V, W, LST, INTT, INT, NEGR, ISEL)
COMMON STØFF, ISTRUK, ISPA, NØUT

Purpose: A macro-group is defined as an energy region
consisting of one group or several groups
adjoining in pairs. MAKRO subdivides the
given energy region (ABN(IM +1), ABN(IL))
into macro-groups meeting the following
conditions:

each macro-group contains

- a) as many groups as possible
- b) NS groups at most
- c) ISD - ISM SGN-meshpoints at most
- d) either groups with only isotropic out-
scattering in the c.m.-system or groups
with only anisotropic outscattering in
the c.m.-system

Arguments: ALFA= $((A-1)/(A+1))^2$ where A is the atomic
weight of the scattering nucleus in neutron
mass units
IL, IM = indices of the first and last group
for which scattering matrices shall
be calculated
ISØ = index of the first group with only
isotropic outscattering in the
c.m.-system
NUEB = scattering width
NS = maximum number of groups per macro-
group
NK = number of basic points in the out-
group scattering region per group

NGR = number of energy groups
NEGR = number of energy group boundaries
(ABN(I), I=1, NEGR) = energy group boundaries
NECU = maximum scattering width
ISM = maximum number of basic points plus
discrete SGNC-energies per macro-
group
ISD = ISM plus maximum number of discrete
SGN-energies per macro-group
STØFF = nucleus in KEDAK notation for which
scattering matrices etc. shall be
calculated
NØUT = printed-output unit
ISEL = pointer for additional output

Auxiliary arrays: V(ISD), W(ISD), LST(2, NGR)

Results: INTT = number of macro-group boundaries
(INT(I), I=1, INTT) = macro-group boundaries
given by group indices
ISDP = increment of ISD if ISD is found
to be too small during execution
comment: the determined macro-groups are
as follows:
1. macro-group: [ABN(INT(2)+1), ABN(IN(1))]
I-th macro-group: [ABN(INT(I+1)+1), ABN(INT(I)+1)]
I = 2, 3, ..., INTT-1

SUBRØUTINE MIXSGT (KT, MAT, DAT, NTT, NTP, NT, ET, ST, EA, EE)

CØMMØN STØFF, ISTRUK, ISPA, NØUT

Purpose: Calculation of the total macroscopic cross
section of a given mixture according to

$$\sum_t(\epsilon) = \sum_{I=1}^{KT} d_I \cdot \sigma_{t,I}(\epsilon)$$

where d_I = density in nuclei per ccm times 10^{-24}
of the I-th nucleus in the mixture and $\sigma_{t,I}$ =
total microscopic cross section of the I-th

nucleus in the mixture; KT see below. E stands for the joint set of discrete energies for which for at least one nucleus in the mixture a σ_t is found on KEDAK.

Arguments: KT = number of different nuclei in the mixture
 (MAT(I), I=1, KT) = nuclei in the mixture
 (DAT(I), I=1, KT) = densities (DAT(I) = d_I)
EA, EE = lower and upper boundary of the energy intervall in which Σ_t shall be calculated
NTT = estimated number of discrete energies of the mixture in the intervall [EA, EE].
NØUT = printed-output unit

Results: NT = number of discrete energies of the mixture in the intervall [EA, EE]
 (ET(I), I=1, NT) = discrete energies in [EA, EE]
 (ST(I), I=1, NT) = total macroscopic cross sections of the mixture corresponding to the ET(I)
NTP = increment of NTT if NTTT is found to be too small during execution

SUBRØUTINE ADD (K, NTT, NTP, NT, ET, ST, NU, D, E, S, EV, SV, EW, SW, EA, EE)

Purpose: Determination of $P_{K, NU+1}(E|ST, ST)$ from $P_{K, NU}(E|ET, ST)$ (definition of $P_{I, J}(E|ET, ST)$ see below)

Arguments: K = index of the nucleus which shall be added to the hitherto calculated mixture of the first K-1 nuclei

EA, EA, NTT see MIXSGT

NT = number of the hitherto determined
discrete energies of the mixture
= number of ET(J) in $P_{K,NU}(E/ET, ST)$
(ET(I), I=1, NT) = hitherto determined discrete
energies
(ST(I), I=1, NT) = hitherto determined total
cross sections
NU = index of the highest energy up to
which the K-th nucleus has been
taken into account
D = density of the K-th nucleus
E, EV = energy points of the polynomial to
be added
S, SV = microscopic total cross sections at
E and EV
EW = ET(NU) before the last modification
SW = ST(NU) before the last modification

Results:

NT = number of discrete energies of the
mixture up to E
(ET(I), I=1, NT) = discrete energies up to E
(ST(I), I=1, NT) = total macroscopic cross
sections up to E
NU = index of the highest energy up to
which the K-th nucleus has been
taken into account
NTP see MIXSGT

Definition of $P(E|ET, ST)$ and $P_{I,J}(E|ET, ST)$:

Be

$$P(E|ET, ST) = P(E|ET(1), ET(2), \dots, ET(NT); ST(1), ST(2), \dots, ST(NT))$$

$$= \begin{cases} 0 & \text{if } E < ET(1) \text{ or } E > ET(NT) \\ ST(J) & \text{if } E = ET(J) \\ ST(J-1) + \frac{ST(J) - ST(J-1)}{ET(J) - ET(J-1)} (E - ET(J-1)) & \text{if } ET(J-1) < E < ET(J) \end{cases}$$

provided $ET(1) < ET(2) < \dots < ET(NT)$.

Be now

$$P_{I,J}(E)_{ET,ST} = \sum_{K=1}^{I-1} P_{K,NT_K}(E)_{ET,ST} + P(E)_{ET_I(1), ET_I(2), \dots, ET_I(J); ST_I(1), ST_I(2), \dots, ST_I(J)}$$

where the $ET_I(K)$ are the discrete σ_t -energies of the I-th nucleus and the $ST_I(K)$ are the corresponding macroscopic cross sections.

SUBROUTINE TRAFØ (IC, NE, ISM, S, XM, AL, A)

Purpose: Transformation of angular distributions for elastic neutron scattering from the c.m.-system into the laboratory system according to equ. (7.7)

Arguments: IC = number of uniformly spaced cosine meshpoints of the angular distributions in the c.m.-system
 NE = number of angular distributions undergoing transformation
 ISM = maximum number of angular distributions ((S(I,J),I=1,IC),J=1,NE) = angular distributions in the c.m.-system
 (A(I),I=1,IC) = uniformly spaced cosine meshpoints of the angular distributions in the c.m.-system with A(1) = -1, A(IC) = 1
 AL = ALFA (see MAKRO)

Results: (A(I),I=1,IC) = cosine meshpoints in the laboratory system
 ((S(I,J),I=1,IC),J=1,NE) = angular distributions in the laboratory system

SUBROUTINE INFØRM (ALFA, NEGR, ABN, R, RSP, IL, IM, INTT, INT, NST,
NSTIS, IR, ISTT, K, NTK, NTT, NØUT)

Purpose: Printing of additional information if ISEL < 0

Arguments: ALFA, NEGR, ABN, IL, IM, INTT, INT, NØUT see
MAKRØ
(R(I),I=1,NEGR-1) = group lethargy widths
(RSP(I),I=1,NEGR-1) = group integrals over the
total weighting spectrum of zeroth
order
(NST(I),I=1,6), NSTIS see LECAL
IR = macro-group index
ISTT = number of energy meshpoints in the
fine-mesh in the IR-th macro-group
K = number of discrete SGN-energies in
the IR-th macro-group.
NTK = pointer for micro-weighting
NTT = number of discrete energies of the
micro-weighting spectrum in the
IR-th macro-group if NTK ≠ 0

SUBROUTINE PUNK (NGR, NEGR, ABN, NA, NE, NEN, ALFA, NK, NR, ISM,
ISMP, ISN, E)

Purpose: Generation of a basic energy mesh uniformly
spaced in lethargy

Arguments: NGR = number of energy groups
NEGR = number of energy group boundaries
(ABN(I),I=1,NEGR) = group boundaries
NA, NE = first and last group in the considered
macro-group
NEN = number of angular distributions in
the considered macro-group
ALFA = $\left(\frac{A-1}{A+1}\right)^2$, A = atomic weight of the
scattering nucleus in neutron mass
units

NK = number of basic points in outgroup scattering region
NR = maximum number of basic points in ingroup scattering region
ISM = maximum number of basic points plus discrete SGNC-energies per macro-group

Results: ISN = calculated number of basic points in the macro-group defined through NA, NE
(E(I), I=1, ISN) = calculated basic points in this macro-group
ISMP = increment of ISM if ISM is too small

SUBROUTINE LØØKO (XMAT, MASSE, CØM, *)

CØMMØN STØFF, ISTRUK, ISPA, NØUT

Purpose: Reading the atomic weight from KEDAK and relating it to the neutron mass as mass unit, calculation of the mass number and of the mean scattering cosine in the laboratory system for isotropic scattering in the c.m.-system

Arguments: STØFF = nucleus in KEDAK notation

Results: XMAT = atomic weight of the nucleus in STØFF in neutron mass units
MASSE = mass number
CØM = $\frac{2}{3 \cdot XMAT}$

Remark: LØØKO calls the Entry NDFLØC of the KEDAK routine NDF

SUBROUTINE LØØK1 (ISD, K, E, SGN, EØ, EE, NT)

CØMMØN STØFF, ISTRUK, ISPA, NØUT

Purpose: Reading of SGN, MUEL and SGT from KEDAK

Arguments: STØFF = nucleus in KEDAK notation for which
data shall be read from KEDAK
ISD = maximum number of data which may
be read with a single call
EØ, EE = lower and upper boundary of the
energy region for which data shall
be read
NT = data-type pointer (NT = 1: SGN are
read; NT = 2: MUEL are read; NT = 3:
SGT are read)
NØUT = printed-output unit

Results: K = number of data (SGN, MUEL or SGT)
in [EØ, EE].
(E(I), I=1, K) = discrete energies at which data
a given
(SGN(I), I=1, K) = data read from KEDAK (SGN,
MUEL or SGT)

Remarks: The conditions $E(1) = EØ$ and $E(K) = EE$ are met
by linear interpolation or constant extra-
polation. LØØK1 calls the Entries NDFLØC
and NDFNXT of the KEDAK routine NDF.

SUBROUTINE LØØK2 (NEN, EA, ICØS, ICØSP, ISM, AR, SGNC)

CØMMØN STØFF, ISTRUK, ISPA, NØUT

Purpose: Reading of SGNC from KEDAK

Arguments: STØFF = nucleus in KEDAK notation
NEN = number of discrete energies at
which SGNC are given

(EA(I),I=1,NEN) = discrete energies for which
SGNC shall be read (they have been
read before with LØØK3)

ICØS = maximum number of SGNC cosine meshpoints

ISM = maximum number of discrete SGNC-energies

NØUT = printed-output unit

Results: ICØSP = 1 if ICØS is too small
= 0 otherwise

(AR(I),I=1,ICØS) = SGNC cosine meshpoints from
KEDAK

((SGNC(I,J),I=1,ICØS),J=1,NEN) = SGNC from KEDAK
corresponding to the cosine mesh AR
and the energy mesh EA

Remarks: The cosine mesh must be the same for all
energies.

LØØK2 calls the Entries NDFLØC and NDFNXT
of the KEDAK routine NDF.

SUBROUTINE LØØK3 (NE, EA, ISØ, NGR, ABN, ISM, ISMP, LST, NEGR, ISEL)

COMMON STØFF, ISTRUK, ISPA, NØUT

Purpose: Reading from KEDAK the entire SGNC-energy-
mesh for one nucleus

Arguments: STØFF, ISM, NØUT see LØØK2
NGR, NEGR, ABN see MAKRØ

Results: ISEL = pointer for additional output
ISMP = estimated increment of ISM if there
are more than ISM angular distributions
on KEDAK for the considered nucleus
NE = number of angular distributions on
KEDAK for the nucleus in question
(EA(I),I=1,NE) = SGNC-energy mesh on KEDAK for
the considered nucleus
(LST(I),I=1,NGR) where LST(I) = number of
energies EA(N) with ABN(I+1)

$$\leq EA(N) < ABN(I)$$

ISØ see MAKRO

Remark: LØØK3 calls the Entry NDFLØC of the KEDAK routine NDF

SUBROUTINE LECAL (MASSE, XMAT, ALFA, ISØT, LEG, NLE, NANF, NEND, NUEB, IL, IM, NEGR, ABN, ICØS, AR, ISN, E, SGNC, NST, NSTIS, ISM, ISD, NECU, ITA, NF, GR, FEKØE, FU, EW, A, H, V, W, F, FEKØ)

Purpose: Control of the calculation of the $w_{\perp}(E \rightarrow g)$ according to equations (7.8) and (7.10)

Arguments:

- MASSE = atomic mass number of the scattering nucleus
- XMAT = atomic weight of the scattering nucleus in neutron mass units
- ALFA = $((XMAT-1)/(XMAT+1))^2$
- ISØT = anisotropy pointer
(ISØT=1: the considered macro-group is within the energy region of anisotropic scattering in the c.m.-system, ISØT = 0 otherwise)
- (LEG(I), I=1,6) = sequence of calculated moments
- NLE = upper Legendre index (input)
- NANF, NEND = indices of the first and the last group in the macro-group for which scattering matrices etc. shall be calculated
- NUEB = scattering width
- IL, IM = indices of the first and the last group for which scattering matrices etc. shall be calculated
- NEGR = number of energy group boundaries
(ABN(I), I=1, NEGR) = energy group boundaries

ICØS = number of cosine meshpoints of the
SGNC on KEDAK
(AR(I),I=1,ICØS) = cosine meshpoints of the
SGNC on KEDAK
ISM = number of basic points plus discrete
SGNC-energies in the macro-group
to be calculated
(E(I),I=1,ISN) = basic points plus discrete
SGNC-energies in the macro-group
to be calculated (coarse mesh)
((SGNC(I,J),I=1,ICØS),J=1,ISN) = angular
distributions in the laboratory
system
(NST(I),I=1,6) = moment-dependent number of
uniformly spaced cosine meshpoints
for μ_0 -integration when ISØT = 1
NSTIS = moment-independent number of uniformly
spaced cosine meshpoints for μ_0 -
integration when ISØT = 0
ISM = maximum number of basic points and
discrete SGNC-energies per macro-group
ISD = ISM plus maximum number of discrete
SGN-energies per macro-group
NECU = maximum scattering width
ITA = pointer for the calculation of
Legendre integrals by means of
LEGINT
(ITA = 0: Legendre integrals are calculated,
ITA = 1: Legendre integrals have already
been calculated)
NF = unit on which the $w_1(E \rightarrow g)$ are stored
temporarily

Auxiliary arrays: GR(NEGR), FEKØE(NEGR), FU(ICØS), EW(ISM), A(ISD),
H(ISD), V(ISD), W(ISD), F(ISD), FEKØ(ISM, NECU)

Results: $w_1(E \rightarrow g)$ on unit NF

SUBROUTINE LEGPOL (NST, A, N, NSTIS)

COMMON /INTEG/ X, F, D

Purpose: Calculation of pointwise Legendre polynomials

Arguments: NST = number of uniformly spaced cosine meshpoints
 N = order up to which Legendre polynomials shall be calculated
 NSTIS = number of uniformly spaced cosine meshpoints if ISOT = 0 (here: maximum length of A)

Results: (A(I), I=1, NST) = uniformly spaced cosine meshpoints
 (X(I), I=1, NST) = uniformly spaced cosine meshpoints
 (F(L, I), I=1, NST), L=1, N+1) = Legendre polynomials pointwise
 $F(L, I) = P_{L-1}(A(I))$
 D = spacing of the A(I):
 $D = A(I) - A(I-1), I = 2, NST$

SUBROUTINE LEGINT (N, NAK, XMAT, ITA)

COMMON /INTEG/ H, GRAL, D

Purpose: Calculation of Legendre integrals

$$B_n(\omega) = \frac{1}{2A} \int_{-1}^{\omega} d\mu_0 \frac{(\mu_0 + \sqrt{\mu_0^2 - 1 + A^2})^2}{\sqrt{\mu_0^2 - 1 + A^2}} P_n(\mu_0)$$

$n = 0, 1, \dots, 5$

Arguments: N = upper Legendre index (is ignored; the $B_n(\omega)$ are always determined for $n = 0, 1, 2, 3, 4, 5$)
 NAK = number of cosine meshpoints
 XMAT = atomic weight in neutron mass units

Results: (H(I), I=1, NAK) = uniformly spaced cosine mesh-
 points, H(1) = -1, H(NAK) = 1
 D = spacing of the H(I)
 (GRAL(L,I), L=1,6), I=1, NAK) where GRAL(L,I)
 = $2 \cdot B_{L-1}(H(I))$
 ITA = pointer (ITA = 1 after calculation)

Remarks: For A = 1, $1 < A \leq 30$, $A > 30$ the above
 integral is treated numerically in different
 ways. For details see ref. 1.

FUNCTION G(N, X, Y)

$$G(N, X, Y) = \frac{X^N - Y^N}{X - Y}$$

is determined according to

$$\frac{X^N - Y^N}{X - Y} = D^{N-1} + \binom{N}{1} D^{N-2} Y + \binom{N}{2} D^{N-3} Y^2 + \dots + \binom{N}{N-1} Y^{N-1}$$

where

$$D = X - Y$$

SUBROUTINE LEGIST (L, NST, IA, IB, NEGR, GR, E)

COMMON /INTEG/ X, GRAL, D

Purpose: Calculation of partial Legendre coefficients
 when scattering is isotropic in the c.m.-
 system according to equ. (7.8)-multiplication
 with 2π is performed in LECAL:

$$E(I) = \frac{GR(I)}{4\pi} \int_{GR(I+1)} d\mu_0 \frac{d\bar{\mu}_0}{d\mu_0} P_{L-1}(\mu_0)$$

using the previously in LEGINT determined
 integrals

$$GRAL(L,I) = 2 \int_{-1}^{X(I)} d\mu_0 \frac{d\bar{\mu}_0}{d\mu_0} P_{L-1}(\mu_0)$$

Arguments: L = Legendre degree plus 1
 NST = number of uniformly spaced cosine mesh-
 points
 (X(I), I=1, NST) = cosine mesh from LEGINT
 D = spacing of the X(I)
 ((GRAL(L, I), L=1, 6), I=1, NST) see above
 NEGR = length of array E
 IA, IB = boundary indices within which the
 E(I) shall be determined
 (GR(I), I=IA, IB) = cosine integration
 boundaries

Results: (E(I), I=IA, IB-1) = partial Legendre
 coefficients (see above)
 E(NEGR) = total Legendre coefficient

$$= \sum_{I=IA}^{IB-1} E(I)$$

SUBROUTINE LEGANS (L, NST, F, IA, IB, NEGR, GR, E, PØLY, A)
 COMMON /INTEG/ X, PØL, D

Purpose: Calculation of partial Legendre coefficients
 when scattering is anisotropic in the c.m.-
 system according to equ. (7.10)-multiplication
 with 2π is performed in LECAL:

$$E(I) = \int_{GR(I+1)}^{GR(I)} d\mu_0 w(E, \mu_0) P_{L-1}(\mu_0)$$

Arguments: L, NST, D, NEGR, IA, IB, GR see LEGIST
 (X(I), I=1, NST) = cosine mesh from LEGPØL
 (F(I), I=1, NST) = pointwise angular distri-
 bution corresponding with the X(I)
 ((PØL(L, I), L=1, 6), I=1, NST) = Legendre polyno-
 mials pointwise from LEGPØL corres-
 ponding with the X(I)

Auxiliary arrays: PØLY(NST), A(NST)

Results: (E(I), I=IA, IB-1) = partial Legendre coefficients
(s. above)

E(NEGR) = total Legendre coefficient

$$= \sum_{I=IA}^{IB-1} E(I)$$

SUBRØUTINE GRUPIN (MASSE, XMAT, ALFA, CØM, ISØT, PM, LEG, NLE, NANF, NEND, NUEB, IL, IM, NGR, NEGR, ABN, ISN, EH, K, H, F, M, G, EN, NTT, ET, ST, MAZ, KSPE, KSPEK, ESP, SPEK, NSPEK, LSPEK, ISM, ISD, NECU, ISCØ, ISEC, NF, E, EW, GR, WA, U, V, W, ISTT, RSP, ELSIG)

Purpose: Weighted energy integration according to (7.2)

Arguments: MASSE, XMAT, ALFA, ISØT, LEG, NLE, NANF, NEND, NUEB, IL, IM, NEGR, ABN, ISN, ISM, ISD, NECU, NF see LECAL

CØM = mean scattering cosine when scattering is isotropic in the c.m.-system
= 2/(3·XMAT)

PM = fraction of MUEL in the corrected first moment

NGR = number of energy groups

(EH(I), I=1, ISN) = basic points plus discrete SGNC-energies in the macro-group to be calculated

K = number of SGN

(H(I), I=1, K) = discrete SGN-energies

(F(I), I=1, K) = SGN corresponding with the H(I)

M = number of MUEL

(G(I), I=1, M) = discrete MUEL-energies

(EN(I), I=1, M) = MUEL corresponding with the G(I)

NTT = micro-weighting spectrum control parameter (NTT = 1: no micro-weighting; NTT > 1: micro-weighting with NTT = number of energy points of the micro-weighting spectrum)

(ET(I),ST(I),I=1,NTT) = micro-weighting
spectrum if NTT > 1
(MAZ(I),I=1,2) = control parameters for the
weighting of higher moments (input)
KSPE = pointer for macro-weighting
(KSPE = 0: weighting with
FUNCTION PHI if MAZ(1) = 0, with
FUNCTION PHI and FUNCTION PHIL
if MAZ(1) = 1;
KSPE = 1: weighting with point-
spectrum in ESP and SPEK according
to MAZ(1)
KSPEK = number of discrete energies of the
macro-weighting spectrum if
KSPE = 1, = 0 otherwise
NSPEK = Max (1, KSPEK)
LSPEK = number of values of the macro-
weighting spectrum if KSPE = 1
= 1 otherwise
(ESP(I),I=1,KSPEK) = macro-spectrum discrete
energies if KSPE = 1; no meaning
otherwise
(SPEK(I),I=1,LSPEK) = macro-spectrum values
if KSPE = 1; no meaning otherwise
ISCØ = dimension of WA (see below and
FLUMMI)
ISEC = dimension of U (see below and
FLUMMI)

Auxiliary arrays: E(ISD), EW(ISM), GR(NEGR), WA(ISCØ), U(ISEC),
V(ISD), W(ISD)

Results: ISTT = number of energy meshpoints in the
fine-mesh consisting of the basic
points (from PUNK), the discrete
SGNC-energies, the discrete SGN-
and MUEL-energies

(RSP(I), I=1, NGR) = energy group integrals of
 the zeroth order weighting spectrum
 (ELSIG(L, I, NN), NN=IL, IM), I=1, NU EB), L=1, NLE+1)
 where with II = NN + I-1

$$\text{ELSIG}(L, I, NN) = \sigma_{L-1}^{NN \rightarrow II}$$

$$= \frac{\int_{NN} dE \sigma(E) w_{L-1}(E \rightarrow II) \bar{\Phi}_{L-1}(E)}{\int_{NN} dE \bar{\Phi}_{L-1}(E)}$$

SUBROUTINE IPOLIN (M, A, B, N, X, Y, HR)

Purpose: Linear interpolation

Arguments: M = number of given points
 (A(I), B(I), I=1, M) = given points
 N = number of given new abszissae
 (X(I), I=1, N) = given new abszissae
 where X(1) = A(1), X(N) = A(M)

Auxiliary array: HR(M)

Results: (Y(I), I=1, N) = new ordinates corresponding
 with the X(I)

SUBROUTINE TRAPEZ (M, E, F, G, W, H, L)

Purpose: Trapezoidal integration according to

$$W = \frac{1}{2} \sum_{I=2}^M (F(I) + F(I-1)) \cdot E(I-1) \quad \text{if } L = 0$$

$$W = \frac{1}{2} \sum_{I=2}^M (F(I) \cdot G(I) + F(I-1) \cdot G(I-1)) \cdot E(I-1) \quad \text{if } L = 1$$

(see chapter 7.2)

Arguments: M = number of meshpoints
 (E(I), I=1, M-1) = weighted intervall lengths
 (F(I), I=1, M) = function to be integrated
 (G(I), I=1, M) = weighting function
 L = pointer for weighting (L = 0: constant
 weighting function; L = 1: weighting
 function not constant in G)

Auxiliary array: H(M)

Results: W(see above)

FUNCTION PHI(E) = 1/E

FUNCTION PHIL(L,E) = 1/E^L

SUBROUTINE MUKØN (ELSIG, ELTØT, NLE, NECU, NUEB, IL, IM, NGR)

Purpose: Determination of σ^{NN} , S_L^{NN+II} and $\bar{\mu}^{NN}$
 from σ_L^{NN+II} according to (7.1), (7.2) and
 (7.3):

$$\sigma^{NN} = \sum_{II \geq NN} \sigma_o^{NN+II},$$

$$S_L^{NN+II} = \sigma_L^{NN+II} / \sigma^{NN}, \quad L = 0, 1, \dots, NLE,$$

$$\bar{\mu}^{NN} = \sum_{II \geq NN} S_1^{NN+II}$$

Arguments: NLE = Legendre degree up to which scattering
 matrices are calculated
 NUEB = scattering width
 IL, IM = first and last calculated energy group
 (((ELSIG(L, I, NN), NN=IL, IM), I=1, NUEB), L=1, NLE+1)
 where (see GRUPIN)
 $ELSIG(L, I, NN) = \sigma_{L-1}^{NN+NN+I-1}$
 NGR = number of energy groups (for dimensioning)
 NECU = maximum scattering width (for
 dimensioning)

Results: $((\text{ELSIG}(L,I,NN), NN=IL,IM), I=1, NUEB), L=1, NLE+1)$

where

$$\text{ELSIG}(L,I,NN) = S_{L-1}^{NN \rightarrow NN+I-1} \quad (\text{see above})$$

$((\text{ELT}\phi\text{T}(I,NN), NN=IL,IM), I=1,2)$ where

$$\text{ELT}\phi\text{T}(1,NN) = \sigma^{NN}$$

$$\text{ELT}\phi\text{T}(2,NN) = \bar{\mu}^{NN} \quad (\text{see above})$$

Remark:

Comparisons with cases which can be treated analytically (see ref. 1: FLUMMI tests) show that the matrixelements $S_{L-1}^{NN \rightarrow NN+I-1}$ are calculated here with a maximum absolute error

$$\Delta S \approx 4 \cdot 10^{-4}$$

Thus quantities smaller than 10^{-5} in absolute value are ignored and set equal to zero.

SUBROUTINE PRINT (ELSIG, ELT ϕ T, NLA, NLE, NECU, NUEB, ISEL, NGR, KSPEK,
MAZ, NTK, KIM)

COMMON ST ϕ FF, ISTRUK, ISPA, N ϕ UT, KPR, NEND, NANF

Purpose: Printing and reserving of final results

Arguments: ELSIG, ELT ϕ T see MUK ϕ N results
NLE, NECU, NUEB, NGR see MUK ϕ N
NLA = first Legendre moment to be calculated
(NLA = 0 always)
NANF, NEND = indices of the first and last
group for which scattering matrices
have been calculated
ST ϕ FF = nucleus for which scattering matrices
have been calculated
MAZ, NTK, KIM see FLUMMI
KSPEK corresponds to KSPE in GRUPIN
ISEL = pointer for additional output
N ϕ UT = printed-output unit
KPR = unit on which all MIGR ϕ S-results are
stored

References

- /1/ H.W.Wiese: FLUMMI - Ein FORTRAN-IV-Programm zur Berechnung von Multigruppenkonstanten aus KEDAK-Daten für höhere Transportnäherungen der Neutronentransportgleichung
(to be published as a KFK-report (in German))



8. Calculation of the fine interval elastic scattering transfer matrices up to P_5 approximation and other quantities required for the REMO-correction. Module 9.

8.1 Basic formulae and notations

a) The fine interval system

The I-th group of the basic broad group system is divided into NFG(I) fine-groups, equidistant in lethargy. Any of these fine groups is divided into NFI(I) fine intervals, equidistant in lethargy. The fine interval width is denoted by ΔE .

b) The fine-interval averaged cross-sections are defined by:

$$\langle \tilde{\sigma} \rangle = \frac{\int_{\Delta E} dE \tilde{\sigma}(E) F(E)}{\int_{\Delta E} dE F(E)} \quad (8.1)$$

The REMO correction procedure requires the total σ_t and the elastic scattering cross-sections σ_e and the averaged cosine which is defined by

$$\langle \mu_e \rangle = \frac{\int_{\Delta E} dE \mu_e(E) \tilde{\sigma}_e(E) F(E)}{\int_{\Delta E} dE \tilde{\sigma}_e(E) F(E)} \quad (8.2)$$

c) The elastic scattering transfer probability matrices are defined by

$$\langle {}_l P(E \rightarrow h) \rangle = \frac{\int_{\Delta E} dE \tilde{\sigma}_e(E) \pi'_l(E \rightarrow h) F(E)}{\sum_h \int_{\Delta E} dE \tilde{\sigma}_e(E) \pi'_o(E \rightarrow h) F(E)} \quad (8.3)$$

where

$$\pi_L(E \rightarrow h) = \frac{(A+1)^2}{4A} \cdot \frac{1}{E} \cdot \int_{\max(\alpha E, E_{h+1})}^{E_h} dE' g(\mu_c) P_L(t) \quad (8.4)$$

where $P_L(t)$ are the Legendre polynomials.

A is the mass of the scatterer, and

$$\alpha = \left(\frac{A-1}{A+1} \right)^2 .$$

The scattering angle in the center-of-mass system

$$\mu_c = 1 - \frac{(A+1)^2}{2A} \cdot \left(1 - \frac{E'}{E} \right) ,$$

and that in the labor system,

$$t = \frac{A+1}{2} \sqrt{\frac{E'}{E}} - \frac{A-1}{2} \sqrt{\frac{E}{E'}} .$$

$g(\mu_c)$ = the angular distribution of the elastic scattering in the center of mass system. It should be normalized i.e.

$$\frac{1}{2} \int_{-1}^{+1} d\mu_c g(\mu_c) = 1$$

($g(\mu_c) \equiv 1$, for isotropic case)

It is practical to transform the integral (8.4) into

$$\pi_L(E \rightarrow h) = \frac{1}{2} \int_{\mu_{h+1}}^{\mu_h} d\mu_c g(\mu_c) P_L(t(\mu_c)) , \quad (8.5)$$

where

$$f(\mu_c) = 0.5 \sqrt{1 - (1 - \mu_c) \frac{2A}{(A+1)^2}} \left\{ A+1 - \frac{A-1}{1 - (1 - \mu_c) \frac{2A}{(A+1)^2}} \right\} \quad (8.6)$$

Let the energy E fall into the group g . The neutron transfer may occur into the groups $g, g+1, \dots, g+k$, where $E_k > \alpha E > E_{k+1}$. As there is an unambiguous relation between the scattering angle μ_c and energy loss, a division of the angle interval $-1 \leq \mu_c \leq 1$ corresponds to the acceptor groups, i.e.

$$-1 \equiv \mu_{k+1} < \mu_k \dots < \mu_1 < 1$$

Thus the neutrons scattered with angle $1 \geq \mu > \mu_1$ remain in the group g , and those with the angle $\mu_j \geq \mu \geq \mu_{j+1}$ will fall into the group $g+j$.

According to our experience, the calculation and the physical interpretation are the simplest when the integration is performed over the center-of-mass scattering angle.

8.2 Method of calculation

According to the formulae (8.1), (8.2), (8.3) and (8.5) the main task is to perform numerical integrations.

For (8.1) and (8.2) a simple trapezoidal rule is used with the mesh points given by the nuclear data file KEDAK. At the end of the integration interval the cross-sections are interpolated.

For the integral (8.3) and (8.5) the Romberg integration method (see app.III) is used. The alternating sign of the Legendre poly-

nomial P_1 and the "ill-behaviour" of the $f(\mu_c)$ at some energy values require special consideration which are described in app. IV.

Before the integral (8.3) is performed, the fine interval ΔE is to be divided into subintervals with the following mesh-points:

- i All energy points for which elastic scattering cross-section are given in the nuclear data file should be a mesh point.
- ii If an energy value E_k/α falls into the fine interval, this also should be a mesh point. The Romberg integration procedure is performed for each subinterval.

8.3 The removal of the inconsistency between the average cosine calculated from the angular distribution and the average cosine available on KEDAK.

There are two types of the above inconsistency:

- a) The angular distribution of the scattered neutrons for the energy range in question are not given, but the retrieved average cosine $\overline{\mu_{ret}}$ is different from $2/3A$. In this case a linear anisotropy is assumed i.e. the following center of mass angular distribution is used

$$f(\mu_c) = 1 + \left(\overline{\mu_{ret}} - \frac{2}{3A} \right) \cdot \frac{3}{1 - \frac{3}{5A^2}} \cdot \mu_c \quad (8.7)$$

- b) The nuclear data file gives center of mass angular distributions for the energy range in question, but the calculated average cosine differs from the retrieved one. This difference can be attributed to the uncertainties in the evaluated angular distribution as well as to the uncertainties in their interpolation. Consequently there exist only ad hoc methods for the removal of this inconsistency. We have applied the following simple treatment.

Let the Legendre expansion of the scattering probability in the labor system be.

$$W_s(\mu) = \sum_{l=0}^{\infty} \frac{2l+1}{2} \cdot b_{s,l} \cdot P_l(\mu), \quad b_{s,0}=1, \quad b_{s,1}=\overline{\mu}_{ang} \quad (8.8)$$

$\overline{\mu}_{ang}$ is the labor average cosine calculated from the given mass angular distribution. The above scattering probability is corrected with the retrieved average cosine i.e.

$$\overline{W}_s(\mu) = W_s(\mu) + \frac{3}{2} (\overline{\mu}_{ret} - \overline{\mu}_{ang}) P_1(\mu) \quad (8.8')$$

and the correction to a transfer matrix element is

$$\Delta \pi_l(E \rightarrow h) = \frac{3}{2} (\overline{\mu}_{ret} - \overline{\mu}_{ang}) \int_{\mu_{h+1}}^{\mu_h} d\mu P_l(\mu) P_1(\mu) \quad (8.9)$$

where (μ_h, μ_{h+1}) is the angle interval in the labor atory system belonging to the scattering from E into the group H.

8.4 Strategy of the extrapolation

It is often the case that for a cross-section or angular distribution the energy interval given in the KEDAK file does not cover the energy range of the group system. In this case an extrapolation is necessary. As the difference between the last energy point and the end of the group energy range is generally not too large, additional energy point coinciding with the end of the group

is taken, with the cross-section value (or angular distribution) equal with the next one. (i.e. the extrapolation is horizontal)

8.5 The description of the program REMO and of the related subroutines

The subroutine REMO is called by the statement

```
CALL REMO (NG, ENG, NSP, F, E, ERR, NLA, NLE, NGRE, NFG, NFI, NJM, NUJM,  
          ISEL, IWORK, WORK, L)
```

NG	The number of energy boundaries
ENG(I),I=1, NG	The energy group boundaries /eV/
NSP	The number of energy points for the neutron spectrum
F(I),I=1, NSP	The neutron spectrum
E(I),I=1, NSP	Energy points for the neutron spectrum
ERR	The error limit for the Romberg integration
NLA, NLE	The first and the last moment to be calculated
NGRE	The number of groups to be calculated
NFG(I),I=1, NGRE	The number of finegroups
NFI(I),I=1, NGRE	The number of fine-intervals per fine-group.
NJM	The maximum number of mesh points for angle integration is 2**NJM+1
NUJM	The maximum number of mesh points for energy integration is 2**NUJM+1
ISEL	= 1:the formula (8.3) will be computed o:the formula to be obtained from (8.3),when $\sigma_s(E) = 1$, will be computed.
IWORK(I),WORK(I),I=1,L	Work-field
L	The length of the work-field.

In the subroutine system related to REMO the dynamical programming method is used, i.e. any dimensioned variable, the dimension of which depends on the task option is dimensioned with variable length. All of such variables are placed in the work field (WORK and IWORK respectively) to be passed by the main control program. This field is generally transferred from one subroutine to the other and it contains a great deal of information to be passed. The first 34 words of WORK contain mainly the control information. (see Table VIII)

The function of the main subroutine REMO is the following:

To settle the control information contained in the Table VIII, and to determine the length of WORK.

To call the subroutine KEDLEC (see below) in order to retrieve the angular distributions in the energy region to be calculated.

REMO calls the subroutine

REMP (ENG, NFG, NFI, F, E, WORK, IWORK, IS) ⁺

IS is the lowest energy group for which angular distributions are given.

The function of REMP: to organize the calculation by group.

It calls the following subroutines: KEDDAT, SINT, ISOFAL, ZWIN, NORM

The subroutine underlying the retrieval of KEDAK data has four entries.

α) KEDDAT (NAMIZ, NAMTYP, NW, NW1, FNW, NDAT, FA, FF, EP, EM, *, *, *)

NAMIZ /8-bytes/ name of the material

NAMTYP /8-bytes/ data type name

NW if there is a further name the sequential number of the required namecombination, otherwise without meaning.

NW1 if there is a further name then the number of retrieved name combination, otherwise without meaning.

FNW if there is a further name then its value, otherwise without meaning.

NDAT the number of retrieved data pairs.

FA(I), I=1,NDAT argument values.

FF(I), I=1,NDAT function values.

EP,EM upper and lower energy boundary [eV], for which the cross-section values are to be retrieved.

If the material and the type are not found then follow RETURN1 and RETURN2 respectively.

⁺ If the parameters of the subroutines are not described by the same symbols as in the program lists, symbols described before or in the input description are used.

b) DATNUM (NAMIZ, NAMTYP, NDAT, NW, EP, EM, *, *, *)

This entry is used when only NDAT is to be determined

c) KEDLEC (NAMIZ, NAMTYP, NW1, EA, IAE, EPW, EMW, FA, FF, FAW, FFW, EP, EM, *, *, *)

This entry is used to retrieve angular distributions in the energy interval $[EPW, EMW]$ [eV]

In this case EP, EM are not energy but cosine of the scattering angle.

EA(I), I=1, NW1 the number of energy points.

IEA(I), I=1, NW1 the number of angle points per angular distribution.

FAW(I), I=1, NW1xNDAT the angle points of angular distributions.

FFW(I), I=1, NW1xNDAT the angular distributions (in one field, one by one)

d) NIVNUM (NAMIZ, NAMTYP, NW1, EPW, EMW, EFR, *, *, *)

EFR is the energy for the first angular distribution. This entry is used for the determination of NW1, and EFR.

SUBROUTINE SINT (ENG(I), NFG(I), NFI(I), F, E, WORK, N, M)

SINT calculates the group averaged total and elastic scattering cross sections and average cosine in the group I.

M = IWORK(7)

For N = IWORK(5) SINT calculates $\langle \sigma_t \rangle$

For N = -IWORK(5) SINT calculates $\langle \sigma_s \rangle$ and $\langle \mu_s \rangle$

Called by REMP.

SUBROUTINE NORM (J2, J1, SGNC, WORK)

NORM normalizes the angular distribution from the J2-th up to the J1-th one.

Called by REMP.

SUBROUTINE ISOFAL (ENG(I), NFG(I), NFI(I), F, E, WORK)

ISOFAL divides the group I into fine intervals and organises the calculation of elastic transfer probabilities for the group I.

Called by REMP.

SUBROUTINE WAHRS (EP, EM, IMAX, WORK(N), ENG(I), NML, NUJM, F, E, TE,
WORK, PER)

N = IWORK (22)

NML = NLE+1

EP, EM are the higher and the lower boundary of the energy interval
to be calculated [eV].

WAHRS divides the interval [EP, EM] further according to the principles
given in 8.2.

Called by ISOFAL

SUBROUTINE LMI (EP, EM, E1, E2, SG1, SG2, AMU1, AMU2, IMAX1, ENG(I),
TE, NML, NUJ, NUJM, F, E, WORK, PER (IMAX+2), IMAX, PL)

SG1, SG2, AMU1, AMU2, E1, E2 elastic scattering cross sections,
retrieved average cosine and the related energy points [eV]

IMAX1 The maximum group change by a scattering in the interval [EP, EM]

Called by WAHRS.

FUNCTION ANINT (BU, BL, K, SG, XL, AZ, ICOS, WORK, T, NJM)

ANINT calculates the integral (8.5) and the correspondence of nota-
tions is the following

K = n

BU = μ_k

BL = μ_{k+1}

XL $\rightarrow \mu_c$

SG $\rightarrow f(\mu_c)$

AZ = $(\bar{\mu}_{ret} - \frac{2}{3A})$

} ICOS number of mesh points

Called by NORM, LMI.

SUBROUTINE AMESH(BU, BL, Z, L2, NX, K)

AMESH prepares the division of angle interval [BU, BL] according to the
roots of Legendre polynomial of order K. (see appendix IV)

Z L2, ... Z NX-1 The roots of the K-th Legendre polynomials in the
interval BU, BL

Called by ANINT.

SUBROUTINE HIDR (BU, BL, K, A, *)

In the case of hydrogen the integral (8.5) can analytically be calculated.
This is done by HIDR.

A = result of the calculation

Called by ANINT.

FUNCTION WINK (XL, SG, KJ, ICOS, AZ, XA)

WINK determines $\int (\mu) \sigma_s$ by means of linear interpolation from the
point-wise angular distribution /XA \rightarrow μ /

Called by ANINT.

FUNCTION FXINT (EP, EM, SGN, ES, F, E, NSP, NDAT)

FXINT calculates the integral $\int_{EM}^{EP} dE \sigma(F) F(E)$ by means of trapezoidal
rule.

SGN(I) cross-section

ES(I) energy mesh-points [eV]

Called by SINT.

FUNCTION BCM (EI, EO, WORK)

BCM gives the cosine of scattering angle in the center of mass system.

EI the energy of neutrons before scattering.

EO the energy of neutron after scattering.

Called by LMI.

FUNCTION ICSOP (E, ENG)

ICSOP gives the group number for energy E.

Called by REMO, LMI, FXINT, MASSIN.

FUNCTION (ANG, XA, A, XM)

ANG gives the cosine of scattering angle in labor system.

XA cosine of scattering angle in center of mass system.

A atomic mass.

XM= $(A+1)^2/2A$.

Called by ANINT.

FUNCTION PTL (N, X)

PTL gives the value $P_N(X)$

Called by ANINT.

SUBROUTINE ZWIN (K, EP, EM, ES, J2, J1)

ES(I), I=1, k energy mesh-points in [eV].

ZWIN gives the values J2, J1. These have the meaning: ES(J2), ..

ES(J1) cover the energy interval EP, EM.

Called by REMP, LMI, FXINT, WAHRS.

SUBROUTINE MASSIN (WORK, ENG)

MASSIN calculates the mass dependent quantities and places them in the corresponding words of WORK /see Table 1./

Called by REMO.

SUBROUTINE INTEN (E, SG, SGNC, EA, ICOS)

INTEN interpolates the angular distribution in energy.

E is the energy for which the angular distribution is to be interpolated.

Called by LMI.

SUBROUTINE SMORN (PI, N, M1, IMAX, WORK)

SMORN normalizes the calculated elastic transfer matrixes in order to get the transfer probabilities (8.3.)

PI(NM1, IMAX) the elastic transfer matrix.

Called by ISOFAL.

SUBROUTINE SEARCH (NFG, NFI, WORK, IWORK, I)

SEARCH organizes the output of $\langle \sigma_t \rangle$, $\langle \sigma_s \rangle$, $\langle \mu_s \rangle$ for the group.

Called by REMO.

SUBROUTINE SUCHM (NFG, NFI, IWORK, WORK, MOM, I, ENG)

SUCHM organizes the output of $\langle \sigma (E_i \rightarrow h) \rangle$ for the group I.

Called by REMO

FUNCTION IWO (N, IWORK)

IWO gives the value IWORK(N).

SUBROUTINE IWIN (L, N, IWORK)

IWIN performs the statement IWORK(L) = N.

In the future program is intended to generalize to the inelastic case. The subroutines INRE, IMPE, EGRENZ serve for this purpose. For elastic scattering calculation these have no any influence.

SUBROUTINE AKOR (BU, BL, L, A)

It calculates the integral in (8.9) for average cosine correction.

BU and BL are the c. of mass scattering angles corresponding to

μ_2 and μ_1 respectively. L is the order of the moment, A is the mass of the scatterer.

Called by LMI.

Table VIII

Control information in the array WØRK

Word	Content	Comment
1	Adress of EA/1/	Energy mesh points for angular distribution
2	Adress of XL/1/	Angle mesh points
3	Adress of SGNC/1/	Angular distributions
4	I	The number of just calculated group
5	Adress of SGN/1/	Elastic scattering or total cross-section
6	Adress of AMU/1/	Retrieved average cosine
7	Adress of ES/1/	Energy mesh points for cross-sections
8	ISEL	
9	Adress of SG/1/	Work-field
10	Adress of TE/1/	Work-field
11	Adress of PER/1/	Work-field
12	Adress of SP/1/	Field for the storage of results
13	Adress of BUF/1/	Buffer-field for output
14	NG	
15	NSP	
16	NLA	
17	NLE	
18	IMAX	Maximum group change+1
19	NJM	
20	NUJM	
21	ICOS	Number of angle mesh points /now it is 21/
22	Address of the first free word of SP	
23	KJ	Control number of the anisotropy
24	NJ	2 NJ+1 number of mesh points to be required for angular integration
25	NDAT	Number of cross section points
26	NIV	Number of energy points for angular distribution
27	ERR	Error for integration procedure
28	A	
29	$((A+1)/(A-1))^2$	

Table VIII (cont.)

ü

Word	Content	Comment
30	$(A+1)^2/2A$	
31	$\log \left(\left(\frac{A+1}{A-1} \right)^2 \right)$	
32	$Q = -0.1$	
33	Address of T/1/	8-bytes work-field
34	Address of PL/1/	8-bytes work-field

9. The calculation of average group cross sections for infinite dilution and of energy resonance self shielding factors from energy point wise data in the resonance region. Module 3.

9.1 The flux weighted energy resonance self shielding factors are calculated for capture, fission, elastic scattering

$$k_f^{x,g}(\sigma_0) = \frac{k \tilde{\sigma}_{x,g}(\sigma_0)}{k \tilde{\sigma}_{x,g}(\sigma_0 \rightarrow \infty)} \quad (9.1)$$

- k = isotope
- x = neutron reaction (n,x)
- (n,e) = elastic scattering, (n,f) = fission,
- (n,c) = total absorption - fission.
- g = energy group
- σ_0 = back ground cross section in barns.

$$k \tilde{\sigma}_{x,g}(\sigma_0) = \frac{\left\langle \frac{k \tilde{\sigma}_x(E) \cdot F(E)}{k \tilde{\sigma}_t(E) + \sigma_0} \right\rangle}{\left\langle \frac{F(E)}{k \tilde{\sigma}_t(E) + \sigma_0} \right\rangle} \quad *$$

(9.2)

is the microscopic effective cross section.

F(E) collision density

$\sigma_t(E)$ is the microscopic total cross section.

Besides that, a flux weighted self shielding factor of the following type is calculated:

$$* \quad \langle y(E) \rangle = \frac{1}{\Delta E_g} \int_{\Delta E_g} y(E) dE$$

$${}_1f_{e,g}(\tilde{\sigma}_o) = \frac{\{ {}^k\tilde{\sigma}_e, {}^k\bar{\mu}_e \}_g(\tilde{\sigma}_o)}{\{ {}^k\tilde{\sigma}_e, {}^k\bar{\mu}_e \}_g(\tilde{\sigma}_o \rightarrow \infty)} \quad (9.3)$$

with

$$\{ {}^k\tilde{\sigma}_e, {}^k\bar{\mu}_e \}_g(\tilde{\sigma}_o) = \frac{\left\langle \frac{{}^k\tilde{\sigma}_e(E) \cdot {}^k\bar{\mu}_e(E) \cdot F(E)}{{}^k\tilde{\sigma}_t(E) + \tilde{\sigma}_o} \right\rangle}{\left\langle \frac{F(E)}{{}^k\tilde{\sigma}_t(E) + \tilde{\sigma}_o} \right\rangle} \quad (9.4)$$

$\bar{\mu}_e(E)$ = average cosine for elastic scattering

The current weighted resonance self shielding factors are calculated for elastic scattering and for the total neutron cross section.

$${}_1f_{t,g}^k(\tilde{\sigma}_o) = \frac{{}_1\tilde{\sigma}_{t,g}^k(\tilde{\sigma}_o)}{{}_1\tilde{\sigma}_{t,g}^k(\tilde{\sigma}_o \rightarrow \infty)} \quad (9.5)$$

$${}_1f_{e,g}^k(\tilde{\sigma}_o) = \frac{{}_1\{ {}^k\tilde{\sigma}_e, {}^k\bar{\mu}_e \}_g(\tilde{\sigma}_o)}{{}_1\{ {}^k\tilde{\sigma}_e, {}^k\bar{\mu}_e \}_g(\tilde{\sigma}_o \rightarrow \infty)} \quad (9.6)$$

with

$${}^k_1\tilde{\sigma}_{t,g}(\tilde{\sigma}_0) = \frac{\left\langle \frac{{}^k\tilde{\sigma}_t(E) \cdot F(E)}{\{ {}^k\tilde{\sigma}_t(E) + \tilde{\sigma}_0 \}^2} \right\rangle}{\left\langle \frac{F(E)}{\{ {}^k\tilde{\sigma}_t(E) + \tilde{\sigma}_0 \}^2} \right\rangle} \quad (9.7)$$

$$\left\{ {}^k\tilde{\sigma}_e, {}^k\bar{\mu}_e \right\}_g(\tilde{\sigma}_0) = \frac{\left\langle \frac{{}^k\tilde{\sigma}_e(E) {}^k\bar{\mu}_e(E) F(E)}{\{ {}^k\tilde{\sigma}_t(E) + \tilde{\sigma}_0 \}^2} \right\rangle}{\left\langle \frac{F(E)}{\{ {}^k\tilde{\sigma}_t(E) + \tilde{\sigma}_0 \}^2} \right\rangle} \quad (9.8)$$

The average group cross sections for infinite dilution are defined:

$${}^k\tilde{\sigma}_{x,g}^\infty = \lim_{\tilde{\sigma}_0 \rightarrow \infty} {}^k\tilde{\sigma}_{x,g}(\tilde{\sigma}_0) = \lim_{\tilde{\sigma}_0 \rightarrow \infty} {}^k_1\tilde{\sigma}_{x,g}(\tilde{\sigma}_0) \quad (9.9)$$

or

$${}^k\tilde{\sigma}_{x,g}^\infty = \frac{\langle {}^k\tilde{\sigma}_x(E) \cdot F(E) \rangle}{\langle F(E) \rangle} \quad (9.10)$$

for capture, fission, elastic scattering and the total neutron reaction, and

$$\left\{ \tilde{\sigma}_e^{k-} / \mu_e^{k-} \right\}_g^\infty = \frac{\langle \tilde{\sigma}_e^{k-}(E), \mu_e^{k-}(E), F(E) \rangle}{\langle F(E) \rangle} \quad (9.11)$$

F(E) is a slowly varying weighting-function, normally a collision density.

9.2 The microscopic cross sections are taken from energy point wise data on the KEDAK-library. The integrations are done by a trapezoidal rule. As integration points the energy points of the KEDAK-tables for $\sigma_t(E)$ and the group boundaries are used. In the averages (9.8) and (9.11) the energy points of the tables for $\sigma_e(E)$ and $\bar{\mu}_e(E)$ are used. If they are tabulated at different energies, they are linearly interpolated. The weighting function is either interpolated from tables or calculated by a function.

9.3 The calculations are performed in the subroutine

FSTRUK (MI, SIGO, NE, ENG, NFE, REFE, EFE, XINTE, ZA, XN, SE, XII,
ZB, XI, MDIM, MDIMP, EMU, XMU, NER, NERP, SN, EN)

The following parameters must be defined:

MI	: number of σ_0 -values.
SIGO	: one-dimensional field containing the σ_0 -values [barn].
NE	: number of group boundaries.
ENG	: one-dimensional field containing the group boundaries [eV].
NFE	: number of energy points of the weighting spectrum.
REFE	: one-dimensional field containing the energy points of the weighting spectrum [eV].

EFE : one-dimensional field containing the weighting function at the energy points of field REFE.

MDIM : length of working fields, should be larger than the number of $\bar{\mu}_e$ -values on KEDAK within one energy group.

NER : length of working fields, should be larger than the maximum number of values for σ_a , σ_e , σ_f on KEDAK within one energy group.

Working fields:

XINTE (NE,4), ZA (MI, NE, 5), XN (NI, NE, 4), SE (MI, NE, 6), XII (NE), ZB (MI, NE), XI (NE), EMU (MDIM), XMU (MDIM), SN (NER, 3), EN (NER, 3).

The following parameters are calculated in the subroutine:

MDIMP : o, if MDIM is large enough .
n, number, by which MDIM should be increased.

NERP : o, if NER is large enough .
n, number, by which NER should be increased.

10. The calculation of group constants in the thermal group, Module 10.

10.1 There is assumed, that the thermal energy range is covered by one energy group. For the cross sections except the elastic scattering cross section a $1/v$ -law, and for the weighting function a Maxwell - distribution is assumed. So the thermal group cross sections are defined as

$${}^k \bar{\sigma}_{x,th} = \frac{\sqrt{k}}{2} \cdot {}^k \sigma_x (0,0253 \text{ eV}) \quad (10.1.)$$

- k = isotope
- x = neutron reaction (n,x) except (n,n) process
- th = thermal

The thermal group cross sections are calculated by formula (10.1.) for all types x of neutron reactions. Furthermore the following group constants are calculated :

- the average elastic scattering cross section

$${}^k \bar{\sigma}_{n,th} = {}^k \bar{\sigma}_n (0,0253 \text{ eV}) \quad (10.2.)$$

- the average cosine of the elastic scattering

$${}^k \bar{\mu}_{e,th} = {}^k \mu_e (0,0253 \text{ eV}) \quad (10.3.)$$

- the average number of secondary neutrons per fission

$${}^k \bar{\nu}_{th} = {}^k \nu (0,0253 \text{ eV}) \quad (10.4.)$$

- the average α - value

$${}^k \bar{\alpha}_{th} = \frac{{}^k \sigma_{\gamma,th}}{{}^k \sigma_{f,th}} \quad (10.5.)$$

- γ = radiative capture
- f = fission

- the average η - value

$${}^k\bar{\eta}_{th} = \frac{{}^k\eta_{th}}{1 + {}^k\alpha_{th}} \quad (10.6.)$$

10.2. The following subroutine is necessary :

THERM (NE, NTY, TYP)

The following parameters are defined by the control program.

NE : number of energy group limits.

NTY : number of cross section types.

TYP (NTY) : one-dimensional field containing the cross section types.

11. The calculation of the average $1/v$ values.
Module 8 .

11.1 The average $1/v$ values are defined as

$$\left(\frac{1}{v}\right)_g = \frac{\int_{(g)} \frac{1}{v} \cdot F(E) dE}{\int_{(g)} F(E) dE} \quad (11.1)$$

11.2 The numerical integration is performed by the following formula:

$$\left(\frac{1}{v}\right)_g = C \cdot \frac{\sum_{i=1}^n \left\{ \frac{F_i}{\sqrt{E_i}} + \frac{F_{i+1}}{\sqrt{E_{i+1}}} \right\} \cdot \{ E_{i+1} - E_i \}}{\sum_{i=1}^n \{ F_i + F_{i+1} \} \cdot \{ E_{i+1} - E_i \}} \quad (11.2)$$

- i : index for the integration points
- E_i : energy in [eV]
- F_i : weighting function
- n : number of integration points in group g.
- c : constant = $7,229286 \cdot 10^{-7}$

If the weighting function is given pointwise, as integration points the energy points of the weighting function are used. Therefore the user must take care of the number of energy points.

If the weighting function is given in an analytic form, ten integration points per groups, equidistant in energy, are used.

In the thermal group the $1/v$ value at 0,0253 is taken.

11.3 The following subroutine is used:

EDV (MM,A, NEF, XS, G, E, V, Y, NES, ES, F)

The following parameters must be defined:

MM : number of group boundaries
A(MM) : group boundaries in [eV].
NEF : number of energy points of the weighting function.
XS(NEF): energy points of the weighting function in [eV]
G(NEF) : weighting function at the energy points XS (NEF)
NES : max (10 x number of groups boundaries, number of energy
points of the weighting function)

E (MM), V(MM), Y(MM), ES(NES), F(NES)
one-dimensional working fields.

12. The calculation of the fission spectrum. Module 7.

12.1 For the calculation of fission spectra the Watt formula is used.

$$\chi(E) = \frac{\exp\left\{-\frac{E_f}{T}\right\}}{\sqrt{\pi \cdot E_f \cdot T}} \cdot \exp\left\{-\frac{E}{T}\right\} \cdot \sinh\left\{\frac{2}{T} \sqrt{E \cdot E_f}\right\} \quad (12.1)$$

E energy in MeV.

E_f fragment kinetic energy per nucleon in MeV.

T Watt fragment nuclear "temperature" in MeV.

The fission spectrum in energy group g is defined as

$$\chi_g = \int_{(g)} \chi(E) dE \quad (12.2)$$

E_g upper group limit in MeV.

E_{g+1} lower group limit in MeV.

Using formula (12.1) one gets

$$\begin{aligned} \chi_g = & \frac{1}{2\sqrt{\pi}} \sqrt{\frac{T}{E_f}} \left\{ \exp(-x_{g+1}^2) - \exp(-x_g^2) - \exp(-y_{g+1}^2) \right. \\ & \left. + \exp(-y_g^2) \right\} - \frac{1}{2} \left\{ \operatorname{erf}(x_{g+1}) - \operatorname{erf}(x_g) \right. \\ & \left. + \operatorname{erf}(y_{g+1}) - \operatorname{erf}(y_g) \right\} \quad (12.3) \end{aligned}$$

$$x_g = \sqrt{\frac{E_g}{T}} - \sqrt{\frac{E_f}{T}} \quad ; \quad x_{g+1} = \sqrt{\frac{E_{g+1}}{T}} - \sqrt{\frac{E_f}{T}}$$

$$y_g = \sqrt{\frac{E_g}{T}} + \sqrt{\frac{E_f}{T}} \quad ; \quad y_{g+1} = \sqrt{\frac{E_{g+1}}{T}} + \sqrt{\frac{E_f}{T}}$$

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$$

12.2 The following subroutine is necessary:

```
SUBROUTINE SPALT (NE, ENG, X)
```

The following parameters are defined by the control program.

NE number of energy group limits,

ENG(NE) one-dimensional field containing the energy group limits
 in [e V].

The following parameter is defined by the subroutine SPALT.

X(NE) one-dimensional field containing the fission spectrum.

Appendix I

A modified version of the formalism of Froelich /1/, /2/, for the calculation of energy resonance self shielding factors in the range of "statistical resonances".*

To take into account overlapping of resonances of the same series, Froelich makes the following assumption /2/, (formula 2.41a):

$${}^s\tilde{\sigma}_{p,eff} + \sum_{r \neq r'} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'} \gg \sum_{r' \neq r} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'} \quad (1)$$

This condition is not fulfilled in the case of strong overlapping, if the resonances are higher than the background cross section ${}^s\sigma_{p,eff}$. Therefore a weaker condition will be used, namely

$${}^s\tilde{\sigma}_{p,eff} + \langle {}^s\tilde{\sigma}_c \rangle + \sum_{r \neq r'} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'} \gg \left| \sum_{r' \neq r} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'} - \langle {}^s\tilde{\sigma}_c \rangle \right| \quad (2)$$

Expression 2.40 in /2/ can be modified to

$$\sum_r \left\langle \frac{{}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'}}{\left[{}^s\tilde{\sigma}_{p,eff} + {}^s\tilde{\sigma}_c + \sum_{r \neq r'} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'} \right] \left[1 - \frac{\sum_{r' \neq r} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'} - \langle {}^s\tilde{\sigma}_c \rangle}{{}^s\tilde{\sigma}_{p,eff} + {}^s\tilde{\sigma}_c + \sum_{r \neq r'} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'}} \right]} \right\rangle \quad (3)$$

Using condition (2), expanding the second term of the denominator of (3) and neglecting quadratic terms in

$$\sum_{r' \neq r} {}^s\tilde{\sigma}_{oc} \cdot {}^s\psi_{r'} - \langle {}^s\tilde{\sigma}_c \rangle \quad , \text{ one gets}$$

* in this appendix the same symbols as in /2/ are used.

the following expression:

$$\begin{aligned}
 & \left\langle \frac{s_{r\sigma_{ox}}}{s_{r\sigma_{oc}}} \left\langle \frac{s_{r\psi}}{\frac{\langle \tilde{\sigma}_t \rangle}{s_{r\sigma_{oc}}} + s_{r\psi}} \right\rangle \right\rangle + \langle s_{\tilde{\sigma}_c} \rangle \left\langle \frac{s_{r\sigma_{ox}}}{\{s_{r\sigma_{oc}}\}^2} \left\langle \frac{s_{r\psi}}{\left\{ \frac{s_{\tilde{\sigma}_c}}{s_{r\sigma_{oc}}} + s_{r\psi} \right\}^2} \right\rangle \right\rangle \\
 & - \left\langle \left\langle \frac{s_{r\sigma_{ox}} \cdot s_{r\sigma_{oc}}}{\{s_{r\sigma_{oc}}\}^2} \left\langle \frac{s_{r\psi} \cdot s_{r\psi}}{\left\{ \frac{\langle \tilde{\sigma}_t \rangle}{s_{r\sigma_{oc}}} + s_{r\psi} \right\}^2} \right\rangle \right\rangle \quad (4)
 \end{aligned}$$

with

$$\langle \tilde{\sigma}_t \rangle = \tilde{\sigma}_{p,eff} + \langle s_{\tilde{\sigma}_c} \rangle$$

The first and the third term of (4) are the same as in expression (2.42) in /2/, when $s_{\sigma_{p,eff}}$ is replaced by $\langle \tilde{\sigma}_t \rangle$.

To calculate the overlapping effect, the assumption is made:

$$\langle \tilde{\sigma}_t \rangle \gg \frac{s_{\sigma_{ox}}}{s_{\sigma_{oc}}} \cdot \frac{s_{\psi}}{r} \quad (5)$$

This condition corresponds to the condition (2.50) and (2.51) in /2/.

With this condition, for the third term of (4) one gets the identically same result as for the second term of (2.42) in /2/.

$$\left\langle \left\langle \frac{s_{r\sigma_{ox}} \cdot s_{r\sigma_{oc}}}{\{s_{r\sigma_{oc}}\}^2} \left\langle \frac{s_{r\psi} \cdot s_{r\psi}}{\left\{ \frac{\langle \tilde{\sigma}_t \rangle}{s_{r\sigma_{oc}}} + s_{r\psi} \right\}^2} \right\rangle \right\rangle \approx \frac{s_D \cdot \langle s_{\tilde{\sigma}_x} \rangle \langle s_{\tilde{\sigma}_c} \rangle}{\Delta \sqrt{2\pi} \langle \tilde{\sigma}_t \rangle^2} \cdot \varepsilon \quad (6)$$

From condition (5), the following approximation for the denominator of the second term of (4) can be deduced:

$$\left\{ \frac{\langle \tilde{\theta}_t \rangle}{\frac{s}{r} \tilde{\theta}_{oc}} + \frac{s\psi}{r} \right\}^2 \cong \frac{\langle \tilde{\theta}_t \rangle}{\left\{ \frac{s}{r} \tilde{\theta}_{oc} \right\}^2} \cdot \left\{ \langle \tilde{\theta}_t \rangle + \frac{s}{r} \tilde{\theta}_{oc} \cdot \frac{s\psi}{r} \right\} \quad (7)$$

Then for the second term of (4) one gets

$$\frac{\langle \frac{s}{r} \tilde{\theta}_c \rangle}{\langle \tilde{\theta}_t \rangle} \cdot \left[\frac{\frac{s}{r} \tilde{\theta}_{ox}}{\frac{s}{r} \tilde{\theta}_{oc}} \left\langle \frac{\frac{s\psi}{r}}{\frac{\langle \tilde{\theta}_t \rangle}{\frac{s}{r} \tilde{\theta}_{oc}} + \frac{s\psi}{r}} \right\rangle \right] \quad (8)$$

Now the first and the second term of (4) can be combined to

$$\left\{ 1 + \frac{\langle \frac{s}{r} \tilde{\theta}_c \rangle}{\langle \tilde{\theta}_t \rangle} \right\} \left[\frac{\frac{s}{r} \tilde{\theta}_{ox}}{\frac{s}{r} \tilde{\theta}_{oc}} \left\langle \frac{\frac{s\psi}{r}}{\frac{\langle \tilde{\theta}_t \rangle}{\frac{s}{r} \tilde{\theta}_{oc}} + \frac{s\psi}{r}} \right\rangle \right] \quad (9)$$

and for this, one gets in the same manner as for the first term of (2.42) in /4/ the following expression.

$$\left\{ 1 + \frac{\langle \frac{s}{r} \tilde{\theta}_c \rangle}{\langle \tilde{\theta}_t \rangle} \right\} \frac{^s \Gamma_x \cdot J(^s \beta, ^s \theta)}{^s D \cdot \cos 2\alpha_l} \quad (10)$$

where now $^s \beta$ is given by

$$^s \beta = \frac{\langle \tilde{\theta}_t \rangle}{\frac{s}{r} \tilde{\theta}_{oc}} \quad (11)$$

So for the effective microscopic cross section one gets the modified result

$${}^s\tilde{\sigma}_{x,g}(E) = {}^s\tilde{\sigma}_{p,eff} \frac{\left[1 + \frac{\langle {}^s\tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma \cdot J({}^s\beta, {}^s\theta)}{{}^sD \cdot \cos 2\delta_l} - \frac{{}^sD \cdot \langle {}^s\tilde{\sigma}_x \rangle \cdot \langle {}^s\tilde{\sigma}_c \rangle}{\Delta \sqrt{2\pi} \cdot \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon}{1 - \left[1 + \frac{\langle {}^s\tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma \cdot J({}^s\beta, {}^s\theta)}{{}^sD} + \frac{{}^sD \cdot \langle {}^s\tilde{\sigma}_c \rangle^2}{\Delta \sqrt{2\pi} \cdot \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon} \quad (12)$$

for capture and fission, and

$${}^s\tilde{\sigma}_{c,g}(E) = {}^s\tilde{\sigma}_{p,eff} \frac{\left[1 + \frac{\langle {}^s\tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma \cdot J({}^s\beta, {}^s\theta)}{{}^sD} - \frac{{}^sD \cdot \langle {}^s\tilde{\sigma}_c \rangle^2}{\Delta \sqrt{2\pi} \cdot \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon}{1 - \left[1 + \frac{\langle {}^s\tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma \cdot J({}^s\beta, {}^s\theta)}{{}^sD} + \frac{{}^sD \cdot \langle {}^s\tilde{\sigma}_c \rangle^2}{\Delta \sqrt{2\pi} \cdot \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon} \quad (13)$$

for the total resonance cross section.

All symbols have the same meaning as in /2/, formula (2.68), except

$${}^s\beta = \frac{\langle \tilde{\sigma}_t \rangle}{{}^s\tilde{\sigma}_{oc}} \quad \text{instead of} \quad \frac{{}^s\tilde{\sigma}_{p,eff}}{{}^s\tilde{\sigma}_{oc}} .$$

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Appendix II

An approximate calculation of current weighted resonance self shielding factors from statistical resonance parameters.

1. Introduction

For the calculation of resonance self shielding factors from statistical resonance parameters a formalism /3/, first developed by R. Froelich in his paper on Doppler theory /2/, was used in the first version of MIGROS, which did not allow the calculation of current weighted self shielding factors. An approximation is proposed, which allows to use the modified formalism for flux weighted self shielding factors of Froelich, given in Appendix I, also for current weighted self shielding factors.

The flux weighted f-factor for the neutrons reaction (n,x) in the energy group g for a resonance isotope may be written as

$$f_{x,g} = \frac{\sum_j \phi_j \cdot \tilde{\sigma}_{x,g}(E_j) \cdot \Delta E_j}{\sum_j \phi_j \cdot \tilde{\sigma}_{x,g}^\infty(E_j) \cdot \Delta E_j} \quad (1)$$

The ΔE_j are energy intervals in the group g, ϕ_j is a weighting function, constant over ΔE_j .

$$\tilde{\sigma}_{x,g}(E) = \frac{\left\langle \frac{\tilde{\sigma}_x(E)}{\tilde{\sigma}_t(E) + \tilde{\sigma}_o} \right\rangle}{\left\langle \frac{1}{\tilde{\sigma}_t(E) + \tilde{\sigma}_o} \right\rangle} \quad (2)$$

is the effective cross section in the energy interval ΔE around E for the neutron reaction (n,x) for the resonance isotope.

$$1) \quad \left\langle y(E) \right\rangle = \frac{1}{\Delta E} \int_{\Delta E} y(E) dE$$

$\tilde{\sigma}_x(E)$ microscopic cross section of the resonance isotope for the neutron reaction (n,x)

$\tilde{\sigma}_t(E)$ microscopic total cross section of the resonance isotope.

$\tilde{\sigma}_0$ total back ground cross section per resonance atom.

$$\tilde{\sigma}_{x,g}^{\infty}(E) = \langle \tilde{\sigma}_x(E) \rangle \quad (3)$$

is the average cross section in ΔE around E. The relation

$$\tilde{\sigma}_{x,g}^{\infty}(E) = \lim_{\tilde{\sigma}_0 \rightarrow \infty} \tilde{\sigma}_{x,g}(E) \quad (4)$$

means that for infinite dilution of the resonance isotope the effective cross section is the same as the average cross section. The interval ΔE is assumed to be so large that many resonances are in it. There is no problem in calculating the average cross section from statistical resonance parameters | 2 |, | 3 |, | 4 |. However, it is not so easy to calculate the effective cross section because of the overlapping of the resonances. The effective cross section is composed of the contributions of the single resonances series s

$$\tilde{\sigma}_{x,g}(E) = \sum_s^s \tilde{\sigma}_{x,g}^s(E) \quad \text{for capture and fission} \quad (5)$$

$$\tilde{\sigma}_{t,g}(E) = \sum_s^s \tilde{\sigma}_{c,g}^s(E) + \tilde{\sigma}_p \quad \text{for the total cross section}$$

$\tilde{\sigma}_{c,g}^s(E)$ is the total resonance cross section of series s
(c = compound, not capture)

$$\tilde{\sigma}_{x,g}^s(E) = \frac{\left\langle \frac{\tilde{\sigma}_x^s(E)}{\tilde{\sigma}_c^s(E) + \sum_{s' \neq s} \tilde{\sigma}_c^{s'}(E) + \tilde{\sigma}_p + \tilde{\sigma}_0} \right\rangle}{\left\langle \frac{1}{\tilde{\sigma}_c^s(E) + \sum_{s' \neq s} \tilde{\sigma}_c^{s'}(E) + \tilde{\sigma}_p + \tilde{\sigma}_0} \right\rangle} \quad (6)$$

${}^s\tilde{\sigma}_c(E)$ is the resonance part of the total cross section of series s.

$\tilde{\sigma}_p$ potential cross section

R. Froelich showed that under the assumption that

$$\frac{\sum_{s' \neq s} [{}^{s'}\tilde{\sigma}_c(E) - \langle {}^{s'}\tilde{\sigma}_c(E) \rangle]}{\tilde{\sigma}_p + \tilde{\sigma}_o + {}^s\tilde{\sigma}_c(E) + \sum_{s' \neq s} {}^{s'}\tilde{\sigma}_c(E)} \ll 1 \quad (7)$$

at the resonances of series s, the effective cross section of series s can be approximated by

$${}^s\tilde{\sigma}_{x,g}(E) = \frac{\left\langle \frac{{}^s\tilde{\sigma}_x(E)}{{}^s\tilde{\sigma}_c(E) + \sum_{s' \neq s} \langle {}^{s'}\tilde{\sigma}_c(E) \rangle + \tilde{\sigma}_p + \tilde{\sigma}_o} \right\rangle}{\left\langle \frac{1}{{}^s\tilde{\sigma}_c(E) + \sum_{s' \neq s} \langle {}^{s'}\tilde{\sigma}_c(E) \rangle + \tilde{\sigma}_p + \tilde{\sigma}_o} \right\rangle} \quad (8)$$

This approximation is correct up to quadratic terms of expression (7).

Introducing

$${}^s\tilde{\sigma}_{p,eff} = \tilde{\sigma}_o + \tilde{\sigma}_p + \sum_{s' \neq s} \langle {}^{s'}\tilde{\sigma}_c(E) \rangle \quad (9)$$

one gets

$${}^s\tilde{\sigma}_{x,g}(E) = \frac{\left\langle \frac{{}^s\tilde{\sigma}_x(E)}{{}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff}} \right\rangle}{\left\langle \frac{1}{{}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff}} \right\rangle} \quad (10)$$

or in a slightly different form

$${}^s\tilde{\sigma}_{x,g}(E) = {}^s\tilde{\sigma}_{p,eff} \cdot \frac{\left\langle \frac{{}^s\tilde{\sigma}_x(E)}{{}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff}} \right\rangle}{1 - \left\langle \frac{{}^s\tilde{\sigma}_c(E)}{{}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff}} \right\rangle} \quad (11)$$

This microscopic cross section of the series s is composed of the contributions of all resonances r of this series in the interval ΔE (if ΔE contains enough resonances, contributions from outside may be neglected).

$${}^s\tilde{\sigma}_{x,g}(E) = {}^s\tilde{\sigma}_{p,eff} \cdot \frac{\sum_r \left\langle \frac{{}^s\tilde{\sigma}_x^r(E)}{{}^s\tilde{\sigma}_c^r(E) + \sum_{r' \neq r} {}^s\tilde{\sigma}_c^{r'}(E) + {}^s\tilde{\sigma}_{p,eff}} \right\rangle}{1 - \sum_r \left\langle \frac{{}^s\tilde{\sigma}_c^r(E)}{{}^s\tilde{\sigma}_c^r(E) + \sum_{r' \neq r} {}^s\tilde{\sigma}_c^{r'}(E) + {}^s\tilde{\sigma}_{p,eff}} \right\rangle} \quad (12)$$

Replacing the sum of all resonances by the statistical average times the average number of resonances in the interval ΔE , the modified formalism of Froelich, given in appendix I, approximates expression (11) by

$${}^s\tilde{\sigma}_{x,g}(E) = {}^s\tilde{\sigma}_{p,eff} \frac{\left[1 + \frac{\langle \tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma_x \cdot J({}^s\beta, {}^s\theta)}{{}^sD \cdot \cos 2\delta_t} - \frac{{}^sD \cdot \langle \tilde{\sigma}_x \rangle \cdot \langle \tilde{\sigma}_c \rangle}{\Delta \sqrt{2\pi} \cdot \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon}{1 - \left[1 + \frac{\langle \tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle} \right] \frac{{}^s\Gamma_x \cdot J({}^s\beta, {}^s\theta)}{{}^sD} + \frac{{}^sD \cdot \langle \tilde{\sigma}_c \rangle^2}{\Delta \sqrt{2\pi} \cdot \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon} \quad (13)$$

for x = fission, capture

$${}^s\tilde{\sigma}_{gg}(E) = {}^s\tilde{\sigma}_{p,eff} \frac{\left[1 + \frac{\langle {}^s\tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle}\right] \frac{{}^s\Gamma \cdot \mathcal{J}({}^s\beta, {}^s\theta)}{{}^s\bar{D}} - \frac{{}^sD \cdot \langle {}^s\tilde{\sigma}_c \rangle^2}{\Delta \sqrt{2\pi} \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon}{1 - \left[1 + \frac{\langle {}^s\tilde{\sigma}_c \rangle}{\langle \tilde{\sigma}_t \rangle}\right] \frac{{}^s\Gamma \cdot \mathcal{J}({}^s\beta, {}^s\theta)}{{}^s\bar{D}} + \frac{{}^sD \cdot \langle {}^s\tilde{\sigma}_c \rangle^2}{\Delta \sqrt{2\pi} \langle \tilde{\sigma}_t \rangle^2} \cdot \epsilon} \quad (14)$$

for the total resonance cross section.

Γ_x width for the neutron reaction (n,X)

Γ total width

$\mathcal{J}(\beta, \theta)$ J-function [5]

$$\theta = \frac{\Gamma}{\Delta}$$

Δ Doppler width [5]

δ_L scattering phase shift

$$\beta = \frac{\langle \tilde{\sigma}_t \rangle}{\tilde{\sigma}_{oc}}$$

$\tilde{\sigma}_{oc}$ total resonance cross section in the peak of the resonance

\bar{D} average level distance

$$\overline{f(\Gamma_n, \Gamma_f)} \equiv \iint f(\Gamma_n, \Gamma_f) \cdot F_n(\Gamma_n) \cdot F_f(\Gamma_f) d\Gamma_n d\Gamma_f$$

$F_n(\Gamma_n)$ distribution of neutron width

$F_f(\Gamma_f)$ distribution of fission width

$$\varepsilon = 2 \cdot \int_0^{\infty} \exp\left\{-\frac{D^2}{2\Delta^2}\right\} \cdot \Omega(D) dD$$

$\Omega(D)$ distribution for level distance.

$$\langle {}^s\tilde{\sigma}_x \rangle = \frac{2\pi}{sD} \cdot \lambda^2 \cdot g \cdot \left[\frac{{}^s\Gamma_n \cdot {}^s\Gamma_x}{s\Gamma} \right]$$

$$\langle {}^s\tilde{\sigma}_c \rangle = \frac{2\pi}{sD} \cdot \lambda^2 \cdot g \cdot \overline{{}^s\Gamma_n} \cdot \cos 2\delta_l$$

$$\langle \tilde{\sigma}_t \rangle = {}^s\tilde{\sigma}_{p,eff} + \langle {}^s\tilde{\sigma}_c \rangle$$

λ neutron wave length

g statistical parameter

The second term in the numerator and the third term in the denominator in (13), (14) are correction terms for overlapping of the resonances. They vanish for separated resonances.

The approximation of R. Froelich for the expression (12) is only valid, if the condition

$${}^s\tilde{\sigma}_{p,eff} + \langle {}^s\tilde{\sigma}_c \rangle + {}^s\tilde{\sigma}_c^r(E) \gg \left| \sum_{r' \neq r} {}^s\tilde{\sigma}_c^{r'} - \langle {}^s\tilde{\sigma}_c \rangle \right| \quad (15)$$

is fulfilled in the range of the resonance r .

2. The extension of the formalism of Froelich to current weighted resonance self shielding factors.

The current weighted total f-factor may be written as

$${}_1f_{t,g} = \frac{\sum_j \phi_j \cdot {}_1\tilde{\sigma}_{t,g}(E_j) \cdot \Delta E_j}{\sum_j \phi_j \cdot \tilde{\sigma}_{t,g}^\infty(E_j) \cdot \Delta E_j} \quad (16)$$

$${}_1\tilde{\sigma}_{t,g}(E) = \frac{\left\langle \frac{\tilde{\sigma}_t(E)}{\{\tilde{\sigma}_t(E) + \tilde{\sigma}_o\}^2} \right\rangle}{\left\langle \frac{1}{\{\tilde{\sigma}_t(E) + \tilde{\sigma}_o\}^2} \right\rangle} \quad (17)$$

is the effective current weighted cross section in the energy interval ΔE around E for the total neutron reaction of the resonance isotope.

The effective total cross section may be composed of several resonance series s

$${}_1\tilde{\sigma}_{t,g}(E) = \sum_s^s {}_1\tilde{\sigma}_{c,g}(E) + \tilde{\sigma}_p \quad (18)$$

width

$${}^s\tilde{\Gamma}_{c,g}(E) = \frac{\left\langle \frac{{}^s\tilde{\Gamma}_c(E)}{\left\{ {}^s\tilde{\Gamma}_c(E) + \sum_{s' \neq s} {}^{s'}\tilde{\Gamma}_c(E) + \tilde{\Gamma}_p + \tilde{\Gamma}_o \right\}^2} \right\rangle}{\left\langle \frac{1}{\left\{ {}^s\tilde{\Gamma}_c(E) + \sum_{s' \neq s} {}^{s'}\tilde{\Gamma}_c(E) + \tilde{\Gamma}_p + \tilde{\Gamma}_o \right\}^2} \right\rangle} \quad (19)$$

The numerator of (19) can also be written as

$$\left\langle \frac{{}^s\tilde{\Gamma}_c(E)}{\left\{ {}^s\tilde{\Gamma}_c(E) + \sum_{s' \neq s} \langle {}^{s'}\tilde{\Gamma}_c(E) \rangle + \tilde{\Gamma}_p + \tilde{\Gamma}_o \right\}^2 \cdot \left\{ 1 + \frac{\sum_{s' \neq s} [{}^{s'}\tilde{\Gamma}_c(E) - \langle {}^{s'}\tilde{\Gamma}_c(E) \rangle]}{\tilde{\Gamma}_p + \tilde{\Gamma}_o + {}^s\tilde{\Gamma}_c(E) + \sum_{s' \neq s} \langle {}^{s'}\tilde{\Gamma}_c(E) \rangle} \right\}^2} \right\rangle$$

Neglecting quadratic terms of

$$\sum_{s' \neq s} \left\{ {}^{s'}\tilde{\Gamma}_c(E) - \langle {}^{s'}\tilde{\Gamma}_c(E) \rangle \right\}$$

in the second bracket, using condition (7) and expression (9),

one gets

$$\left\langle \frac{{}^s\tilde{\Gamma}_c(E)}{\left\{ {}^s\tilde{\Gamma}_c(E) + {}^s\tilde{\Gamma}_{p,eff} \right\}^2} \right\rangle - 2 \cdot \left\langle \frac{{}^s\tilde{\Gamma}_c(E) \cdot \sum_{s' \neq s} \left\{ {}^{s'}\tilde{\Gamma}_c(E) - \langle {}^{s'}\tilde{\Gamma}_c(E) \rangle \right\}}{\left\{ {}^s\tilde{\Gamma}_c(E) + {}^s\tilde{\Gamma}_{p,eff} \right\}^3} \right\rangle$$

Because the resonances, which belong to different series, are not correlated, the second term may be written as

$$2 \cdot \left\langle \frac{{}^s\tilde{\Gamma}_c(E)}{\left\{ {}^s\tilde{\Gamma}_c(E) + {}^s\tilde{\Gamma}_{p,eff} \right\}^3} \right\rangle \cdot \left\langle \frac{\sum_{s' \neq s} \left\{ {}^{s'}\tilde{\Gamma}_c(E) - \langle {}^{s'}\tilde{\Gamma}_c(E) \rangle \right\}}{\left\{ {}^s\tilde{\Gamma}_c(E) + {}^s\tilde{\Gamma}_{p,eff} \right\}^3} \right\rangle$$

The second term is identically zero.

The same can be done for the denominator of (19), so that one gets

$${}^s\tilde{\sigma}_{c,ig}(E) = \frac{\left\langle \frac{{}^s\tilde{\sigma}_c(E)}{\left\{ {}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff} \right\}^2} \right\rangle}{\left\langle \frac{1}{\left\{ {}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff} \right\}^2} \right\rangle} \quad (20)$$

To calculate the current weighted effective cross section (20), one can use the formalism of Froelich described in chapter I, if the following condition is fulfilled:

$$\left| {}^s\tilde{\sigma}_c(E) - \langle {}^s\tilde{\sigma}_c(E) \rangle \right| \ll \sum_{s'} \langle {}^{s'}\tilde{\sigma}_c(E) \rangle + \tilde{\sigma}_p + \tilde{\sigma}_o \quad (21)$$

The numerator of (20) may be written as

$$\left\langle \frac{{}^s\tilde{\sigma}_c(E)}{\left\{ \left[{}^s\tilde{\sigma}_c(E) - \langle {}^s\tilde{\sigma}_c(E) \rangle \right] + \left[{}^s\tilde{\sigma}_{p,eff} + \langle {}^s\tilde{\sigma}_c(E) \rangle \right] \right\}^2} \right\rangle$$

Using condition (21) and neglecting quadratic terms in

$$\left| {}^s\tilde{\sigma}_c(E) - \langle {}^s\tilde{\sigma}_c(E) \rangle \right|$$

one gets

$$\frac{1}{2 \cdot \left\{ {}^s\tilde{\sigma}_{p,eff} + \langle {}^s\tilde{\sigma}_c(E) \rangle \right\}} \cdot \left\langle \frac{{}^s\tilde{\sigma}_c(E)}{{}^s\tilde{\sigma}_c(E) + 0.5 \left\{ {}^s\tilde{\sigma}_{p,eff} - \langle {}^s\tilde{\sigma}_c(E) \rangle \right\}} \right\rangle$$

For the denominator of (20), one gets

$$\frac{1}{2 \cdot \{ {}^s\tilde{\sigma}_{p,eff} + \langle {}^s\tilde{\sigma}_c(E) \rangle \}} \left\langle \frac{1}{{}^s\tilde{\sigma}_c(E) + 0.5 \{ {}^s\tilde{\sigma}_{p,eff} - \langle {}^s\tilde{\sigma}_c(E) \rangle \}} \right\rangle$$

Therefore, under condition (21) the current weighted effective cross section can be approximated by a flux weighted effective cross section of the form (11)

$${}^s\tilde{\sigma}_{c,g}(E) = {}^s\tilde{\sigma}_{p,eff}^* \frac{\left\langle \frac{{}^s\tilde{\sigma}_c(E)}{{}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff}^*} \right\rangle}{1 - \left\langle \frac{{}^s\tilde{\sigma}_c(E)}{{}^s\tilde{\sigma}_c(E) + {}^s\tilde{\sigma}_{p,eff}^*} \right\rangle} \quad (22a)$$

and

$${}^s\tilde{\sigma}_{p,eff}^* = 0.5 \{ {}^s\tilde{\sigma}_{p,eff} - \langle {}^s\tilde{\sigma}_c(E) \rangle \} \quad (22b)$$

Now the formalism of Froelich can be used yielding the result (14), where ${}^s\tilde{\sigma}_{p,eff}$ is replaced by ${}^s\tilde{\sigma}_{p,eff}^*$.

3. Limits of the approximation

The limits of this approximation are given by condition (21) and the modified condition (15), when ${}^s\tilde{\sigma}_{p,eff}$ is replaced by ${}^s\tilde{\sigma}_{p,eff}^*$. Condition (15) then takes the form

$$\frac{{}^s\tilde{\sigma}_{p,eff} + \langle {}^s\tilde{\sigma}_c(E) \rangle + \langle {}^s\tilde{\sigma}_c^r(E) \rangle}{2} \left| \sum_{r \neq r'} {}^s\tilde{\sigma}_c^{r'}(E) - \langle {}^s\tilde{\sigma}_c \rangle \right| \quad (23)$$

Equation (23) is certainly fulfilled for separated resonances, as long as

$${}^s\tilde{\sigma}_{p,eff} > \langle {}^s\tilde{\sigma}_c(E) \rangle \quad (24)$$

Equation (24) can also be written as

$${}^s\tilde{\sigma}_{p,eff}^* > 0. \quad (25)$$

This condition is necessary, because (22a) is not defined for negative ${}^s\tilde{\sigma}_{p,eff}^*$. For strong overlapping, condition (23) is more sensible than condition (15).

In the energy range of "statistical resonances" condition (21) cannot be proved. To get some information on the influence of condition (21) on the approximation (22), some calculations in the region of resolved resonance parameters can be done. The flux weighted and the current weighted effective cross sections (10) and (20) may be calculated for a certain energy interval

exactly. In the resolved resonance region it is not necessary to distinguish the single resonance series. Then one gets for the total resonance cross section in the energy group g

$$\tilde{\sigma}_{c,g}(\tilde{\sigma}_o^*) = \frac{\left\langle \frac{\tilde{\sigma}_c(E)}{\tilde{\sigma}_c(E) + \tilde{\sigma}_{p,eff}^*} \right\rangle}{\left\langle \frac{1}{\tilde{\sigma}_c(E) + \tilde{\sigma}_{p,eff}^*} \right\rangle} \quad (26a)$$

$${}_1\tilde{\sigma}_{c,g}(\tilde{\sigma}_o) = \frac{\left\langle \frac{\tilde{\sigma}_c(E)}{\{\tilde{\sigma}_c(E) + \tilde{\sigma}_{p,eff}\}^2} \right\rangle}{\left\langle \frac{1}{\{\tilde{\sigma}_c(E) + \tilde{\sigma}_{p,eff}\}^2} \right\rangle} \quad (26b)$$

$$\tilde{\sigma}_{p,eff} = \tilde{\sigma}_p + \tilde{\sigma}_o \quad (26c)$$

$$\tilde{\sigma}_{p,eff}^* = 0.5 \cdot \{ \tilde{\sigma}_{p,eff} - \langle \tilde{\sigma}_c(E) \rangle \} \quad (26d)$$

$$\tilde{\sigma}_o^* = \tilde{\sigma}_{p,eff}^* - \tilde{\sigma}_p = 0.5 \{ \tilde{\sigma}_o - (\tilde{\sigma}_p + \langle \tilde{\sigma}_c(E) \rangle) \} \quad (26e)$$

The averages $\langle \rangle$ are performed over the energy group g .

For the corresponding resonance self shielding factors one gets

$$f_{t,g}(\tilde{\sigma}_o^*) = \frac{\tilde{\sigma}_{c,g}(\tilde{\sigma}_o^*) + \tilde{\sigma}_p}{\tilde{\sigma}_{c,g}^\infty + \tilde{\sigma}_p} \quad (27a)$$

and

$$1/f_{t,g}(\tilde{\sigma}_0) = \frac{1/\tilde{\sigma}_{c,g}(\tilde{\sigma}_0) + \tilde{\sigma}_p}{\tilde{\sigma}_{c,g}^\infty + \tilde{\sigma}_p} \quad (27b)$$

neglecting the macroscopic weighting.

In fig. 1 and fig. 2 the relative errors of the approximated effective resonance cross section and the approximated total resonance self shielding factor for U 235 in the energy range from 50 eV to 100 eV for several temperatures are plotted. As can be seen from these figures, the errors of the total resonance self shielding factors are smaller than those of the effective resonance cross sections because of the contribution of the potential cross section. The lowest $\tilde{\sigma}_0$ -value, for which the approximation can be used is given by condition (24). The average resonance cross section for this example is $\langle \tilde{\sigma}_c(E) \rangle = 49,9$ barns; the potential cross section is $\tilde{\sigma}_p = 10,3$ barns. So the lowest possible background cross section is $= 39,6$ barns. This restriction of the approximation to $\tilde{\sigma}_0$ -values greater than 50 barns is unimportant, because the $\tilde{\sigma}_0$ -values that correspond to the concentration of U 235 in reactors are in the range of several hundred barns. From fig. 1 one can see that for $\tilde{\sigma}_0$ -values greater than 500 barns the approximation is better than 1 % for all three temperatures. Since with increasing energy the resonances become smaller and the overlapping becomes more important, one can expect, that in the range of the unresolved resonances, condition (21) is fulfilled better than in the resolved resonance region. Therefore the errors of the approximation should be smaller than in fig. 1.

In fig. 3 and fig. 4 the same results for U 238 in the energy range from 2 keV to 3,3, keV are plotted. In contrast to U 235 the errors of the approximation for the effective resonance cross section (fig. 3) of U 238 are much larger, because of the

higher resonances. On the other hand, the effective resonance cross section U 238 is smaller than the potential cross section, so that the errors in the total resonance self shielding factors are much smaller. In fig. 4 one can see that the errors are smaller than 2 - 3 % in the whole range of $\tilde{\sigma}_0$. The average resonance cross section for this example is $\langle \tilde{\sigma}_c(E) \rangle = 5,9$ barns, the potential cross section is $\tilde{\sigma}_p = 10,7$ barns, so that condition (24) can be fulfilled for all $\tilde{\sigma}_0$ -values.

4. Conclusions

From calculations in the resolved resonance region one can draw the conclusion that the proposed approximation for the current weighted total resonance self shielding factors yields results with relative errors of a few percent. This is sufficient for reactor calculations.

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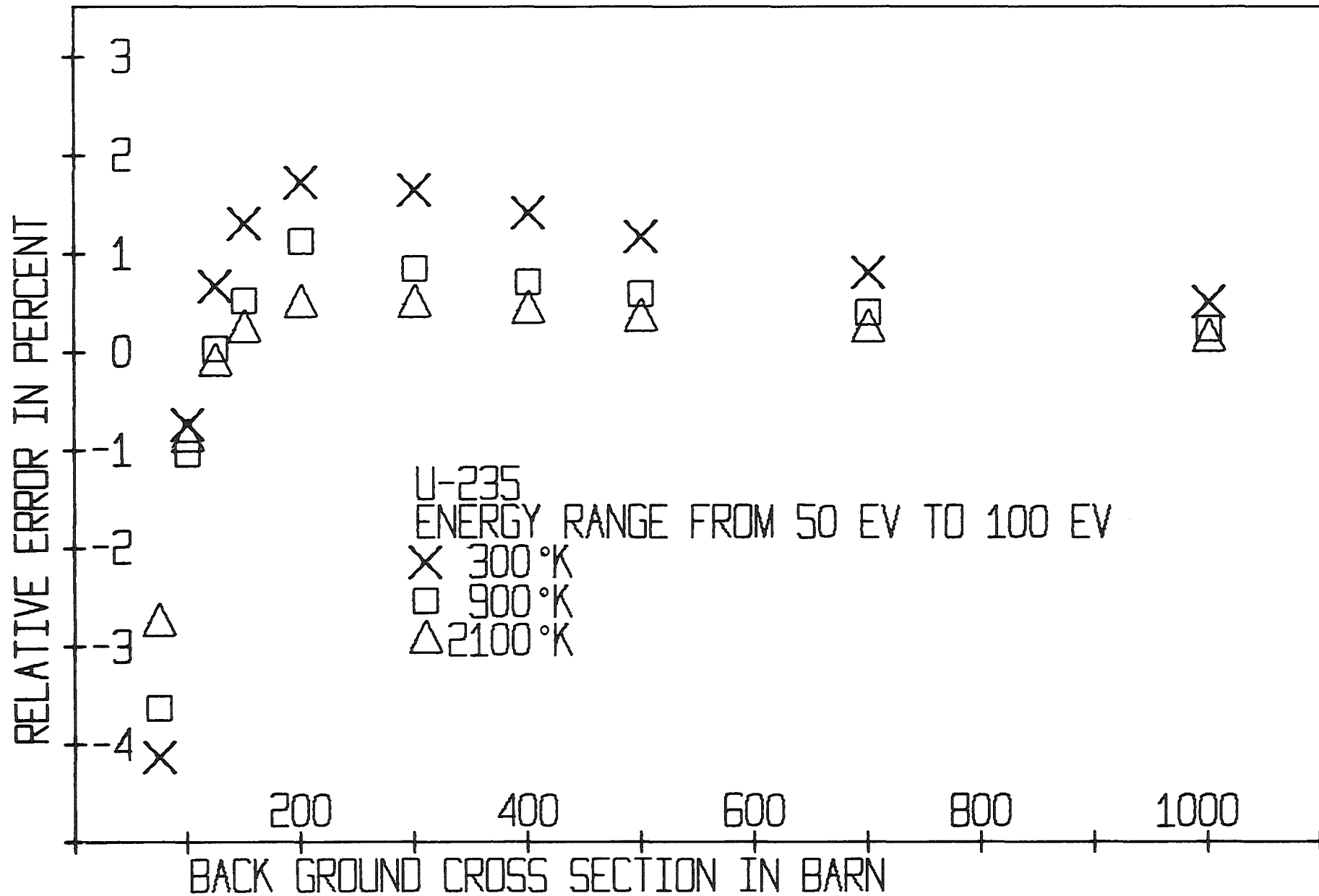


Fig. 1 : $[\sigma_c(\sigma_0^*) - \sigma_c(\sigma_0)] / [0.01 \sigma_c(\sigma_0)]$ over σ_0 .

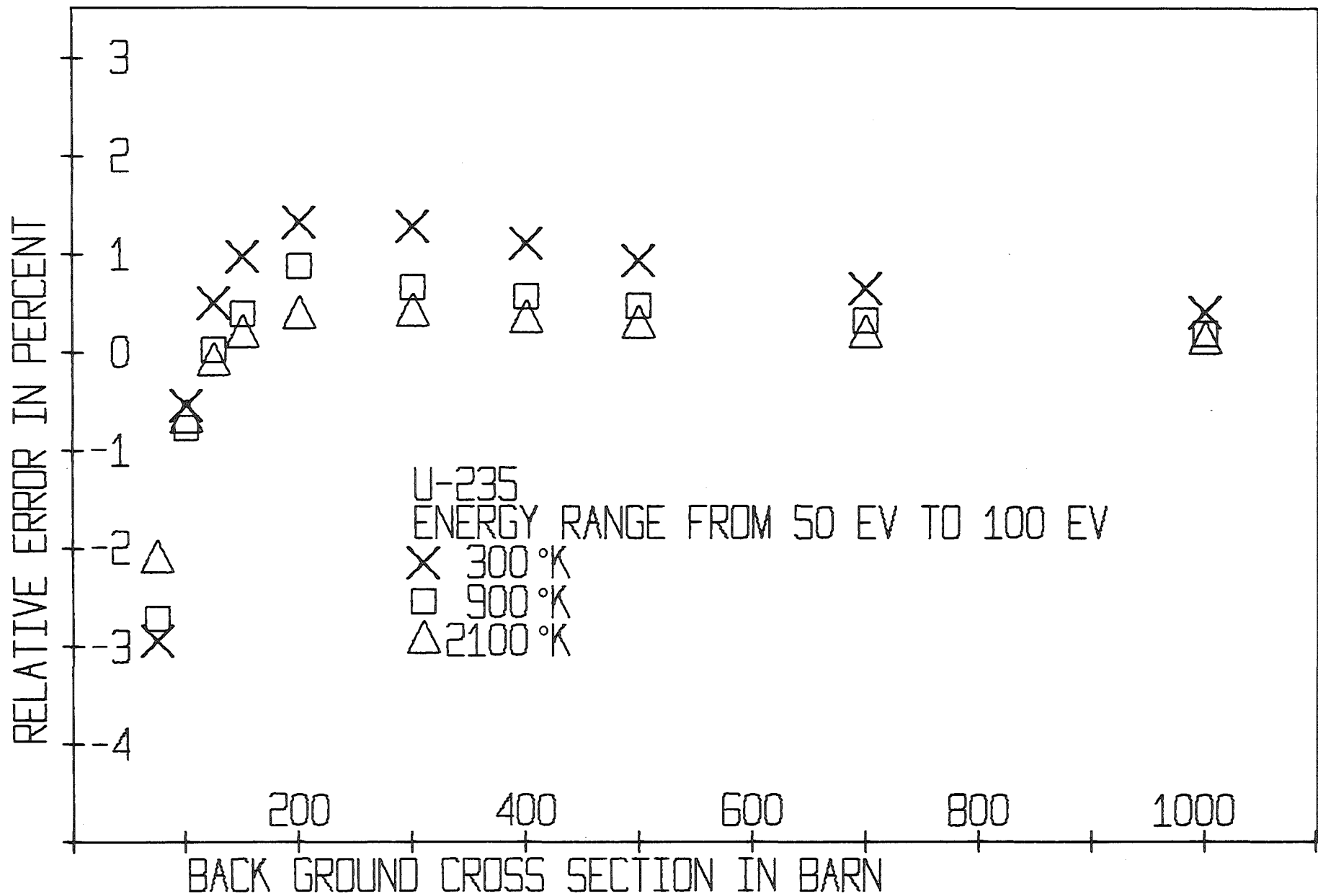


Fig.2: $[f_t(\sigma_0^*) - f_t(\sigma_0)] / [0.01 f_t(\sigma_0)]$ over σ_0 .

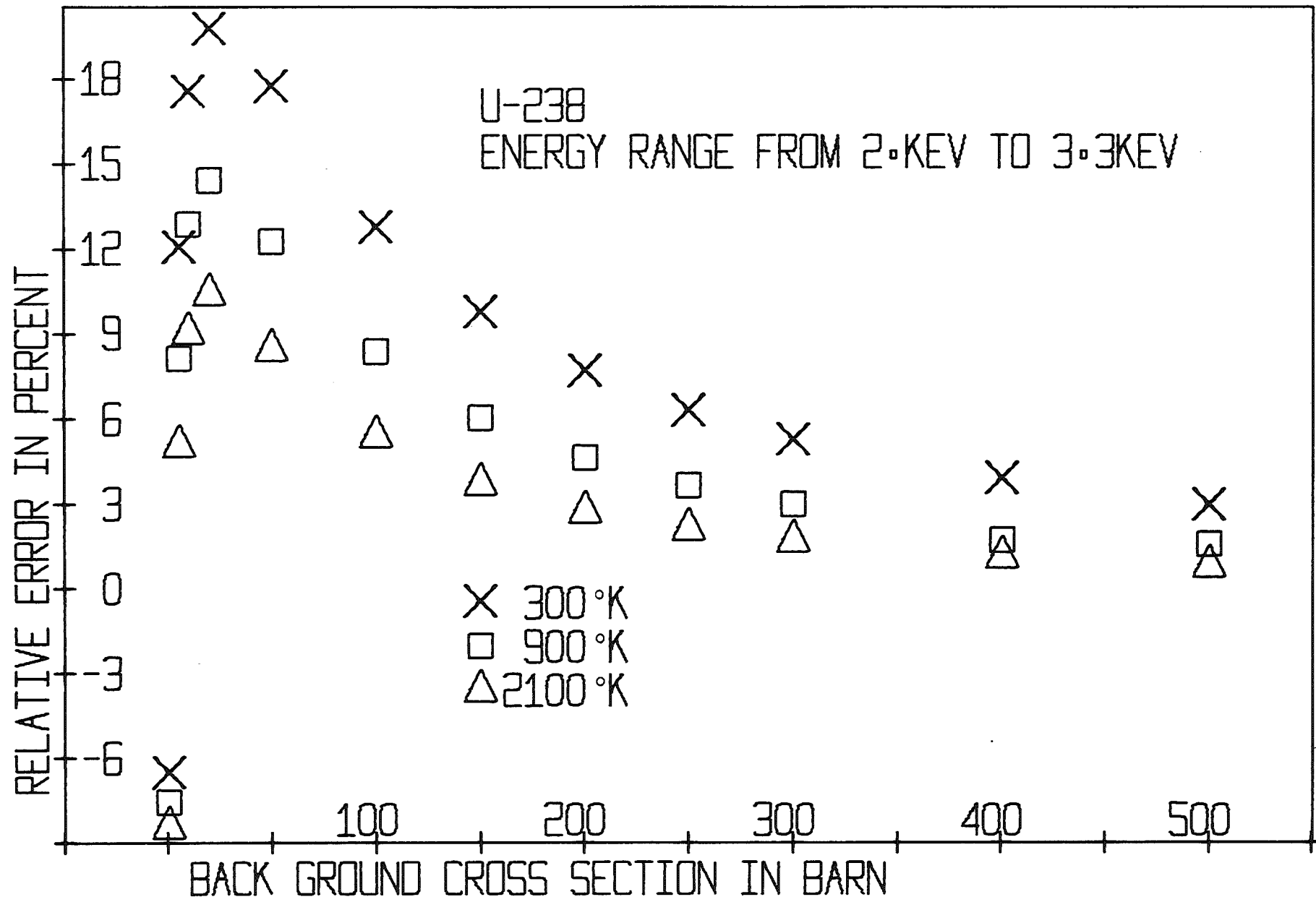


Fig 3.: $[\sigma_c(\sigma_0^*) - \sigma_c(\sigma_0)] / [0.01 \sigma_c(\sigma_0)]$ over σ_0 .

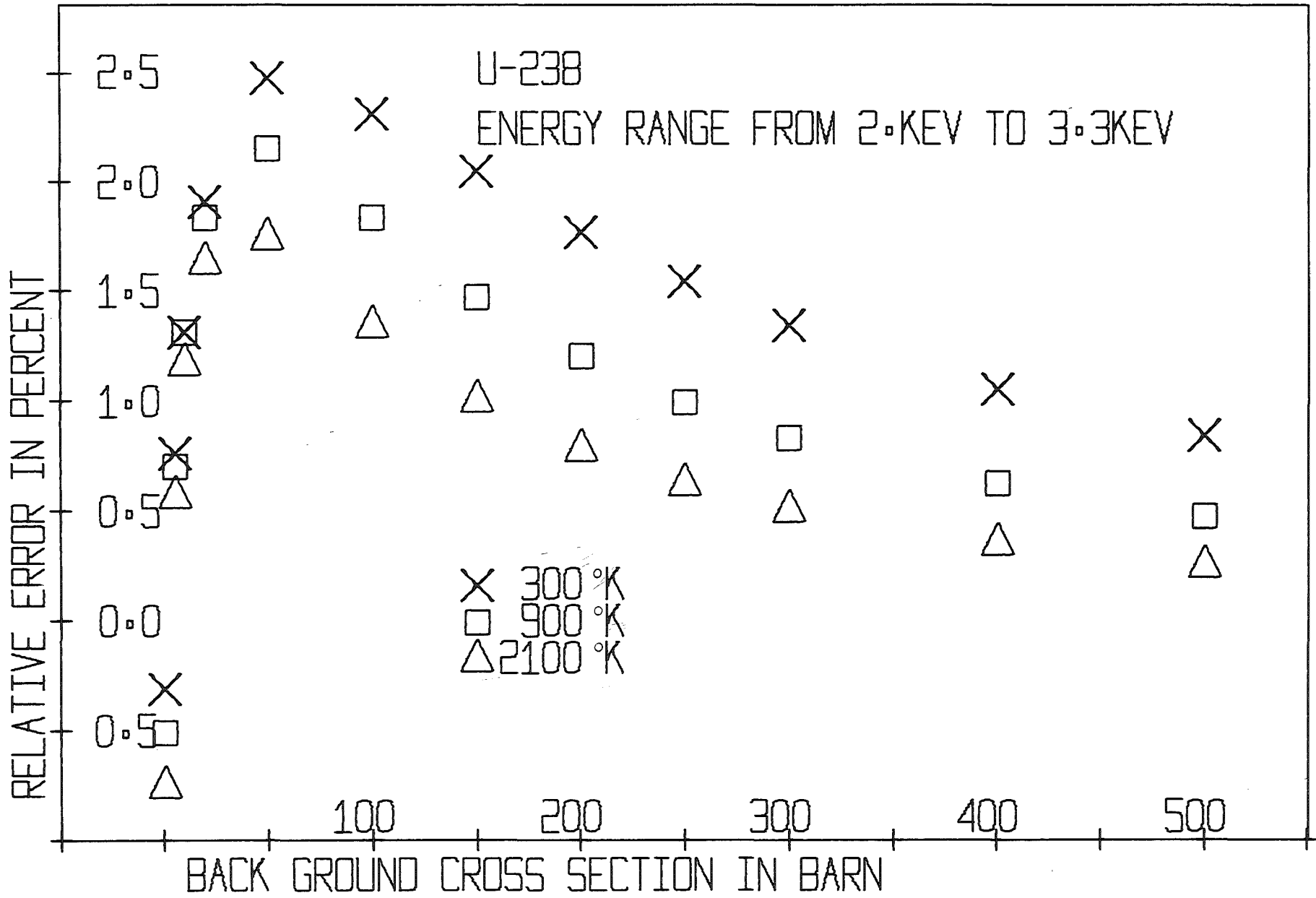


Fig. 4 : $[f_t(\sigma_0^*) - f_t(\sigma_0)] / [0.01 f_t(\sigma_0)]$ over σ_0 .

Appendix III Method of integration /Romberg procedure/

The Romberg method which has been used for numerical integration is briefly the following.

Let us have a function $f(x)$ and the integral

$$I = \int_a^b dx f(x) \quad (1)$$

is to be calculated.

First an equidistant mesh with 2^{n+1} mesh points / including the end points of the interval $| a, b |$ are taken and the sum

$$I_{0n} = \left\{ 0.5 [f(a) + f(b)] + \sum_{k=1}^{2^n-1} f(x_k^n) \right\} \Delta X_n \quad (2)$$

is calculated, where

$$\Delta X_n = \frac{b-a}{2^n} \quad ; \quad x_k^n = a + k \cdot \Delta X_n \quad (k=1, \dots, 2^n-1)$$

Evidently

$$I_{0, n+1} = 0.5 \cdot I_{0n} + \Delta X_{n+1} \sum_{k=1,3,\dots}^{2^{n+1}-1} f(x_k^{n+1}) \quad (3)$$

where

$$\Delta X_{n+1} = 0.5 \Delta X_n, \quad x_k^{n+1} = a + k \Delta X_{n+1} \quad (k=1,3,5,\dots)$$

The second summation is taken over the new mesh points of the higher order mesh.

Let us take the expression

$$I_{mn} = \frac{2^{2m} \cdot I_{m-1, n} - I_{m-1, n-1}}{2^{2m} - 1} \quad (4)$$

It is easy to prove that the deviation of I_{mn} from I is of the order $(\Delta X_n)^m$ i.e. I_{nn} is the best approximation for the given n .

The I_{mn} form a triangular matrix.

$$\begin{array}{cccc}
 I_{00} & & & \\
 I_{01} & I_{11} & & \\
 \cdot & & & \\
 \cdot & & & \\
 \cdot & & & \\
 I_{0n-1} & I_{1n-1} & \dots & I_{n-1n-1} \\
 I_{0n} & I_{1n} & \dots & I_{n-1n} & I_{nn}
 \end{array}$$

The programming of the method is quite straightforward. A row of the above matrix is always to be stored and by means of the formula (3), (4) the next row can be determined. The process is finished if the quantity

$$\left| \frac{I_{nn} - I_{n-1, n-1}}{I_{nn}} \right| < \varepsilon$$

where ε is the required accuracy.

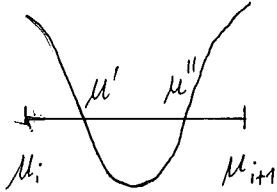
The following should be noted.

1. Only one ε /ERR/ is given. This used for the energy integration. For the angular integration $\varepsilon/10$ is prescribed.
2. Some of the higher moments are very small as far as heavier elements are concerned. In order to avoid a "bad" convergence, the energy integration process for the moment in question is finished when $|I_{nn}| < 10^{-10}$
3. In order to avoid a waste of computing time the integration will be finished if n reaches a maximal value / for energy integration NUJM, for angle integration NJM /. Reaching these values in the course integration an error message will be generated and the calculation is continued. From the error message the user can decide whether the results are acceptable.

Appendix IV The subintervals in the angular integration

The integrand of (8.5) is a product of two functions. In order to save computing time the actual behaviour of these functions. in the integration interval $[\mu_i, \mu_{i+1}]$ should be taken into account.

a) The polynomial $P_n(t(\mu))$ has n zeros. E.g. let us assume that



two zeros are in the interval $[\mu_i, \mu_{i+1}]$ and then the integration interval is divided into three subinterval: $[\mu_i, \mu']$, $[\mu', \mu'']$, $[\mu'', \mu_{i+1}]$.

However sometimes it occurs that the intervals (μ_i, μ') or (μ'', μ_{i+1}) are too small. This can give rise to difficulties in the Romberg integration method. Therefore the width of

the extreme intervals are investigated and if one of them is less than a given limit (now this limit is $|\mu'' - \mu'|/10$) then it will be joined with the next one.

b) The function $\gamma(\mu)$ has often shown a "nasty" character. This "nasty" character can already be detected at the normalization of angular distribution: N passes beyond NJM. If in a energy group a "nasty" character is detected then the integration interval is divided into subintervals defined by the angle point of the angular distribution given in the nuclear data file.

Appendix V. Sample problem1. Job-Control Cards

```

JOB ORIGIN FROM LOCAL DEVICE=RD2      ,020.
//INR017MI JOB (0017,1)1,P6M1B),KRIEG,CLASS=A,TIME=10,REGION=340K
// EXEC PGM=MIGROS2,PARM=170000
//STEPLIB DD DSN=LOAD.NUSYS,VOL=SER=NUSICE,UNIT=3330,DISP=SHR
//FT08F001 DD UNIT=SYSDA,SPACE=(TRK,100)
//FT06F001 DD SYSCUT=A,
// DCB=(RECFM=FBA,LRECL=133,BLKSIZE=931),SPACE=(TRK,25)
//FT01F001 DD UNIT=2314,VOL=SER=NUSYS0,DSN=KNDF,DISP=SHR
//FT03F001 DD UNIT=SYSDA,SPACE=(TRK,20)
//FT10F001 DD UNIT=SYSDA,SPACE=(TRK,200)
//FT05F001 DD *
//

```

2. Input

```

@U 235@ 1
@BLOC@ 4
7 1 18 18 2 14 14 4 14 14 5 5 3 7 25 1 8 25 1 10 26 26
@ENDE@ 17
@D 16@ 0
@BLOC@ 4
5 3 4 4 6 3 3 6 12 12 9 3 3 9 12 12
@ENDE@ 17
@ENDE@ 18

```

On the following pages the output of the results is given. Besides the output of the results for every MIGROS-2-job, a German input-description with an input sample and the print out of the actual input is provided. It has been omitted here.

PROGRAMM KENNZIFFER 1
 PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTABSCHIRMFAKTOREN VON AUFGELOESTEN RESONANZPARAMETERN

MATERIAL		TEMPERATUR	GRUPPE	GRENZEN	
U 235		300.00	18	0.4650E 02	0.1000E 03
SIGMA G		SIGMA N	SIGMA F	SIGMAN1	SIGMAT1
0.23004425E 02		0.12802895E 02	0.28783478E 02	0.12802895E 02	0.64590790E 02
SIGMA 0	FG	FN	FF	FN1	FT1
0.0	0.41953827E 00	0.94163764E 00	0.43603367E 00	0.92344695E 00	0.36389023E 00
0.10000000E 02	0.48846396E 00	0.94783056E 00	0.50455952E 00	0.92906642E 00	0.41785538E 00
0.10000000E 03	0.71442896E 00	0.96913391E 00	0.72282106E 00	0.95165300E 00	0.62696862E 00
0.10000000E 04	0.93719780E 00	0.99258685E 00	0.93785942E 00	0.98616427E 00	0.90250689E 00
0.10000000E 05	0.99270159E 00	0.99910563E 00	0.99279928E 00	0.99824953E 00	0.98808670E 00
0.10000000E 06	0.99930418E 00	0.99990535E 00	0.99939227E 00	0.99981219E 00	0.99876916E 00
0.10000000E 07	0.99991238E 00	0.99999189E 00	0.99992269E 00	0.99998176E 00	0.99987227E 00

MATERIAL		TEMPERATUR	GRUPPE	GRENZEN	
U 235		900.00	18	0.4650E 02	0.1000E 03
SIGMA G		SIGMA N	SIGMA F	SIGMAN1	SIGMAT1
0.22994293E 02		0.12800570E 02	0.28795914E 02	0.12800570E 02	0.64590775E 02
SIGMA 0	FG	FN	FF	FN1	FT1
0.0	0.51503837E 00	0.95064503E 00	0.52545762E 00	0.92998606E 00	0.41661143E 00
0.10000000E 02	0.58327061E 00	0.95673221E 00	0.59176672E 00	0.93649042E 00	0.48082870E 00
0.10000000E 03	0.78867930E 00	0.97656202E 00	0.78884345E 00	0.96075338E 00	0.70187891E 00
0.10000000E 04	0.95903468E 00	0.99510008E 00	0.95769149E 00	0.99061841E 00	0.93352538E 00
0.10000000E 05	0.99543375E 00	0.99943763E 00	0.99527752E 00	0.99889314E 00	0.99233168E 00
0.10000000E 06	0.99956995E 00	0.99993879E 00	0.99959248E 00	0.99988109E 00	0.99922240E 00
0.10000000E 07	0.99995285E 00	0.99999326E 00	0.99993795E 00	0.99998796E 00	0.99991435E 00

MATERIAL		TEMPERATUR	GRUPPE	GRENZEN	
U 235		2100.00	18	0.4650E 02	0.1000E 03
SIGMA G		SIGMA N	SIGMA F	SIGMAN1	SIGMAT1
0.22947205E 02		0.12792241E 02	0.28690109E 02	0.12792241E 02	0.64429550E 02
SIGMA 0	FG	FN	FF	FN1	FT1
0.0	0.60038620E 00	0.95793915E 00	0.60191047E 00	0.93366927E 00	0.47596300E 00
0.10000000E 02	0.66217357E 00	0.96402717E 00	0.66151232E 00	0.94240171E 00	0.54420239E 00
0.10000000E 03	0.83948672E 00	0.98195404E 00	0.83347678E 00	0.96827495E 00	0.75921273E 00
0.10000000E 04	0.97122353E 00	0.99655139E 00	0.96887636E 00	0.99331534E 00	0.95158803E 00
0.10000000E 05	0.99685490E 00	0.99961168E 00	0.99658918E 00	0.99923521E 00	0.99457383E 00
0.10000000E 06	0.99969882E 00	0.99995881E 00	0.99969935E 00	0.99991757E 00	0.99944764E 00
0.10000000E 07	0.99996769E 00	0.99999511E 00	0.99995691E 00	0.99999112E 00	0.99994034E 00

PROGRAMM KENNZIFFER 2
 PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTABSCHIRMFAKTOREN VON STATISTISCHEN RESONANZPARAMETERN

MATERIAL TEMPERATUR GRUPPE GRENZEN
 U 235 300.00 14 0.1000E 04 0.2150E 04

SIGMA G SIGMA N SIGMA F SIGMAN1 SIGMAT1
 0.33934994E 01 0.12306777E 02 0.74021378E 01 0.12306777E 02 0.23102356E 02

SIGMA 0 FG FN FF FN1 FT1

0.0 0.77354968E 00 0.98292530E 00 0.77513117E 00 0.96743643E 00 0.74214095E 00
 0.10000000E 02 0.84917933E 00 0.98786384E 00 0.85143816E 00 0.97653657E 00 0.83018833E 00
 0.10000000E 03 0.95035458E 00 0.99311745E 00 0.94899106E 00 0.99301809E 00 0.96027339E 00
 0.10000000E 04 0.99418187E 00 0.99908733E 00 0.99393761E 00 0.99815631E 00 0.99331576E 00
 0.10000000E 05 0.99942076E 00 0.99990880E 00 0.99939656E 00 0.99981672E 00 0.99934429E 00
 0.10000000E 06 0.99994653E 00 0.99999177E 00 0.99994475E 00 0.99998331E 00 0.99993974E 00
 0.10000000E 07 0.10000000E 01 0.10000000E 01 0.10000000E 01 0.10000000E 01 0.10000000E 01

MATERIAL TEMPERATUR GRUPPE GRENZEN
 U 235 900.00 14 0.1000E 04 0.2150E 04

SIGMA G SIGMA N SIGMA F SIGMAN1 SIGMAT1
 0.33934994E 01 0.12306777E 02 0.74021378E 01 0.12306777E 02 0.23102356E 02

SIGMA 0 FG FN FF FN1 FT1

0.0 0.80794090E 00 0.98593742E 00 0.80581999E 00 0.96952617E 00 0.75254667E 00
 0.10000000E 02 0.87867934E 00 0.99032903E 00 0.87774932E 00 0.97942162E 00 0.84590864E 00
 0.10000000E 03 0.96768624E 00 0.99529058E 00 0.96637660E 00 0.99207813E 00 0.96425921E 00
 0.10000000E 04 0.99666643E 00 0.99947429E 00 0.99652636E 00 0.99893630E 00 0.99615276E 00
 0.10000000E 05 0.99966973E 00 0.99994755E 00 0.99965602E 00 0.99989378E 00 0.99962395E 00
 0.10000000E 06 0.99996924E 00 0.99999493E 00 0.99996859E 00 0.99998987E 00 0.99996555E 00
 0.10000000E 07 0.10000000E 01 0.10000000E 01 0.10000000E 01 0.10000000E 01 0.10000000E 01

MATERIAL TEMPERATUR GRUPPE GRENZEN
 U 235 2100.00 14 0.1000E 04 0.2150E 04

SIGMA G SIGMA N SIGMA F SIGMAN1 SIGMAT1
 0.33934994E 01 0.12306777E 02 0.74021378E 01 0.12306777E 02 0.23102356E 02

SIGMA 0 FG FN FF FN1 FT1

0.0 0.83612275E 00 0.98812699E 00 0.83280033E 00 0.97210991E 00 0.76919627E 00
 0.10000000E 02 0.89682126E 00 0.99082017E 00 0.89438051E 00 0.98215550E 00 0.86282295E 00
 0.10000000E 03 0.97812480E 00 0.99687904E 00 0.97724724E 00 0.99318576E 00 0.97129792E 00
 0.10000000E 04 0.99780726E 00 0.99965549E 00 0.99771470E 00 0.99930209E 00 0.99746072E 00
 0.10000000E 05 0.99978393E 00 0.99996555E 00 0.99977404E 00 0.99993122E 00 0.99975389E 00
 0.10000000E 06 0.99997985E 00 0.99999696E 00 0.99997896E 00 0.99999368E 00 0.99997759E 00
 0.10000000E 07 0.10000000E 01 0.10000000E 01 0.10000000E 01 0.10000000E 01 0.10000000E 01

PROGRAMM KENNZIFFER 7
PROGRAMM ZUR BERECHNUNG DES SPALTSPEKTRUMS

U 235 EINFALLSENERGIE = 0.0 VON GRUPPE 25 BIS 1
CHI 0
0.10703186E-09 0.32318170E-09 0.99962483E-09 0.31642213E-08 0.98985815E-08 0.31139621E-07 0.99362353E-07
0.31199102E-06 0.98313831E-06 0.31392656E-05 0.98584596E-05 0.31055373E-04 0.99062676E-04 0.31036860E-03
0.97276270E-03 0.30689863E-02 0.93898997E-02 0.23881126E-01 0.61028074E-01 0.14057308E 00 0.20236573E 00
0.26991755E 00 0.18339217E 00 0.98322520E-01 0.15635300E-01

PROGRAMM KENNZIFFER 8
PROGRAMM ZUR BERECHNUNG DES 1/V- GRUPPENMITTELWERTES

GRUPPE	1/V
1	0.25519209E-09
2	0.32105696E-09
3	0.40761239E-09
4	0.53069216E-09
5	0.70539552E-09
6	0.96682840E-09
7	0.13673027E-08
8	0.19336588E-08
9	0.27883824E-08
10	0.40953445E-08
11	0.60131669E-08
12	0.88176435E-08
13	0.12950633E-07
14	0.19015328E-07
15	0.27883853E-07
16	0.40953452E-07
17	0.60131697E-07
18	0.88176307E-07
19	0.12950636E-06
20	0.19015351E-06
21	0.27883834E-06
22	0.40953523E-06
23	0.60131714E-06
24	0.88176398E-06
25	0.12950613E-05
26	0.40279137E-05

PROGRAMM KENNZIFFER 10
PROGRAMM ZUR BERECHNUNG DER THERMISCHEN QUERSCHNITTE

U 235	MUEL	26
0.28400000E-02		
U 235	NUE	26
0.24299994E 01		
U 235	SGA	26
0.60553198E 03		
U 235	SGF	26
0.51578394E 03		
U 235	SGC	26
0.89748047E 02		
U 235	SGI	26
0.0		
U 235	SGN	26
0.16599991E 02		
U 235	SG2N	26
0.0		

PROGRAMM KENNZIFFER 3
 PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTABSCHIRMFAKTOREN VON PUNKTWEISE GEGEBENEN WIRKUNGSQUERSCHNITTEN

***WARNING 3. 5 : IN THE ENERGY GROUP 4 THE SELFSHIELDING FACTORS ARE SET TO 1, BECAUSE THE CROSS SECTIONS FOUND ON KEDAK FOR THIS GROUP ALL ARE ZERO

0 16 GRUPPE = 4 GRUPPENGRENZEN 0.14000000E 07 0.25000000E 07

	SIGMA A	SIGMA N	SIGMA N01	SIGMA N1	SIGMA T1
0.0	0.0	0.17582741E 01	0.19087684E 00	0.19087684E 00	0.17582703E 01
SIG0	FA	FN	FN01	FN1	FT1
0.0	0.10000000E 01	0.81050962E 00	0.95290893E 00	0.91999471E 00	0.59960687E 00
0.10000000E 02	0.10000000E 01	0.97468221E 00	0.98883152E 00	0.97875798E 00	0.95077115E 00
0.10000000E 03	0.10000000E 01	0.99691004E 00	0.99860436E 00	0.99722087E 00	0.99385238E 00
0.10000000E 04	0.10000000E 01	0.99968928E 00	0.99986470E 00	0.99971133E 00	0.99936974E 00
0.10000000E 05	0.10000000E 01	0.99997497E 00	0.99999422E 00	0.99997669E 00	0.99994141E 00
0.10000000E 06	0.10000000E 01	0.10000029E 01	0.10000010E 01	0.10000048E 01	0.10000029E 01
0.10000000E 07	0.10000000E 01	0.10000000E 01	0.10000000E 01	0.10000000E 01	0.10000000E 01

PROGRAMM KENNZIFFER 6
 PROGRAMM ZUR BERECHNUNG ELASTISCHER STREUMATRIZEN

***WARNING NDF. 2 : THE FURTHER NAME 0.15830016E 08 IS GREATER THAN THE GREATEST FURTHER NAME INCLUDED IN THE KEDAK LIBRARY FOR 0 16 SGNC

***WARNING NDF. 2 : THE FURTHER NAME 0.15830016E 08 IS GREATER THAN THE GREATEST FURTHER NAME INCLUDED IN THE KEDAK LIBRARY FOR 0 16 SGNC

***MESSAGE 6.1 : WARNING NDF. 2 MAY BE IGNORED

PROGRAMM KENNZIFFER 6
 PROGRAMM ZUR BERECHNUNG ELASTISCHER STREUMATRIZEN

***WARNING NDF. 2 : THE FURTHER NAME 0.15830016E 08 IS GREATER THAN THE GREATEST FURTHER NAME INCLUDED IN THE KEDAK LIBRARY FOR 0 16 SGNC

***WARNING NDF. 2 : THE FURTHER NAME 0.15830016E 08 IS GREATER THAN THE GREATEST FURTHER NAME INCLUDED IN THE KEDAK LIBRARY FOR 0 16 SGNC

***MESSAGE 6.1 : WARNING NDF. 2 MAY BE IGNORED

SGNC0 ELASTISCHE STREUMATRIX 0. ORDNUNG FUER 0 16
3. GRUPPE
0.88370
0.11629

SGNC1 ELASTISCHE STREUMATRIX 1. ORDNUNG FUER 0 16
3. GRUPPE
0.29687
-0.02758

SGNC2 ELASTISCHE STREUMATRIX 2. ORDNUNG FUER 0 16
3. GRUPPE
0.19188
0.00079

SGNC3 ELASTISCHE STREUMATRIX 3. ORDNUNG FUER 0 16
3. GRUPPE
0.08567
-0.00520

SGNC4 ELASTISCHE STREUMATRIX 4. ORDNUNG FUER 0 16
3. GRUPPE
0.01674
-0.00050

SGNC5 ELASTISCHE STREUMATRIX 5. ORDNUNG FUER 0 16
3. GRUPPE
-0.00131
0.00036

ERLAEUTERUNG
L-TE ORDNUNG, G-TE GRUPPE, I-TE ZEILE: MATRIXELEMENT L-TER ORDNUNG
FUER STREUUNG AUS DER G-TEN GRUPPE IN DIE (G+I-1)-TE GRUPPE, BEZOGEN
AUF DEN TOTALEN ELASTISCHEN QUERSCHNITT (TOTALES \varnothing -MOMENT) DER
G-TEN GRUPPE

M A K R O W I C H T U N G

ALLE MOMENTE WIE DAS \varnothing -MOMENT MIT
 $F(0,E)$ (STANDARD $F(0,E) = 1/E$)

M I K R O W I C H T U N G (FEINSTRUKTUR)

ALLE MOMENTE MIT $FS(0,E) = 1$
(KEINE FEINSTRUKTURWICHTUNG)

TOTALE ELASTISCHE STREUQUERSCHNITTE SGN UND STREUKOSINUS MUEL FUER 0 16

3.GRUPPE
2.040E 00
2.693E-01

ERLAEUTERUNG
G-TE GRUPPE, 1. ZEILE: TOTALER ELASTISCHER QUERSCHNITT DER G-TEN
GRUPPE
G-TE GRUPPE, 2. ZEILE: MITTLERER STREUKOSINUS DER G-TEN GRUPPE

PROGRAMM KENNZIFFER 6
PROGRAMM ZUR BERECHNUNG ELASTISCHER STREUMATRIZEN

***WARNING NDF. 2 : THE FURTHER NAME 0.15830016E 08 IS GREATER THAN THE GREATEST FURTHER NAME
INCLUDED IN THE KEDAK LIBRARY FOR 0 16 SGNC

***WARNING NDF. 2 : THE FURTHER NAME 0.15830016E 08 IS GREATER THAN THE GREATEST FURTHER NAME
INCLUDED IN THE KEDAK LIBRARY FOR 0 16 SGNC

***MESSAGE 6.1 : WARNING NDF. 2 MAY BE IGNORED

SGNC0 ELASTISCHE STREUMATRIX 0. ORDNUNG FUER 0 16
12.GRUPPE
0.84206
0.15793

SGNC1 ELASTISCHE STREUMATRIX 1. ORDNUNG FUER 0 16
12.GRUPPE
0.09014
-0.04811

SGNC2 ELASTISCHE STREUMATRIX 2. ORDNUNG FUER 0 16
12.GRUPPE
0.00343
-0.00263

SGNC3 ELASTISCHE STREUMATRIX 3. ORDNUNG FUER 0 16
12.GRUPPE
0.00004
-0.00004

SGNC4 ELASTISCHE STREUMATRIX 4. ORDNUNG FUER 0 16
12.GRUPPE
0.0
0.0

SGNC5 ELASTISCHE STREUMATRIX 5. ORDNUNG FUER 0 16
12.GRUPPE
-0.00005
0.00005

ERLAEUTERUNG
L-TE ORDNUNG, G-TE GRUPPE, I-TE ZEILE: MATRIXELEMENT L-TER ORDNUNG
FUER STREUUNG AUS DER G-TEN GRUPPE IN DIE (G+I-1)-TE GRUPPE, BEZOGEN
AUF DEN TOTALEN ELASTISCHEN QUERSCHNITT (TOTALES 0.MOMENT) DER
G-TEN GRUPPE

M A K R O W I C H T U N G

ALLE MOMENTE WIE DAS 0. MOMENT MIT
 $F(0,E)$ (STANDARD $F(0,E) = 1/E$)

M I K R O W I C H T U N G (FEINSTRUKTUR)

ALLE MOMENTE MIT $FS(0,E) = 1$
(KEINE FEINSTRUKTURWICHTUNG)

TOTALE ELASTISCHE STREUQUERSCHNITTE SGN UND STREUKOSINUS MUEL FUER 0 16

12.GRUPPE
3.700E 00
4.203E-02

ERLAEUTERUNG
G-TE GRUPPE, 1. ZEILE: TOTALER ELASTISCHER QUERSCHNITT DER G-TEN
GRUPPE
G-TE GRUPPE, 2. ZEILE: MITTLERER STREUKOSINUS DER G-TEN GRUPPE


```

ELASTIC SCATTERING MATRIX  SGNC0  FOR 0 16  GROUP= 3
FROM GROUP 3 TO GROUP 3
 2.8227E-02  7.8761E-02  1.2239E-01  1.6345E-01  2.1659E-01  2.7287E-01
 3.2734E-01  3.7757E-01  4.2379E-01  4.6571E-01  5.0372E-01  5.3808E-01
 5.6899E-01  5.9703E-01  6.2213E-01  6.3841E-01  6.5196E-01  6.6610E-01
 6.8074E-01  6.9571E-01  7.1089E-01  7.2614E-01  7.4286E-01  7.6173E-01
 7.7919E-01  7.7882E-01  7.7721E-01  7.8409E-01  7.9834E-01  8.1127E-01
 8.2497E-01  8.5349E-01  8.8278E-01  9.0954E-01  9.3579E-01  9.5933E-01
 9.7833E-01  9.9636E-01  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
FROM GROUP 3 TO GROUP 4
 9.7177E-01  9.2124E-01  8.7761E-01  8.3655E-01  7.8341E-01  7.2713E-01
 6.7266E-01  6.2243E-01  5.7621E-01  5.3429E-01  4.9628E-01  4.6192E-01
 4.3101E-01  4.0297E-01  3.7787E-01  3.6159E-01  3.4804E-01  3.3390E-01
 3.1926E-01  3.0429E-01  2.8911E-01  2.7386E-01  2.5714E-01  2.3827E-01
 2.2081E-01  2.2118E-01  2.2279E-01  2.1591E-01  2.0166E-01  1.8873E-01
 1.7503E-01  1.4651E-01  1.1722E-01  9.0457E-02  6.4213E-02  4.0672E-02
 2.1672E-02  3.6398E-03  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0

```

```

ELASTIC SCATTERING MATRIX  SGNC1  FOR 0 16  GROUP= 3
FROM GROUP 3 TO GROUP 3
 2.7734E-02  7.5561E-02  1.1431E-01  1.4845E-01  1.9148E-01  2.3499E-01
 2.7504E-01  3.1030E-01  3.4129E-01  3.6826E-01  3.9177E-01  4.1230E-01
 4.3019E-01  4.4603E-01  4.5948E-01  4.6035E-01  4.5594E-01  4.5095E-01
 4.4530E-01  4.3874E-01  4.3129E-01  4.2293E-01  4.1456E-01  4.0657E-01
 3.9628E-01  3.6824E-01  3.3286E-01  3.1342E-01  3.2575E-01  3.3829E-01
 3.4885E-01  3.4520E-01  3.3349E-01  3.1575E-01  2.9724E-01  2.7649E-01
 2.4648E-01  2.1389E-01  1.9446E-01  1.7864E-01  1.6288E-01  1.5741E-01
 1.6541E-01  1.7411E-01  1.8281E-01  1.9164E-01  2.0459E-01  2.2056E-01
 2.4255E-01  2.6603E-01  2.8650E-01  3.0190E-01  3.1725E-01  3.3260E-01
 3.4599E-01  3.5241E-01  3.5740E-01  3.6040E-01  3.6180E-01  3.5350E-01
 3.0990E-01  2.6659E-01  2.5582E-01  3.3799E-01  3.7447E-01  3.6958E-01
 3.6843E-01  3.7472E-01  3.9501E-01  4.2433E-01
FROM GROUP 3 TO GROUP 4
 2.3338E-01  1.7569E-01  1.2704E-01  8.6316E-02  4.8585E-02  1.2328E-02
 -2.0506E-02  -4.8652E-02  -7.2483E-02  -9.2315E-02  -1.0870E-01  -1.2212E-01
 -1.3295E-01  -1.4172E-01  -1.4866E-01  -1.5337E-01  -1.5682E-01  -1.5903E-01
 -1.6013E-01  -1.6034E-01  -1.5965E-01  -1.5806E-01  -1.5469E-01  -1.4906E-01
 -1.4352E-01  -1.4886E-01  -1.5519E-01  -1.5658E-01  -1.5317E-01  -1.4929E-01
 -1.4343E-01  -1.2335E-01  -1.0140E-01  -8.0464E-02  -5.8776E-02  -3.8308E-02
 -2.0987E-02  -3.6007E-03  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0
 0.0  0.0  0.0  0.0  0.0  0.0

```

ELASTIC SCATTERING MATRIX SGNC2 FOR 0 16 GROUP= 3

FROM GROUP 3 TO GROUP 3

2.6766E-02	6.9427E-02	9.9230E-02	1.2126E-01	1.4720E-01	1.7011E-01
1.8788E-01	2.0108E-01	2.1076E-01	2.1789E-01	2.2328E-01	2.2758E-01
2.3134E-01	2.3489E-01	2.3776E-01	2.2621E-01	2.3853E-01	1.9059E-01
1.7245E-01	1.5412E-01	1.3577E-01	1.1762E-01	1.0012E-01	8.3564E-02
6.7289E-02	4.7667E-02	2.6716E-02	2.3535E-02	5.2308E-02	8.3559E-02
1.1489E-01	1.2767E-01	1.3727E-01	1.4783E-01	1.6092E-01	1.7487E-01
1.8748E-01	2.0201E-01	2.0390E-01	2.0228E-01	2.0096E-01	2.3267E-01
2.8650E-01	3.3315E-01	2.9590E-01	2.4299E-01	2.3671E-01	2.5272E-01
2.6846E-01	2.8442E-01	2.8420E-01	2.5902E-01	2.2789E-01	1.8492E-01
1.4501E-01	1.7260E-01	2.1578E-01	2.2436E-01	2.3190E-01	2.4786E-01
2.4956E-01	1.9408E-01	1.4298E-01	1.3944E-01	1.4051E-01	1.4157E-01
1.4177E-01	1.3337E-01	1.2280E-01	1.1221E-01		

FROM GROUP 3 TO GROUP 4

3.0559E-02	-2.2521E-02	-6.2757E-02	-9.2215E-02	-1.0496E-01	-1.1080E-01
-1.1154E-01	-1.0771E-01	-1.0039E-01	-9.0533E-02	-7.8945E-02	-6.6286E-02
-5.3117E-02	-3.9740E-02	-2.6679E-02	-1.8461E-02	-1.1604E-02	-4.4953E-03
2.8189E-03	1.0335E-02	1.7854E-02	2.5172E-02	3.1843E-02	3.7576E-02
4.2971E-02	5.1106E-02	6.0393E-02	7.0159E-02	7.9630E-02	8.7197E-02
9.1914E-02	8.4465E-02	7.4059E-02	6.2687E-02	4.8797E-02	3.3856E-02
1.9663E-02	3.5235E-03	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0

ELASTIC SCATTERING MATRIX SGNC3 FOR 0 16 GROUP= 3

FROM GROUP 3 TO GROUP 3

2.5362E-02	6.0881E-02	7.9167E-02	8.6817E-02	9.3899E-02	9.5909E-02
9.3285E-02	8.8295E-02	8.2518E-02	7.7178E-02	7.3001E-02	7.0453E-02
6.9835E-02	7.1111E-02	7.3614E-02	6.5249E-02	5.3199E-02	4.2406E-02
3.2903E-02	2.4729E-02	1.7893E-02	1.2341E-02	8.0041E-03	4.7971E-03
3.2926E-03	1.1762E-02	2.5130E-02	4.3321E-02	6.5863E-02	8.7285E-02
1.0672E-01	1.1678E-01	1.2388E-01	1.2897E-01	1.3047E-01	1.2671E-01
1.1345E-01	9.5190E-02	8.7900E-02	8.4020E-02	8.0139E-02	7.6330E-02
7.2562E-02	6.8435E-02	6.0340E-02	5.1719E-02	5.2188E-02	5.6995E-02
6.1722E-02	6.6513E-02	7.2391E-02	7.9836E-02	8.4337E-02	8.2694E-02
8.1630E-02	9.7462E-02	1.2062E-01	1.4744E-01	1.6920E-01	1.6948E-01
1.6070E-01	1.2653E-01	9.5032E-02	8.9697E-02	8.6919E-02	8.4150E-02
8.1361E-02	7.8374E-02	7.5339E-02	7.2302E-02		

FROM GROUP 3 TO GROUP 4

-2.1149E-02	-5.8151E-02	-7.7923E-02	-8.5943E-02	-8.5758E-02	-7.9057E-02
-6.7736E-02	-5.4067E-02	-3.9597E-02	-2.5588E-02	-1.2753E-02	-1.5550E-03
7.6880E-03	1.5046E-02	2.0595E-02	2.5728E-02	3.0158E-02	3.3334E-02
3.5209E-02	3.5753E-02	3.4958E-02	3.2880E-02	2.9583E-02	2.5159E-02
1.9983E-02	1.6188E-02	1.0800E-02	6.8687E-04	-1.3722E-02	-2.7010E-02
-3.8428E-02	-4.1735E-02	-4.2329E-02	-4.0933E-02	-3.5941E-02	-2.7825E-02
-1.7787E-02	-3.4100E-03	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0

ELASTIC SCATTERING MATRIX SGNC4 FOR 0 16 GROUP= 3

FROM GROUP 3 TO GROUP 3

2.3577E-02	5.0616E-02	5.6672E-02	5.1047E-02	4.2891E-02	3.0842E-02
1.7509E-02	5.8715E-03	-2.9387E-03	-8.2105E-03	-1.0031E-02	-8.7210E-03
-4.7690E-03	1.2609E-03	8.5238E-03	1.0897E-02	1.2070E-02	1.4060E-02
1.6559E-02	1.9294E-02	2.1858E-02	2.3859E-02	2.4942E-02	2.4802E-02
2.3806E-02	2.7629E-02	3.1281E-02	3.0953E-02	2.6293E-02	2.0377E-02
1.3148E-02	2.4266E-03	-8.9155E-03	-1.9021E-02	-2.6037E-02	-2.8643E-02
-2.5830E-02	-1.6090E-02	-1.5306E-02	-1.7784E-02	-2.0165E-02	-1.1739E-02
3.9115E-03	1.6794E-02	6.4191E-03	-9.0435E-03	-5.3530E-03	7.4612E-03
2.0062E-02	3.2832E-02	3.7516E-02	2.9630E-02	2.2722E-02	1.8113E-02
1.4322E-02	2.2889E-02	3.6789E-02	5.3304E-02	6.6725E-02	6.6888E-02
6.2156E-02	4.4344E-02	2.7793E-02	2.4347E-02	2.2187E-02	2.0034E-02
1.8061E-02	1.7856E-02	1.8101E-02	1.8348E-02		

FROM GROUP 3 TO GROUP 4

-1.9280E-02	-4.7367E-02	-5.4472E-02	-4.9475E-02	-3.9058E-02	-2.4184E-02
-8.0334E-03	6.4495E-03	1.8038E-02	2.6119E-02	3.0737E-02	3.2227E-02
3.1071E-02	2.7838E-02	2.3144E-02	1.8920E-02	1.4174E-02	8.6146E-03
2.5447E-03	-3.7595E-03	-9.8970E-03	-1.5470E-02	-2.0127E-02	-2.3557E-02
-2.5856E-02	-2.9678E-02	-3.2368E-02	-2.9979E-02	-2.1376E-02	-1.1483E-02
-1.0080E-03	6.8100E-03	1.3802E-02	1.9570E-02	2.2244E-02	2.0893E-02
1.5480E-02	3.2629E-03	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0

ELASTIC SCATTERING MATRIX SGNC5 FOR 0 16 GROUP= 3

FROM GROUP 3 TO GROUP 3

2.1479E-02	3.9439E-02	3.4446E-02	1.9551E-02	3.5546E-03	-1.2091E-02
-2.4171E-02	-3.0887E-02	-3.2541E-02	-3.0025E-02	-2.4597E-02	-1.7517E-02
-9.9262E-03	-2.5902E-03	3.8581E-03	9.3175E-03	1.3829E-02	1.7008E-02
1.8640E-02	1.9620E-02	1.6962E-02	1.3829E-02	9.4772E-03	4.2549E-03
-1.5461E-03	-9.7245E-03	-1.9033E-02	-2.5900E-02	-2.7145E-02	-2.5613E-02
-2.1620E-02	-1.8142E-02	-1.3216E-02	-7.2960E-03	-2.3016E-03	1.0929E-04
-1.9805E-04	-9.1969E-03	-1.1478E-02	-1.0673E-02	-9.8508E-03	-7.1093E-03
-3.0953E-03	5.7997E-04	1.1216E-03	9.4173E-04	1.1144E-03	1.4553E-03
1.7890E-03	2.1276E-03	1.8221E-03	5.2818E-04	-1.1974E-03	-3.9112E-03
-6.2858E-03	-2.1162E-03	4.1595E-03	9.2182E-03	1.3064E-02	1.2546E-02
1.0472E-02	4.8066E-03	-5.9652E-04	-2.1279E-03	-3.2757E-03	-4.4197E-03
-5.2295E-03	-2.7488E-03	5.6752E-04	3.8876E-03		

FROM GROUP 3 TO GROUP 4

-1.7354E-02	-3.6147E-02	-3.1979E-02	-1.7839E-02	-2.0927E-03	1.3420E-02
2.5361E-02	3.1952E-02	3.3451E-02	3.0795E-02	2.5232E-02	1.8011E-02
1.0282E-02	2.8107E-03	-3.7705E-03	-9.2399E-03	-1.3713E-02	-1.6852E-02
-1.8446E-02	-1.8391E-02	-1.6694E-02	-1.3522E-02	-9.1282E-03	-3.8677E-03
1.8404E-03	8.3427E-03	1.5514E-02	2.1221E-02	2.3019E-02	2.2071E-02
1.8461E-02	1.3109E-02	5.8558E-03	-2.3907E-03	-9.7082E-03	-1.3798E-02
-1.2887E-02	-3.0857E-03	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0


```

ELASTIC SCATTERING MATRIX  SGNC0  FOR 0 16  GROUP= 12
FROM GROUP 12 TO GROUP 12
 2.4398E-02  7.2843E-02  1.2077E-01  1.6817E-01  2.1506E-01  2.6145E-01
 3.0734E-01  3.5274E-01  3.9763E-01  4.4205E-01  4.8600E-01  5.2946E-01
 5.7245E-01  6.1498E-01  6.5705E-01  6.9868E-01  7.3984E-01  7.8057E-01
 8.2085E-01  8.6070E-01  9.0012E-01  9.3912E-01  9.7769E-01  9.9987E-01
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
 1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00  1.0000E 00
FROM GROUP 12 TO GROUP 13
 9.7560E-01  9.2716E-01  8.7923E-01  8.3183E-01  7.8494E-01  7.3855E-01
 6.9266E-01  6.4726E-01  6.0237E-01  5.5795E-01  5.1400E-01  4.7054E-01
 4.2755E-01  3.8502E-01  3.4295E-01  3.0132E-01  2.6016E-01  2.1943E-01
 1.7915E-01  1.3930E-01  9.9878E-02  6.0885E-02  2.2314E-02  1.3161E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0

```

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ELASTIC SCATTERING MATRIX  SGNC1  FOR 0 16  GROUP= 12
FROM GROUP 12 TO GROUP 12
 2.3694E-02  6.7944E-02  1.0756E-01  1.4264E-01  1.7328E-01  1.9957E-01
 2.2160E-01  2.3946E-01  2.5323E-01  2.6301E-01  2.6889E-01  2.7093E-01
 2.6923E-01  2.6387E-01  2.5492E-01  2.4246E-01  2.2657E-01  2.0733E-01
 1.8481E-01  1.5907E-01  1.3019E-01  9.8251E-02  6.3310E-02  4.1831E-02
 4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02
 4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02
 4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02
 4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02
 4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02
 4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02  4.1700E-02
FROM GROUP 12 TO GROUP 13
 1.8007E-02  -2.6244E-02  -6.5862E-02  -1.0094E-01  -1.3158E-01  -1.5787E-01
 -1.7990E-01  -1.9776E-01  -2.1153E-01  -2.2131E-01  -2.2719E-01  -2.2923E-01
 -2.2753E-01  -2.2217E-01  -2.1322E-01  -2.0076E-01  -1.8487E-01  -1.6563E-01
 -1.4311E-01  -1.1737E-01  -8.8494E-02  -5.6551E-02  -2.1610E-02  -1.3129E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0

```



```

ELASTIC SCATTERING MATRIX   SGNC2   FOR 0 16   GROUP= 12
FROM GROUP   12 TO GROUP   12
 2.2331E-02  5.8828E-02  8.4085E-02  9.9409E-02  1.0608E-01  1.0531E-01
 9.8319E-02  8.6247E-02  7.0222E-02  5.1329E-02  3.0616E-02  9.1103E-03
-1.2205E-02 -3.2376E-02 -5.0477E-02 -6.5617E-02 -7.6927E-02 -8.3573E-02
-8.4746E-02 -7.9666E-02 -6.7575E-02 -4.7748E-02 -1.9481E-02  6.6218E-04
 7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04
 7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04
 7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04
 7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04
 7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04
 7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04
 7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04  7.9485E-04
FROM GROUP   12 TO GROUP   13
-2.1560E-02 -5.8058E-02 -8.3314E-02 -9.8639E-02 -1.0531E-01 -1.0454E-01
-9.7548E-02 -8.5477E-02 -6.9452E-02 -5.0558E-02 -2.9845E-02 -8.3398E-03
 1.2976E-02  3.3146E-02  5.1247E-02  6.6388E-02  7.7698E-02  8.4343E-02
 8.5516E-02  8.0437E-02  6.8345E-02  4.8518E-02  2.0251E-02  1.3066E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0

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ELASTIC SCATTERING MATRIX   SGNC3   FOR 0 16   GROUP= 12
FROM GROUP   12 TO GROUP   12
 2.0397E-02  4.6736E-02  5.5371E-02  5.1107E-02  3.8133E-02  2.0051E-02
-9.3706E-05 -1.9804E-02 -3.7097E-02 -5.0493E-02 -5.8985E-02 -6.2022E-02
-5.9495E-02 -5.1713E-02 -3.9386E-02 -2.3605E-02 -5.8383E-03  1.2104E-02
 2.8073E-02  3.9599E-02  4.3910E-02  3.7935E-02  1.8332E-02  1.2961E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
FROM GROUP   12 TO GROUP   13
-2.0398E-02 -4.6737E-02 -5.5373E-02 -5.1108E-02 -3.8134E-02 -2.0052E-02
 9.2736E-05  1.9803E-02  3.7096E-02  5.0492E-02  5.8983E-02  6.2020E-02
 5.9493E-02  5.1711E-02  3.9385E-02  2.3604E-02  5.8370E-03 -1.2105E-02
-2.8074E-02 -3.9601E-02 -4.3912E-02 -3.7937E-02 -1.8333E-02 -1.2972E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0

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ELASTIC SCATTERING MATRIX  SGNC4  FOR 0 16      GROUP= 12
FROM GROUP 12 TO GROUP 12
 1.8015E-02  3.3250E-02  2.7181E-02  1.0437E-02 -9.2209E-03 -2.6488E-02
-3.8141E-02 -4.2691E-02 -4.0035E-02 -3.1138E-02 -1.7719E-02 -1.9804E-03
 1.3692E-02  2.6975E-02  3.5856E-02  3.8871E-02  3.5321E-02  2.5484E-02
 1.0813E-02 -5.8810E-03 -2.0256E-02 -2.6259E-02 -1.5987E-02 -1.2872E-04
-2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07
-2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07
-2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07
-2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07
-2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07
-2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07
-2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07 -2.5071E-07
FROM GROUP 12 TO GROUP 13
-1.8015E-02  3.3250E-02  2.7182E-02 -1.0438E-02  9.2206E-03  2.6488E-02
 3.8141E-02  4.2691E-02  4.0035E-02  3.1137E-02  1.7719E-02  1.9802E-03
-1.3692E-02 -2.6975E-02 -3.5856E-02 -3.8872E-02 -3.5321E-02 -2.5485E-02
-1.0813E-02  5.8808E-03  2.0256E-02  2.6259E-02  1.5987E-02  1.2847E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
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 0.0 0.0 0.0 0.0 0.0 0.0

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ELASTIC SCATTERING MATRIX  SGNC5  FOR 0 16      GROUP= 12
FROM GROUP 12 TO GROUP 12
 1.5329E-02  2.0021E-02  4.4840E-03 -1.4239E-02 -2.6976E-02 -3.0232E-02
-2.4354E-02 -1.1998E-02  3.1234E-03  1.7222E-02  2.7149E-02  3.0877E-02
 2.7769E-02  1.8614E-02  5.4863E-03 -8.5982E-03 -2.0172E-02 -2.6006E-02
-2.3915E-02 -1.3722E-02  1.6574E-03  1.4978E-02  1.3366E-02  1.2693E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
FROM GROUP 12 TO GROUP 13
-1.5329E-02 -2.0021E-02 -4.4840E-03  1.4239E-02  2.6976E-02  3.0232E-02
 2.4354E-02  1.1998E-02 -3.1235E-03 -1.7222E-02 -2.7149E-02 -3.0877E-02
-2.7769E-02 -1.8614E-02 -5.4864E-03  8.5979E-03  2.0172E-02  2.6006E-02
 2.3915E-02  1.3722E-02 -1.6574E-03 -1.4978E-02 -1.3366E-02 -1.2692E-04
 0.0 0.0 0.0 0.0 0.0 0.0
 0.0 0.0 0.0 0.0 0.0 0.0
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**DATAMAIN' STELLT FUER EINE FORTRAN-SUBROUTINE 'MAIN' SPEICHERFELDER,
**DEREN LAEANGEN IM EXEC-PARM-FELD ANGEGEBEN SIND,BEREIT UND GIBT SIE
**ALS ARGUMENTE AN 'MAIN' NACH FOLGENDEM SCHEMA WEITER.
* // EXEC FHG,PARM.G=(X,Y),REGION.G=RK X,Y SIND 2 FELDLAENGEN
*AUSFUEHRUNG VON 'DATAMAIN' LIEFERT ANFANGSADRESSEN DER ZWEI FELDER IN
*DER TASK-REGION.
* N > 0 GIBT ANZAHL DER FELDER AN (NORMALFALL)
* DURCH 'SUBROUTINE MAIN(N,X,JX,Y,JY)' KOENNEN DIE 2 FELDER IN 'MAIN' SO-
* WOHL ALS FESTKOMMA-ALS AUCH ALS GLEITKOMMA-FELDER BENUTZT WERDEN.
* UNTER X(1)=JX(1) BZW. Y(1)=JY(1) KOENNEN DIE LAENGEN XX,YY IN 'MAIN'
* ABGEHOLT WERDEN.
* N = 0 IST FEHLERFALL.HINWEIS AUF FEHLERART UNTER JX(1)-JX(35).
* MAXIMALLAENGE DES PARM-FELDES IST 40 EBCDIC-ZEICHEN.
CATAMAIN START 0 ANFANG VON 'DATAMAIN'
      EXTRN MAIN
      BC 15,12(0,15) SPRUNG UM KONSTANTEN
      DC C'DATAMAIN' PROG.NAME
      STM 14,12,12(13) RETTEN DER REGISTER 14,15,0-12
      BALR 10,0 LADE BASIS-REG.10
      USING *,10
      ST 13,AREA+4 SYSTEM SA-ADR. NACH PROG.SA
      LA 4,AREA PROG.SA-ADRESSE
      ST 4,8(0,13) NACH SYSTEM-SA
      L 4,0(0,1) /4/=ADR.DES PARM-FELDS
      LH 3,0(0,4) /3/=PARM-FELD-LAENGE
      LH 2,=X'0001' /2/=INKREMENT=1
      AR 4,2 /4/=1/4+1
      AR 3,4 /3/=COMPARAND=ENDADR.DES PARM-FELDES
      LH 5,=X'0000' /5/=I=PARAMETERZAEHLER=0
      LPR 1,4 /1/=ANFANG DES NEUEN PARAMETERS-1
      CHARLOOP BXH 4,2,LOOPEND /4/=1/4+1,IF(1/4-1/3)*+1,*+1,LOOPEND
      CLI 0(4),X'6B' VERGLEICH AUF KOMMA
      BNE CHARLOOP
      CVBLOOP LR 7,4 /7/=ANZAHL DER
      SR 7,1 EBCDIC-ZEICHEN
      LPR 7,7
      SH 7,=X'0C02' -1 OHNE ,
      BL NOPARM
      LH 6,=X'0C05'
      SR 6,7
      BL ERREXIT
      AH 7,=X'0C70' /7/=7*16+7/=PACK-LAENGEN
      STC 7,PACK+1 PACK-MODIFIKATION
      BL PARMLOOP IF(7)PARMLOOP,PACK,PACK
      PACK PACK X,1(8,1) /X/=UNSIGNIERTE DEZIMALLAENGE
      NI X+7,X'FE' /X/=+X/
      CVB 6,X DEZIMAL-BINAER-WANDLUNG
      LTR 6,6
      BNH NOPARM
      ST 6,LNG(5) LNG(I)=/6/=BINAERLAENGE
      AH 5,=X'0C04'
      SR 4,2
      BXH 4,2,GTMN KVT02901 P
      B PARMLOOP GOTO PARMLOOP
      LOOPEND LR 6,5 /6/=1/5/ AND GO TO CVBLOOP
      B CVBLOOP
      GTMN MVI AREAS,X'00' LOESCHEN DES
      MVC AREAS+1(83),AREAS AREAS-FELDES
      SH 5,=X'0C04' LETZTER PARAMETER

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IC 6,=X'80'
STC 6,LNG(5)
LA 13,AREA
GETMAIN LC,LA=LNG,A=AREAS+4
LTR 6,15 IF(1/15)GTMNCOND,**2,GTMNCOND
BNE GTMNCOND /15/=0 ODER 4, COND-CODE NACH GETMAIN
IC 6,=X'00'
STC 6,LNG(5)
AH 5,=X'0C04'
LR 8,5 /8/=K=I
LR 6,5 /6/=J=I
AR 5,5 /5/=I=I+I
      EPLOOP SH 6,=X'0004' J=J-1
      BL EPILOG IF(J)EPILOG,*+1,*+1
      SH 5,=X'0004' I=I-1
      L 7,AREAS+4(6) /7/=AREAS(J+1)/
      ST 7,AREAS+4(5) AREAS(I)=AREAS(J+1)
      SH 5,=X'0004' I=I-1
      ST 7,AREAS+4(5) AREAS(I)=AREAS(J+1)=AREAS(I+1)
      L 9,LNG(6) /9/=LNG(J)
      ST 9,0(7) /77/=1/9/,1.WORT DES FELDES/=FELDLAENGE
      B EPLOOP
      EPILOG SRA 8,2(0) N=K=/8/*4**=-1
      ST 8,N
      LA 8,N
      ST 8,AREAS
      RUFIBCOM L 13,AREA+4 /13/=PROG.SA-ADR
      L 15,=V(IBCUM#) AUFRUF IBCOM
      BAL 14,64(15) (OPEN)
      RUFMAIN LA 13,AREA /13/=PROG.SA-ADR
      LA 1,AREAS /1/=ARG.LIST-ADR
      L 15,MAINAD AUFRUF
      BALR 14,15 MAIN
      CNOP 0,4
      RETURNOS L 13,AREA+4
      STOP L 15,=V(IBCUM#)
      BAL 14,52(15)
      DC X'0540404040F00000'
      GTMNCOND LA 6,MSG3
      B ERREX
      NOPARM LA 6,MSG2
      B ERREX
      ERREXIT LA 6,MSG1
      ERREX ST 6,AREAS+4
      ST 6,AREAS+8
      CONDEXIT B CONDEXIT
      LH 8,=X'0C00'
      B EPILOG
      MAINAD DC A(MAIN)
      MSG1 DC C' FELDLAENGE HAT MEHR ALS 6 ZIFFERN. '
      MSG2 DC C' EXEC-PARM-FEHLER,Z.B.NICHT>0. '
      MSG3 DC C' TASK-REGION IST ZU KLEIN. '
      X DS 1D
      N DS 1F
      AREA CS 18F
      LNG CS 20F
      AREAS CS 41F
      END DATAMAIN

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870
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1010
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1100
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1120
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SUBROUTINE FSPIE
C FSPIE IS A SPECIAL ERROR-DETECTING SUBRCUTINE ,WHICH IN CASE
C OF AN ABNORMAL END DETERMINS THE PSW AND PRINTS THIS PSW +
C A TRACE-BACK + THE REGISTER CONTENTS + THE SYSTEM COMPLETION
C CODE.. FOR FSPIE IS INSTALATION DEPENDENT ITS CODE IS NOT
C DISTRIBUTED HERE
RETURN
END

```

```

C PROGRAMM ZUR STEUERUNG DES MIGROS-SYSTEMS
C

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```

SUBROUTINE MAIN (N,XL,NXL)
REAL*8 MAT1,N1,N2
DIMENSION XL(1),NXL(1)
COMMON MAT1,K(2),NOUT,LIZ
DATA N1/'BLOC '/,N2/'ENDE '/,SGC/'SGC '/
NINP=5
NOUT=6
KDE=1
LIZ=3
NF=8
IF(N.NE.0) GO TO 18
WRITE (NOUT,19) (XL(I),I=1,9)
19 FORMAT(1H ,9A4)
STOP
18 CALL FSPIE
IADR=NXL(1)/4
CALL EING(NOUT)
CALL FREEFO (NINP,NF,NOUT,XL,XL,XL)
CALL NDFOPN (KDE,NXL,NUDAT,0)
WRITE (NOUT,1) NUDAT
1 FORMAT (1H1/' STAND DER KERNDATENBIBLIOTHEK VOM',110)
IZI=0
2 READ (NF)
16 READ (NF) MAT1,MAT2
IZI=IZI+2
IF(MAT1.EQ.N1.AND.MAT2.LT.2) GO TO 2
IF(MAT1.EQ.N1.AND.MAT2.EQ.2) GO TO 3
NE=26
GO TO 4
3 READ (NF) NE
READ (NF) MAT1,MAT2
IZI=IZI+2
4 IF(MAT1.EQ.N1.AND.MAT2.EQ.3) GO TO 5
NFE=1
LSPEC=1
GO TO 6
5 READ (NF) LSPEC,NFE
READ (NF) MAT1,MAT2
IZI=IZI+2
6 IF(MAT1.EQ.N1.AND.MAT2.NE.4) GO TO 7
READ (NF) NA

```

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10 READ (NF) MAT1,MAT2
IZI=IZI+2
20 GO TO 9
30 7 WRITE (NOUT,8)
40 8 FORMAT(1H0/' ***ERROR 0. 1 : THE INPUT CF THE MODULE NUMBERS WAS N
50 10T FOUND')
60 STOP
70 9 IF(MAT1.EQ.N1.AND.MAT2.EQ.5) GO TO 10
80 NT=3
GO TO 11
10 READ (NF) NT
READ (NF) MAT1,MAT2
IZI=IZI+2
11 IF(MAT1.EQ.N1.AND.MAT2.EQ.6) GO TO 12
MI=7
GO TO 13
12 READ (NF) MI
READ (NF) MAT1,MAT2
IZI=IZI+2
13 IF(MAT1.EQ.N1.AND.MAT2.EQ.7) GO TO 14
NTYP=7
GO TO 15
14 READ (NF) NTYP,(XL(I),XL(I+200),I=1,NTYP)
24 READ (NF) MAT1,MAT2
IZI=IZI+2
DO 26 I=1,NTYP
IF(XL(I).EQ.SGC) NTYP=NTYP+1
26 CCNTINUE
15 IF(MAT1.EQ.N1.AND.MAT2.LT.12) GO TO 22
IF(MAT1.EQ.N1.AND.MAT2.GT.12) GO TO 23
IF(MAT1.EQ.N2) GC TO 23
READ (NF) NMAT
IZI=IZI+1
GO TO 25
22 READ (NF)
READ (NF) MAT1,MAT2
IZI=IZI+2
GC TO 15
23 NMAT=1
25 DG 17 I=1,IZI
17 BACKSPACE NF
NM =NTYP*2+1
NZ=NM+NMAT*2
NENG=NZ+NMAT
NSPEC=NENG+NE
NARB=NSPEC+NFE*{(LSPEC+1)}
NTEMP=NARB+NA*3
NSI=NTEMP+NT
NG=NSI+MI+1
NMI=NG+NA
NM2=NMI+NA*NE
NFG=NM2+NA*NE
NFI=NFG+NA*NE
NGG=NFI+NA*NE
NII=NGG+NE
NFR=NII+NE

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```

I=NFR/2*2
IF(I.EQ.NFR) NFR=NFR+1
IF(NFR.LT.IADR) GO TO 20
WRITE (NOUT,21)
21 FORMAT(1H0/' ***ERROR 0. 2 : THE PARM.G - FIELD IS ALREADY TOO SHO
IRT FOR THE CONTROL MODULE')
STOP
20 NX=NE+1
CALL INPUT (XL(1),XL(NM),XL(NZ),
1 XL(NENG),XL(NSPEC),XL(NARB),XL(NTEMP),XL(NSI),
1XL(NG),XL(NN1),XL(NN2),XL(NFG),XL(NFI),XL(NGG),XL(NII),
3NMAT,MI,NE,LSPEC+1,NFE,NA,NT,NTYP,NFR,XL,NF,NX,IADR)
IZI=0
READ (NF) MAT1
IF(MAT1.NE.N2) GO TO 16
I=0
WRITE (LIZ) I,N2
RETURN
END

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1000
1010
1020
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1060
1070
1080
1090
1100
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1120
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1170
1180

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C EINGABEBESCHREIBUNG FUER DAS PROGRAMMSYSTEM MIGROS 10
C SUBROUTINE EING(NOUT) 20
WRITE (NOUT,1) 30
1 FORMAT(1H1/' EINGABEBESCHREIBUNG FUER DAS PROGRAMMSYSTEM MIGROS'// 40
1' *****' // 50
2' EINLEITUNG' // 60
3' DAS PROGRAMMSYSTEM MIGROS ERLAUET DIE BERECHNUNG VON MIKROS-' / 80
4' KOPISCHEN GRUPPENKONSTANTEN VON KERNDATEN. ALS KERNDATEN-' / 90
5' BIBLIOTHEK WIRD KEDAK BENOETIGT. DAS PROGRAMMSYSTEM ENTHAELT' / 100
6' FOLGENDE PROGRAMME : ' // 110
7' KENNZIFFER ARBEITSPROGRAMM' // 120
8' 1 PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTAB-' / 130
9' SCHIRMFAKTOREN VON AUFGELCESTEN RESONANZPARAME-' / 140
A' TERN (12400)' / 150
B' 2 PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTAB-' / 160
C' SCHIRMFAKTOREN VON STATISTISCHEN RESONANZ-' / 170
D' PARAMETERN (01787)' / 180
E' 3 PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTAB-' / 190
F' SCHIRMFAKTOREN VON PUNKTWEISE GEGEBENEN WIRKUNGS-' / 200
G' QUERSCHNITTEN (01700)' / 210
WRITE (NOUT,14) 220
14 FORMAT( 230
1' 4 PROGRAMM ZUR BERECHNUNG VON GRUPPENWIRKUNGSQUER-' / 240
2' SCHNITTEN BEI UNENDLICHER VERDUENNUNG (01797)' / 250
3' 5 PROGRAMM ZUR BERECHNUNG INELASTISCHER STREU-' / 260
4' MATRIZEN (26640)' / 270
5' 6 PROGRAMM ZUR BERECHNUNG ELASTISCHER STREUMATRIZEN' / 280
6' (23871)' / 290
7' 7 PROGRAMM ZUR BERECHNUNG DES SPALTSPEKTRUMS' / 300
8' (01780)' / 310
9' 8 PROGRAMM ZUR BERECHNUNG DES 1/√- GRUPPENMITTEL-' / 320

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A' WERTES (10242)' / 330
B' 9 PROGRAMM ZUR BERECHNUNG DER ELASTISCHEN STREU-' / 340
C' MATRIZEN FUER DIE REMG- KORREKTUR' / 350
D' 10 PROGRAMM ZUR BERECHNUNG DER THERMISCHEN' / 360
E' QUERSCHNITTE (01722)' // 370
F' DIE NACHFOLGENDE EINGABE IST JEWEILS FUER EIN MATERIAL ZU' / 380
G' ERSTELLEN. SOLLEN GRUPPENKONSTANTEN FUER MEHRERE MATERIALIEN' / 390
H' IN EINEM LAUF BERECHNET WERDEN, SO IST DIESE EINGABE' / 400
G' ENTSPRECHEND OFT ZU WIEDERHOLEN')
WRITE(NOUT,30) 420
30 FORMAT(1H , 430
1' DIE EINGABE IST DEN FREEFO - KONVENTIONEN ENTSPRECHEND' / 440
2' FORMATFREI ZU ERSTELLEN') 450
WRITE (NOUT,13) 460
13 FORMAT(1H1/1H0/1H0/ 470
2' 1. KARTE' / 480
3' MAT MATERIALNAME IN KEDAK - NCMENKLATUR' / 490
4' ISTRUK 0: DAS MATERIAL WIRD ALS STRUKTURMATERIAL BEHANDELT' / 500
5' 1: DAS MATERIAL WIRD ALS SCHWERES ISOTOP BEHANDELT'// 510
6' FALLS DIE INFORMATION DARUEBER, OB ES SICH UM EIN NICHT SPALT-' / 520
7' BARES ODER UM EIN SPALTBARES MATERIAL HANDELT NICHT DEM' / 530
8' LETZTEN ZEICHEN DES MATERIALNAMENS ENTCNMEN WERDEN SOLL' / 540
9' - WOBEI EINE GERADE ZIFFER ALS LETZTES ZEICHEN EIN NICHT' / 550
A' SPALTBARES MATERIAL UND EINE UNGERADE ZIFFER ALS LETZTES' / 560
B' ZEICHEN EIN SPALTBARES MATERIAL BEZEICHNET - FOLGEN DIE BEIDEN' / 570
C' NAECHSTEN KARTEN :'/) 580
WRITE (NOUT,2) 590
2 FORMAT(' 2. KARTE'// 600
1' @BLOC@ 1 KONSTANTE' / 610
2' 3. KARTE' / 620
4' ISPA 0: DAS MATERIAL IST NICHT SPALTBAR' / 630
5' 1: DAS MATERIAL IST SPALTBAR' // 640
6' FALLS DIE ENERGIEGRUPPENGRENZEN NICHT DIE DES 26-GRUPPEN' / 650
7' ABN-SATZES SEIN SOLLEN, FOLGEN DIE BEIDEN NAECHSTEN KARTEN : ' // 660
8' 4. KARTE' / 670
9' @BLOC@ 2 KONSTANTE' / 680
A' 5. KARTE' / 690
B' NE ANZAHL DER ENERGIEGRUPPENGRENZEN' / 700
C' (ENG(I),I=1,NE) ENERGIEGRUPPENGRENZEN IN EV IN AUFSTEIGEN-' / 710
D' DER REIHENFOLGE GEORDET' / 720
E' DAS PROGRAMMSYSTEM MIGROS KANN EINE BELIEBIGE ANZAHL VON ' / 730
F' ENERGIEGRUPPEN BEHANDELN'//) 740
WRITE (NOUT,3) 750
3 FORMAT( 760
1' FALLS EIN PUNKTWEISE VORGEGEBENES STOSSDICHTESPEKTRUM VERWEN-' / 770
2' DET WERDEN SOLL, FOLGEN DIE BEIDEN NAECHSTEN KARTEN : ' // 780
3' 6. KARTE' / 790
4' @BLOC@ 3 KONSTANTE' / 800
5' 7. KARTE' / 810
6' NSPEC ANZAHL DER SPEKTREN' / 820
6' NFE ANZAHL DER SPEKTRUMSSTUETZPUNKTE' / 830
7' (EF(I),I=1,NFE),{(F(I,J),I=1,NFE),J=1,NSPEC)' / 840
8' EF : ENERGIESTUETZPUNKTE DER STOSSDICHTESPEKTREN IN EV' / 850
8' F : ZUGEOERIGE WERTE DER STOSSDICHTESPEKTREN' / 860
A' IST DIE ZAHL DER EINGELESENEN SPEKTREN GROESSER 1, SO' / 870
B' WERDEN DIE SPEKTREN 2, 3, ... NUR IM ARBEITSPROGRAMM NR. 6' / 880

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C* INTERPRETIERT UND BEDEUTEN DORT DIE MAKROSKOPISCHE WICH- / 890
D* TUNGSFUNKTION FUER DAS 1., 2.,...MOMENT DER ELASTISCHEN / 900
E* STREUMATRIX. ES IST DARAUFG ZU ACHTEN, DASS DIE ZAHL DER / 910
F* SPEKTREN MIT DER ZAHL DER ZU BERECHNENDEN MOMENTE UEBEREIN- / 920
G* STIMMT. WIRD NUR EIN SPEKTRUM ANGEGEBEN, SO WIRD DIESES / 930
H* FUER ALLE MOMENTE DER ELASTISCHEN STREUMATRIX VERWENDET.) / 940
WRITE(NOUT,31) / 950
31 FORMAT( / 960
1* DIE ANZAHL DER STUETZPUNKTE IST BELIEBIG, ES IST JEDOCH ZU / 970
2* BEACHTEN, DASS DIE SPEKTREN DEN GESAMTEN BETRACHTETEN / 980
3* ENERGIEBEREICH UEBERSTREICHEN SOLLTEN. WENN DIE KARTEN 6 / 990
4* UND 7 NICHT ANGEGEBEN WERDEN, BERECHNET MIGROS DIE SPEK- / 1000
5* TRUMSWERTE F(I,1) MITTELS DER FUNKTION PHI(E) UND DIE / 1010
6* SPEKTRUMSWERTE (F(I,L),L=2,NLE+1) MITTELS DER FUNKTION / 1020
7* PHI(L,E). NLE SIEHE KARTE 23. STANDARDFUNKTIONEN SIND ://) / 1030
WRITE(NOUT,43) / 1040
43 FORMAT(1H1/1H0/1H0/ / 1050
E* FUNCTION PHI(E) FUNCTION PHI(L,E) / 1060
F* PHI=1./E PHI=1./E**L / 1070
1* RETURN RETURN / 1080
2* END END / 1090
3* VERGL. AUCH DIE KARTE 27 : MAZ(1) / 1100
WRITE(NOUT,4) / 1110
4 FORMAT( / 1120
1* 8. KARTE / 1130
2* @BLOC@ 4 KONSTANTE / 1140
3* 9. KARTE / 1150
4* NA ANZAHL DER ANZULAUFENDEN ARBEITSPROGRAMME / 1160
5* ((NR(I,J),I=1,3),J=1,NA) / 1170
6* HIERBEI BEDEUTET: / 1180
7* NR(1,J) KENNZIFFER DER ARBEITSPROGRAMME / 1190
8* NR(2,J) NUMMER DER ENERGIEGRUPPE, AB DER DAS / 1200
9* PROGRAMM NR(1,J) RECHNEN SOLL / 1210
A* NR(3,J) NUMMER DER ENERGIEGRUPPE, BIS ZU DER / 1220
B* EINSCHLIESSLICH DAS PRGGRAMM NR(1,J) / 1230
C* RECHNEN SOLL. DIESE GRUPPENNUMMERN WERDEN / 1240
D* IM ENERGETISCH AUFSTIEGENDEN SINNE ANGEGEBEN. / 1250
E* NR(2,J).GE.NR(3,J) DIE NUMERIERUNG DER / 1260
F* ENERGIEGRUPPEN ERFOLGT MIT FALLENDER ENERGIE, / 1270
G* D.H. DIE ENERGIEGRUPPE, DIE ZU DEN HOECHSTEN / 1280
H* ENERGIEGRENZEN IN @BLOC @ 2 GEHOERT, / 1290
I* ERHAELT DIE NUMMER 1. / 1300
WRITE(NOUT,15) / 1310
15 FORMAT( / 1320
1* KOMMENTAR : / 1330
2* EIN ARBEITSPROGRAMM KANN MHRFACH AUFGERUFEN WERDEN. DIES IST / 1340
3* NOTWENDIG, FALLS MEHRERE NICHT ZUSAMMENHAENGENDE BEREICHE VON / 1350
4* ENERGIEGRUPPEN VORGEGBEN SIND. FUER JEDEN ZUSAMMENHAENGENDEN / 1360
5* BEREICH VON ENERGIEGRUPPEN HAT EIN AUFRUF EINES ARBEITSPRO- / 1370
6* GRAMMS ZU ERFOLGEN. / 1380
7* NR(2,J), NR(3,J) BEDEUTEN DIE NUMMERN DER AUSSTREUGRUPPEN IM / 1390
8* FALLE VON STREUMATRIZEN. GERECHNET WERDEN ALLE EINSTREUGRUPPEN) / 1400
WRITE(NOUT,5) / 1410
5 FORMAT( / 1420
1* WENN DIE PROGRAMME ZUR BERECHNUNG VON SELBSTABSCHIRMFAKTOREN / 1430
2* AUS RESONANZ- BZW STATISTISCHEN DATEN ANGELAUFEN WERDEN / 1440
2* (PROGRAMMKENNZIFFER 1 UND 2) UND IM / 1450
3* FALLE EINES SCHWEREN ISOTOPIS (ISTRUK=1) NICHT MIT DEN 3 STAN- / 1460
4* DARDTEMPERATUREN 300, 900 UND 2100 GRAD KELVIN UND IM FALLE / 1470
5* EINES STRUKTURMATERIALS (ISTRUK=0) NICHT MIT DER STANDARDTEM- / 1480
6* PERATUR 0 GRAD KELVIN GERECHNET WERDEN SOLL, FOLGEN DIE BEIDEN / 1490
7* KARTEN : / 1500
8* 10. KARTE / 1510
9* @BLOC@ 5 KONSTANTE / 1520
A* 11. KARTE / 1530
B* NT ANZAHL DER TEMPERATUREN / 1540
C* (TEMP(I),I=1,NT) TEMPERATUREN IN GRAD KELVIN / 1550
D* DIE ANZAHL DER TEMPERATUREN IST BELIEBIG) / 1560
WRITE(NOUT,6) / 1570
6 FORMAT(1H0, / 1580
1* WENN EINES DER PROGRAMME ZUR BERECHNUNG VON SELBSTABSCHIRMFAK- / 1590
2* TORN ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 1 UND 2 UND 3) / 1600
2* UND NICHT MIT DEN 7 STANDARD / 1610
3* VERDUENNUNGSGRADEN DES ABN-SATZES ALS C, 10, 10**2, 10**3, / 1620
4* 10**4, 10**5, 10**6 GERECHNET WERDEN SOLL, FOLGEN DIE KARTEN: // / 1630
5* 12. KARTE / 1640
6* @BLOC@ 6 KONSTANTE // / 1650
WRITE(NOUT,7) / 1660
7 FORMAT(1H1/1H0/1H0/ / 1670
7* 13. KARTE / 1680
8* MI ANZAHL DER VERDUENNUNGSGRADEN / 1690
9* (SIGO(I),I=1,MI) VERDUENNUNGSGRADEN / 1700
A* DIE ANZAHL DER VERDUENNUNGSGRADEN IST BELIEBIG / 1710
1* WENN DAS PROGRAMM ZUR BERECHNUNG VON GRUPPENWIRKUNGSQUER- / 1720
2* SCHNITTEN BEI UNENDLICHER VERDUENNUNG ANGELAUFEN WIRD / 1730
2* (PROGRAMMKENNZIFFER 4 UND 10) UND / 1740
3* NICHT DIE 7 STANDARDTYPEN MUEL,NUE,SGA,SGF,SGI,SGN UND SG2N // / 1750
4* BERUECKSICHTIGT WERDEN SOLL, FOLGEN DIE NAECHSTEN BEIDEN / 1760
5* KARTEN : / 1770
6* 14. KARTE / 1780
7* @BLOC@ 7 KONSTANTE / 1790
8* 15. KARTE / 1800
9* NTP ANZAHL DER WIRKUNGSQUERSCHNITTSTYPEN / 1810
A* (ITYP(I),I=1,NTP) WIRKUNGSQUERSCHNITTSTYPEN / 1820
B* DIE ANZAHL DER QUERSCHNITTSTYPEN IST BELIEBIG // / 1830
WRITE(NOUT,20) / 1840
20 FORMAT( / 1850
1* KOMMENTAR: ES KOENNEN FUER BELIEBIGE WIRKUNGSQUERSCHNITTE / 1860
2* GRUPPENMITTELWERTE BERECHNET WERDEN, SOFERN DIESE AUF KEDAK / 1870
3* VORHANDEN SIND. DIES SIND ZUR ZEIT: SGA, SGALP, SGF, SGG, SGI, / 1880
4* SGN, SGP, SGT, SGTR, SGX, SG2N. FERNER WERDEN GEEIGNET / 1890
5* DEFINIERT GRUPPENMITTELWERTE DER GROESSEN ALPHA, ETA, MUEL, / 1900
6* NUE BERECHNET, SOWIE DER WIRKUNGSQUERSCHNITT SGC = SGA - SGF / 1910
6* (DEFINITION DER OBEN GENANNTEN GROESSEN SIEHE / 1920
7* KEDAK-NOTIZ NR. 17, INR-NOTIZ NR. 279/71.) FALLS DIE QUER- / 1930
8* SCHNITTSTYPEN SGA UND SGF BERECHNET WERDEN, WIRD AUTOMATISCH / 1940
9* AUCH DER TYP SGC = SGA - SGF BERECHNET //) / 1950
WRITE(NOUT,8) / 1960
8 FORMAT( / 1970
1* WENN DAS PROGRAMM ZUR BERECHNUNG DER SELBSTABSCHIRMFAKTOREN / 1980
2* AUS RESONANZDATEN ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 1) UND / 1990
2* DIE ANZAHL DER PRO ENERGIESTUETZSTELLE MITZUNEHMENDEN RESONAN- / 2000

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3* ZEN NICHT 10 BEIDSEITS DER STUETZSTELLE SEIN SOLL, FOLGEN DIE* / 2010
4* BEIDEN NAECHSTEN KARTEN :* // 2020
6* 16. KARTE* / 2030
7* @BLOC@ 8 KONSTANTE* / 2040
8* 17. KARTE* / 2050
9* NRES ANZAHL DER RESONANZEN* // 2060
A* WENN DAS PROGRAMM ZUR BERECHNUNG VON SELBSTABSCHIRMFAKTOREN* / 2070
B* AUS RESONANZDATEN ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 1)* / 2080
C* UND DER ZULAESSIGE INTEGRATIONSFEHLER NICHT 0.05 SEIN SOLL,* / 2090
D* FOLGEN DIE BEIDEN KARTEN :* // 2100
D* 18. KARTE* / 2110
E* @BLOC@ 9 KONSTANTE* / 2120
F* 19. KARTE* / 2130
G* ERROR ZULAESSIGER INTEGRATIONSFEHLER*/) 2140
WRITE (NOUT,9) 2150
9 FORMAT( 2160
1* WENN DAS PROGRAMM ZUR BERECHNUNG INELASTISCHER STREUMATRIZEN* / 2170
2* ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 5) UND DER PARAMETER XNUE* / 2180
3* FUER DAS WEISSKOPF*SCHE VERDAMPFUNGSMODELL NICHT 0.16 SEIN * / 2190
4* SOLL (KERNTemperatur THETA(E)=SQRT(E/(XNUE*A))), WOBEI A DAS * / 2200
5* ATOMGEWICHT IST UND E IN MEV ANGENOMMEN WIRD), FOLGEN DIE* / 2210
6* BEIDEN NAECHSTEN KARTEN :*/IHL/IHO/IHC/ 2220
7* 20. KARTE* / 2230
8* @BLOC@ 10 KONSTANTE* / 2240
9* 21. KARTE* / 2250
A* XNUE PARAMETER FUER DAS WEISSKOPF*SCHE* / 2260
B* VERDAMPFUNGSMODELL* // 2270
C* WENN DAS PROGRAMM ZUR BERECHNUNG DER ELASTISCHEN STREUMATRIZEN*/ 2280
D* ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 6) UND NICHT MIT DER* / 2290
E* FOLGENDEN STANDARDEINGABE GERECHNET WERDEN SOLL:* / 2300
F* ISEL = 2 WENN ISTRUK=0* / 2310
G* ISEL = 1 WENN ISTRUK=1* / 2320
H* NLA = 0* / 2330
I* NLE = 5* / 2340
J* WERDEN DIE FOLGENDEN KARTEN BENUTZT :*/) 2350
WRITE (NOUT,10) 2360
10 FORMAT( 2370
1* 22. KARTE* / 2380
2* @BLOC@ 11 KONSTANTE* / 2390
3* 23. KARTE* / 2400
4* ISEL* / 2410
5* |ISEL| =2: FUER LEICHTE UND MITTELSCHWERE ELEMENTE* / 2420
6* (TOT. ELAST. STREUQUERSCHNITT UNGEAENDERT)* / 2430
7* =1: FUER SCHWERE ELEMENTE* / 2440
8* (TOT. ELAST. STREUQUERSCHNITT = 1 GESETZT)* / 2450
9* ISEL)0: ZUSAETZLICHER PAPIERAUSDRUCK VON ZWISCHEN-*/ 2460
10* INFORMATIEN* / 2470
11* NLA UNTERER LEGENDRE INDEX* / 2480
12* NLA GEHT UEBER IN 0 (WEGEN NORMIERUNG)* / 2490
13* NLE OBERER LEGENDRE INDEX (MAXIMAL 5)* / 2500
14* NLE GEHT UEBER IN MAX(1,NLE) (WEGEN MUEL-ANPASSUNG)* // 2510
15* WENN DAS PROGRAMM ZUR BERECHNUNG DER ELASTISCHEN STREUMATRIZEN*/ 2520
16* ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 6) UND EINE FEINSTRUKTUR-*/ 2530
17* RICHTUNG BERUECKSICHTIGT WERDEN SOLL, FOLGEN DIE KARTEN :* // 2540
18* 24. KARTE*) 2550
WRITE (NOUT,32) 2560

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32 FORMAT( 2570
D* @BLOC@ 12* / 2580
E* 25. KARTE* / 2590
F* NMAT ANZAHL DER MATERIALIEN IN DER* / 2600
G* MISCHEUNG* / 2610
H* (NAME(I),TZ(I),I=1,NMAT) NAME : MATERIALNAME* / 2620
I* TZ : TEILCHENDICHTE * 1.E-24* // 2630
1* WENN DAS PROGRAMM ZUR BERECHNUNG DER ELASTISCHEN STREUMATRIZEN*/ 2640
2* ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 6) UND NICHT ALLE MOMENTE* / 2650
3* MIT DEM SPEKTRUM DES 0. MOMENTES GEWICHTET WERDEN SOLLEN,* / 2660
4* FOLGEN DIE KARTEN :* // 2670
5* 26. KARTE / 2680
6* @BLOC@ 13* / 2690
7* 27. KARTE*) 2700
WRITE (NOUT,33) 2710
33 FORMAT( 2720
8* MAZ(1) 0 : ALLE MOMENTE WERDEN MIT DEM MAKROSPEKTRUM* / 2730
9* DES 0. MOMENTES GEWICHTET* / 2740
10* 1 : JEDES MOMENT HAT EIN EIGENES MAKROSPEKTRUM* / 2750
11* MAZ(2) 0 : ALLE MOMENTE WERDEN MIT DEM MIKROSPEKTRUM* / 2760
12* DES 0. MOMENTES GEWICHTET* / 2770
13* 1 : JEDES MOMENT HAT EIN EIGENES MIKROSPEKTRUM* / 2780
14* STANDARD : 1/SGT**((L+1) , L=0,...NLE*/) 2790
WRITE (NOUT,41) 2800
41 FORMAT(IHL/IHO/IHC/ 2810
1* WENN DAS PROGRAMM ZUR BERECHNUNG DER ELASTISCHEN STREUMATRIZEN*/ 2820
2* ANGELAUFEN WIRD (PROGRAMMKENNZIFFER 6) UND DIE MAXIMALE ANZAHL*/ 2830
3* VON GRUPPEN IN EINER MAKRO - GRUPPE NICHT 4 , DIE ANZAHL DER* / 2840
4* GRUNDSTUETZPUNKTE IN DEM ENERGIEBEREICH EINER GRUPPE, AUS DEM* / 2850
5* IN NACHBARGRUPPEN STREUUNG MOEGLICH IST,NICHT GLEICH 70 IST* / 2860
6* UND DIE ANZAHL DER GRUNDSTUETZPUNKTE IN DEM ENERGIEBEREICH* / 2870
7* EINER GRUPPE, AUS DEM KEINE STREUUNG IN NACHBARGRUPPEN* / 2880
8* MOEGLICH IST, NICHT GLEICH 16 IST, FOLGEN DIE KARTEN :* // 2890
9* 28. KARTE* / 2900
10* @BLOC@ 14 KONSTANTE* / 2910
11* 29. KARTE* / 2920
12* NS MAXIMALE ANZAHL VON GRUPPEN IN EINER MAKRO - GRUPPE* / 2930
13* NK ANZAHL DER GRUNDSTUETZPUNKTE IN DEM ENERGIEBEREICH* / 2940
14* EINER GRUPPE, AUS DEM IN NACHBARGRUPPEN STREUUNG* / 2950
15* MOEGLICH IST* / 2960
16* NR ANZAHL DER GRUNDSTUETZPUNKTE IN DEM ENERGIEBEREICH* / 2970
17* EINER GRUPPE, AUS DEM KEINE STREUUNG IN NACHBARGRUP-*/ 2980
18* PEN MOEGLICH IST*/) 2990
WRITE (NOUT,42) 3000
42 FORMAT( 3010
1* WENN DAS PROGRAMM ZUR BERECHNUNG DER REMO-DATEN ANGELAUFEN* / 3020
2* WIRD (PROGRAMMKENNZIFFER 9) UND DER ZULAESSIGE INTEGRATIONS-*/ 3030
3* FEHLER NICHT 0.05, NJM NICHT GLEICH 6 SEIN SOLL, WOBEI* / 3040
4* 2**NJM+1 DIE ANZAHL DER WINKELSTUETZPUNKTE FUER DIE WINKEL-*/ 3050
5* INTEGRATION IST, UND NUJM NICHT GLEICH 10 SEIN SOLL, WOBEI* / 3060
6* 2**NUJM+1 DIE MAXIMAL ZULAESSIGE ANZAHL VON ENERGIESTUETZPUNK-*/ 3070
7* TEN FUER DIE ENERGIEINTEGRATION IST, FOLGEN DIE BEIDEN* / 3080
8* KARTEN :* // 3090
9* 30. KARTE* / 3100
10* @BLOC@ 15 KONSTANTE* / 3110
11* 31. KARTE* / 3120

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B*   ERR   ZULAESSIGER INTEGRATIONSFEHLER* / 3130
C*   NJM   WOBEI 2**NJM+1 DIE ANZAHL DER WINKELSTUETZPUNKTE* / 3140
D*   FUER DIE WINKELINTEGRATION IST* / 3150
E*   NUJM  WOBEI 2**NUJM+1 DIE MAXIMAL ZULAESSIGE ANZAHL VON* / 3160
F*   ENERGIESTUETZPUNKTEN FUER DIE ENERGIEINTEGRATION IST* / 3170
G)
WRITE (NOUT,11) / 3190
11 FORMAT( / 3200
1* WENN DAS PROGRAMM ZUR BERECHNUNG DER REMO-DATEN ANGELAUFEN* / 3210
2* WIRD (PROGRAMMKENNZIFFER 9) UND NICHT MIT DER FOLGENDEN* / 3220
3* STANDARDEINGABE GERECHNET WERDEN SOLL : * / 3230
4* ISELR = 1 WENN ISTRUK=0* / 3240
4* ISELR = 0 WENN ISTRUK=1 / 3250
4* NLRA=0 UNTERER LEGENDRE INDEX* / 3260
5* NLRE=5 OBERER LEGENDRE INDEX* / 3270
6* NCALL=0 ALLE ENERGIEGRUPPEN WERDEN IN ALLEN ANLAEUFEN DES* / 3280
7* PROGRAMMS IN 14 UNTERGRUPPEN ZU JE 5 ENERGIESTUETZ- / 3290
8* PUNKTEN UNTERTEILT* / 3300
9* NFG=14 ANZAHL DER UNTERGRUPPEN PRO ENERGIEGRUPPE* / 3310
A* NFI=5 ANZAHL DER FEININTERVALLÄ PRO UNTERGRUPPE* / 3320
B* WERDEN DIE FOLGENDEN KARTEN BENOETIGT : * // 3330
1* 32. KARTE* / 3340
2* @BLOC@ 16 KONSTANTE* // 3350
WRITE (NOUT,12) / 3360
12 FORMAT(IH1/IH0/IH0/ / 3370
3* 33. KARTE* / 3380
4* ISELR=1 FUER LEICHTE UND MITTELSCHWERE ELEMENTE* / 3390
4* ISELR=0 FUER SCHWERE ELEMENTE* / 3400
4* NLRA UNTERER LEGENDRE INDEX* / 3410
5* NLRE OBERER LEGENDRE INDEX* / 3420
6* NCALL 0: ALLE ENERGIEGRUPPEN WERDEN IN ALLEN ANLAEUFEN* / 3430
7* DES PROGRAMMS IN 14 UNTERGRUPPEN ZU JE 5 ENER- / 3440
8* GIESTUETZPUNKTEN UNTERTEILT* / 3450
9* >0: ANZAHL DER PROGRAMMAUFRUFE (MUSS MIT DER ZAHL * / 3460
9* DER AUFRUFE DES ARBEITSPRCGRAMMS 9 IN KARTE 9* / 3470
9* UEBEREINSTIMMEN)* // 3480
A* WENN NCALL.GT.0 , FOLGT NCALL-MAL DIE KARTE : * // 3490
B* 34. KARTE* / 3500
C* NGRE ANZAHL DER BEREICHE VON ENERGIEGRUPPEN, DIE * / 3510
C* JEWEILS AUF GLEICHE WEISE ZU UNTERTEILEN SIND* / 3520
D* (N1(I), NUMMER DER ERSTEN ZU UNTERTEILENDEN ENERGIEGRUPPE* / 3530
E* N2(I), NUMMER DER LETZTEN ZU UNTERTEILENDEN GRUPPE*) / 3540
WRITE (NOUT,16) / 3550
16 FORMAT( / 3560
F* NFG(I), ANZAHL DER UNTERGRUPPEN PRO ENERGIEGRUPPE* / 3570
G* NFI(I), I=1,NGRE) ANZAHL DER ENERGIESTUETZPUNKTE PRO* / 3580
H* UNTERGRUPPE* // 3590
I* DIE LETZTE KARTE LAUTET : * // 3600
1* 35. KARTE* / 3610
2* @ENDE@ 17 KONSTANTE* // 3620
3* DIE KARTEN 1 BIS 35 SIND FUER JEDES ZU BEHANDELNDE MATERIAL* / 3630
4* ZU WIEDERHOLEN. DIE LETZTE EINGABEKARTE DES JOBS LAUTET : * // 3640
5* @ENDE@ 18* // 3650
WRITE (NOUT,21) / 3660
21 FORMAT( / 3670
1* DAS PROGRAMMSYSTEM MIGROS BENOETIGT DD-KARTEN FUER * / 3680

2* FOLGENDE EXTERNE SPEICHER * // 3690
3* 8 EINHEIT, AUF DIE FREEFO DIE ENTSCHLUESSELTE EINGABE* / 3700
4* SCHREIBT* / 3710
5* 1 EINHEIT, AUF DER DIE KERNDATENBIBLIOTHEK STEHT* / 3720
6* 3 EINHEIT, AUF DIE ALLE ARBEITSPRCGRAMME IHRE INTERNE* / 3730
7* AUSGABE SCHREIBEN* / 3740
8* 10 ZWISCHENEINHEIT, DIE NUR VOM PRGGRAMM MIT DER KENNZIFFER 6* // 3750
9* BENOETIGT WIRD* // 3760
A* AUF DER EXEC - KARTE IST DIE ANGABE DES PARM.G - PARAMETERS* / 3770
B* ERFORDERLICH. DIE LAENGE DES BENOETIGTEN PARM.G - FELDES IST* / 3780
C* VON DER EINGABE UND DEN GEWUENSCHTEN ARBEITSPROGRAMMEN* / 3790
D* ABHAENGIG* / 3800
1* IM FOLGENDEN WERDEN DIE SYMBOLE DER EINGABEBESCHREIBUNG* / 3810
1* VERWENDET* // 3820
1* DIE STEUERPHASE BENOETIGT* // 3830
2* 2*NTYP+3*NMAT+NFE*(NSPEC+1)+NT+MI+1+3*NE+4*NA*(1+NE) WORTE* // 3840
WRITE (NOUT,22) / 3850
22 FORMAT(IH1/IH0/IH0/ / 3860
2* WOBEI FUER DEN FALL, DASS EINE DIESER GROESSEN NICHT IN DER* / 3870
1* EINGABE GESETZT IST, DIE WERTE DER STANDARDEINGABE BENUTZT* / 3880
2* WERDEN, ALSO: NTYP=7 , NFE=1 , NSPEC=1 , NT=3 , MI=7 , NMAT=1* // 3890
3* HIERZU IST DER PLATZBEDARF DES ANGELAUFENEN ARBEITSPROGRAMMS* / 3900
4* MIT DEM GROESSTEN PLATZBEDARF ZU ADDIEREN* // 3910
5* PLATZBEDARF DER EINZELNEN ARBEITSPROGRAMME* // 3920
6* KENNZIFFER PLATZBEDARF* // 3930
7* 1 7*MI+3*NE+7*MAX(300,ANZAHL DER RESONANZENERGIEN* / 3940
8* AUF KEDAK)+5*MAX(400,MAXIMALE ANZAHL DER ENERGIE- / 3950
8* STUETZPUNKTE, DIE FUER DIE INTEGRATION UEBER EINE* / 3960
8* ENERGIEGRUPPE BENOETIGT WIRD) WORTE*) / 3970
WRITE (NOUT,24) / 3980
24 FORMAT( / 3990
E* KOMMENTAR:* // 4000
1* DIE ANZAHL DER STUETZSTELLEN FUER DIE INTEGRATION* / 4010
2* KANN VOM BENUTZER NICHT IM VORAUS BESTIMMT WERDEN.* / 4020
3* FALLS DIE VORGEGEBENEN 400 SPEICHERWORTE NICHT* / 4030
4* AUSREICHEN, WIRD DIE DIMENSION DES SPEICHERPLATZES* / 4040
5* ERHOEHET, FALLS ES DAS PARM.G-FELD ERLAUBT. WENN DAS* // 4050
6* PARM.G-FELD NICHT AUSREICHEND GROSS IST, WIRD DAS* / 4060
8* ARBEITSPROGRAMM UEBERSPRUNGEN. EMPFOHLEN WIRD IN* / 4070
9* DIESEM FALLE EINE VERDOPPELUNG DES FUER DIE* / 4080
A* INTEGRATION VORGEGEHENEN FELDES IM PARM.G-FELD* // 4090
A* 2 35*MI WORTE* // 4100
B* 3 6*NE+16*MI*NE+2*MAX(1500,ANZAHL DER MUEL-WERTE AUF* / 4110
C* KEDAK IM GEWUENSCHTEN ENERGIEBEREICH)+6*MAX(1500,* / 4120
D* ANZAHL DER SGA-WERTE, ANZAHL DER SGN-WERTE,* / 4130
B* ANZAHL DER SGF-WERTE AUF KEDAK IM GEWUENSCHTEN* / 4140
E* ENERGIEBEREICH) WORTE* // 4150
F* 4 6*NE+2*MAX(1500,ANZAHL DER SGN-WERTE,ANZAHL DER* / 4160
G* SGF-WERTE AUF KEDAK IM GEWUENSCHTEN ENERGIEBEREICH)* // 4170
H* WORTE* // 4180
WRITE (NOUT,23) / 4190
23 FORMAT( / 4200
1* 5 12*NE+3*(NE+1)+3*MAX(700,ANZAHL DER SGI-WERTE AUF* / 4210
2* KEDAK OBERHALB DER SCHWELLE FUER INELASTISCHE* / 4220
3* STREUUNG)+3*MAX(25,ANZAHL DER INELASTISCHEN ANRE- / 4230
4* GUNGSNIVEAUS)+MAX(25*400,ANZAHL DER INELASTISCHEN* / 4240

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5* ANREGUNGSNIVEAUS*ANZAHL DER ENERGIESTUETZPUNKTE DES* / 4250
6* ANREGUNGSQUERSCHNITTES ZUM 1. NIVEAU)+MAX(400,' / 4260
7* ANZAHL DER SGI-WERTE AUF KEDAK OBERHALB DER* / 4270
8* SCHWELLE FUER INELASTISCHE STREUUNG) WORTE* // 4280
9* 6 5*MAX(400,ISM)+7*MAX(1000,ISD)+4*(NE+1)+4*NE+* / 4290
A* MAX(400*21,ISM*ICOS)+MAX(400*2,ISM*NECU)+6*MAX(2,* / 4300
B* NECU)*NE+2*ICOS+2*MAX(4000,NTT) WORTE* / 4310
C* DABEI BEDEUTET : * / 4320
D* ISM 300+MAXIMALZAHL VON ENERGIESTUETZSTELLEN DER* / 4330
E* SGNC AUF KEDAK IN VIER BENACHBARTEN ENERGIE- / 4340
G* ISD ISM+MAXIMALZAHL VON ENERGIESTUETZSTELLEN DER* / 4350
H* SGN ODER MUEL AUF KEDAK PRO ENERGIEGRUPPE* / 4360
F* GRUPPEN* / 4370
I* ICOS ANZAHL DER WINKELSTUETZSTELLEN DER SGNC AUF* / 4380
J* KEDAK* / 4390
WRITE (NOUT,25) / 4400
25 FORMAT(1H1/1H0/1H0/ / 4410
1* NECU MAXIMALE STREUBREITE IN ENERGIEGRUPPEN (=2* / 4420
2* FUER EINFACHE UEBERSTREUUNG)* // 4430
1* NTT ANZAHL DER SGT - STUETZPUNKTE FUER DIE GEGEBE- / 4440
2* NE MISCHUNG IN 4 BENACHBARTEN ENERGIEGRUPPEN* / 4450
3* 7 NE WORTE* // 4460
4* 8 3*NE+2*MAX(10*NE,NFE) WORTE* // 4470
5* 9 39+NMI*IMAX*(3+NUJM)+ICOS*(2*NIV)+NIV+IMAX* / 4480
6* +2*NJM+3*NDAT+NE+NFIN*IZV+BUF* / 4490
7* DABEI BEDEUTET: * / 4500
8* NMI=NLRE+1* / 4510
9* NFIN MAXIMALE ANZAHL DER FEININTERVALLE PRO* / 4520
A* ENERGIEGRUPPE*) / 4530
WRITE (NOUT,28) / 4540
28 FORMAT( / 4550
1* IZV = MAX(4,IMAX*NMI)* / 4560
B* NDAT MAXIMAL MOEGLICHE ANZAHL VON ENERGIESTUETZ- / 4570
C* PUNKTEN FUER TOTALE ODER ELASTISCHE QUER- / 4580
D* SCHNITTE IN EINER ENERGIEGRUPPE* / 4590
E* NIV ANZAHL DER ENERGIESTUETZSTELLEN DER SGNC AUF* / 4600
F* KEDAK* / 4610
G* ICOS ANZAHL DER WINKELSTUETZSTELLEN DER SGNC AUF* / 4620
H* KEDAK* / 4630
I* IMAX MAXIMALE STREUBREITE IN ENERGIEGRUPPEN (=2* / 4640
J* FUER EINFACHE UEBERSTREUUNG)*) / 4650
WRITE(NOUT,29) / 4660
29 FORMAT( / 4670
5* BUF = MAX(NDAT,2+IMAX*NFIN)* // 4680
J* 10 ES IST KEIN SPEICHERPLATZ ERFORDERLICH* /// 4690
K* DIE ANGABE DER REGION AUF DER JOB - KARTE SETZT SICH * / 4700
L* ZUSAMMEN AUS DEM PARM.G - FELD UMGERECHNET IN K - BYTES* / 4710
M* + 170K - BYTES FUER DAS MIGROS - PROGRAMM UND DIE PUFFER* /// 4720
1* EINGABEBEISPIEL* // 4730
2* //INR017MI JOB (0017,101,P6M1B),KRIEG,CLASS=A,TIME=10,* / 4740
3* // REGION=340K* / 4750
4* // EXEC PGM=MIGROS2,PARM=170000* / 4760
4* //STEPLIB DD DSN=LOAD.NUSYS,UNIT=3330,VOL=SER=NUSICE,DISP=SHR* / 4770
5* //FT08F001 DD UNIT=SYSDA,SPACE=(TRK,10)* / 4780
5* //FT06F001 DD SYSOUT=A,DCB=(RECFM=FBA,LRECL=133,BLKSIZE=931),* / 4790
5* // SPACE=(TRK,80,RLSE)* / 4800

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6* //FT01F001 DD UNIT=2314,VOL=SER=NUSYS,CSN=KNDF,DISP=SHR* / 4810
7* //FT03F001 DD UNIT=3330,VOL=SER=SCRCIC,DISP=(NEW,KEEP),* / 4820
8* // SPACE=(TRK,20,RLSE),DSN=GROUP.CONSTANT.OUTPUT,* / 4830
7* // DCB=(RECFM=VS,LRECL=32768,BLKSIZE=800)* / 4840
9* //FT10F001 DD UNIT=SYSDA,SPACE=(TRK,100)* / 4850
WRITE (NOUT,26) / 4860
26 FORMAT( / 4870
1* //FT05F001 DD * / 4880
F* @U 235@ 1* / 4890
G* @BLOC@ 4* / 4900
1* 7 1 18 18 2 14 14 4 14 14 5 5 3 7 25 1 8 25 1 10 26 26* / 4910
3* @ENDE@ 17* / 4920
2* @D 16@ 0* / 4930
6* @BLOC@ 4* / 4940
7* 5 3 4 4 6 3 3 6 12 12 9 3 3 9 12 12* / 4950
E* @ENDE@ 17* / 4960
4* @ENDE@ 18* /// 4970
5* ENDE DER EINGABEBESCHREIBUNG* // 4980
6* *****'/1H1) / 4990
RETURN / 5000
END / 5010

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C / 10
C ROUTINE FOR PRODUCING AN UNFORMATTED INPUT-FILE / 20
C / 30
C / 40
C SUBROUTINE FREEFO (INP,NFI,NFO,LF,F,NF) / 50
DIMENSION LF(1),F(1),NF(1),JZ(2) / 60
REAL*8 N9,NV8/5HNLFIN/,VC / 70
LCGICAL*1 JF(8),JX(2) / 80
INTEGER*2 NFE(80),LV(18),JY(4),LL,KJFE,STERN/2H* / 90
EQUIVALENCE (JZ(1),JF(1),JY(1),N8),(LL,JX(1)) / 100
DATA LV(1)/1H /,LV(2)/1H0/,LV(3)/1H1/,LV(4)/1H2/,LV(5)/1H3/, / 110
1LV(6)/1H4/,LV(7)/1H5/,LV(8)/1H6/,LV(9)/1H7/,LV(10)/1H8/, / 120
2LV(11)/1H9/,LV(12)/1H+/,LV(13)/1H+/,LV(14)/1H-/,LV(15)/1H./, / 130
3LV(16)/1HE/,LV(17)/1H@/,LV(18)/1H*/,LE/4HHEXA/,LFC/4HFORM/ / 140
4,LSPE/4HSPEC/,LNO/4HNORM/ / 150
C / 160
C / 170
C / 180
IY=80 / 190
GOTO 9111 / 200
C / 210
C / 220
ENTRY FREE72 (INP,NFI,NFO,LF,F,NF) / 230
IY=72 / 240
C / 250
C / 260
9111 V=1. / 270
MV=1 / 280
LPP=0 / 290
NF(1)=0 / 300

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```

LSU=0
LS=0
LP=0
NS=0
LO=0
N=0
LL=LV(1)/256
KSPNO=0
KOUT=0
C
C
C
33 IF(NF(1).EQ.LE) GOTO 2
   IF(NF(1).EQ.LFO) GOTO 2
   GOTO 201
200 KOUT=1
   GOTO 12
201 JZ(1)=NF(1)
   JZ(2)=NF(2)
   IF(N8.EQ.NV8) GOTO 200
   READ (INP,1,END=200,ERR=3) (NFE(I),I=1,80)
   1 FORMAT(80A1)
   GO TO 4
   2 IF(NFI)203,203,202
202 ENDFILE NFI
   REWIND NFI
203 RETURN
   3 WRITE (NFO,5)
   5 FORMAT(1H0/48H ERROR-CONDITION IN DATA TRANSFER OR INPUT-ERROR)
   STOP
   4 IF (IY.EQ.80) GOTO 6667
   JKFE=NFE(73)
   NFE(73)=STERN
6667 WRITE (NFO,6) (NFE(I),I=1,80)
   6 FORMAT(1X,80A1)
   IF (IY.EQ.80) GOTO 6668
   NFE(73)=JKFE
C
C
C
6668 IF(NF(1).EQ.LNO) GOTO 500
   IF(NF(1).EQ.LSPE) GOTO 501
   GOTO 502
500 KSPNO=0
   GOTO 11
501 KSPNO=1
   GOTO 11
502 IF(NFE(1).EQ.LV(1)) GOTO 10
   IF(N)11,11,12
   12 IF(NFI)13,13,144
144 IF(KSPNO)145,145,14
   14 WRITE (NFI) N,(NF(I),I=1,N)
111 IF(KOUT)11,11,2
145 WRITE (NFI) (NF(I),I=1,N)
   GOTO 111
13 NS=NS+1

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310 LF(NS)=N
320 N1=NS+1
330 N2=NS+N
340 N=0
350 DO 15 I=N1,N2
360 N=N+1
370 15 LF(I)=NF(N)
380 NS=N2
390 GOTO 111
400 11 N=0
410 J=0
420 GO TO 16
430 10 J=1
440 16 J=J+1
C
C
C
97 DO 20 K=1,18
   IF(NFE(J).EQ.LV(K)) GO TO 21
20 CONTINUE
   GO TO 3
C
C
C
21 IF(K-1)30,30,22
30 IF(LS)31,31,32
31 IF(J-IY)16,33,33
C
C
C
32 IF(LPP)40,40,41
40 N=N+1
   NF(N)=LSU*MV
47 LSU=0
   LS=0
   LO=0
   MV=1
   V=1.
   GOTO 31
C
C
C
41 M=LP-LS
   IF(LS-9)42,43,43
43 LSU=LSUR
42 IF(M)44,45,46
44 IF(78+M)3,3,45
46 IF(75-M)3,3,45
45 N=N+1
   VC=V
   F(N)=DFLOAT(LSU)*VC*10.**M
   LP=0
   LPP=0
   GO TO 47
C
C

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1010
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B 14

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C
22 IF(K-11)50,50,23
50 LS=LS+1
   LSU=10*LSU+K-2
   IF(LS-9)511,52,511
52 L SUR=LSU
511 IF(LPP)51,51,883
883 LC=-1
   J=J+1
   IF(J-IY)884,884,32
51 IF(J-IY)16,32,32

C
C
C
23 IF(K-14)60,60,24
60 IF(L0)61,61,3
61 L0=1
   IF(K-14)62,63,63
63 V=-1.
   MV=-1
62 IF(J-IY)64,3,3
64 J=J+1
   DC 65 K=2,11
   IF(NFE(J).EQ.LV(K)) GO TO 50
65 CONTINUE
   IF(NFE(J).EQ.LV(15)) GO TO 70
   GO TO 3

C
C
C
24 IF(K-15)70,70,25
70 IF(LP)71,71,3
71 LP=LS
   LPP=1
   IF(J-IY)72,73,73
73 IF(LS)3,3,41
72 J=J+1
   DO 74 K=2,11
   IF(NFE(J).EQ.LV(K)) GO TO 50
74 CONTINUE
   IF(NFE(J).EQ.LV(1)) GO TO 73
   IF(NFE(J).EQ.LV(16)) GO TO 81
   LC=0
884 LA=0
   LV1=1
   LP1=0
   IF(J-IY)882,882,3

C
C
C
25 IF(K-16)80,80,26
80 IF(LPP)3,3,81
81 LA=0
   LC=1
   LV1=1
   LP1=0

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1980

IF(J-IY)82,3,3
82 J=J+1
   IF(NFE(J).EQ.LV(1)) GO TO 83
882 IF(NFE(J).EQ.LV(12)) GO TO 83
   IF(NFE(J).EQ.LV(13)) GO TO 83
   IF(NFE(J).EQ.LV(14)) GO TO 84
   IF(LC)97,3,85
84 LV1=-1
83 IF(J-IY)86,3,3
86 J=J+1
85 DC 87 K=2,11
   IF(NFE(J).EQ.LV(K)) GO TO 88
87 CONTINUE
   IF(NFE(J).EQ.LV(1)) GO TO 89
   GO TO 3
89 IF(LA)3,3,90
88 LA=1
   LP1=10*LP1+K-2
   IF(J-IY)86,90,90
90 LP=LP+LP1*LV1
   GO TO 41

C
C
C
26 IF(K-17)300,300,301
300 M=5
   K7=17
   GO TO 117
301 M=4
   K7=18
117 LC=0
116 LA=0
   DO 100 L=1,4
100 JY(L)=LV(1)
110 J=J+1
   IF(J-IY)101,102,102
102 IF(NFE(J).EQ.LV(K7)) GOTO 120
   LC=0
   GOTO 121
120 J=J-1
121 IF(LC)33,3,112
101 IF(NFE(J).EQ.LV(K7)) GO TO 106
   GO TO 107
106 IF(LC)105,3,102
107 LA=LA+1
   LC=1
   LL=NFE(J)
   JF(LA)=JX(1)
   IF(LA-M)110,112,112
112 N=N+1
   NF(N)=JZ(1)
   IF(K-17)433,433,434
433 N=N+1
   NF(N)=JZ(2)
434 LC=-1
   IF(NFE(J+1).EQ.LV(K7)) GOTO 110

```

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GOTO 116
105 IF(NFE(J+1).EQ.LV(1)) GOTO 16
GO TO 3
END

```

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2560
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```

DC 9 I=2,NE
IF(ENG(I).LT.ENG(I-1)) GO TO 10
9 CONTINUE
GC TO 14
10 WRITE (NOUT,12) (ENG(I),I=1,NE)
12 FORMAT(1H0/' ***ERROR 0. 3 : THE ENERGY GROUP BOUNDARIES'/(8E16.8)
1)
WRITE (NOUT,13)
13 FORMAT(' ARE NOT GIVEN IN AN INCREASING ORDER')
STOP
8 NE=26
ENG(1)=0.215
ENG(2)=0.465
ENG(3)=1.
DO 15 K=1,5
DC 15 J=1,3
MO=3*(K-1)+J
15 ENG(MO+3)=ENG(MO)*10.
ENG(19)=0.2E6
ENG(20)=0.4E6
ENG(21)=0.8E6
ENG(22)=1.4E6
ENG(23)=2.5E6
ENG(24)=4.0E6
ENG(25)=6.5E6
ENG(26)=1.0E7
14 IF(MAT2-3)1,16,17
16 READ (NF) LSP,NFE,((EF(I,J),I=1,NFE),J=1,NSPEC)
READ (NF) MAT1,MAT2
DO 19 I=2,NFE
IF(EF(I,1).LE.EF(I-1,1)) GO TO 19
18 CCNTINUE
GO TO 20
19 WRITE (NOUT,21)
21 FORMAT(1H0/' ***ERROR 0. 4 : THE ENERGY PCINTS OF THE MACROSCOPIC
1WEIGHTING FUNCTION ARE NOT GIVEN IN AN INCREASING ORDER')
STOP
17 NFE=1
LSP=0
20 IF(MAT2-4)1,22,23
22 READ (NF) NA,(NR(1,I),NR(2,I),NR(3,I),I=1,NA)
DO 11 I=1,NA
IF(NR(2,I).GE.NR(3,I).AND.NR(2,I).LE.(NE-1)) GO TO 11
IF(NR(1,I).EQ.10.AND.NR(3,I).LE.NE) GO TO 11
WRITE (NOUT,54) NR(1,I)
54 FORMAT(1H0/' ***ERROR 0. 5 : THE CHOICE OF ENERGY GROUP BOUNDARIES
1 FOR PROGRAM',16,' IS NOT VALID')
STOP
11 CONTINUE
IF(NFE.EQ.1) GO TO 114
MIN=1
MAX=NE-1
DO 112 I=1,NA
IF(NR(2,I).GT.MIN) MIN=NR(2,I)
IF(NR(3,I).LT.MAX) MAX=NR(3,I)
112 CCNTINUE

```

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810
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830
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1010
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PRGRAMM ZUM LESEN UND PRUEFEN DER EINGABE FUER DAS MIGROS-SYSTEM
SUBROUTINE INPUT (ITYP,NAME,TZ,ENG,EF,NR,TEMP,SIGC,NRE,NN1,NN2,
1NFG,NFI,NG,NI,NMAT,MI,NE,NSPEC,NFE,NA,NT,NTYP,NFR,XL,NF,NX,IADR)
REAL*8 MAT1, MAT,ITYP,N2,NAME,SGC,SGA,SGF
DIMENSION ENG(NE),EF(NFE,NSPEC),NAME(NMAT),TZ(NMAT),MAZ(2),
1 NR(3,NA),TEMP(NT),SIGO(MI),ITYP(NTYP),
1 INZ(5),IPM(2), NRE(NA),NN1(NA,NE),NN2(NA,NE),NFG(NA,NE),NFI(NA,NE)
2,XL(1),NG(NE),NI(NE)
COMMON MAT,ISTRUK,ISPA,NOUT,LIZ,NANF,NEND,KL
EQUIVALENCE (MAT,IPM(1))
DATA INZ/'0','2','4','6','8',/N2/'ENDE' /
1,N3/' ',/SGC/'SGC' '/',SGA/'SGA' '/',SGF/'SGF' '/'
CALL FSPIE
MDIM=1500
NER=1500
LDIM=1500
IRE=300
ISTE=400
IWE=400
ISG=400
NET=700
NAE=25
NTT=4000
NMAX=5
ICOS=21
ISM=400
ISD = 1000
NECU = 2
ISCO=ISM*ICOS
ISEC=ISM*NECU
IREMO=0
READ (NF) MAT,ISTRUK
READ (NF) MAT1,MAT2
IF(MAT2-1)1,2,3
2 READ (NF) ISPA
READ (NF) MAT1,MAT2
GO TO 4
3 DO 5 I=1,5
IF(IPM(2).EQ.INZ(I)) GO TO 6
5 CONTINUE
ISPA=1
GO TO 4
6 ISPA=0
4 IF(MAT2-2)1,7,8
7 READ (NF) NE,(ENG(I),I=1,NE)
READ (NF) MAT1,MAT2

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AN=ENG(NE-MIN)	1040	35 DO 39 I=1,NA	1600
END=ENG(NE-MAX+1)	1050	IF(NR(1,I).EQ.1.OR.NR(1,I).EQ.2.OR.NR(1,I).EQ.3) GO TO 40	1610
IF(EF(1,1).GE.AN) GO TO 113	1060	39 CONTINUE	1620
IF(EF(NFE,1).GT.END) GO TO 114	1070	MI=0	1630
113 WRITE (NOUT,115)	1080	GO TO 37	1640
115 FORMAT(1H0/' ***ERROR 0. 6 : THE MACROSCOPIC WEIGHTING FUNCTION DO	1090	40 MI=7	1650
IES NOT INCLUDE ALL ENERGY GROUPS DESIRED')	1100	SIGO(1)=0.	1660
STCP	1110	DO 41 I=1,6	1670
114 DO 116 I=1,NA	1120	41 SIGO(I+1)=10.**I	1680
IF(NR(1,I).EQ.3.AND.IPM(2).EQ.N3) GO TO 117	1130	37 IF(MAT2-7)1,42,43	1690
GO TO 116	1140	42 READ (NF) NTYP,(ITYP(I),I=1,NTYP)	1700
117 WRITE (NOUT,118) MAT	1150	READ (NF)MAT1,MAT2	1710
118 FORMAT(1H0/' ***WARNING 0. 1 : THE SELFSHIELDING FACTORS FOR ',A6,	1160	DO 149 I=1,NTYP	1720
1' CAN ONLY BE CALCULATED FROM RESONANCE DATA/' THE MODULE NUMBER	1170	IF(ITYP(I).EQ.SGC) GO TO 150	1730
2WAS MODIFIED APPROPRIATE')	1180	149 CCNTINUE	1740
NR(1,I)=1	1190	GO TO 151	1750
116 CONTINUE	1200	150 DO 152 J=1,NTYP	1760
READ (NF) MAT1,MAT2	1210	IF(ITYP(J).EQ.SGA) GO TO 153	1770
GO TO 24	1220	152 CCNTINUE	1780
23 WRITE (NOUT,25)	1230	ITYP(I)=SGA	1790
25 FORMAT(1H0/' ***ERROR 0. 7 : THE IDENTIFICATION NUMBERS OF MODULES	1240	153 DO 154 J=1,NTYP	1800
1 TO BE CALLED CAN NOT BE FOUND')	1250	IF(ITYP(J).EQ.SGF) GO TO 155	1810
STOP	1260	154 CONTINUE	1820
24 IF(MAT2-5)1,26,27	1270	IF(ITYP(I).EQ.SGA) GO TO 156	1830
26 READ (NF) MT,(TEMP(I),I=1,MT)	1280	ITYP(I)=SGF	1840
READ (NF) MAT1,MAT2	1290	GO TO 151	1850
DO 28 I=1,NA	1300	156 NTYP=NTYP+1	1860
IF(NR(1,I).EQ.1.OR.NR(1,I).EQ.2) GO TO 25	1310	ITYP(NTYP)=SGF	1870
28 CONTINUE	1320	GC TO 151	1880
WRITE (NOUT,30)	1330	155 IF(ITYP(I).EQ.SGA) GO TO 151	1890
30 FORMAT(1H0/' ***WARNING 0. 2 : THE DECLARATION OF THE TEMPERATURE	1340	NTYP=NTYP-1	1900
1IS NOT NECESSARY AND IS IGNORED')	1350	DO 157 J=I,NTYP	1910
GO TO 29	1360	157 ITYP(J)=ITYP(J+1)	1920
27 DO 31 I=1,NA	1370	151 DO 44 I=1,NA	1930
IF(NR(1,I).EQ.1.OR.NR(1,I).EQ.2) GO TO 32	1380	IF(NR(1,I).EQ.4) GO TO 45	1940
31 CONTINUE	1390	44 CONTINUE	1950
NT=0	1400	WRITE (NOUT,46)	1960
GO TO 29	1410	46 FORMAT(1H0/' ***WARNING 0. 4 : THE DECLARATION OF THE CROSS SECTION	1970
32 IF(ISTRUK.EQ.0) GO TO 34	1420	IN TYPES IS NOT NECESSARY AND IS IGNORED')	1980
NT=3	1430	GC TO 45	1990
TEMP(1)=300.	1440	43 DO 47 I=1,NA	2000
TEMP(2)=900.	1450	IF(NR(1,I).EQ.4.OR.NR(1,I).EQ.10) GO TO 48	2010
TEMP(3)=2100.	1460	47 CONTINUE	2020
GC TO 29	1470	NTYP=0	2030
3→ NT=1	1480	GC TO 45	2040
TEMP(1)=0.	1490	48 NTYP=7	2050
29 IF(MAT2-6)1,33,35	1500	CALL DOPW (8HMUEL ,ITYP(1))	2060
33 READ (NF) MI,(SIGO(I),I=1,MI)	1510	CALL DOPW (8HNUE ,ITYP(2))	2070
READ (NF) MAT1,MAT2	1520	CALL DOPW (8HSGA ,ITYP(3))	2080
DO 36 I=1,NA	1530	CALL DOPW (8HSGF ,ITYP(4))	2090
IF(NR(1,I).EQ.1.OR.NR(1,I).EQ.2.OR.NR(1,I).EQ.3) GO TO 37	1540	CALL DOPW (8HSGI ,ITYP(5))	2100
36 CONTINUE	1550	CALL DOPW (8HSGN ,ITYP(6))	2110
WRITE (NOUT,38)	1560	CALL DOPW (8HSG2N ,ITYP(7))	2120
38 FORMAT(1H0/' ***WARNING 0. 3 : THE DECLARATION OF THE BACKGROUND C	1570	45 IF(MAT2-8)1,49,50	2130
CROSS SECTIONS IS NOT NECESSARY AND IS IGNORED')	1580	49 READ (NF) NRES	2140
GO TO 37	1590	READ (NF) MAT1,MAT2	2150

DO 51 I=1,NA	2160	133 IF(IP.EQ.1) GO TC 135	2720
IF(NR(1,I).EQ.1) GO TO 52	2170	WRITE (NOUT,83)	2730
51 CCNTINUE	2180	135 READ (NF) NMAT,(NAME(I),TZ(I),I=1,NMAT)	2740
WRITE (NOUT,53)	2190	READ (NF) MAT1,MAT2	2750
53 FORMAT(1H0/' ***WARNING 0. 5 : THE DECLARATION OF THE NUMBER OF RE	2200	NTK=1	2760
ISONANCES IS NOT NECESSARY AND IS IGNCREd')	2210	GO TO 136	2770
GO TO 52	2220	134 NTK=0	2780
50 NRES=10	2230	136 IF(MAT2-13)1,137,140	2790
52 IF(MAT2-9)1,56,57	2240	137 IF(IP.EQ.1) GO TO 139	2800
56 READ (NF) ERROR	2250	WRITE(NOUT,83)	2810
READ (NF) MAT1,MAT2	2260	139 READ (NF) MAZ(1),MAZ(2)	2820
DC 58 I=1,NA	2270	READ (NF) MAT1,MAT2	2830
IF(NR(1,I).EQ.1) GO TO 59	2280	IF(MAZ(1).EQ.0) GO TO 138	2840
58 CONTINUE	2290	IF(LSP.EQ.0) GO TO 138	2850
WRITE (NOUT,60)	2300	IF(LSP-(NLE+1))142,138,138	2860
60 FORMAT(1H0/' ***WARNING 0. 6 : THE DECLARATION OF THE ALLOWED INTE	2310	142 LP=NLE+1	2870
GRATION ERROR IS NOT NECESSARY AND IS IGNCREd')	2320	WRITE (NOUT,143) LP	2880
GO TO 59	2330	143 FORMAT(1H0/' ***ERRDR 0. 8 : THE NUMBER OF WEIGHTING FUNCTIONS IS	2890
57 ERROR=0.05	2340	'LESS THAN THE NUMBER OF LEGENDRE MOMENTS.'/15,' NUMBER OF WEIGHTIN	2900
59 IF(MAT2-10)1,63,64	2350	2G FUNCTIONS ARE NEEDED')	2910
63 READ (NF) XNUE	2360	STOP	2920
READ (NF) MAT1,MAT2	2370	140 MAZ(1)=0	2930
DO 65 I=1,NA	2380	MAZ(2)=0	2940
IF(NR(1,I).EQ.5) GO TO 66	2390	138 IF(MAT2-14)1,145,146	2950
65 CONTINUE	2400	145 IF(IP.EQ.1) GO TO 147	2960
WRITE (NOUT,67)	2410	WRITE(NOUT,83)	2970
67 FORMAT(1H0/' ***WARNING 0. 7 : THE DECLARATION OF THE PARAMETER FO	2420	147 READ(NF)NS,NK,NRR	2980
1R THE WEISSKOPF EVAPORATION MODEL IS NOT NECESSARY AND IS IGNORED'	2430	READ(NF)MAT1,MAT2	2990
2)	2440	GO TO 148	3000
GO TO 56	2450	146 NS=4	3010
64 XNUE=0.16	2460	NK=70	3020
66 DO 77 I=1,NA	2470	NRR=16	3030
IF(NR(1,I).EQ.6) GO TO 78	2480	148 DO 68 I=1,NA	3040
77 CONTINUE	2490	IF(NR(1,I).EQ.9) GO TO 69	3050
IP=0	2500	68 CONTINUE	3060
GO TO 79	2510	IP=0	3070
78 IP=1	2520	GO TO 70	3080
LSP=NF*LSPE	2530	69 IP=1	3090
79 IF(MAT2-11)1,80,81	2540	70 IF(MAT2-15)1,71,72	3100
80 IF(IP.EQ.1) GO TO 82	2550	71 IF(IP.EQ.1) GO TO 73	3110
WRITE (NOUT,83)	2560	WRITE (NOUT,74)	3120
83 FORMAT(1H0/' ***WARNING 0. 8 : THE INPUT FOR MODULE 6 IS NOT NECES	2570	74 FCRMAT(1H0/' ***WARNING 0. 9 : THE INPUT FOR MODULE 9 IS NOT NECES	3130
SARY')	2580	ISARY')	3140
82 READ (NF) ISEL,NLA,NLE	2590	73 READ (NF) ERR,NJM,NUJM	3150
READ (NF) MAT1,MAT2	2600	READ (NF) MAT1,MAT2	3160
NLA=0	2610	GO TO 75	3170
IF(NLE.LT.1) NLE=1	2620	72 IF(IP.EQ.0) GO TO 75	3180
IF(NLE.GT.5) NLE=5	2630	ERR=0.05	3190
GO TO 84	2640	NJM=6	3200
81 NLA=0	2650	NUJM=10	3210
NLE=5	2660	75 IF(MAT2-16)1,76,103	3220
IF(ISTRUK.EQ.0) ISEL=2	2670	76 IF(IP.EQ.1) GO TO 104	3230
IF(ISTRUK.EQ.1) ISEL=1	2680	WRITE (NOUT,74)	3240
84 IF(IP.EQ.0) GO TO 141	2690	104 READ (NF) ISELR,NLRA,NLRE,NCALL	3250
IF(NFE.GT.0.AND.EF(1,1).EQ.0) EF(1,1)=1.E-3	2700	IF(NCALL.NE.0) GC TO 158	3260
141 IF(MAT2-12)1,133,134	2710	READ(NF)MAT1,MAT2	3270

GO TO 105	3280	CALL FGEM(MI,SIGO,NE,ENG,NFE,EF(1,1),EF(1,2),NT,TEMP,ERROR,NRES,	3840
158 IS=0	3290	1NGR,XL(NFR),XL(NFR+7*MI),IRE,IREF,XL(NFR+7*MI+3*NE),XL(NFR+7*MI+	3850
DO 109 I=1,NA	3300	23*NE+IRE),XL(NFR+7*MI+3*NE+IRE*2),XL(NFR+7*MI+3*NE+3*IRE),	3860
IF(NR(1,I).EQ.9) IS=IS+1	3310	3XL(NFR+7*MI+3*NE+4*IRE),XL(NFR+7*MI+3*NE+5*IRE),	3870
109 CONTINUE	3320	4XL(NFR+7*MI+3*NE+6*IRE),ISTE,ISTEP,XL(NFR+7*MI+3*NE+7*IRE))	3880
IF(IS.EQ.NCALL) GO TO 110	3330	IF(IREF.GT.0.OR.ISTEP.GT.0) GO TO 126	3890
WRITE(NOUT,111) IS,NCALL	3340	IF(NGR.EQ.0) GO TO 101	3900
111 FORMAT(1H0/' ***ERROR 0.9 : MODULE 9 IS CALLED ',I3,' TIMES, IN T	3350	NANF=NE-NGR	3910
HE INPUT BLOC 16 NCALL IS SET TO',I3)	3360	KL=KL-1	3920
STOP	3370	GO TO 91	3930
110 DO 106 I=1,NCALL	3380	101 IF(KL.LE.NA) GO TO 100	3940
READ(NF)NGRE,(NN1(I,J),NN2(I,J),NFG(I,J),NFI(I,J),J=1,NGRE)	3390	GO TO 102	3950
NRE(I)=NGRE	3400	91 IF(NFR+35*MI.GT.IADR) GO TO 121	3960
DO 106 J=1,NGRE	3410	CALL FSTAT(MI,SIGO,NE,ENG,NFE,EF(1,1),EF(1,2),NT,TEMP,XL(NFR),	3970
IF(NN1(I,J).GE.NN2(I,J)) GO TO 106	3420	1XL(NFR+15*MI),XL(NFR+30*MI))	3980
WRITE(NOUT,107) I	3430	NGR=0	3990
107 FORMAT(1H0/' ***ERROR 0.10 : IN BLOC 16 THE INPUT OF THE ENERGY GR	3440	IF(KL.LE.NA) GO TO 100	4000
10UPS IS NOT VALID FOR THE ',I3,',' PROGRAM CALL')	3450	GO TO 102	4010
STOP	3460	125 MDIM=MDIM+MDIMP	4020
106 CONTINUE	3470	NER=NER+NERP	4030
READ(NF)MAT1,MAT2	3480	92 ISIG=0	4040
GO TO 108	3490	IF(SIGO(MI).GE.1.E6) GO TO 144	4050
103 NLRA=0	3500	ISIG=1	4060
NLRE=5	3510	MI=MI+1	4070
ISELR=1	3520	SIGO(MI)=1.E6	4080
NCALL=0	3530	144 IF(NFR+6*NE+16*MI*NE+2*MDIM+6*NER.GT.IADR) GO TO 121	4090
105 NFG(1,1)=14	3540	CALL FSTRUK(MI,SIGO,NE,ENG,NFE,EF(1,1),EF(1,2),XL(NFR),XL(NFR+4*NE	4100
NFI(1,1)=5	3550	1),XL(NFR+4*NE+5*MI*NE),XL(NFR+4*NE+9*MI*NE),XL(NFR+4*NE+15*MI*NE),	4110
NRE(1)=1	3560	2XL(NFR+5*NE+15*MI*NE),	4120
108 IF(MAT1.EQ.N2) GO TO 88	3570	3XL(NFR+5*NE+16*MI*NE),MDIM,MDIMP,XL(NFR+6*NE+16*MI*NE),	4130
1 WRITE(NOUT,89)	3580	4XL(NFR+6*NE+16*MI*NE+MDIM),NER,NERP,XL(NFR+6*NE+16*MI*NE+2*MDIM),	4140
89 FORMAT(1H0/' ***ERROR 0.11 : THE INPUT IS NOT ORDERED IN INCREASIN	3590	5XL(NFR+6*NE+16*MI*NE+2*MDIM+3*NER))	4150
IG ARGUMENTS')	3600	IF(ISIG.EQ.1) MI=MI-1	4160
STOP	3610	IF(MDIMP.GT.0.OR.NERP.GT.0) GO TO 125	4170
88 NGR=0	3620	IF(KL.LE.NA) GO TO 100	4180
KL=1	3630	GO TO 102	4190
100 NANF=NR(2,KL)	3640	124 LDIM=LDIM+LDIMP	4200
NEND=NR(3,KL)	3650	93 IF(NFR+6*NE+2*LDIM.GT.IADR) GO TO 121	4210
IF(NR(1,KL).EQ.1) GO TO 90	3660	CALL SUND(NE,ENG,NFE,EF(1,1),EF(1,2),NTYP,ITYP,XL(NFR),	4220
IF(NR(1,KL).EQ.2) GO TO 91	3670	1XL(NFR+NE),XL(NFR+2*NE),XL(NFR+3*NE),XL(NFR+4*NE),XL(NFR+5*NE),	4230
IF(NR(1,KL).EQ.3) GO TO 92	3680	2LDIM,LDIMP,XL(NFR+6*NE),XL(NFR+6*NE+LDIM))	4240
IF(NR(1,KL).EQ.4) GO TO 93	3690	IF(LDIMP.GT.0) GO TO 124	4250
IF(NR(1,KL).EQ.5) GO TO 94	3700	IF(KL.LE.NA) GO TO 100	4260
IF(NR(1,KL).EQ.6) GO TO 95	3710	GO TO 102	4270
IF(NR(1,KL).EQ.7) GO TO 96	3720	127 NET=NET+NETP	4280
IF(NR(1,KL).EQ.8) GO TO 97	3730	NAE=NAE+NAEP	4290
IF(NR(1,KL).EQ.9) GO TO 119	3740	ISG=ISG+ISGP	4300
IF(NR(1,KL).EQ.10) GO TO 120	3750	IWE=IWE+IWP	4310
WRITE(NOUT,98) NR(1,KL)	3760	94 L=NFR+8*NE+2*NX+2*NET+2*NAE	4320
98 FORMAT(1H0/' ***ERROR 0.12 : THE DESIRED MODULE',I6,' IS NOT CONTA	3770	IF(L+4*NE+NX+NET+NAE*(ISG+1)+IWE.GT.IADR) GO TO 121	4330
INED IN MIGROS')	3780	CALL SCAT(NE,ENG, XNUE,NFE,EF(1,1),EF(1,2),NX,XL(NFR),	4340
STOP	3790	1XL(NFR+2*NE),XL(NFR+4*NE),XL(NFR+6*NE),XL(NFR+8*NE),NET,	4350
126 IRE=IRE+IREP	3800	2XL(NFR+8*NE+2*NX),NAE,XL(NFR+8*NE+2*NX+2*NET),XL(L),XL(L+NE),	4360
ISTE=ISTE+ISTEP	3810	3XL(L+2*NE),XL(L+3*NE),XL(L+4*NE),NETP,XL(L+4*NE+NX),NAEP,	4370
90 NGR=0	3820	4XL(L+4*NE+NX+NET),ISG,ISGP,XL(L+4*NE+NX+NET+NAE),IWE,IWP,	4380
IF(NFR+7*MI+3*NE+7*IRE+5*ISTE.GT.IADR) GO TO 121	3830	5XL(L+4*NE+NX+NET+NAE+NAE*ISG))	4390


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4 FORMAT(1H0/' ***ERROR NDF. 1 : THE DD-CARD FOR UNIT 1 DOES NOT CHA 310
  1RACTERIZE A VALID KEDAK LIBRARY') 320
  STOP 330
1 CONTINUE 340
  IFD=ISATZ(4) 350
  IAD(1)=ISATZ(5) 360
  IAD(2)=ISATZ(8) 370
  IAD(3)=ISATZ(11) 380
  K=4 390
  IR(1)=ISATZ(6) 400
  IR(2)=ISATZ(9) 410
  IR(3)=ISATZ(12) 420
  IW(1)=ISATZ(7) 430
  IW(2)=ISATZ(10) 440
  IW(3)=ISATZ(13) 450
  DO 3 J=1,3 460
  N=IR(J) 470
  IWJ=IW(J) 480
  IF(IS=N-1)5,6,5 490
5 READ (LBN*N) (ISATZ(II),II=1,NSZ) 500
  IS=N+1 510
6 IF(J-3)10,11,11 520
10 L=3 530
  GO TO 326 540
11 L=4 550
326 IMP=IAD(J)*L+IWJ-1 560
  8 IF(IMP-NSZ)13,14,15 570
14 N=1 580
  GO TO 16 590
15 N=2 600
16 DO 12 L=IWJ,NSZ 610
  IAD(K)=ISATZ(L) 620
12 K=K+1 630
  GO TO (3,17),N 640
17 IMP=IMP-NSZ 650
  IWJ=1 660
  READ (LBN*IS) (ISATZ(II),II=1,NSZ) 670
  IS=IS+1 680
  GO TO 8 690
13 DO318 L=IWJ,IMP 700
  IAD(K)=ISATZ(L) 710
318 K=K+1 720
  3 CONTINUE 730
  RETURN 740
C 750
C 760
C SUBROUTINEN NDFLOC , LDFLOC , IDFLOC 770
C 780
  ENTRY NDFLOC (KONTR,NNAM,DAT1,ID,KC) 790
  IF(NSZ.NE.880) GO TO 1004 800
  ID=0 810
  KC=0 820
  DO 403 LS=1,2 830
  IF (DAT1(LS).NE.MNAM(LS)) GO TO 104 840
403 CONTINUE 850
  MGL=2 860

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  GOTO 107 870
104 MGL=1 880
107 J=NNAM(1)*2 890
  DO 51 N=1,J 900
  51 Z(N)=DAT1(N) 910
  L=1 920
  LS=1 930
  DO 18 N=1,J 940
  IF(N-4)19,218,20 950
  19 IF(N-3)20,218,218 960
  20 MNAM(L)=KDAT1(N) 970
  L=L+1 980
  GOTO 18 990
213 IBEST(LS)=KDAT1(N) 1000
  LS=LS+1 1010
  18 CCNTINUE 1020
  KK=1 1030
  GOTO427 1040
C 1050
  ENTRY LDFLOC (KONTR,NNAM,INAM,DAT2) 1060
  IF(NSZ.NE.880) GO TO 1004 1070
  DO 127 LS=1,2 1080
  IF (INAM(LS).NE.MNAM(LS)) GO TO 128 1090
127 CCNTINUE 1100
  MGL=2 1110
  GOTO 129 1120
128 MGL=1 1130
129 J=NNAM(1)*2 1140
  DO 21 N=1,J 1150
  21 MNAM(N)=INAM(N) 1160
  KK=2 1170
427 K=4 1180
  IWJ=0 1190
  DO 22 LS=1,3,2 1200
  IWJ=IWJ+1 1210
  N=IAD(IWJ)*3 1220
  DO 23 M=1,N 1230
  IF (MNAM(LS).NE.IAD(K)) GO TO 23 1240
  IF (MNAM(LS+1).EQ.IAD(K+1)) GO TO 26 1250
  23 K=K+3 1260
  WRITE (NOUTP,2000) (MNAM(M),M=1,4) 1270
2000 FORMAT(1H0/' ***WARNING NDF. 1 : THE DATA FOR ',2A4,1X,2A4,' ARE N 1280
  10T INCLUDED IN THE'/ ' CONVERSIONTABLE OF THE KEDAK LIBRARY') 1290
  MNAM(1)=1 1300
  KONTR=0 1310
  IF(JJ-1)98,98,96 1320
  26 NLNA(IWJ)=IAD(K+2) 1330
  22 K=N+4 1340
  GO TO (227,228),MGL 1350
C 1360
  ENTRY IDFLOC (KONTR,NNAM,INAM,DAT2) 1370
  IF(NSZ.NE.880) GO TO 1004 1380
  KK=3 1390
  IF(INAM(1)-NUNA(1))130,131,130 1400
  131 MGL=2 1410
  GOTO 132 1420

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130	MGL=1	1430	49	NNK=NWP	1990
132	NUNA(1)=INAM(1)	1440		NWR=IDR	2000
	NUNA(2)=INAM(2)	1450		NWW=IDW	2010
	IF(NNAM(1)-2)27,27,52	1460		IF(IS=NWR-1)44,45,44	2020
52	J=NNAM(1)	1470	44	READ(LBN*NWR)(IWNA(II),II=1,NSZ)	2030
	DO 53 LS=3,J	1480		NWR=NWR+1	2040
53	MNAM(2*LS-1)=INAM(LS)	1490		GO TO 62	2050
27	GO TO (227,228),MGL	1500	45	DO 61 LS=1,NSZ	2060
227	N=IAD(3)*4	1510	61	IWNA(LS)=ISATZ(LS)	2070
	IWJ=(IAD(1)+IAD(2))*3+4	1520		NWR=NWR+1	2080
	DO 28 LS=1,N,4	1530	62	DC 68 N=1,NNK	2090
	IF(NUNA(1)-IAD(IWJ))28,29,28	1540		NKC=N	2100
28	IWJ=IWJ+4	1550		NW=NWW	2110
	KONTR=0	1560		KP=5	2120
	GO TO 24	1570		JD=3	2130
29	NT=IAD(IWJ+1)	1580		DO 54 LS=1,NWN	2140
	JR=IAD(IWJ+2)	1590		IF(IWNA(NWW)-MNAM(KP))58,55,56	2150
	JW=IAD(IWJ+3)	1600	56	KONTR=0	2160
228	KR=JR	1610		IVY=1	2170
	KW=JW	1620		IF(KK-2)57,59,859	2180
	IVY=0	1630	59	INAM(KP)=IWNA(NWW)	2190
	IF(IS-KR-1)30,31,30	1640		GOTO 60	2200
30	READ(LBN*KR)(ISATZ(II),II=1,NSZ)	1650	859	INAM(JD)=IWNA(NWW)	2210
	IS=KR+1	1660		GO TO 60	2220
31	DO 32 LS=1,NT	1670	57	DAT1(KP+2)=XWNA(NWW)	2230
	IF(NUNA(2)-ISATZ(KW))33,34,33	1680		GOTO 60	2240
33	KW=KW+7	1690	55	KONTR=1	2250
	IF(KW-NSZ)32,32,35	1700	60	KP=KP+2	2260
35	READ(LBN*IS)(ISATZ(II),II=1,NSZ)	1710		JD=JD+1	2270
	IS=IS+1	1720		NWW=NWW+1	2280
	KW=KW-NSZ	1730		IF(NWW-NSZ)54,54,64	2290
32	CONTINUE	1740	64	READ(LBN*NWR)(IWNA(II),II=1,NSZ)	2300
	KONTR=0	1750		NWR=NWR+1	2310
	GO TO 24	1760		NWW=1	2320
34	KW=KW+1	1770	54	CONTINUE	2330
	DO 36 LS=1,6	1780		GO TO 74	2340
	IF(KW-NSZ)37,37,38	1790	58	NWW=NW+NWN+3	2350
38	READ(LBN*IS)(ISATZ(II),II=1,NSZ)	1800		IF(NWW-NSZ)68,68,70	2360
	IS=IS+1	1810	70	READ(LBN*NWR)(IWNA(II),II=1,NSZ)	2370
	KW=1	1820		NWR=NWR+1	2380
37	GO TO (39,40,40,41,42,43),LS	1830	73	NWW=NWW-NSZ	2390
39	NWN=ISATZ(KW)	1840	68	CONTINUE	2400
	GOTO 36	1850		KONTR=0	2410
40	NNAM(LS)=ISATZ(KW)	1860		IF(KK.EQ.3) GO TO 580	2420
	GOTO 36	1870		WRITE(NOUTP,2002) MNAM(5),(MNAM(II),II=1,4)	2430
41	NWP=ISATZ(KW)	1880	2002	FORMAT(1H0/' ***WARNING NDF. 2 : THE FURTHER NAME ',E16.8,' IS GRE	2440
	GO TO 36	1890		1ATER THAN THE GREATEST FURTHER NAME '/' INCLUDED IN THE KEDAK LIBR	2450
42	IDR=ISATZ(KW)	1900		2ARY FOR',1X,2A4,1X,2A4)	2460
	GO TO 36	1910		GO TO 98	2470
43	IDW=ISATZ(KW)	1920	580	WRITE(NOUTP,581) MNAM(5),(INAM(II),II=1,2)	2480
36	KW=KW+1	1930		581 FORMAT(1H0/' ***WARNING NDF. 2 : THE FURTHER NAME ',E16.8,' IS GREA	2490
	IF(KK-1)46,46,47	1940		1TER THAN THE GREATEST FURTHER NAME '/' INCLUDED IN THE KEDAK LIBRAR	2500
46	NNAM(4)=0	1950		2Y FOR',2I10)	2510
47	DO 80 LS=1,4	1960		GO TO 98	2520
80	NN(LS)=NNAM(LS)	1970	1004	WRITE(NOUTP,1005)	2530
	IF(NWN)384,384,49	1980	1005	FORMAT(1H0/' ***ERROR NDF. 2 : AT FIRST THE --OPN - ROUTINE MUST B	2540

IE CALLED*)	2550	93 RETURN	3110
STOP	2560	C	3120
74 NWP=IWNA(NWW)	2570	C	3130
DO 75 LS=1,2	2580	C	3140
NWW=NWW+1	2590	C	3150
IF(NWW-NSZ)76,76,78	2600	ENTRY NDFNXT (KONTR,NNAM,DAT1,ID,KC)	3160
78 READ (LBN*NWR) (IWNA(II),II=1,NSZ)	2610	ID=0	3170
NWR=NWR+1	2620	KC=0	3180
NWW=1	2630	LL=1	3190
76 GO TO (81,82),LS	2640	GO TO 101	3200
81 IDR=IWNA(NWW)	2650	C	3210
GOTO 75	2660	ENTRY LDFNXT(KONTR,NNAM,INAM,DAT2)	3220
82 IDW=IWNA(NWW)	2670	LL=2	3230
75 CONTINUE	2680	GCTO 101	3240
NWW=NWW+1	2690	C	3250
48 IF(NWN)384,384,383	2700	ENTRY IDFNXT(KONTR,NNAM,INAM,DAT2)	3260
384 IF(IS-IDR-1)83,385,83	2710	LL=3	3270
383 IF(NWR-IDR-1)83,84,83	2720	101 NPA=NPA+1	3280
385 DO 386 L=1,NSZ	2730	IVY=0	3290
386 JDAT(L)=ISATZ(L)	2740	IF(NPA-NWP)102,102,103	3300
IDR=IDR+1	2750	103 KONTR=0	3310
GO TO 388	2760	IF (NWN)387,387,389	3320
84 DO 85 L=1,NSZ	2770	389 NKO=NKO+1	3330
85 JDAT(L)=IWNA(L)	2780	IF(NKO-NNK)391,391,387	3340
IDR=IDR+1	2790	387 RETURN	3350
GO TO 388	2800	102 KCNTR=1	3360
83 READ (LBN*IDR) (JDAT(II),II=1,NSZ)	2810	391 DO 304 LS=1,3	3370
IDR=IDR+1	2820	304 NNAM(LS)=NN(LS)	3380
388 NPA=1	2830	IF(LL-2)105,124,106	3390
86 JD=NNAM(2)+NNAM(3)	2840	105 NNAM(4)=0	3400
IF(KK-2)87,88,88	2850	KDAT1(1)=MNAM(1)	3410
87 LS=NNAM(1)*2+1	2860	KDAT1(2)=MNAM(2)	3420
GO TO 89	2870	KDAT1(3)=IBEST(1)	3430
88 LS=1	2880	KDAT1(4)=IBEST(2)	3440
89 DO 90 L=1,JD	2890	KDAT1(5)=MNAM(3)	3450
GO TO (91,92,92),KK	2900	KDAT1(6)=MNAM(4)	3460
91 DAT1(LS)=XJDAT(IDW)	2910	DO 401 II=1,6	3470
DAT1(LS+1)=0.	2920	401 DAT1(II)=Z(II)	3480
LS=LS+2	2930	GO TO 125	3490
GO TO 93	2940	124 DO 126 LS=1,4	3500
92 DAT2(LS)=XJDAT(IDW)	2950	126 INAM(LS)=MNAM(LS)	3510
LS=LS+1	2960	GO TO 125	3520
93 IDW=IDW+1	2970	106 DC 327 LS=1,2	3530
IF(IDW-NSZ) 90,90,94	2980	327 INAM(LS)=NUNA(LS)	3540
94 READ(LBN*IDR) (JDAT(II),II=1,NSZ)	2990	125 IF(NWN)86,86,390	3550
IDR=IDR+1	3000	390 IF(KONTR)108,108,86	3560
IDW=1	3010	108 KP=NWN+3	3570
90 CONTINUE	3020	IVY=1	3580
IF(IVY)1003,1003,58	3030	L=5	3590
1003 KONTR=1	3040	DO 109 LS=1,KP	3600
GO TO 98	3050	IF(NWW-NSZ)110,110,111	3610
24 WRITE(NDUTP,97)NUNA(1),NUNA(2)	3060	111 READ (LBN*NWR) (IWNA(II),II=1,NSZ)	3620
97 FORMAT(1H0/' ***WARNING NDF. 3 : THE DATA FOR',2I9,' ARE NOT INCLU	3070	NWR=NWR+1	3630
IDED IN THE KEDAK LIBRARY')	3080	NWW=1	3640
IF(JJ-1)98,98,96	3090	110 IF(LS-NWN)112,112,113	3650
96 RETURN 1	3100	112 IF(LL-2)114,115,116	3660

114	L=L+2	3670	—WRITE (NOUTP,9001)	120
	KDAT1(L)=IWNA(NWW)	3680	9001 FORMAT(' PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTABSCHIRMFAKTORE	130
	KDAT1(L+1)=0	3690	IN VON AUFGELÖSTEN RESONANZPARAMETERN'/)	140
	DAT1(L)=Z(L)	3700	CALL FSP1E	150
	DAT1(L+1)=Z(L+1)	3710	CALL DOPW (8HBEST ,NN(1))	160
	GO TO 109	3720	CALL DOPW (8HISOT1 ,NN(2))	170
115	INAM(L)=IWNA(NWW)	3730	CALL DOPW (8HISOT2 ,NN(3))	180
	L=L+2	3740	CALL DOPW (8HRES ,NN(4))	190
	GO TO 109	3750	CALL DOPW (8HMIGR ,MM)	200
116	L=LS+2	3760	IREP=0	210
	INAM(L)=IWNA(NWW)	3770	IREF=0	220
113	L=LS-NWN	3780	ISTEP=0	230
	GO TO (119,118,117),L	3790	ISTEP=0	240
119	NWP=IWNA(NWW)	3800	CALL EXPPX(X,X,X,X)	250
	GOTO 109	3810	NEFE=NEF	260
118	IDR=IWNA(NWW)	3820	IF(NEF.NE.1) GOTO 411	270
	GOTO 109	3830	NEFE=0	280
117	IDW=IWNA(NWW)	3840	411 PI=3.14159	290
109	NWW=NWW+1	3850	NE1=NE-1	300
	GO TO 388	3860	NMIN=NMR	310
	END	3870	12 IZAHL(1)=3	320
			IFEST(1)=MATN	330
			IFEST(2)=NN(1)	340
			IFEST(3)=NN(2)	350
			CALL NDFLOC(IDD, IZAHL, FEST, IDAT, ICOD)	360
			IF(IDD)110,111,110	370
			110 A=FEST(4)	380
			IFEST(3)=NN(3)	390
			CALL NDFLOC(IDD, IZAHL, FEST, IDAT, ICOD)	400
			IF(IDD)115,111,115	410
			115 RLA=FEST(4)	420
			R=FEST(5)	430
			SIGP=4.*PI*R**2	440
			IFEST(3)=NN(4)	450
			NR=1	460
			CALL NDFLOC(IDD, IZAHL, FEST, IDAT, ICOD)	470
			IF(IDD)112,111,112	480
			112 ER(NR)=FEST(4)	490
			L(NR)=IFEST(5)	500
			GJ(NR)=FEST(7)	510
			GAT(NR)=FEST(8)	520
			GAN(NR)=FEST(9)	530
			GAG(NR)=FEST(10)+FEST(12)+FEST(13)	540
			GAF(NR)=FEST(11)	550
			NR=NR+1	560
			IF(NR.LE.IRE) GO TO 200	570
			NR=1	580
			IREF=IREF+1	590
			200 CALL NDFNXT(IDD, IZAHL, FEST, IDAT, ICOD)	600
			IF(IDD)112,113,112	610
			113 NR=NR-1	620
			IF(IREF.EQ.0) GO TO 5	630
			IREP=(IREF-1)*IRE+NR+1	640
			RETURN	650
			5 DO 63 IE=1,NE1	660
			IF(ENG(IE+1)-ER(NR))63,63,64	670
	FUNCTION PHI(E)	10		
	PHI=1./E	20		
	RETURN	30		
	END	40		
	SUBROUTINE DOPW(I,J)	10		
	REAL*8 I,J	20		
	J=I	30		
	RETURN	40		
	END	50		
	SUBROUTINE FGEMINS,SIGO,NE,ENG,NEF,ES,F,NT,TEMP,PR,NMR,NFST,SUM,	10		
	1SUO,IRE,IREP,ER,L,GJ,GAT,GAN,GAG,GAF,ISTE,ISTEP,STE)	20		
	REAL*8 MATN,FEST,IFEST,NN,MM,III	30		
	DIMENSION ENG(NE),ES(NEF),F(NEF),TEMP(NT),SIGO(NS),IFEST(20),	40		
	1FEST(20),	50		
	2 NN(4), IZAHL(4), SUMO(3), SUM(NS,7), STE(5, ISTE), SUO(NE,3),	60		
	3ER(IRE),L(IRE),GJ(IRE),GAT(IRE),GAN(IRE),GAG(IRE),GAF(IRE)	70		
	COMMON MATN,ISTRUK,ISPA,NOUTP,LZWF,IR(2),KL	80		
	EQUIVALENCE (FEST(1),IFEST(1))	90		
	WRITE (NOUTP,9000)	100		
9000	FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 1')	110		

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63 CONTINUE
   NLAST=NE1
   GOTO 66
64 NLAST=IE-1
   IF(NLAST)65,65,66
65 NFST=1
   GCTO 3
66 NFST =NLAST+1
   4 NANF=NE-IR(1)
   IF(NANF-NLAST)510,510,3
510 NEND=NE-IR(2)
502 IF(NEND-NLAST)503,503,505
503 NLAST=NEND
   NFST=0
505 DO 60 IT=1,NT
   T=TEMP(IT)
   NR1=1
   DO 10 IE=NANF,NLAST
   DC 36 ISO=1,3
36 SUMO(ISO)=.0
   DO 37 ISO=1,NS
   DO 37 J1=1,7
37 SUM(ISO,J1)=.0
   SUME=.0
   K=1
   STE(1,K)=ENG(IE)
   DO 15 INR=NR1,NR
   IF(ENG(IE)-ER(INR))16,15,15
16 IF(ENG(IE+1)-ER(INR))18,18,17
17 K=K+1
   IF(K.LE.ISTE) GO TO 201
   K=1
   ISTE=ISTEF+1
201 STE(1,K)=ER(INR)
15 CONTINUE
18 K=K+1
   IF(K.LE.ISTE) GO TO 202
   K=1
   ISTE=ISTEF+1
202 STE(1,K)=ENG(IE+1)
   KMAX=K
   IF(ISTEF.EQ.0) GO TO 250
   ISTEP=(ISTEF-1)*ISTE+K+50
   RETURN
250 NR1=INR
   DO 19 K=1,KMAX
   CALL WIRQU(NR1,STE(1,K),STE(2,K),STE(3,K),STE(5,K),NR,IRE,ER,GJ,
   IGAN,GAT,GAF,L,GAG,T,A,R,RLA,NMIN)
19 STE(4,K)=STE(5,K)-STE(3,K)-STE(2,K)
22 KMX=KMAX+1
   KMX1=KMAX-1
   IF(KMX1)33,34,33
33 IF(KMX.LE.ISTE) GO TO 833
   ISTEP=50
   RETURN
833 DO 20 K=1,KMX1

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680 K1=KMX-K
690 DO 20 J1=1,5
700 STE(J1,K1+1)=STE(J1,K1)
710 KMAX=KMAX+1
720 35 STE(1,2)=(STE(1,3)+STE(1,1))/2.
730 CALL WIRQU(NR1,STE(1,2),STE(2,2),STE(3,2),STE(5,2),NR,IRE,ER,GJ,
740 IGAN,
750 IGAN,GAF,L,GAG,T,A,R,RLA,NMIN)
760 STE(4,2)=STE(5,2)-STE(3,2)-STE(2,2)
770 DST=STE(1,2)-STE(1,1)
780 CALL STOSS(STE(1,1),STE(1,2),STE(1,3),F1,F2,F3,NEF,ES,F)
790 DO 21 J1=2,4
800 FLC1=(3.*STE(J1,1)*F1+STE(J1,3)*F3)*DST*.25
810 FLC2=(STE(J1,1)*F1+STE(J1,2)*F2)*DST*.5
820 IF(FLC1)45,21,45
830 45 AFL=ABS((FLC1-FLC2)/FLC1)
840 IF(AFL-PR )21,21,22
850 21 CONTINUE
860 DO 24 J1=1,3
870 SUMO(J1)=(STE(J1+1,1)*F1+STE(J1+1,2)*F2)*DST*.5+SUMO(J1)
880 SUME=SUME+(F1+F2)*DST*.5
890 DO 25 ISO=1,NS
900 SIG1=STE(5,1)+SIGO(ISO)
910 SIG2=STE(5,2)+SIGO(ISO)
920 SIG3=SIG1*SIG1
930 SIG4=SIG2*SIG2
940 SUM(ISO,5)=(F1/SIG3+F2/SIG4)*DST*.5+SUM(ISO,5)
950 SUM(ISO,1)=(F1/SIG1+F2/SIG2)*DST*.5+SUM(ISO,1)
960 DO 325 J1=6,7
970 325 SUM(ISO,J1)=(STE(J1-2,1)*F1/SIG3+STE(J1-2,2)*F2/SIG4)*DST*.5
980 1+SUM(ISO,J1)
990 DO 25 J1=2,4
1000 25 SUM(ISO,J1)=(STE(J1,1)*F1/SIG1+STE(J1,2)*F2/SIG2)*DST*.5+SUM(ISO,J
1010 11)
1020 DO 27 J1=2,4
1030 FLC1=(STE(J1,1)*F1+3.*STE(J1,3)*F3)*DST*.25
1040 FLC2=(STE(J1,2)*F2+STE(J1,3)*F3)*DST*.5
1050 IF(FLC1)46,27,46
1060 46 AFL=ABS((FLC1-FLC2)/FLC1)
1070 IF(AFL-PR )27,27,28
1080 27 CONTINUE
1090 DO 29 J1=1,3
1100 29 SUMO(J1)=(STE(J1+1,2)*F2+STE(J1+1,3)*F3)*DST*.5+SUMO(J1)
1110 SUME=SUME+(F2+F3)*DST*.5
1120 DO 30 ISO=1,NS
1130 SIG1=STE(5,2)+SIGO(ISO)
1140 SIG2=STE(5,3)+SIGO(ISO)
1150 SIG3=SIG1*SIG1
1160 SIG4=SIG2*SIG2
1170 SUM(ISO,5)=SUM(ISO,5)+(F2/SIG3+F3/SIG4)*DST*.5
1180 SUM(ISO,1)=SUM(ISO,1)+(F2/SIG1+F3/SIG2)*.5*DST
1190 DO 330 J1=6,7
1200 330 SUM(ISO,J1)=SUM(ISO,J1)+(STE(J1-2,2)*F2/SIG3+STE(J1-2,3)*F3/
1210 1SIG4)*DST*.5
1220 DO 30 J1=2,4
1230 30 SUM(ISO,J1)=SUM(ISO,J1)+(STE(J1,2)*F2/SIG1+STE(J1,3)*F3/SIG2)*DST*
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1.5
DO 31 K=3,KMAX
DO 31 J1=1,5
31 STE(J1,K-2)=STE(J1,K)
KMAX=KMAX-2
GOTO 22
28 DO 32 J1=1,5
32 STE(J1,1)=STE(J1,2)
GOTO 35
34 IF(IT.EQ.1.AND.IE.EQ.NANF) GO TO 935
GO TO 934
935 INDEX=0
WRITE (LZWF) INDEX,MM
934 INDEX=6
WRITE(LZWF) INDEX,MATN ,TEMP(IT),IE,ENG(IE),ENG(IE+1)
DO 38 ISO=1,3
38 SUMO(ISO)=SUMO(ISC)/SUME
INDEX=5
SGT1=SUMO(1)+SUMO(2)+SUMO(3)
WRITE(LZWF) INDEX,SUMO(2),SUMO(3),SUMO(1),SUMO(3),SGT1
INDEX=6
WRITE(NOUTP,957)
957 FORMAT(1H0,18X,8HMATERIAL,3X,10HTEMPERATUR,2X,6HGRUPPE,7X,7HGRENZE
IN)
IV=NE-IE
WRITE(NOUTP,958) MATN ,TEMP(IT),IV,ENG(IE),ENG(IE+1)
958 FORMAT(1H ,18X,A5,F14.2,6X,I2,1X,2E12.4/1X)
WRITE(NOUTP,950)
950 FCRMAT(1H0,21X,7HSIGMA G,9X,7HSIGMA N,9X,7HSIGMA F,9X,7HSIGMAN1,
19X,7HSIGMAT1)
WRITE (NOUTP,951) SUMO(2),SUMO(3),SUMO(1),SUMO(3),SGT1
951 FORMAT(1H ,15X,5E16.8)
WRITE(NOUTP,952)
952 FORMAT(1H0,6X,7HSIGMA C,11X,2HFG,14X,2HFN,14X,2HFF,14X,3HFN1,
113X,3HFT1/1X)
DO 39 ISO=1,NS
DO 40 J1=2,4
IF(SUMO(J1-1))420,421,420
421 SUM(ISO,J1)=1.
GOTO 40
420 SUM(ISO,J1)=SUM(ISO,J1)/(SUM(ISO,1)*SUMO(J1-1))
40 CONTINUE
SUM(ISO,6)=SUM(ISO,6)/(SUM(ISO,5)*SUMC(3))
SUM(ISO,7)=SUM(ISO,7)/(SUM(ISO,5)*(SUMO(1)+SUMO(2)+SUMO(3)))
WRITE (NOUTP,953) SIGO(ISO),SUM(ISO,3),SUM(ISC,4),SUM(ISO,2)
1,SUM(ISC,6),SUM(ISO,7)
39 WRITE(LZWF) INDEX,SIGO(ISO),SUM(ISC,3),SUM(ISC,4),SUM(ISO,2)
1,SUM(ISO,6),SUM(ISO,7)
953 FORMAT(6E16.8)
10 CONTINUE
60 CGNTINUE
3 CONTINUE
111 KL=KL+1
RETURN
END

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SUBROUTINE WIRQ (INR,E,SF,SG,ST,IRE,
1 ER,GJ,GAN,GAT,GAF,L,GAG,T,A,R,RLA)
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1 REAL*8 MAT
DIMENSION ER(IRE),GJ(IRE),GAN(IRE),GAT(IRE),GAF(IRE),L(IRE),GAG(IR
1E)
COMMON MAT,ISTRUK,ISPA,NGUTP
PI=3.14159
XKO=8.6165E-5
ER=ER(INR)
IF(ERI)20,20,21
20 ERI=1.
GOTO 21
21 ELA=RLA/SQRT(ERI)
SOC=4.*PI*ELA**2*GJ(INR)*GAN(INR)/GAT(INR)
XLEX=FLOAT(L(INR))-0.5
SOC=SOC*(E/ERI)**XLEX
X=2.*(E-ER(INR))/GAT(INR)
IT=T
IF(IT)4,5,4
5 PSI=1./(1.+X*X)
CHI=X*PSI
GOTO 6
4 DELTA=SQRT(4.*XKO*T*ERI/A)
TETA=(DELTA/GAT(INR))**2
CALL PSIXI(X,TETA,PSI,CHI)
6 IF(L(INR)-1)1,1,2
2 SG=.0
SF=.0
ST=.0
GOTO 3
1 SG=SOC*GAG(INR)*PSI/GAT(INR)
SF=SOC*GAF(INR)*PSI/GAT(INR)
ELA=RLA/SQRT(E)
DELTA1=2.*(R/ELA-ATAN(R*FLOAT(L(INR))/ELA))
ST=SOC*(PSI*COS(DELTA1)+CHI*SIN(DELTA1))
SIGP=4.*PI*R**2
IF(ST+SIGP)10,10,3
10 WRITE(NOUTP,900)ER(INR),E,ST,SIGP
900 FORMAT(1H0,77H***WARNING 1.01 : THE CONTRIBUTION OF THE RESONANCE
1WITH THE RESONANCE ENERGY,E13.6,3H EV/1H ,18X,13HAT THE ENERGY,E14
2.6,13H IS NEGATIVE,,E14.6,22H BARNS,AND ITS ABSOLUT/1H ,18X,46HVAL
3UE LARGER THAN THE POTENTIAL CRCSS SECTION,E14.0,7H BARNS.)
3 RETURN
END
SUBROUTINE WIRQU(NR1,E,SIFG,SIGG,SIGT,NR,IRE,
1 ER,GJ,GAN,GAT,GAF,L,GAG,
1T,A,R,RLA,NMIN)
1 REAL*8 MAT
DIMENSION ER(IRE),GJ(IRE),GAN(IRE),GAT(IRE),GAF(IRE),L(IRE),GAG(IR
1E)
COMMON MAT,ISTRUK,ISPA,NOUTP
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PI=3.14159
SIGP=4.*PI*R**2
NR2=NR1+1
DO 1 I=1,NR1
NR3=NR2-I
IF(E-ER(NR3))1,2,2
1 CONTINUE
2 SIFG=0.
SIGG=.0
SIGT=SIGP
IA=NR3-NMIN
IF(IA)3,3,4
3 IA=1
4 IE=NR3+NMIN
IF(IE-NR)5,6,6
6 IE=NR
DO 7 INR=IA,IE
CALL WIRQ(INR,E,SF,SG,ST,IRE,
1 ER,GJ,GAN,GAT,GAF,L,GAG,T,A,R,RLA)
SIFG=SIFG+SF
SIGG=SIGG+SG
SIGT=SIGT+ST
7 CONTINUE
IF(SIGT)10,10,11
10 WRITE(NOUTP,900) SIGT,E
11 CONTINUE
900 FORMAT(1H0,'***WARNING 1.02 : THE TOTAL CROSS SECTION IS NEGATIVE,
1',E14.6,' BARNs, AT THE ENERGY'/19X,E14.8,' EV. THE TOTAL CROSS SE
2CTION HAS BEEN SET EQUAL TO 1. BARN')
SING=SIGT-SIGG-SIFG
IF(SING-1.0)12,13,13
12 SING=1.0
SIGT=SING+SIGG+SIFG
13 RETURN
END

SUBROUTINE STOSS(E1,E2,E3,F1,F2,F3,NFE,ES,F)
DIMENSION F(NFE),ES(NFE)
NEFE=NFE
IF(NFE.EQ.1) NEFE=0
IF(NEFE)1,2,1
2 F1=PHI(E1)
F2=PHI(E2)
F3=PHI(E3)
GOTO 3
1 E=E1
IND=1
9 DO 4 I=1,NEFE
IF(E-ES(I))5,5,4
4 CONTINUE
I=NEFE
5 FF=F(I)-(ES(I)-E)*(F(I)-F(I-1))/(ES(I)-ES(I-1))

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GOTO(6,7,8),IND
6 F1=FF
IND=2
E=E2
GOTO 9
7 F2=FF
IND=3
E=E3
GOTO 9
8 F3=FF
3 RETURN
END

SUBROUTINE PSIXI(X,T,U,V)
C
C GENAU EINMAL IST IN EINEM PRGGRAMM, DAS PSIXI AUFRUFT,
C VOR DIESEM AUFRUF DAS STATEMENT
C
C CALL EXPPX(X,X,X,X)
C
C W SETZEN
C
C DIMENSION D(8),C(8)
XX=X*X
BB=1.+XX
U=HORP/BB
V=X*U
DO 1 N=1,8
AA=4.*T*C(N)
CC=2.*HORP*D(N)
ALPHA=1.-XX+AA
DIV=1./(ALPHA*ALPHA+4.*XX)
U=U+CC*(BB+AA)*DIV
1 V=V+CC*X*(BB-AA)*DIV
IF(T.LE.8)GOTO 2
RT=SQRT(T)
ETA=PI/(H*RT)
EXPO=(XX-1.)/(4.*T)+ETA
IF(EXPO.GT.25.)GOTO 2
XI=X/(2.*T)
ZETA=X*ETA
ACCS=COS(XI)
BSIN=SIN(XI)
CCOS=EXP(-ETA)-COS(ZETA)
DSIN=SIN(ZETA)
BOT=1./(CCOS*CCOS+DSIN*DSIN)
REAL=BOT*(ACOS*CCOS-BSIN*DSIN)
AMAG=BOT*(BSIN*CCOS+ACOS*DSIN)
EXIT=RT*PI*EXP(-EXPO)/RT
U=U+EXIT*REAL
V=V-EXIT*AMAG
2 RETURN

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ENTRY EXPPX
PI=3.141593
RTPI=SQRT(PI)
H=0.7
HORP=H/RTPI
B=H*H/(4.*PI*PI)
DO 3 N=1,3
ENN=N*N
C(N)=ENN*H*H
3 D(N)=EXP(-C(N))
RETURN
END

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SUBROUTINE FSTAT (MI, SIGO, J, ENG, NFE, REFE, EFE, MT, TEMP, SE, SM,
1XEUGZ)
REAL*8 X, IX, MAT, NN
1, MM
DIMENSION ENG(J), EFE(NFE), REFE(NFE), TEMP(MT), SIGO(MI), NADAT(2),
2EY(100), NN(5), NAME(4), X(100), IX(100), CHI(25,4), GF(10,100),
3 GN(10), XL(10), DN(10), GI(10), XNYN(10), FXNYN(10), GG(10),
4SE(3,5,MI), SM(3,5,MI), E(3), XEUGZ(5,MI), XSUG1(5), SU(5)
5, EF(3)
COMMON MAT, ISTRUK, ISPA, NOUTP, LIZ, NANF, NENC, KL
EQUIVALENCE (X(1), IX(1))
DATA CHI/8.383E-4, 5.8607E-3, 1.595668E-2, 3.122396E-2, 5.182214E-2,
X7.796678E-2, 0.1110434, 0.14875186, 0.19382632, 0.2462014, 0.3066044,
X0.37595912, 0.45541448, 0.5464511, 0.65097026, 0.7714919, 0.9114051,
X1.075423, 1.2703559, 1.5065244, 1.8006327, 2.1822634, 2.7109872,
X3.5798648, 5.9721596, 2.0273E-2, 6.195225E-2, 0.1054415, 0.15091425,
X0.198549, 0.2485715, 0.301227, 0.35680875, 0.41567075, 0.47820725,
X0.544925, 0.61641525, 0.693415, 0.77684375, 0.867878, 0.968046,
X1.0793867, 1.2047147, 1.3480617, 1.5155082, 1.716864, 1.9695355,
X2.309333, 2.832582, 4.2188757, 5.9001333E-2, 0.1333242, 0.19497146,
X0.251645, 0.3074142, 0.36262126, 0.4181491, 0.4746072, 0.53251576,
X0.5965109, 0.65162876, 0.7165988, 0.78492293, 0.86824213, 0.94071066,
X1.0257787, 1.118885, 1.2221354, 1.3385186, 1.4725188, 1.6312598,
X1.8271452, 2.0844432, 2.4887964, 3.497654, 0.10176025, 0.196785,
X0.26520225, 0.32682625, 0.3842565, 0.43967925, 0.49512475, 0.545537,
X0.6059465, 0.659715, 0.717307, 0.77698, 0.83930475, 0.90493675,
X0.97467075, 1.049486, 1.1306795, 1.2199627, 1.319785, 1.4337535,
X1.567579, 1.7310672, 1.9429085, 2.2809085, 3.089838/
WRITE (NOUTP, 9000)
9000 FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 2')
WRITE (NOUTP, 9001)
9001 FORMAT (' PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTABSCHIRMFAKTORE
IN VON STATISTISCHEN RESONANZPARAMETERN'/)
NEFE=NFE
IF(NFE.EQ.1) NEFE=0
ISIG=0
IF(SIGO(MI).GE.1.E6) GO TO 260
MI=MI+1
SIGO(MI)=1.E6

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ISIG=1
260 CALL DOPW (8HFSTAT ,MM)
CALL DOPW (8HBEST ,NN(1))
CALL DOPW (8HISOT1 ,NN(2))
CALL DOPW (8HISOT2 ,NN(3))
CALL DOPW (8HST ,NN(4))
CALL DOPW (8HSTGF ,NN(5))
I=0
WRITE (LIZ) I, MM
NGR=J-NANF
NEN=J+1-NEND
IX(1)=MAT
IX(2)=NN(1)
IX(3)=NN(2)
NAME(1)=3
CALL NDFLOC (KONTR, NAME, X, NUDAT, IS)
IF (KONTR) 51, 51, 62
51 WRITE (NOUTP, 153) KCNTR
153 FORMAT(1H0, 3X, 7HKONTR =I3)
GC TO 1000
62 DEL=344.489/X(4)
SPIN=DABS(X(6))
IX(3)=NN(3)
CALL NDFLOC (KONTR, NAME, X, NUDAT, IS)
IF (KONTR) 51, 51, 63
63 XA=X(4)*X(4)
RQU=X(5)*X(5)
EBI=X(6)
IX(3)=NN(4)
IT=1
CALL NDFLOC (KONTR, NAME, X, NUDAT, IS)
IF (KONTR) 51, 51, 28
28 GG(IT)=X(6)*1.E3
XL(IT)=IX(4)
IF (XL(IT)-5.E-5.LE.1.) GO TO 3000
WRITE (NOUTP, 3001)
3001 FCRMAT(1H0, '***ERROR 2.01 : A CALCULATION OF RESONANCE SELFSHIELDI
ING FACTORS FOR L.GT.1 (L=ANGULAR MOMENTUM) IS NOT POSSIBLE')
GO TO 1000
3000 DN(IT)=X(7)*1.E3
GN(IT)=X(8)*1.E3
GI(IT)=(2.*X(5)+1.)/(4.*SPIN+2.)
FXNYN(IT)=IX(10)
XNYN(IT)=IX(11)
CALL NDFNXT (KONTR, NAME, X, NUDAT, IS)
IF (KONTR) 65, 65, 66
66 IT=IT+1
GO TO 28
65 IF (ISPA) 1000, 201, 404
404 NAME(1)=3
IX(1)=MAT
IX(2)=NN(1)
IX(3)=NN(5)
IY=1
KI=1
CALL NDFLOC (KONTR, NAME, X, NUDAT, IS)

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93 IF(KONTR)51,51,74	960	DC 12 KK=L,NEFE	1520
74 IF(X(4).LT.ENG(NGR+1)) GO TO 75	970	IF(REFE(KK)-E(IST))12,38,213	1530
WRITE(NOUTP,9002)ENG(NGR),ENG(NGR+1)	980	12 CONTINUE	1540
9002 FORMAT(1H0,37H***ERROR 2.02 : THE ENERGY GRUP FROM,E16.8,6H EV TO	990	KK=NEFE	1550
1,E16.8,22H EV IS NOT POSSIBLE IN/19X,67HTHIS MODULE,BECAUSE NO STA	1000	213 EF(IST)=EFE(KK-1)+(EFE(KK)-EFE(KK-1))/(REFE(KK)-REFE(KK-1))*(E(IST	1560
2TISTICAL INFORMATION IS AVAILABLE IN THIS/19X,21HENERGYRANGE ON KE	1010	1)-REFE(KK-1))	1570
3DAK.)	1020	GO TO 26	1580
NGR=NGR+1	1030	38 EF(IST)=EFE(KK)	1590
IF(NGR-NEN)74,1000,1000	1040	26 L=KK	1600
75 IF(X(4).LE.ENG(NGR)) GO TO 76	1050	5 CONTINUE	1610
EY(IY)=ENG(NGR)	1060	27 H=(E(3)-E(2))*3.333333E-1	1620
GF(KI,IY)=X(8)*1.E3	1070	DO 29 KI=1,MI	1630
EY(IY+1)=X(4)	1080	DO 29 IS=1,LIJ	1640
GF(KI,IY+1)=X(8)*1.E3	1090	IF(SE(1,IS,KI).NE.-99999..AND.SE(2,IS,KI).NE.-99999..AND.SE(3,IS,	1650
78 KI=KI+1	1100	KI).NE.-99999.) GO TO 2000	1660
IF(KI.GT.IT) GO TO 92	1110	XEUGZ(IS,KI)=-99999.	1670
CALL NDFNXT(KONTR,NAME,X,NUDAT,IS)	1120	GO TO 29	1680
GC TO 93	1130	2 000 XEUGZ(IS,KI)=H*(SE(1,IS,KI)*EF(1)+4.*SE(2,IS,KI)*EF(2)+SE(3,IS,KI)	1690
92 KI=1	1140	1*EF(3))	1700
IY=IY+1	1150	29 CONTINUE	1710
NEY=IY-1	1160	DO 30 IS=1,LIJ	1720
79 CALL NDFNXT(KONTR,NAME,X,NUDAT,IS)	1170	IF(SM(1,IS,1).NE.-99999..AND.SM(2,IS,1).NE.-99999..AND.SM(3,IS,1)	1730
IF(KONTR)201,201,91	1180	1.NE.-99999.) GO TO 2001	1740
76 EY(IY)=X(4)	1190	SL(IS)=-99999.	1750
GF(KI,IY)=X(8)*1.E3	1200	GO TO 2002	1760
GO TO 78	1210	2 001 SU(IS)=H*(SM(1,IS,1)*EF(1)+4.*SM(2,IS,1)*EF(2)+SM(3,IS,1)*EF(3))	1770
91 EY(IY)=X(4)	1220	2 002 IF(SE(1,IS,MI).NE.-99999..AND.SE(2,IS,MI).NE.-99999..AND.SE(3,IS,	1780
GF(KI,IY)=X(8)*1.E3	1230	MI).NE.-99999.) GO TO 2004	1790
KI=KI+1	1240	XSUG1(IS)=-99999.	1800
IF(KI.GT.IT) GO TO 92	1250	GO TO 30	1810
GO TO 79	1260	2 004 XSUG1(IS)=H*(SE(1,IS,MI)*EF(1)+4.*SE(2,IS,MI)*EF(2)+SE(3,IS,MI)*EF	1820
201 DO 7 I=1,MT	1270	1(3))	1830
TT=TEMP(I)	1280	30 CONTINUE	1840
IF(TT.EQ.0) GO TO 700	1290	XNE=H*(EF(1)+4.*EF(2)+EF(3))	1850
K=NGR-1	1300	DO 32 KI=1,MI	1860
N=1	1310	DO 32 IS=1,LIJ	1870
IF(K+1-NEN)4,1000,1000	1320	IF(XEUGZ(IS,KI).NE.-99999..AND.XSUG1(IS).NE.-99999.) GO TO 2003	1880
4 K=K+1	1330	XEUGZ(IS,KI)=-99999.	1890
E(N)=ENG(K)	1340	GC TO 32	1900
16 CALL QUER(E(N),TT,N,SE,SM,XNYN,FXNYN,DEL,RQU,CHI,	1350	2 003 XEUGZ(IS,KI)=XEUGZ(IS,KI)/XSUG1(IS)	1910
1DN,EBI,XA,XL,GN,GG,NEY,EY,GF,GI,IT,MI,SIGO)	1360	32 CONTINUE	1920
21 GO TO (23,24,25),N	1370	DO 31 IS=1,LIJ	1930
23 IF(K+1-NEN)6,6,7	1380	IF(SU(IS).EQ.-99999.) GO TO 31	1940
6 IF(EY(NEY).LT.ENG(K+1)) GO TO 7	1390	SU(IS)=SU(IS)/XNE	1950
N=2	1400	31 CONTINUE	1960
E(N)=(ENG(K)+ENG(K+1))*0.5	1410	KI=K-1	1970
GO TO 16	1420	IS=J-KI	1980
24 N=3	1430	WRITE(NOUTP,130) MAT,TT,IS,ENG(KI),ENG(K)	1990
GO TO 4	1440	130 FORMAT(1H0,16X,8HMATERIAL,2X,10HTEMPERATUR,2X,6HGRUPPE,8X,7HGRENZE	2000
25 LIJ=5	1450	1N/1H,17X,A5,4X,F8.2,5X,I3,2E12.4)	2010
128 IF(NEFE)300,301,300	1460	IN=6	2020
301 DO 1 IST=1,3	1470	WRITE(LIJ)IN,MAT,TT,KI,ENG(KI),ENG(K)	2030
1 EF(IST)=PHI(E(IST))	1480	132 WRITE(NOUTP,37)(SU(IS),IS=1,LIJ)	2040
GO TO 27	1490	37 FORMAT(1H0,20X,7HSIGMA G,9X,7HSIGMA N,9X,7HSIGMA F,9X,7HSIGMAN1,	2050
300 L=2	1500	19X,7HSIGMAT1/1H,14X,5E16.8)	2060
DO 5 IST=1,3	1510	IN=LIJ	2070

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WRITE(LIZ)IN,(SU(IS),IS=1,LIJ)
WRITE(NOUTP,138)
138 FORMAT(1H0,5X,7HSIGMA 0,11X,2HFG,14X,2HFN,14X,2HFF,14X,3HFN1,13X,
13HFT1/1X)
DO 139 KK=1,MI
WRITE(NOUTP,36)SIGO(KK),(XEUGZ(IS,KK),IS=1,LIJ)
36 FORMAT(1X,E15.8,5E16.8)
IN=LIJ+1
139 WRITE(LIZ)IN,SIGO(KK),(XEUGZ(IS,KK),IS=1,LIJ)
40 DO 70 KK=1,MI
DO 70 IN=1,LIJ
SE(1,IN,KK)=SE(3,IN,KK)
70 SM(1,IN,KK)=SM(3,IN,KK)
E(1)=E(3)
GO TO 23
700 WRITE(NOUTP,701)
701 FORMAT(1H0,93H***ERROR 2.03 : A CALCULATION OF RESONANCE SELF SHIE
LDING FACTORS AT ZERO TEMPERATURE (DEGREE/17X,39HKELVIN) IS NOT PO
SSIBLE IN THIS MODULE,)
7 CONTINUE
1000 IF(ISIG.EQ.1) MI=MI-1
BACKSPACE LIZ
READ(LIZ) I,NN(1)
IF(I.EQ.0.AND.NN(1).EQ.MM) BACKSPACE LIZ
KL=KL+1
RETURN
END

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ST=0.
DO 1 I=1,IT
AC(I)=SIGC(E,I,XA,GI,DN,EBI,GN,RQU,XL)
PH(I)=PHASE(E,I,XA,RQU,XL)
ACI(I)=AC(I)*PH(I)
1 ST=ST+ACI(I)
ST=ST+SIGO(K)+SE(N,2,K)
DO 2 I=1,IT
ISE=0
IJ=1
KJ=3
CALL EZZ(E,I,EZG,EZC,EZF,DRSG,DRFS, CHI,XNYN,FXNY,GN,RQU,XL,XA,
1GG,NEY,EY,GFA)
D=DMIT(E,I,DN,EBI)
DL=DELTA(E,TT,I,DEL)
A(1)=AC(I)*DRSG
G=GAMN(E,I,GN,RQU,XL,XA)+GAMG(E,I,GG)+GAF*(E,I,NEY,EY,GFA)
IF(G/DL-C)22,22,922
922 TETA=G/DL
WRITE(NOUTP,923)TETA,I,TT,E
923 FORMAT(1H0,31H***WARNING 2.01 : GAMMA/DELTA =,E14.6,11H FOR SERIES
1,I,13,19H AT THE TEMPERATURE,E14.6/19X,31HDEGREE KELVIN,AND AT THE E
2NERGY,E14.6,4H EV.)
22 XY=D*ACI(I)*EPSI(E,TT,I,DN,EBI,DEL)/(2.506628*DL*ST**2)
SPI=ST
IF((ST-ACI(I)).GT.0) GO TO 23
DO 26 L=IJ,5
SE(N,L,K)=-99999.
26 SM(N,L,K)=-99999.
WRITE(NOUTP,42)SIGO(K),I
42 FORMAT(1H0,28H***WARNING 2.02 : FOR SIGO =,E12.5,33H SPEFF* IS NEG
1ATIVE IN THE SERIES,I4)
GO TO 25
23 A(3)=0.
B(3)=0.
AA=0.398942*ACI(I)*D*EZC/(DL*SPI)
IF(AA-TEST)3,3,4
3 BT=(1.-AA)*D*ACI(I)/SPI
B(1)=(1.-AA*EZG/EZC)*D*A(1)/SPI
IF(ISPA)5,5,6
6 A(3)=AC(I)*DRFS
B(3)=(1.-AA*EZF/EZC)*D*A(3)/SPI
5 A(2)=ACI(I)-A(3)-A(1)
B(2)=BT-B(1)-B(3)
QF=1.+ACI(I)/(SPI+ACI(I))
BT=QF*BT
B(1)=QF*B(1)
B(2)=QF*B(2)
B(3)=QF*B(3)
IF(ISE.EQ.0) GO TO 27
A(4)=A(2)
A(5)=ACI(I)
B(4)=B(2)
B(5)=BT
27 XNEN=1./(1.-BT/D+XY*ACI(I))
IF(XNEN.GT.0) GO TO 202

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SUBROUTINE QUER (E,TT,N,SE,SM,XNYN,FXNY,DEL,RQU,CHI, DN,EBI,XA, 10
1XL,GN,GG,NEY,EY,GFA,GI,IT,MI,SIGO) 20
REAL*8 MAT 30
DIMENSION SIGO(MI),A(5),B(5),ACI(10), 40
1PH(10),AC(10),GFF(25),TERM(5),CHI(25,4),XNYN(10),FXNY(10),GNN(25), 50
2SE(3,5,MI),SM(3,5,MI),DN(10),XL(10),GN(10),GG(10),EY(100), 60
3GFA(10,100),GI(10) 70
COMMON MAT,ISTRUK,ISPA,NOUTP,LIZ,NANF,NENC,KL 80
C=0.5 90
TEST=0.1 100
DO 25 K=1,MI 110
SE(N,1,K)=0. 120
SE(N,2,K)=12.566371*RQU 130
IF(ISPA .EQ.1) GO TO 41 140
SE(N,3,K)=1. 150
GO TO 40 160
41 SE(N,3,K)=0. 170
40 SE(N,4,K)=SE(N,2,K) 180
SE(N,5,K)=SE(N,2,K) 190
SM(N,1,K)=0. 200
SM(N,2,K)=SE(N,2,K) 210
SM(N,3,K)=0. 220
SM(N,4,K)=SE(N,2,K) 230
SM(N,5,K)=SE(N,2,K) 240

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WRITE (NOUTP,203)
203 FORMAT(1H0,52H***WARNING 2.03 : INACCURATE CVERLAPPING CORRECTION.
1)
202 CCONTINUE
DO 57 L=IJ,KJ
TERM(L)=(ST-ACI(I))*XNEN*(B(L)/D-XY*A(L))
SE(N,L,K)=SE(N,L,K)+TERM(L)
7 SM(N,L,K)=SM(N,L,K)+A(L)
57 CONTINUE
GO TO 20
4 GO=GAMN(E,I,GN,RQU,XL,XA)
CI=SPI*GO*1.5707963/(ACI(I)*D)
IF(ISPA)8,8,9
8 GF=0.
NF=1
NN=25
GFF(1)=0.
NU=XNYN(I)
DC 10 M=1,NN
10 GNN(M)=GO*CHI(M,NU)
11 GAX=GAMG(E,I,GG)
BT=0.
B(1)=0.
B(2)=0.
LG=1
DO 12 M=1,NN
GAX=GNN(M)
DO 12 NI=1,NF
GFX=GFF(NI)
GTX=GAX+GNX+GFX
XK=1.442695*ALOG(1.E5*CI*GTX/GNX)
XT=GTX/DL
CALL TAB (XT,XK,LG,DSJ,XABCJ,DJK)
B(3)=B(3)+GFX*XABCJ
B(1)=B(1)+XABCJ
12 BT=BT+GTX*XABCJ
AZ=NN
BZ=NF
FK=AZ*BZ
B(1)=B(1)*GAX/(PH(I)*FK)
B(3)=B(3)/(PH(I)*FK)
BT=BT/FK
GO TO 5
9 NU=XNYN(I)
A(3)=AC(I)*DRFS
NN=25
NF=25
DC 200 M=1,NN
200 GNN(M)=GO*CHI(M,NU)
GF=GAFM(E,I,NEY,EY,GFA)
NU=FXNY(I)
DO 201 M=1,NF
201 GFF(M)=GF*CHI(M,NU)
GO TO 11
20 IF(ISE.EQ.1) GO TO 21
ISE=1

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FT=ST
ST=ST*0.5
IJ=4
KJ=5
GO TO 22
21 ST=FT
2 CONTINUE
25 CONTINUE
RETURN
END

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SUBROUTINE EZZ (E,NS,EZG,EZC,EZF,DRSG,DRFS, CHI,XNYN,FXNYN,GN,
IRQU,XL,XA,GG,NEY,EY,GF)
REAL*8 MAT
DIMENSION CHI(25,4),XNYN(10),FXNYN(10),GN(10),XL(10),GG(10)
1, EY(100),GF(10,100),X(2),Y(2)
COMMON MAT,ISTRUK,ISPA,NOUTP
GNN=GAMN(E,NS,GN,RQU,XL,XA)
GNF=GAFM(E,NS,NEY,EY,GF)
GNG=GAMG(E,NS,GG)
NY=XNYN(NS)
NFY=FXNYN(NS)
DO 1 I=1,2
X(I)=.0
1 Y(I)=.0
IF(ISPA)6,6,5
5 NN=25
NF=25
XFN=625.
7 DO 2 I=1,NN
DO 2 J=1,NF
X1=((GNN*CHI(I,NY))**2)/(GNN*CHI(I,NY)+GNG+GNF*CHI(J,NFY))
Y1=GNN*CHI(I,NY)/(GNN*CHI(I,NY)+GNG+GNF*CHI(J,NFY))
X(1)=X(1)+X1
Y(1)=Y(1)+Y1
X(2)=X(2)+X1*GNF*CHI(J,NFY)
2 Y(2)=Y(2)+Y1*GNF*CHI(J,NFY)
EZG=X(1)/(Y(1)*GNN)
EZC=1.+2./XNYN(NS)
EZF=X(2)/(Y(2)*GNN)
DRSG=Y(1)*GNG/(GNN*XFN)
DRFS=Y(2)/(GNN*XFN)
GO TO 10
6 NN=25
XFN=25.
DO 3 I=1,NN
X1=((GNN*CHI(I,NY))**2)/(GNN*CHI(I,NY)+GNG)
Y1=GNN*CHI(I,NY)/(GNN*CHI(I,NY)+GNG)
X(1)=X(1)+X1
3 Y(1)=Y(1)+Y1
EZG=X(1)/(Y(1)*GNN)
EZC=1.+2./XNYN(NS)

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EZF=.0	420	GAMN=GN(NS)*SQRT(E)*(E#RQU+(1.-XL(NS))*XA)/(E#RQU+XA)	30
DRSG=Y(1)*GNG/(GNN*XFN)	430	RETURN	40
DRFS=.0	440	END	50
10 CONTINUE	450		
RETURN	460		
END	470		
		FUNCTION GAMG(E,NS,GG)	10
		DIMENSION GG(10)	20
		GAMG=GG(NS)	30
		RETURN	40
		END	50
FUNCTION DMIT(E,NS,DN,EBI)	10		
DIMENSION DN(10)	20		
DMIT=DN(NS)*((EBI+E)/EBI)**2*EXP(-SQRT(.89.72E-6)*(SQRT(EBI+E)-	30		
1 SQRT(EBI)))	40		
RETURN	50		
END	60		
		FUNCTION GAFM(E,NS,NEU,EY,GF)	10
		REAL*8 MAT	20
		DIMENSION EY(100),GF(10,100),X1(2),FX1(2)	30
		COMMON MAT,ISTRUK,ISPA	40
		IF(ISPA)2,2,3	50
		3 CALL SUCH (E,NSU,NEU,EY)	60
		DO 1 J=1,2	70
		NSU1=NSU+J	80
		X1(J)=EY(NSU1)	90
		1 FX1(J)=GF(NS,NSU1)	100
		GAFM=POL(E,X1,FX1)	110
		GO TO 4	120
		2 GAFM=0.	130
		4 RETURN	140
		END	150
		FUNCTION POL(X,X1,FX1)	10
		DIMENSION X1(2),FX1(2)	20
		IF(X1(2)-X1(1))1,2,1	30
		2 POL=FX1(1)	40
		GO TO 3	50
		1 POL=FX1(1)+(FX1(2)-FX1(1))*(X-X1(1))/(X1(2)-X1(1))	60
		3 RETURN	70
		END	80
		SUBROUTINE SUCH (E,NSU,NEU,EY)	10
		DIMENSION EY(100)	20
		DO 14 K=2,NEU	30
		IF(EY(K)-E)14,13,13	40
		14 CONTINUE	50
FUNCTION DELTA(E,T,NS,DEL)	10		
DELTA=SQRT(DEL*T*E)	20		
RETURN	30		
END	40		
		FUNCTION EPSI(E,T,NS,DN,EBI,DEL)	10
		DIMENSION EP(21),EPP(21),X1(2),FX1(2),DN(10)	20
		DATA EP/0.,1.,1.5,2.,2.5,3.,3.5,4.,5.,7.,10.,14.,20.,30.,40.,	30
		X50.,60.,70.,80.,90.,100./,EPP/2.5066,1.7241,1.3705,1.0661,	40
		X0.82,0.6291,0.4845,0.3758,0.23261,0.10002,0.03577,0.012276,	50
		X0.003647,8.523E-4,2.8817E-4,1.2683E-4,5.979E-5,3.694E-5,2.5643E-5,	60
		X1.8679E-5,8.733E-6/	70
		X=DMIT(E,NS,DN,EBI)/DELTA(E,T,NS,DEL)	80
		DO 1 I=2,21	90
		IF(X-EP (I))4,4,1	100
		1 CONTINUE	110
		4 DO 5 K=1,2	120
		I1=I+K-2	130
		X1(K)=EP (I1)	140
		5 FX1(K)=EPP(I1)	150
		EPSI=POL(X,X1,FX1)/X	160
		RETURN	170
		END	180
FUNCTION GAMN(E,NS,GN,RQU,XL,XA)	10		
DIMENSION GN(10),XL(10)	20		

13 NSU=K-2	60	X0.12,0.13,0.14,0.15,0.16,0.17,0.18,0.19,0.20,0.26,0.32,0.38,0.44,	270
RETURN	70	X0.50,0.56,0.62,0.68,0.74,0.80,0.86,0.92,0.98/	280
END	80	DATA T01/348.3270,329.0620,309.3460,289.0880,268.2430,246.8240,	290
		X224.91800,202.69300,180.40300,158.38700,137.04200,116.78800,	300
		X 98.01930, 81.05280, 66.09140, 53.20820, 42.35590, 33.39130,	310
		X 26.11050, 20.28100/	320
FUNCTION SIGC(E,NS,XA,GI,DN,EBI,GN,RQU,XL)	10	DATA T02/205.0000,194.1000,183.5000,173.0000,162.5000,151.9000,	330
DIMENSION GI(10),DN(10),GN(10),XL(10)	20	X141.20000,130.20000,119.10000,107.80000, 96.43000, 85.17000,	340
SIGC=19.73921*XA*GI(NS)*GAMN(E,NS,GN,RQU,XL,XA)/(E*DMIT(E,NS,DN,	30	X 74.17000, 63.64000, 53.78000, 44.78000, 36.74000, 29.75000,	350
EBI))	40	X 23.80000, 18.85000/	360
RETURN	50	DATA T03/154.4000,145.3000,136.9000,128.9000,121.2000,113.7000,	370
END	60	X106.30000, 98.80000, 91.28000, 83.67000, 75.99000, 68.28000,	380
		X 60.61000, 53.09000, 45.83000, 38.98000, 32.65000, 26.95000,	390
		X 21.95000, 17.65000/	400
		DATA T04/129.5430,120.8780,113.1000,105.9960, 99.3851, 93.1199,	410
		X 87.07990, 81.16910, 75.31440, 69.46550, 63.59470, 57.69990,	420
		X 51.80430, 45.95800, 40.23460, 34.72510, 29.52760, 24.73430,	430
		X 20.41850, 16.62540/	440
FUNCTION PHASE(E,NS,XA,RQU,XL)	10	DATA T05/115.3000,106.6000, 99.0000, 92.2100, 86.0500, 80.3700,	450
DIMENSION XL(10)	20	X 75.03000, 69.93000, 64.97000, 60.10000, 55.26000, 50.43000,	460
B=SQRT(RQU*E/XA)	30	X 45.60000, 40.81000, 36.08000, 31.47000, 27.06000, 22.93000,	470
PHASE=COS(2.*(B-XL(NS)*ATAN(B)))	40	X 19.14000, 15.75000/	480
RETURN	50	DATA T06/106.4000, 97.5700, 89.9000, 83.1800, 77.2000, 71.8000,	490
END	60	X 66.84000, 62.19000, 57.76000, 53.47000, 49.27000, 45.13000,	500
		X 41.00000, 36.91000, 32.86000, 28.89000, 25.06000, 21.43000,	510
		X 18.05000, 14.98000/	520
		DATA T07/100.5080, 91.4730, 83.6906, 76.9337, 71.0030, 65.7266,	530
SUBROUTINE TAB(XT,XK,LG,DSJ,XABCJ,DJK)	10	X 60.95750, 56.57200, 52.46730, 48.56010, 44.78500, 41.09360,	540
DIMENSION XXK(20),XJT(33),XJ(660)	20	X 37.45480, 33.85470, 30.29690, 26.80240, 23.40750, 20.15950,	550
DIMENSION T01(20),T02(20),T03(20),T04(20),T05(20),T06(20),T07(20),	30	X 17.11060, 14.31030/	560
T08(20),T09(20),T10(20),T11(20),T12(20),T13(20),T14(20),T15(20),	40	DATA T08/ 96.4200, 87.1900, 79.2700, 72.4300, 66.4800, 61.2400,	570
T16(20),T17(20),T18(20),T19(20),T20(20),T21(20),T22(20),T23(20),	50	X 56.58000, 52.35000, 48.45000, 44.79000, 41.30000, 37.93000,	580
T24(20),T25(20),T26(20),T27(20),T28(20),T29(20),T30(20),T31(20),	60	X 34.64000, 31.40000, 28.21000, 25.07000, 22.02000, 19.08000,	590
T32(20),T33(20)	70	X 16.29000, 13.72000/	600
EQUIVALENCE (XXK(1),R),(XJT(1),T),(XJ(1),U),	80	DATA T09/ 93.4900, 84.0900, 76.0200, 69.0800, 63.0800, 57.8400,	610
1 (A,C,Y),(B,D,Y1,VO),(T2,Y5,P,EAXG,XK05,XXXX),	90	X 53.22000, 49.07000, 45.30000, 41.82000, 38.53000, 35.39000,	620
2(T4,Y6,XK06,J),(DK,RI),(XK1,SJ0),(XK2,SJ1),(DKA,R1,V1),(DKB,R2,V2)	100	X 32.36000, 29.39000, 26.48000, 23.62000, 20.83000, 18.14000,	630
3,(DKC,R3,V3),(SV,N),(AB1,TD,M),(AB2,AXG),(SY1,XXXT,S1),	110	X 15.58000, 13.19000/	640
4(XK02,XJA2,S2),(SXM,XK03,XJA3,S3),(XK4,XJB1),(R1K,XJB2),	120	DATA T10/ 91.3415, 81.7729, 73.5745, 66.5357, 60.4667, 55.1985,	650
5(R2K,XJB3),(R3K,XJC1),(BETA,XJC3),(K,SG,SJ3),(Y2,ZZK,XJ1),(Y3,	130	X 50.58150, 46.48530, 42.79780, 39.42410, 36.28530, 33.31770,	660
6WG),(XJ3,Y4),(SY,I),(XJC2,SJ2),(XJA1,SJ4)	140	X 30.47200, 27.71280, 25.01900, 22.38310, 19.81150, 17.32270,	670
EQUIVALENCE (XJ(1),T01(1)),(XJ(21),T02(1)),(XJ(41),T03(1)),	150	X 14.94500, 12.71180/	680
1(XJ(61),T04(1)),(XJ(81),T05(1)),(XJ(101),T06(1)),(XJ(121),T07(1)),	160	DATA T11/ 89.7100, 80.0100, 71.6900, 64.5600, 58.4200, 53.1100,	690
2(XJ(141),T08(1)),(XJ(161),T09(1)),(XJ(181),T10(1)),(XJ(201),T11(1)	170	X 48.48000, 44.40000, 40.76000, 37.47000, 34.43000, 31.59000,	700
3),(XJ(221),T12(1)),(XJ(241),T13(1)),(XJ(261),T14(1)),(XJ(281),T15(180	X 28.89000, 26.30000, 23.78000, 21.32000, 18.92000, 16.60000,	710
41)),(XJ(301),T16(1)),(XJ(321),T17(1)),(XJ(341),T18(1)),(XJ(361),T1	190	X 14.38000, 12.28000/	720
59(1)),(XJ(381),T20(1)),(XJ(401),T21(1)),(XJ(421),T22(1)),	200	DATA T12/ 88.4600, 78.6400, 70.2200, 62.9900, 56.7800, 51.4200,	730
6(XJ(441),T23(1)),(XJ(461),T24(1)),(XJ(481),T25(1)),(XJ(501),T26(1)	210	X 46.77000, 42.70000, 39.09000, 35.84000, 32.88000, 30.14000,	740
7),(XJ(521),T27(1)),(XJ(541),T28(1)),(XJ(561),T29(1)),(XJ(581),T30(220	X 27.55000, 25.09000, 22.70000, 20.39000, 18.14000, 15.97000,	750
81)),(XJ(601),T31(1)),(XJ(621),T32(1)),(XJ(641),T33(1))	230	X 13.88000, 11.90000/	760
DATA XXK/ 5.2, 5.6, 6.0, 6.4, 6.8, 7.2, 7.6, 8.0, 8.4, 8.8, 9.2,	240	DATA T13/ 87.4785, 77.5623, 69.0480, 61.7399, 55.4611, 50.0523,	770
X 9.6,10.0,10.4,10.8,11.2,11.6,12.0,12.4,12.8/	250	X 45.37000, 41.28650, 37.68900, 34.47960, 31.57410, 28.90220,	780
DATA XJT/0.01,0.02,0.03,0.04,0.05,0.06,0.07,0.08,0.09,0.10,0.11,	260	X 26.40610, 24.04060, 21.77230, 19.58010, 17.45450, 15.39730,	790
		X 13.62070, 11.54430/	800
		DATA T14/ 86.7000, 76.7000, 68.1000, 60.7200, 54.3800, 48.9200,	810
		X 44.21000, 40.11000, 36.51000, 33.32000, 30.46000, 27.84000,	820

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X 25.41000, 23.13000, 20.96000, 18.86000, 16.84000, 14.89000,
X 13.01000, 11.22000/
DATA T15/ 86.0600, 75.9900, 67.3300, 59.8900, 53.4900, 47.9800,
X 43.23000, 39.11000, 35.51000, 32.34000, 29.49000, 26.92000,
X 24.55000, 22.33000, 20.24000, 18.23000, 16.29000, 14.43000,
X 12.64000, 10.93000/
DATA T16/ 85.5500, 75.4200, 66.6900, 59.1900, 52.7400, 47.1900,
X 42.41000, 38.26000, 34.66000, 31.48000, 28.66000, 26.12000,
X 23.79000, 21.63000, 19.60000, 17.66000, 15.80000, 14.01000,
X 12.29000, 10.66000/
DATA T17/ 85.1184, 74.9383, 66.1636, 58.6102, 52.1123, 46.5214,
X 41.70360, 37.53840, 33.91760, 30.74480, 27.93540, 25.41610,
X 23.12450, 21.09880, 19.02780, 17.15060, 15.35620, 13.63380,
X 11.98190, 10.40740/
DATA T18/ 84.7600, 74.5400, 65.7200, 58.1200, 51.5800, 45.9500,
X 41.10000, 36.91000, 33.28000, 30.10000, 27.30000, 24.80000,
X 22.53000, 20.45000, 18.52000, 16.69000, 14.95000, 13.29000,
X 11.70000, 10.18000/
DATA T19/ 84.4600, 74.2000, 65.3400, 57.7000, 51.1200, 45.4600,
X 40.58000, 36.37000, 32.72000, 29.53000, 26.73000, 24.25000,
X 22.00000, 19.96000, 18.06000, 16.28000, 14.59000, 12.97000,
X 11.43000, 9.96300/
DATA T20/ 84.2023, 73.9074, 65.0142, 57.3417, 50.7292, 45.0336,
X 40.12650, 35.89180, 32.22580, 29.03470, 26.23560, 23.75540,
X 21.53060, 19.50740, 17.64160, 15.89810, 14.25100, 12.68360,
X 11.18790, 9.76410/
DATA T21/ 83.2433, 72.8209, 63.7912, 55.9766, 49.2205, 43.3851,
X 38.34820, 34.00070, 30.24480, 26.99280, 24.16570, 21.69310,
X 19.51260, 17.56960, 15.81730, 14.21660, 12.73650, 11.35310,
X 10.05070, 8.82059/
DATA T22/ 82.7808, 72.2943, 63.1948, 55.3048, 48.4693, 42.5527,
X 37.43530, 33.01150, 29.18750, 25.87920, 23.01150, 20.51690,
X 18.33520, 16.41300, 14.70340, 13.16620, 11.76730, 10.48140,
X 9.28662, 8.16990/
DATA T23/ 82.5236, 72.0012, 62.8617, 54.5281, 48.0456, 42.0794,
X 36.91100, 32.43670, 28.56450, 25.21300, 22.30950, 19.78900,
X 17.59340, 15.67080, 13.97510, 12.46640, 11.11010, 9.87782,
X 8.74691, 7.70062/
DATA T24/ 82.3662, 71.8216, 62.6574, 54.6965, 47.7842, 41.7860,
X 36.58420, 32.07560, 28.16950, 24.78600, 21.85410, 19.31000,
X 17.09880, 15.16850, 13.47470, 11.97790, 10.64360, 9.44274,
X 8.35127, 7.35063/
DATA T25/ 82.2628, 71.7037, 62.5233, 54.5442, 47.6121, 41.5923,
X 36.36750, 31.83500, 27.90470, 24.49750, 21.54360, 18.98080,
X 16.75420, 14.81430, 13.11730, 11.62440, 10.30150, 9.11911,
X 8.05277, 7.08266/
DATA T26/ 82.1914, 71.6222, 62.4304, 54.4388, 47.4928, 41.4579,
X 36.21670, 31.66690, 27.71890, 24.29400, 21.32300, 18.74490,
X 16.50530, 14.55600, 12.85400, 11.36100, 10.04360, 8.87225,
X 7.82223, 6.87307/
DATA T27/ 82.1399, 71.5634, 62.3636, 54.3629, 47.4068, 41.3608,
X 36.10770, 31.54510, 27.58380, 24.14540, 21.16120, 18.57060,
X 16.32020, 14.36230, 12.65480, 11.16000, 9.84471, 8.67993,
X 7.64071, 6.70620/
DATA T28/ 82.1016, 71.5197, 62.3138, 54.3063, 47.3427, 41.2885,
X 36.02630, 31.45410, 27.48250, 24.03370, 21.03900, 18.43850,

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X 16.17910, 14.21370, 12.50080, 11.00330, 9.68843, 8.52743,
X 7.49539, 6.57129/
DATA T29/ 82.0723, 71.4863, 62.2758, 54.2631, 47.2937, 41.2331,
X 35.96400, 31.38430, 27.40480, 23.94770, 20.94470, 18.33610,
X 16.06920, 14.09740, 12.37950, 10.87900, 9.56356, 8.40460,
X 7.37737, 6.46077/
DATA T30/ 82.0494, 71.4602, 62.2461, 54.2293, 47.2554, 41.1898,
X 35.91530, 31.32960, 27.34380, 23.83010, 20.87040, 18.25520,
X 15.98200, 14.00470, 12.28240, 10.77890, 9.46232, 8.30433,
X 7.28033, 6.36920/
DATA T31/ 82.0311, 71.4394, 62.2224, 54.2024, 47.2249, 41.1553,
X 35.87640, 31.28600, 27.29510, 23.82600, 20.81080, 18.19010,
X 15.91170, 13.92970, 12.20340, 10.69720, 9.37921, 8.22151,
X 7.19965, 6.29253/
DATA T32/ 82.0164, 71.4226, 62.2032, 54.1806, 47.2002, 41.1274,
X 35.84490, 31.25060, 27.25560, 23.78210, 20.76240, 18.13710,
X 15.85430, 13.86820, 12.13850, 10.62960, 9.31019, 8.15236,
X 7.13192, 6.22775/
DATA T33/ 82.0042, 71.4088, 62.1875, 54.1628, 47.1800, 41.1045,
X 35.81900, 31.22160, 27.22310, 23.74590, 20.72250, 18.09340,
X 15.80680, 13.81720, 12.08440, 10.57320, 9.25227, 8.09408,
X 7.07449, 6.17256/
N=20
M=33
DG 1 I=1,N
IF(XK-XXK(I))2,3,1
1 CONTINUE
3 IF(I-2)9,33,23
23 IF(I-N+1)33,33,9
33 XXXK=XK
GC TO 12
2 IF(I-2)9, 5,24
24 IF(I-N+1)32,32,9
32 A=ABS(XK-XXK(I))
B=XK-XXK(I-1)
IF(B-A)4,4,5
4 XXXK=XXK(I-1)
I=I-1
GC TO 12
5 XXXK=XXK(I)
12 IF(XT-XXT(M-1))499,499,9
499 DO 6 J=1,M
IF(XT-XXT(J))7,8,6
6 CONTINUE
8 IF(J-2)9,18,28
28 IF(J-M+1)18,18,9
13 XXXT=XT
GC TO 13
7 IF(J-2)9,11,29
29 IF(J-M+1)17,17,9
17 C=ABS(XT-XXT(J))
D=XT-XXT(J-1)
IF(D-C)10,10,11
10 XXXT=XXT(J-1)
J=J-1
GC TO 13

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VI 28

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11	XXXT=XXT(J)	1950	4449	BETA=2.**XK*1.E-5	2510
13	K=(J-1)*N+I	1960		Y=XA/BETA	2520
	XJB2=XJ(K)	1970		IF(LG)995,995,50	2530
	XJA2=XJ(K-1)	1980	50	IF(LG-1)995,700,995	2540
	XJC2=XJ(K+1)	1990	995	SV=2.*TD	2550
	K=K-N	2000		T2=1./SV	2560
	XJB1=XJ(K)	2010		T4=T2*T2	2570
	XJA1=XJ(K-1)	2020		SXA=(-T4-T2)*XA+T4	2580
	XJC1=XJ(K+1)	2030		SXM=(1-2.)*T4-T2)*XM+2.*T4	2590
	K=K+2*N	2040		DK=SV*XA	2600
	XJB3=XJ(K)	2050		SG=(0.5*TD*DK*(SXA+2.*XA+SV*SXA)-((1.+SV)*XA-1.)*(DK+TD*TD*SXA))/	2610
	XJA3=XJ(K-1)	2060		1(SV*DK*DK*TD)	2620
	XJC3=XJ(K+1)	2070		XK01=SXA*0.5/XK04	2630
	DK=XXXK-XXK(I-1)	2080		XK02=(XA*SXM-SXA*XM)/(XA*XA)	2640
	AB1=XXT(J+1)-XXXT	2090		XK03=(SXA*WG-XA*SG*0.5/WG)/(WG*WG)	2650
	AB2=XXXT-XXT(J-1)	2100		R1=(-17.)*XK01-16.*XK02+XK03	2660
	AB1=AB1/AB2	2110		R2=5.3333333*(8.*XK01+11.*XK02-XK03)	2670
	IF(LG-2)157,800,157	2120		R3=5.3333333*(-5.)*XK01-8.*XK02+XK03	2680
157	IF(XK-XXXK)159,159,70	2130		XK05=XM/XA	2690
159	XK1=(XXXK-XK)/DK	2140		XK06=XA/WG	2700
	XK2=(XK-XXK(I-1))/DK	2150		R1K=32.-17.*XK04-16.*XK05+XK06	2710
	XJ1=XK1*XJA1+XK2*XJB1	2160		R2K=5.3333333*(-18.+8.*XK04+11.*XK05-XK06)	2720
	XJ2=XK1*XJA2+XK2*XJB2	2170		R3K=5.3333333*(12.-5.*XK04-8.*XK05+XK06)	2730
	XJ3=XK1*XJA3+XK2*XJB3	2180		SY=SXA/BETA	2740
	GO TD 71	2190		IF(Y-0.3)900,25,25	2750
70	XK1=(XXX(I+1)-XK)/DK	2200	25	SV=1./(1.+Y)	2760
	XK2=(XK-XXXK)/DK	2210		XJ0=SQRT(SV)	2770
	XJ1=XK1*XJB1+XK2*XJC1	2220		Y1=1./Y	2780
	XJ2=XK1*XJB2+XK2*XJC2	2230		XJ1=(1.-XJ0)*Y1	2790
	XJ3=XK1*XJB3+XK2*XJC3	2240		XJ2=(0.5-XJ1)*Y1	2800
71	DSJ=(AB1+1.)/(AB1*(XT-XXT(J-1))/(XJ3-XJ2)+(XT-XXT(J+1))/(XJ1-XJ2))	2250		XJ3=(0.375-XJ2)*Y1	2810
	IF(LG)222,111,222	2260		XJ4=(0.3125-XJ3)*Y1	2820
222	XABCJ=XJ2+(XT-XXXT)*D SJ	2270		SJC=SY*XJ0*SV*(-0.5)	2830
	IF(LG)800,800,221	2280		SY1=SY*Y1	2840
221	IF(LG-1)111,111,800	2290		SJ1=(-SJ0)*Y1-XJ1*SY1	2850
800	XK1=XT-XXT(J-1)	2300		SJ2=(-SJ1)*Y1-XJ2*SY1	2860
	XK2=XT-XXT(J+1)	2310		SJ3=(-SJ2)*Y1-XJ3*SY1	2870
	XJ1=AB1/(AB1+1.)	2320		SJ4=(-SJ3)*Y1-XJ4*SY1	2880
	DKA=XJ1*(XK1/(XJA3-XJA2)+XK2/(XJA1-XJA2))	2330		GO TD 1000	2890
	DKB=XJ1*(XK1/(XJB3-XJB2)+XK2/(XJB1-XJB2))	2340	900	XJ5=(((((1.444644*Y-0.1494459)*Y+0.1549810)*Y-0.1611802)*Y	2900
	DKC=XJ1*(XK1/(XJC3-XJC2)+XK2/(XJC1-XJC2))	2350		1+0.1681880)*Y-0.1761970)*Y+0.1854705)*Y-0.1963806)*Y+0.2094726)*Y	2910
	XK1=XT-XXXT	2360		2-0.2255859)*Y+0.2460937	2920
	XJ1=XJA2+XK1/DKA	2370		XJ4=0.2734375-Y*XJ5	2930
	XJ2=XJB2+XK1/DKB	2380		XJ3=0.3125-Y*XJ4	2940
	XJ3=XJC2+XK1/DKC	2390		XJ2=0.375-Y*XJ3	2950
	DJK=(0.5*(XJ3-XJ1)+(XK-XXXK)/DK*(XJ3-2.*XJ2+XJ1))/DK	2400		XJ1=0.5-Y*XJ2	2960
	GO TD 111	2410		XJ0=1.-Y*XJ1	2970
9	AXG=XT*XT	2420		SJ5=(((((1.444644*Y-1.3450131)*Y+1.239848)*Y-1.1282614)*Y	2980
	IF(LG)19,1119,19	2430		1+1.009128)*Y-0.830985)*Y+0.741882)*Y-0.5851418)*Y+0.4189452)*Y-	2990
19	IF(XT-6.)1119,2229,2229	2440		20.2255859)*SY	3000
1119	TD=1./AXG	2450		SJ4=- (SY*XJ5+Y*SJ5)	3010
	EAXG=EXP(0.25*AXG)	2460		SJ3=- (SY*XJ4+Y*SJ4)	3020
	XA=XT*0.8862265*(1.-ERF(0.5*XT))*EAXG	2470		SJ2=- (SY*XJ3+Y*SJ3)	3030
	XM=1.253314*XT*(1.-ERF(7.071067E-1*XT))*EAXG*EAXG	2480		SJ1=- (SY*XJ2+Y*SJ2)	3040
	XK04=SQRT(XA)	2490		SJ0=- (SY*XJ1+Y*SJ1)	3050
	WG=SQRT(((1.+2.*TD)*XA-1.)/(8.*XA*TD*TD))	2500	1000	DSJ=(-3.1415926)/(BETA*AXG *XT)*(XK01*XJC+R1*(-XJ1+2.*XJ2)+XK04	3060

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1*SJO+R1K*(-SJ1+2.*SJ2)+R2*((-3.)*XJ2+4.*XJ3)+R2K*((-3.)*SJ2+4.*
2SJ3)+R3*((-5.)*XJ3+6.*XJ4)+R3K*((-5.)*SJ3+6.*SJ4)
IF(LG)60,111,49
49 IF(LG-1)111,111,60
60 IF(XT-6.)700,220,220
700 ZZK=XM/XA
S1=1.-XK04
S2=4.-2.*XK04-2.*ZZK
S3=1.+0.08333333*(XA/WG-5.*XK04-8.*ZZK)
6669 RI=1.+Y
V0=1./SQRT(RI)
Y2=Y*Y
Y3=Y2*Y
Y4=Y2*Y2
IF(Y-0.2)55,66,66
55 V1=(-Y+8.)*Y/(16.*RI)+(Y4-8.*Y3)*(Y4-8.*Y3)/(4096.*((Y+2.)*Y+1.
1*Y2)
GO TO 77
66 V1=(8.-((-Y+4.)*Y+8.)*V0)/Y2
77 IF(Y-2.)88,99,99
88 SV=RI*R1
Y5=2.+Y
Y6=Y5*Y5*Y5
V2=Y/(2.*Y5*RI)-Y3*Y2/(128.*Y6*SV)
P=(3.*Y+16.)*Y+16.
V3=3.*Y2/(RI*P)-3.*Y4*Y4/(16.*SV*P*P)
GO TO 44
99 V2=(-8.-4.*Y+((Y+8.)*Y+8.)*V0)*4./Y3
V3=(((-6.)*Y-32.)*Y-32.+(((Y+18.)*Y+48.)*Y+32.)*V0)*12./Y4
44 XABCJ=(V0-S1*V1+S2*V2+S3*V3)*1.5707963/BETA
IF(LG)220,111,1111
1111 IF(LG-2)111,220,111
220 DJK=-0.69315*XABCJ-1.088797*Y/SXA*(XKC4*SJO+R1K*(-SJ1+2.*SJ2)
1+R2K*(-3.*SJ2+4.*SJ3)+R3K*(-5.*SJ3+6.*SJ4))
GO TO 111
2229 BETA=2.**XK*1.E-5
XABCJ=1.5707963/(BETA*SQRT(1.+1./BETA-6./AXG))
IF(LG)1119,111,34
34 IF(LG-1)111,111,1119
111 RETURN
END
SUBROUTINE FSTRUK(MI,SIGO,NE,ENG,NFE,REFE,EFE,XINTE,ZA,XN,SE,XII, 10
1 ZB,XI,MDIM,MDIMP,EMU,XMU,NER,NERP,SN,EN) 20
REAL*8 MAT,NTYP,FEST,NFEST,NN 30
DIMENSION ENG(NE),REFE(NFE),EFE(NFE),SIGO(MI),SN(NER,3),EN(NER,3 40
1),FEST(10),ISTM(3) 50
2,NFEST(10),NADAT(2),NSATZ(4),E(3),FLUSS(2),ST(2),NTYP(6), 60
3XINTE(NE,4),ZA(MI,NE,5),XN(MI,NE,4),SE(MI,NE,6),XII(NE), 70
4 EMU(MDIM),XMU(MDIM), ZB(MI,NE),XI(NE) 80
COMMON MAT,ISTRUK,ISPA,NOUTP,JA,NANF,NEND,KL 90
EQUIVALENCE (FEST(1),NFEST(1)) 100

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3070 MDIMF=0 110
3080 MDIMP=0 120
3090 NERF=0 130
3100 NERP=0 140
3110 NEFE=NFE 150
3120 IF(NFE.EQ.1) NEFE=0 160
3130 CALL DOPW(8HSTRK,NN) 170
3140 CALL DOPW(8HSGA,NTYP(1)) 180
3150 CALL DOPW(8HSGN,NTYP(2)) 190
3160 CALL DOPW(8HSGF,NTYP(3)) 200
3170 CALL DOPW(8HSGT,NTYP(4)) 210
3180 CALL DOPW(8HBEST,NTYP(5)) 220
3190 CALL DOPW(8HMUEL,NTYP(6)) 230
3200 WRITE(NOUTP,9000) 240
3210 9000 FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 3') 250
3220 WRITE(NOUTP,9001) 260
3230 9001 FORMAT(' PROGRAMM ZUR BERECHNUNG VON RESONANZSELBSTABSCHIRMFAKTORE 270
IN VON PUNKTWEISE GEGEBENEN WIRKUNGSQUERSCHNITTEN'/) 280
IY=NE-NANF 290
NEN=NE+1-NEND 300
NSATZ(1)=3 310
NFEST(1)=MAT 320
NFEST(2)=NTYP(5) 330
JI=NEN-1 340
DO 3 LI=IY,JI 350
DO 3 LA=1,MI 360
ZB(LA,LI)=0. 370
DO 300 N=1,5 380
300 ZA(LA,LI,N)=0. 390
DO 3 N=1,4 400
XN(LA,LI,N)=0. 410
IF(ISTRUK.EQ.0.AND.N.EQ.3) XN(LA,LI,N)=1. 420
3 CONTINUE 430
NFEST(3)=NTYP(6) 440
CALL NDFLOC(NSUCH1,NSATZ,FEST,NUDAT,NC) 450
IF(NSUCH1)85,85,1001 460
85 WRITE(NOUTP,87){NFEST(K),K=1,3} 470
87 FORMAT(1H0/' ***ERROR 3. 1 : THE DATA FOR ',3A6,' COULD NOT BE FOU 480
IND IN THE KEDAK LIBRARY') 490
GO TO 100 500
1001 EMU(1)=FEST(4) 510
XMU(1)=FEST(5) 520
M=2 530
1005 IF(EMU(1)-ENG(IY))1002,1002,1004 540
1004 WRITE(NOUTP,11) ENG(IY),(NFEST(K),K=1,3),EMU(1) 550
IY=IY+1 560
IF(IY-JI)1005,1005,1000 570
1002 CALL NDFNXT(NSUCH1,NSATZ,FEST,NUDAT,NC) 580
IF(NSUCH1)1010,1010,1006 590
1006 IF(FEST(4)-ENG(IY))1003,1003,1007 600
1003 EMU(1)=FEST(4) 610
XMU(1)=FEST(5) 620
GO TO 1002 630
1007 IF(FEST(4)-ENG(NEN))1008,1014,1014 640
1008 EMU(M)=FEST(4) 650
XMU(M)=FEST(5) 660

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M=M+1
IF(M-MDIM)1002,1002,2000
2000 M=1
MDIMF=MDIMF+1
GO TO 1002
1014 EMU(M)=FEST(4)
XMU(M)=FEST(5)
GO TO 1009
1010 M=M-1
DO 1011 LI=IY,NEN
IF(EMU(M)-ENG(LI))1013,1012,1011
1011 CONTINUE
1012 NEN=LI
LI=NE-NEN+1
IF(NEN-1.LT.JI) WRITE (NOUTP,1015) LI
1015 FORMAT(1H0/' ***WARNING 3. 1 : IT IS ONLY POSSIBLE TO CALCULATE UN
1TIL GROUP ',I4,' BECAUSE KEDAK DOES NOT CONTAIN FURTHER MUEL VALUE
2S')
JI=NEN-1
GO TO 1009
1013 NEN=LI-1
LI=NE-NEN+1
WRITE (NOUTP,1015) LI
JI=NEN-1
1009 MM=M
IF(MDIMF.EQ.0) GO TO 2001
MDIMP=(MDIMF-1)*MDIM+M+1
RETURN
2001 DO 2222 I=1,4
IF(I.EQ.3.AND.ISTRUK.EQ.0) GO TO 2222
IF(IY-JI)400,400,1000
400 L=IY
LY=2
NFEST(3)=NTYP(I)
CALL NDFLOC(NSUCH1,NSATZ,FEST,NUDAT,NC)
IF(NSUCH1)85,85,86
86 E(1)=FEST(4)
SIG1=FEST(5)
IF(I-3)121,121,120
121 EN(1,I)=FEST(4)
SN(1,I)=FEST(5)
IST=1
IF(I.NE.2) GO TO 120
DO 3012 KK=2,MM
IF(EMU(KK)-E(1))3012,3013,3014
3012 CONTINUE
GO TO 3007
3013 XMU1=XMU(KK)
GO TO 120
3014 XMU1=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(1)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
120 IX=1
66 IF(E(1)-ENG(L))8,29,10
10 WRITE(NOUTP,11) ENG(L),(NFEST(K),K=1,3),E(1)
11 FORMAT(1H0/' ***WARNING 3. 2 : THE LOWER ENERGY GROUP BOUNDARY ',
1E16.8,'EV IS NOT IN THE AVAILABLE ENERGY RANGE IN KEDAK FOR ',3A5/

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670 2' THE LOWER ENERGY GROUP BOUNDARY HAS BEEN MODYFIED TO ',E16.8,'EV 1230
680 3') 1240
690 IY=IY+1 1250
700 GO TO 201 1260
710 8 CALL NDFNXT(NSUCH1,NSATZ,FEST,NUDAT,NC) 1270
720 IF(NSUCH1)12,12,13 1280
730 12 WRITE(NOUTP,14)NFEST(1),NFEST(2),NTYP(I),E(1) 1290
740 14 FORMAT(1H0/' ***WARNING 3. 3 : THE DESIRED DATA FOR ',3A5,' ARE ON 1300
750 1LY AVAILABLE IN KEDAK UNTIL ',E16.8,'EV') 1310
760 JI=L 1320
770 IF(NERF.EQ.0) GO TO 34 1330
780 NERP=(NERF-1)*NER+IST+1 1340
790 RETURN 1350
800 13 IF(FEST(4)-ENG(L))18,19,15 1360
810 19 IJ=1 1370
820 GO TO 20 1380
830 18 IJ=2 1390
840 20 E(1)=FEST(4) 1400
850 SIG1=FEST(5) 1410
860 IF(I-3)123,123,122 1420
870 123 EN(1,I)=FEST(4) 1430
880 SN(1,I)=FEST(5) 1440
890 IST=1 1450
900 IF(I.NE.2) GO TO 122 1460
910 DO 3000 KK=2,MM 1470
920 IF(EMU(KK)-E(1))3000,3001,3002 1480
930 3000 CONTINUE 1490
940 3007 WRITE (NOUTP,3003) NTYP(I),NTYP(I) 1500
950 3003 FORMAT(1H0/' ***WARNING 3. 4 : THE TYPE ',A6,' CAN NOT BE CALCULAT 1510
960 1ED'/' BECAUSE THE ENERGY POINTS FOR MUEL ARE ALL LESS THAN THE FIR 1520
970 2ST ENERGY POINT OF ',A6) 1530
980 GO TO 2 1540
990 3001 XMU1=XMU(KK) 1550
1000 GO TO 122 1560
1010 3002 XMU1=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(1)-EMU(KK-1))/(EMU(KK)- 1570
1020 1EMU(KK-1)) 1580
1030 122 IX=1 1590
1040 GO TO (29,8),IJ 1600
1050 15 E(2)=FEST(4) 1610
1060 SIG2=FEST(5) 1620
1070 SIG1=SIG1+(SIG2-SIG1)*(ENG(L)-E(1))/(E(2)-E(1)) 1630
1080 E(1)=ENG(L) 1640
1090 IF(I-3)125,125,124 1650
1100 125 EN(1,I)=E(1) 1660
1110 SN(1,I)=SIG1 1670
1120 EN(2,I)=E(2) 1680
1130 SN(2,I)=SIG2 1690
1140 IST=2 1700
1150 IF(I.NE.2) GO TO 124 1710
1160 IK=2 1720
1170 IF(KK.GT.2) IK=KK-1 1730
1180 DO 3004 KK=IK,MM 1740
1190 IF(EMU(KK)-E(1))3004,3005,3006 1750
1200 3004 CONTINUE 1760
1210 GO TO 3007 1770
1220 3005 XMU1=XMU(KK) 1780

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GC TO 3008
3006 XMU1=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(1)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
3009 IF(KK.GT.2) IK=KK-1
DO 3009 KK=IK,MM
IF(EMU(KK)-E(2))3009,3010,3011
3009 CONTINUE
GC TO 3007
3010 XMU2=XMU(KK)
GO TO 124
3011 XMU2=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(2)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
124 IX=2
29 IF(NEFE)210,101,210
101 FLUSS(1)=PHI(E(1))
GO TO (9,16),IX
210 DO 102 LR=2,NEFE
IF(REFE(LR)-E(1))102,103,104
102 CONTINUE
104 FLUSS(1)=EFE(LR-1)+(EFE(LR)-EFE(LR-1))/(REFE(LR)-REFE(LR-1))*(E(1)
1-REFE(LR-1))
GO TO (9,16),IX
103 FLUSS(1)=EFE(LR)
GO TO (9,16),IX
9 CALL NDFNXT(NSUCH1,NSATZ,FEST,NUDAT,NC)
IF(NSUCH1)12,12,22
22 E(2)=FEST(4)
SIG2=FEST(5)
IF(E(2)-ENG(NEN))3015,3015,16
3015 IF(I.NE.2) GO TO 16
IK=KK-1
DO 3016 KK=IK,MM
IF(EMU(KK)-E(2))3016,3017,3018
3016 CONTINUE
GO TO 3007
3017 XMU2=XMU(KK)
GO TO 16
3018 XMU2=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(2)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
16 IF(E(2)-ENG(L+1))26,26,25
25 E(3)=E(2)
SIG3=SIG2
IX=-1
SIG2=SIG1+(SIG3-SIG1)*(ENG(L+1)-E(1))/(E(3)-E(1))
E(2)=ENG(L+1)
IF(I-3)127,127,23
127 IST=IST+1
ISTM(I)=IST
IF(IST.LE.NER) GO TO 2002
IST=1
NERF=NERF+1
2002 EN(IST,I)=E(2)
SN(IST,I)=SIG2
IF(I.NE.2) GO TO 23
IK=KK-1
DO 3020 KK=IK,MM

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IF(EMU(KK)-E(2))3020,3021,3022
3020 CONTINUE
GO TO 3007
3021 XMU2=XMU(KK)
GO TO 23
3022 XMU2=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(2)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
GO TO 23
26 IX=0
128 IF(I-3)126,126,23
126 IST=IST+1
ISTM(I)=IST
IF(IST.LE.NER) GO TO 2003
IST=1
NERF=NERF+1
2003 EN(IST,I)=E(2)
SN(IST,I)=SIG2
IF(I.NE.2) GO TO 23
IK=KK-1
DO 3024 KK=IK,MM
IF(EMU(KK)-E(2))3024,3025,3026
3024 CONTINUE
GO TO 3007
3025 XMU2=XMU(KK)
GO TO 23
3026 XMU2=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(2)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
23 IF(NEFE)27,28,27
28 FLUSS(2)=PHI(E(2))
GO TO 30
27 DO 31 LR=LY,NEFE
IF(REFE(LR)-E(2))31,32,33
31 CONTINUE
LR=NEFE
33 FLUSS(2)=EFE(LR-1)+(EFE(LR)-EFE(LR-1))/(REFE(LR)-REFE(LR-1))*(E(2)
1-REFE(LR-1))
GO TO 36
32 FLUSS(2)=EFE(LR)
36 LY=LR
30 IF(I-4)49,601,49
601 KK=3
LRJ=2
IF(ISTRUK.EQ.1) LRJ=3
DO 60 JAN=1,LRJ
IST=ISTM(JAN)
M=1
50 DO 51 JJ=2,IST
IF(EN(JJ,JAN)-E(M))51,52,53
51 CONTINUE
JJ=IST
53 ST(M)=SN(JJ-1,JAN)+(SN(JJ,JAN)-SN(JJ-1,JAN))/(EN(JJ,JAN)-
1EN(JJ-1,JAN))*(E(M)-EN(JJ-1,JAN))
GO TO 54
52 ST(M)=SN(JJ,JAN)
54 IF(M-1)57,57,58
57 M=2

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GO TO 50
58 EV=0.5*(E(2)-E(1))
IF(KK.GT.2) IK=KK-1
DO 4000 KK=IK,MM
IF(EMU(KK)-E(1))4000,4001,4002
4000 CONTINUE
GO TO 3007
4001 XMU1=XMU(KK)
GO TO 4003
4002 XMU1=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(1)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
4003 IF(KK.GT.2) IK=KK-1
DO 4004 KK=IK,MM
IF(EMU(KK)-E(2))4004,4005,4006
4004 CONTINUE
GO TO 3007
4005 XMU2=XMU(KK)
GO TO 4007
4006 XMU2=XMU(KK-1)+(XMU(KK)-XMU(KK-1))*(E(2)-EMU(KK-1))/(EMU(KK)-
1EMU(KK-1))
4007 DO 60 LI=1,MI
VA=FLUSS(1)/(SIG1 +SIGO(LI))
VB=FLUSS(2)/(SIG2 +SIGO(LI))
IF(JAN.NE.2) GO TO 3028
ZB(LI,L )=ZB(LI,L )+EV*(ST(1)*XMU1 *VA+ST(2)*XMU2*VB)
3028 ZA(LI,L,JAN)=ZA(LI,L,JAN)+EV*(ST(1)*VA+ST(2)*VB)
3029 XN(LI,L,JAN)=XN(LI,L,JAN)+EV*(VA+VB)
IF(JAN.NE.2) GO TO 60
VA=FLUSS(1)/(SIG1+SIGO(LI))*2
VB=FLUSS(2)/(SIG2+SIGO(LI))*2
ZA(LI,L,4)=ZA(LI,L,4)+EV*(ST(1)*XMU1*VA+ST(2)*XMU2*VB)
3031 XN(LI,L,4)=XN(LI,L,4)+EV*(VA+VB)
ZA(LI,L,5)=ZA(LI,L,5)+EV*(SIG1*VA+SIG2*VB)
60 CONTINUE
ST(1)=ST(2)
49 E(1)=E(2)
SIG1=SIG2
FLUSS(1)=FLUSS(2)
IF(I.EQ.2) XMU1=XMU2
IF(E(2)-ENG(L+1))9,34,34
34 IF(I-3)69,69,62
62 JFAK=0
DO 4 LA=1,LRJ
4 XINTE(L,LA)=ZA(MI,L,LA)/XN(MI,L,LA)
XINTE(L,4)=ZA(MI,L,5)/XN(MI,L,4)
XI(L)=ZB(MI,L)/XN(MI,L,2)
XII(L)=ZA(MI,L,4)/XN(MI,L,4)
DO 63 LI=1,MI
DO 363 LA=1,LRJ
IF(XINTE(L,LA).EQ.0.OR.XN(LI,L,LA).EQ.0) GO TO 366
IF(ISTRUK.EQ.1.AND.LA.EQ.1) GO TO 5301
SE(LI,L,LA)=ZA(LI,L,LA)/(XN(LI,L,LA)*XINTE(L,LA))
GO TO 363
5301 SE(LI,L,LA)=(ZA(LI,L,LA)/XN(LI,L,LA)-ZA(LI,L,3)/XN(LI,L,3))/(
IXINTE(L,LA)-XINTE(L,3))
GC TO 363

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366 SE(LI,L,LA)=1.
JFAK=1
363 CONTINUE
IF(XI(L).NE.0) GO TO 367
SE(LI,L,4)=1.
JFAK=1
GO TO 900
367 SE(LI,L,4)=ZB(LI,L)/(XN(LI,L,2)*XI(L))
900 IF(XII(L).NE.0) GO TO 467
SE(LI,L,5)=1.
JFAK=1
GO TO 901
467 SE(LI,L,5)=ZA(LI,L,4)/(XN(LI,L,4)*XII(L))
901 IF(XINTE(L,4).EQ.0) GO TO 368
SE(LI,L,6)=ZA(LI,L,5)/(XN(LI,L,4)*XINTE(L,4))
GO TO 63
368 SE(LI,L,6)=1.
JFAK=1
63 CONTINUE
IF(JFAK.EQ.0) GO TO 69
LZ=NE-L
WRITE (NOUTP,5300) LZ
5300 FORMAT(1HO/' ***WARNING 3. 5 : IN THE ENERGY GROUP ',I4,' THE SELF
1SHIELDING FACTORS ARE SET TO 1, BECAUSE THE CROSS SECTIONS '/' FOUN
2D ON KEDAK FOR THIS GROUP ALL ARE ZERC')
69 L=L+1
IF(L-JI)64,64,2
64 IF(IX)166,9,9
166 IF(E(3)-ENG(L+1))67,67,68
68 E(2)=ENG(L+1)
SIG2=SIG1+(SIG3-SIG1)*(E(2)-E(1))/(E(3)-E(1))
GO TO 128
67 E(2)=E(3)
SIG2=SIG3
GO TO 26
201 L=L+1
IF(L-JI)66,66,2
2 IF(NERF.EQ.0) GO TO 2222
NERP=(NERF-1)*NER+IST+1
RETURN
2222 CONTINUE
N=0
WRITE (JA) N,NN
DO 70 LI=IY,JI
IN=5
WRITE (JA) IN,MAT ,LI,ENG(LI),ENG(LI+1)
LZ=NE-LI
WRITE(NOUTP,35) MAT ,LZ,ENG(LI),ENG(LI+1)
35 FORMAT(1HO,A5,9H GRUPPE =I3,3X,15HGRUPPENGRENZEN 2E16.8)
IF(ISTRUK.EQ.1) GO TO 5201
WRITE (NOUTP,5200)
5200 FCRMAT(1HO,22X,7HSIGMA A,11X,7HSIGMA N,10X,9HSIGMA NO1,9X,
18HSIGMA N1,10X,8HSIGMA T1/)
GC TO 5203
5201 WRITE (NOUTP,5202)
5202 FCRMAT(1HO,22X,'SIGMA C',11X,'SIGMA N',11X,'SIGMA F',10X,

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1'SIGMA N01',9X,'SIGMA N1',10X,'SIGMA T1'//
XINTE(LI,1)=XINTE(LI,1)-XINTE(LI,3)
5203 IN=LRJ+3
WRITE(JA) IN,(XINTE(LI,LA),LA=1,LRJ),XI(LI),XI(LI),XINTE(LI,4)
WRITE(NOUTP,37)(XINTE(LI,LA),LA=1,LRJ),XI(LI),XI(LI),XINTE(LI,4)
37 FORMAT(16X,6E18.8)
IF(ISTRUK.EQ.1) GO TO 5204
WRITE(NOUTP,39)
39 FORMAT(1H0,5X,4HSGO,15X,2HFA,16X,2HFN,16X,4HFN01,14X,3HFN1,15X,
13HFT1/)
GO TO 5206
5204 WRITE(NOUTP,5205)
5205 FORMAT(1H0,5X,4HSGO,15X,2HFC,16X,2HFN,16X,2HFF,14X,4HFN01,14X,
13HFN1,15X,3HFT1/)
5206 IN=LRJ+4
DO 70 I=1,MI
WRITE(JA) IN,SIGO(I),(SE(I,LI,LA),LA=1,LRJ),(SE(I,LI,LA),LA=4,6)
70 WRITE(NOUTP,38) SIGO(I),(SE(I,LI,LA),LA=1,LRJ),(SE(I,LI,LA),LA=4,6
1)
38 FORMAT(E16.8,6E18.8)
1000 CONTINUE
71 KL=KL+1
RETURN
100 STOP
END

SUBROUTINE SUND (MM,ENG,NFE,REFE,EFE,ITYP,ITNAM,SGC,DUE,XINTE,
1ZINT,XNEN,STREU,LDIM,LDIMP,SE,FSE)
INTEGER*2 IHC(2),IMP(2)
REAL*8 MAT,ITNAM,FEST,NFEST,NN,MMM
DIMENSION ENG(MM),REFE(NFE),EFE(NFE),ITNAM(ITYP),SGC(MM),E(3),
1FLUSS(2),DUE(MM),XINTE(MM),ZINT(MM),XNEN(MM),STREU(MM),NADAT(2),
2NSATZ(4),FEST(10),NFEST(10),NN(11),SE(LDIM),FSE(LDIM),SFN(2)
COMMON MAT,ISTRUK,ISPA,NOUTP,JA,NANF,NEND,KL
EQUIVALENCE (FEST(1),NFEST(1)),(IHC(1),FM)
DATA IMP/'PU','U '/
WRITE(NOUTP,9000)
9000 FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 4')
WRITE(NOUTP,9001)
9001 FORMAT(' PROGRAMM ZUR BERECHNUNG VON GRUPPENWIRKUNGSQUERSCHNITTEN
1BEI UNENDLICHER VERDUENNUNG')
CALL FSPIE
LDIF=0
LDIMP=0
NEFE=NFE
IF(NFE.EQ.1) NEFE=0
NGR=MM-NANF
NEN=MM+1-NEND
CALL DOPW (8HBEST ,NFEST(2))
CALL DOPW (8HSGA ,NN(1))
CALL DOPW (8HSGF ,NN(2))
CALL DOPW (8HSGN ,NN(3))

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CALL DOPW (8HMUEL ,NN(4))
CALL DOPW (8HNUE ,NN(5))
CALL DOPW (8HSGC ,NN(6))
CALL DOPW (8HSTR ,NN(7))
CALL DOPW (8HH 01 ,NN(8))
CALL DOPW (8HALPHA ,NN(9))
CALL DOPW (8HETA ,NN(10))
CALL DOPW (8HSGG ,NN(11))
CALL DOPW (8HSGKE ,MMM)
INR=0
LL=0
HM=MAT
IF(IHC(1).EQ.IMP(1)) GO TO 802
IF(IHC(1).NE.IMP(2)) GO TO 822
802 LAR=1
DO 803 JJ=1,ITYP
IF(ITNAM(JJ).EQ.NN(1)) GO TO 805
IF(ITNAM(JJ).NE.NN(2)) GO TO 803
805 LL=LL+1
IF(LL-2)803,801,803
803 CONTINUE
IF(LL-2)807,801,807
807 WRITE(NOUTP,808) MAT
808 FORMAT(1H0/' ***WARNING 4. 1 : THE GRUOP CROSS SECTION SGC CAN NOT
1 BE CALCULATED FOR ',A5/' BECAUSE THE REACTION TYPES SGF AND SGA A
2RE NOT SPECIFIED IN THE INPUT')
GO TO 801
822 LAR=0
DO 821 JJ=1,ITYP
IF(ITNAM(JJ).EQ.NN(1)) GO TO 801
821 CONTINUE
GO TO 807
801 NFEST(1)=MAT
IGRUP=NEN-1
DO 870 JJ=NGR,IGRUP
870 SGC(JJ)=0.
JJJ=1
DO 3 JJ=1,ITYP
KGRU=0
NSATZ(1)=3
KSIK=0
IF(ITNAM(JJ).NE.NN(9).OR.ITNAM(JJ).NE.NN(10)) GO TO 871
IF(LAR.EQ.0) GO TO 875
KSIK=1
NFEST(3)=NN(11)
GO TO 2003
875 WRITE(NOUTP,876)MAT,ITNAM(JJ)
876 FORMAT(1H0/' ***WARNING 4. 2 : THE GRUOP CROSS SECTION FOR ',2A8,
1' CAN NOT BE CALCULATED'/' BECAUSE THE VALUES OF SGF ARE ZERO')
GO TO 3
874 KSIK=2
NFEST(3)=NN(2)
GO TO 2003
879 KSIK=3
GO TO 2002
871 IF(ITNAM(JJ).EQ.NN(4)) GO TO 2001

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IF(ITNAM(JJ).EQ.NN(5)) GO TO 2002	830	5 WRITE(NOUTP,7) (NFEST(K),K=1,3)	1390
GO TO 2003	840	7 FORMAT(1H0/' ***WARNING 4. 4 : DATA FOR ',3A5,' COULD NOT BE FOUND	1400
2001 NFEST(3)=NN(3)	850	1 IN THE KEDAK LIBRARY')	1410
GO TO 2004	860	GC TO 300	1420
2002 NFEST(3)=NN(2)	870	6 E(1)=FEST(4)	1430
2004 CALL NDFLOC(NSUCH1,NSATZ,FEST,NUDAT,NC)	880	SIG1=FEST(5)	1440
IF(NSUCH1)2005,2005,2006	890	IF(ITNAM(JJ).EQ.NN(4).OR.ITNAM(JJ).EQ.NN(5)) GO TO 2023	1450
2005 SE(1)=ENG(NGR)	900	GO TO 66	1460
FSE(1)=0.	910	2023 DO 2016 KK=2,MMM	1470
SE(2)=ENG(NEN)	920	IK=KK	1480
FSE(2)=0.	930	IF(SE(KK)-E(1))2016,2017,2018	1490
M=2	940	2016 CCNTINUE	1500
WRITE(NOUTP,2070) FEST(3),ITNAM(JJ)	950	2022 WRITE (NOUTP,2019) ITNAM(JJ)	1510
2070 FORMAT(1H0/' ***WARNING 4. 3 : TYPE ',A6,' IS SET TO ZERO. THIS TY	960	2019 FORMAT(1H0/' ***WARNING 4. 5 : TYPE ',A6,' CAN NOT BE CALCULATED B	1520
1PE IS USED TO CALCULATE THE GROUP CROSS SECTION ',A6)	970	1ECAUSE THE ENERGY SCALE OF'/' THIS TYPE BEGINS AT A HIGHER ENERGY	1530
GO TO 2003	980	2THAN THE HIGHEST ENERGY GROUP BOUNDARY')	1540
2006 SE(1)=FEST(4)	990	GC TO 3	1550
FSE(1)=FEST(5)	1000	2017 SFN(1)=FSE(KK)	1560
IF(SE(1)-ENG(NGR))2008,2008,2009	1010	GO TO 66	1570
2009 WRITE (NOUTP,11) ENG(NGR),(NFEST(K),K=1,3),SE(1)	1020	2018 SFN(1)=FSE(KK-1)+(FSE(KK)-FSE(KK-1))*(E(1)-SE(KK-1))/(SE(KK)-SE(1580
2008 M=2	1030	1KK-1))	1590
2013 CALL NDFNXT(NSUCH1,NSATZ,FEST,NUDAT,NC)	1040	66 IF(E(1)-ENG(L))8,9,10	1600
IF(NSUCH1)2223,2223,2010	1050	10 WRITE (NOUTP,11) ENG(L),(NFEST(K),K=1,3),E(1)	1610
2010 IF(FEST(4)-ENG(NGR))2011,2011,2012	1060	11 FORMAT(1H0/' ***WARNING 4. 6 : THE LOWER ENERGY GROUP BOUNDARY ',	1620
2011 SE(1)=FEST(4)	1070	1E16.8,'EV IS NOT IN THE AVAILABLE ENERGY RANGE IN KEDAK FOR '/A,	1630
FSE(1)=FEST(5)	1080	23A5,' THE LOWER ENERGY GROUP BOUNDARY HAS BEEN MODYFIED TO ',E16.8	1640
GO TO 2013	1090	3,'EV')	1650
2012 IF(FEST(4)-ENG(NEN))2014,2015,2015	1100	IF(E(1)-ENG(L+1))9,200,200	1660
2014 SE(M)=FEST(4)	1110	8 CALL NDFNXT(NSUCH1,NSATZ,FEST,NUDAT,NC)	1670
FSE(M)=FEST(5)	1120	IF(NSUCH1)12,12,13	1680
M=M+1	1130	12 WRITE(NOUTP,14)NFEST(1),NFEST(2),ITNAM(JJ),E(1)	1690
IF(M-LDIM)2013,2013,2090	1140	14 FORMAT(1H0/' ***WARNING 4. 7 : THE DESIRED DATA FOR ',3A5,' ARE ON	1700
2090 M=1	1150	1LY AVAILABLE IN KEDAK UNTIL ',E16.8,'EV')	1710
LDIF=LDIF+1	1160	IZLM=-1	1720
GO TO 2013	1170	GC TO 50	1730
2223 M=M-1	1180	13 IZLM=1	1740
GO TO 2003	1190	IF(FEST(4)-ENG(L))18,19,15	1750
2015 SE(M)=FEST(4)	1200	19 IJ=1	1760
FSE(M)=FEST(5)	1210	GC TO 20	1770
2003 LY=2	1220	18 IJ=2	1780
IF(LDIF.EQ.0) GO TO 2091	1230	20 E(1)=FEST(4)	1790
LDIMP=(LDIF-1)*LDIM+M+1	1240	SIG1=FEST(5)	1800
RETURN	1250	IF(ITNAM(JJ).EQ.NN(4).OR.ITNAM(JJ).EQ.NN(5)) GO TO 2024	1810
2091 MMM=M	1260	GO TO (9,8),IJ	1820
L=NGR	1270	2024 IF(KK.GT.2) KK=KK-1	1830
M=L	1280	DO 2059 IK=KK,MMM	1840
DO 4 LI=L,IGRUP	1290	IF(SE(IK)-E(1))2059,2020,2021	1850
ZINT(LI)=0.	1300	2059 CONTINUE	1860
DUE(LI)=0.	1310	GO TO 2022	1870
4 XNEN(LI)=0.	1320	2020 SFN(1)=FSE(IK)	1880
IF(KS IK.NE.3) GO TO 971	1330	GO TO (9,8),IJ	1890
NFEST(3)=NN(5)	1340	2021 SFN(1)=FSE(IK-1)+(FSE(IK)-FSE(IK-1))*(E(1)-SE(IK-1))/(SE(IK)-SE(1900
GO TO 970	1350	1IK-1))	1910
971 IF(KS IK.EQ.0) NFEST(3)=ITNAM(JJ)	1360	GC TO (9,8),IJ	1920
970 CALL NDFLOC(NSUCH1,NSATZ,FEST,NUDAT,NC)	1370	15 E(2)=FEST(4)	1930
IF(NSUCH1)5,5,6	1380	SIG2=FEST(5)	1940

SIG1=SIG1+(SIG2-SIG1)*(ENG(L)-E(1))/(E(2)-E(1))	1950	GO TO 2022	2510
E(1)=ENG(L)	1960	2031 SFN(2)=FSE(IK)	2520
IF(ITNAM(JJ).EQ.NN(4).OR.ITNAM(JJ).EQ.NN(5)) GO TO 2037	1970	GO TO 23	2530
GO TO 16	1980	2032 SFN(2)=FSE(IK-1)+(FSE(IK)-FSE(IK-1))*(E(2)-SE(IK-1))/(SE(IK)-SE(IK-1))	2540
2037 IF(IK.GT.2) IK=IK-1	1990	1(IK-1))	2550
KK=IK	2000	GC TO 23	2560
DO 2038 IK=KK,MMMM	2010	26 IX=0	2570
IF(SE(IK)-E(1))2038,2039,2040	2020	23 IF(ITNAM(JJ).EQ.NN(4)) GO TO 72	2580
2038 CONTINUE	2030	IF(ITNAM(JJ).NE.NN(5)) GO TO 70	2590
GO TO 2022	2040	72 IFL=1	2600
2039 SFN(1)=FSE(IK)	2050	GO TO 573	2610
GO TO 2041	2060	972 ZINT(L)=ZINT(L)+0.5*(SIG1*SFN(1)*FLUSS(1)+SIG2*SFN(2)*FLUSS(2))*	2620
2040 SFN(1)=FSE(IK-1)+(FSE(IK)-FSE(IK-1))*(E(1)-SE(IK-1))/(SE(IK)-SE(IK-1))	2070	1(E(2)-E(1))	2630
1(IK-1))	2080	XNEN(L)=XNEN(L)+0.5*(SFN(1)*FLUSS(1)+SFN(2)*FLUSS(2))*(E(2)-E(1))	2640
2041 IF(IK.GT.2) IK=IK-1	2090	IJ=0	2650
KK=IK	2100	GO TO 46	2660
DO 2042 IK=KK,MMMM	2110	70 IJ=1	2670
IF(SE(IK)-E(2))2042,2043,2044	2120	IFL=2	2680
2042 CONTINUE	2130	573 IF (NEFE) 27,28,27	2690
IF(SE(MMMM)-ENG(L+1))2042,2043,2044	2140	28 FLUSS(1)=PHI(E(1))	2700
2043 SFN(2)=FSE(IK)	2150	FLUSS(2)=PHI(E(2))	2710
GO TO 16	2160	GO TO (972,30),IFL	2720
2044 SFN(2)=FSE(IK-1)+(FSE(IK)-FSE(IK-1))*(E(2)-SE(IK-1))/(SE(IK)-SE(IK-1))	2170	27 J=1	2730
1(IK-1))	2180	GO TO 35	2740
GO TO 16	2190	34 J=2	2750
9 CALL NDFNXT(NSUCH1,NSATZ,FEST,NUDAT,NC)	2200	LY=LR	2760
IF(NSUCH1)12,12,22	2210	35 DO 31 LR=LY,NEFE	2770
22 E(2)=FEST(4)	2220	IF(REFE(LR)-E(J))31,32,33	2780
SIG2=FEST(5)	2230	31 CONTINUE	2790
IF(E(2)-ENG(NEN))2071,2071,16	2240	33 FLUSS(J)=EFE(LR-1)+(EFE(LR)-EFE(LR-1))/(REFE(LR)-REFE(LR-1))*(E(J)-	2800
2071 IZLM=1	2250	1-REFE(LR-1))	2810
IF(ITNAM(JJ).EQ.NN(4).OR.ITNAM(JJ).EQ.NN(5)) GO TO 2033	2260	GO TO (34,36),J	2820
GO TO 16	2270	32 FLUSS(J)=EFE(LR)	2830
2033 IF(IK.GT.2) IK=IK-1	2280	GO TO (34,36),J	2840
KK=IK	2290	36 LY=LR	2850
DO 2034 IK=KK,MMMM	2300	GO TO (972,30),IFL	2860
IF(SE(IK)-E(2))2034,2035,2036	2310	30 ZINT(L)=ZINT(L)+0.5*(SIG1*FLUSS(1)+SIG2*FLUSS(2))*(E(2)-E(1))	2870
2034 CONTINUE	2320	IF(MAT .NE.NN(8)) GO TO 120	2880
GO TO 2022	2330	IF(ITNAM(JJ).NE.NN(3)) GO TO 120	2890
2035 SFN(2)=FSE(IK)	2340	121 EN=ALOG(E(2)/E(1))	2900
GO TO 16	2350	SF1=SIG1*FLUSS(1)	2910
2036 SFN(2)=FSE(IK-1)+(FSE(IK)-FSE(IK-1))*(E(2)-SE(IK-1))/(SE(IK)-SE(IK-1))	2360	SF2=SIG2*FLUSS(2)	2920
1(IK-1))	2370	DUE(L)=DUE(L)+SF1*EN+SF2-SF1-(SF2-SF1)/(E(2)-E(1))*E(1)*EN	2930
16 IF(E(2)-ENG(L+1))26,26,25	2380	120 XNEN(L)=XNEN(L)+0.5*(FLUSS(1)+FLUSS(2))*(E(2)-E(1))	2940
25 E(3)=E(2)	2390	46 IF(E(2)-ENG(L+1))47,50,50	2950
SIG3=SIG2	2400	50 IF(ITNAM(JJ).EQ.NN(5).OR.ITNAM(JJ).EQ.NN(9)) GO TO 503	2960
IX=-1	2410	GO TO 500	2970
SIG2=SIG1+(SIG3-SIG1)*(ENG(L+1)-E(1))/(E(3)-E(1))	2420	503 IF(XNEN(L).NE.0) GO TO 500	2980
E(2)=ENG(L+1)	2430	M=M+1	2990
IF(ITNAM(JJ).EQ.NN(4).OR.ITNAM(JJ).EQ.NN(5)) GO TO 2029	2440	KGRU=1	3000
GO TO 23	2450	GO TO 122	3010
2029 IF(IK.GT.2) IK=IK-1	2460	500 XINTE(L)=ZINT(L)/XNEN(L)	3020
KK=IK	2470	IF(MAT .NE.NN(8)) GO TO 122	3030
DO 2030 IK=KK,MMMM	2480	IF(ITNAM(JJ).NE.NN(3)) GO TO 122	3040
IF(SE(IK)-E(2))2030,2031,2032	2490	STREU(L)=DUE(L)/XNEN(L)	3050
2030 CONTINUE	2500	122 L=L+1	3060

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IF (IZLM) 300, 300, 452
452 E(1)=E(2)
SIG1=SIG2
SFN(1)=SFN(2)
IF(L-IGRUP) 51, 51, 52
51 IF(IX) 53, 9, 9
53 IF(E(3)-ENG(L+1)) 107, 107, 108
108 E(2)=ENG(L+1)
SIG2=SIG1+(SIG3-SIG1)*(E(2)-E(1))/(E(3)-E(1))
GO TO 23 IF(ITNAM(JJ).EQ.NN(4).OR.ITNAM(JJ).EQ.NN(5)) GO TO 2045
107 E(2)=E(3)
SIG2=SIG3
IF(ITNAM(JJ).EQ.NN(4).OR.ITNAM(JJ).EQ.NN(5)) GO TO 2025
GO TO 26
2025 IF(IK.GT.2) IK=IK-1 KK=IK
DO 2026 IK=KK, MMMM
IF(SE(IK)-E(3)) 2026, 2027, 2028
2026 CONTINUE
IF(SE(MMMM)-ENG(L+1)) 2022, 2027, 2028
2027 SFN(2)=FSE(IK)
GO TO 26
2028 SFN(2)=FSE(IK-1)+(FSE(IK)-FSE(IK-1))*(E(3)-SE(IK-1))/(SE(IK)-SE(IK-1))
GO TO 26
47 E(1)=E(2)
SIG1=SIG2
SFN(1)=SFN(2)
GO TO 9
52 NNA=NEN-1
IF(KGRU.EQ.0) GO TO 501
WRITE(NOUTP, 502) MAT, ITNAM(JJ)
502 FORMAT(/' ***WARNING 4. 8 : THE GROUP CROSS SECTION FOR', 2A9, ' CAN
INOT BE CALCULATED IN ALL ENERGY GROUPS REQUIRED'/' BECAUSE THE FIS
2SIGN CROSS SECTIONS ON KEDAK ARE EQUAL TO ZERO')
IF(MM-M.LT.MM-NNN) GO TO 3
501 IF(JJJ-1) 54, 56, 54
56 N=0
WRITE(JA) N, MMM
JJJ=2
54 N=2
WRITE(JA) N, M, NNN
J=4
N=NEN-M
IF(KSIK.EQ.3) GO TO 881
IF(KSIK.EQ.2) GO TO 877
IF(KSIK.EQ.0) GO TO 872
DO 873 J=M, IGRUP
873 STREU(J)=XINTE(J)
GO TO 874
877 DO 878 J=M, IGRUP
878 STREU(J)=STREU(J)/XINTE(J)
IF(ITNAM(JJ).EQ.NN(10)) GO TO 879
DO 880 J=M, IGRUP
880 XINTE(J)=STREU(J)
GO TO 872
881 DO 882 J=M, IGRUP

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** 2045 IF(IK.GT.2) IK=IK-1
KK=IK
DO 2046 IK=KK, MMMM
IF(SE(IK)-E(2)) 2046, 2047, 2048
2046 CONTINUE
GO TO 2022
2047 SFN(2)=FSE(IK)
GO TO 23
2048 SFN(2)=FSE(IK-1)+(FSE(IK)-FSE(IK-1))*
(FSE(IK-1))/(E(2)-SE(IK-1))
2(SE(IK)-SE(IK-1))
GO TO 23

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882 XINTE(J)=XINTE(J)/(1.+STREU(J))
872 WRITE(JA) J, MAT, ITNAM(JJ)
WRITE(JA) N, (XINTE(J), J=M, IGRUP)
IST=MM-M
NI=MM-NNN
WRITE(NOUTP, 910) MAT, ITNAM(JJ), IST, NI
910 FORMAT(IH0, 2A7, 9H GRUPPEI3, 4H BISI3)
WRITE(NOUTP, 909) (XINTE(J), J=M, IGRUP)
909 FORMAT(IH, 7E16.8/(1X, 7E16.8))
IF(ITNAM(JJ).EQ.NN(1)) GO TO 810
IF(ITNAM(JJ).EQ.NN(2)) GO TO 812
GO TO 816
810 DO 813 J=M, IGRUP
813 SGC(J)=SGC(J)+XINTE(J)
IF(INR.EQ.1) GO TO 818
INR=1
GO TO 3
812 DO 814 J=M, IGRUP
814 SGC(J)=SGC(J)-XINTE(J)
IF(INR.EQ.1) GO TO 818
INR=1
GO TO 3
818 J=2
WRITE(JA) J, M, NNN
J=4
WRITE(JA) J, MAT, NN(6)
WRITE(JA) N, (SGC(J), J=M, IGRUP)
WRITE(NOUTP, 910) MAT, NN(6), IST, NI
WRITE(NOUTP, 909) (SGC(J), J=M, IGRUP)
GO TO 3
816 IF(MAT.NE.NN(8)) GO TO 3
IF(ITNAM(JJ).NE.NN(3)) GO TO 3
J=2
WRITE(JA) J, M, NNN
J=4
WRITE(JA) J, MAT, NN(7)
WRITE(JA) N, (STREU(J), J=M, IGRUP)
WRITE(NOUTP, 910) MAT, NN(7), IST, NI
WRITE(NOUTP, 909) (STREU(J), J=M, IGRUP)
GO TO 3
300 IF(L-IGRUP) 600, 600, 52
600 DO 57 K=L, IGRUP
57 XINTE(K)=0.
GO TO 52
200 XINTE(L)=0.
L=L+1
IF(L-IGRUP) 66, 66, 52
3 CONTINUE
100 KL=KL+1
63 RETURN
END

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VI 37

C	CALCULATION OF THE INELASTIC SCATTERING PROBABILITIES	10	KI=2	570
C		20	CALL AKED(E, ET,SGIP,JMAT,KMAX,LMAX,SU,QUOT,WERT,AE,EMAX,EMIN,	580
	SUBROUTINE SCAT(NX,EG, XNUE,NFE,EF,FI,NE27,WAHR,PROB,VW,QUER,	30	1IA,NAB,NFE,EF,FI,NE,NX,NET,NETF,NETP,SGIT,NAE,NAEF,NAEP,ISG,ISGF,	590
	IE,NET,ET,NAE,AE,SU,QUOT,LBA,WEIN,AG,NETP,SGIT,NAEP,KMAX,ISG,ISGP,	40	2ISGP,IWE,IWF,IWP)	600
	2SGIP,IWE,IWP,WERT)	50	IF(NETF.GT.0.OR.ISGF.GT.0.OR.NAEF.GT.0.OR.IWF.GT.0) RETURN	610
	REAL*8 NAM,ISOT,NAM4,EO,EU,DIFF,ET,AE,EMIN,EMAX,E,SUM,QUER,WAHR,	60	IF(IA.GE.IANF) GOTO 205	620
	IPROB,WAI,VW,UWA,EH,AB	70	NEA=IANF	630
	DIMENSION AE(NAE),KMAX(NAE),ET(NET),SGIP(NAE,ISG),WERT(IWE),	80	GOTO 206	640
	1EG(NX),EF(NFE),FI(NFE),NN(4),SU(NX),QUOT(NX),WAHR(NX),PROB(NX),	90	205 NEA=IA	650
	2LBA(NX),VW(NX),QUER(NX),E(NE27),WEIN(NX),AG(NE27),SGIT(NET),	100	IF(IA.LE.IE) GOTO 206	660
	3ISOT(2)	110	WRITE(NA,179)	670
	COMMON NAM,ISTRUK,ISPA,NA,LIZ,NANF,NEND,KL	120	179 FORMAT (' ***ERROR 5.04 : WITH INCIDENT NEUTRON ENERGIES OUT OF	680
	DATA ISOT/'BEST ', 'ISOT1'/	130	1THE ENRGY GROUPS REQUESTED IN THE INPUT INELASTIC SCATTERING	690
	WRITE (NA ,9000)	140	2CANNOT OCCUR,SINCE THE FIRST ENERGY AT WHICH INELASTIC SCATTERING	700
9000	FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 5')	150	3CROSS SECTIONS ARE STORED ON KEDAK IS HIGHER THAN THE UPPER	710
	WRITE (NA ,9001)	160	4ENERGY LIMIT OF THE HIGHEST ENERGY GROUP (LOWEST GROUP NUMBER	720
9001	FORMAT(' PROGRAMM ZUR BERECHNUNG INELASTISCHER STREUMATRIZEN'/)	170	5ACCORDING TO THE ABN-SET')	730
	WRITE(NA,9002)	180	GO TO 999	740
9002	FORMAT(1H0/' PROBSGI°/1X)	190	206 KD=IE-NEA+1	750
	CALL FSPIE	200	WRITE(NA,172) KI,NAM,KD	760
	NN(1)=3	210	172 FORMAT(1H0,3X,1I,3X,A5,3X,I2)	770
	AE(1)=NAM	220	MZ=0	780
	AE(2)=ISOT(1)	230	CALL DOUB (8HSMTOT ,NAM4)	790
	AE(3)=ISOT(2)	240	26 WRITE(LIZ) MZ,NAM4	800
	CALL NDFLOC (I,NN,AE,K,K)	250	IF(ET(1).LT.E(IA+1)) GOTO 39	810
	IF(I.EQ.0) GO TO 1001	260	204 KD=KD-1	820
	A=AE(4)	270	39 KZ=KI+1	830
	NE=NX	280	WRITE (LIZ) KZ,NAM,KD	840
	IF(EG(1)-1.E-3)305,301,307	290	C	850
305	WRITE (NA,304)	300	LCCP OVER ALL OUTSCATTERING GROUPS TO BE CONSIDERED	860
304	FORMAT (' ***ERROR 5.01 : THE LOWER BOUNDARY OF THE LOWEST ENERGY	310	DO 1 I=NEA,IE	870
	1GROUP HAS TO BE LARGER OR EQUAL TO 1.E-3 WHICH IS THE LOWEST	320	NEI=NE-I	880
	2ENERGY ON KEDAK FOR ALL DATA TYPES. THIS IS NOT FULFILLED.')	330	IF(E(I+1).LE.ET(1)) GOTO 1	890
	GO TO 999	340	10 GES=0.	900
307	NE=NE27	350	IF(I-NAB)83,85,1002	910
	E(1)=0.	360	85 EH=EMAX	920
	E(1)=1.E-3	370	GO TO 89	930
	AG(1)=1.E-3	380	83 EH=E(I+1)	940
	DO 302 I=2,NE27	390	89 NEI=NE-1	950
	E(I)=0.	400	C	960
	AG(I)=EG(I-1)	410	LCCP OVER ALL INSCATTERING GROUPS	970
302	E(I)=EG(I-1)	420	DO 2 K=1,I	980
	GO TO 303	430	SUM=DBLE(0.0)	990
301	DO 300 I=1,NE	440	C	1000
	E(I)=0.	450	LCCP OVER ALL INELASTIC EXCITATION LEVELS OF THE MATERIAL CONSIDER	1010
	AG(I)=EG(I)	460	DO 3 J=1,JMAT	1020
	E(I)=EG(I)	470	IF(J.NE.1) GOTO 6	1030
300	CONTINUE	480	7 KMIN=0	1040
C	DETERMINATION OF THE EFFECTIVE INITIAL ENERGY GROUP OUT OF WHICH	490	GO TO 8	1050
C	INELASTIC SCATTERING OCCURS AND OF THE LAST GROUP UP TO WHICH	500	6 KMIN=KMAX(1)-KMAX(J)	1060
C	THE TRANSITION PROBABILITIES ARE REQUIRED	510	C	1070
303	NEI=NE-1	520	DETERMINATION OF THE LIMITS EU,EO FOR INTEGRATION OVER THE	1080
	DO 37 IM=1,NE1	530	OUTSCATTERING GROUP	1090
37	QUOT(IM)=0.0	540	8 IF((E(K)+AE(J)).LT.E(I)) GOTC 9	1100
	IANF=NE-NANF	550	17 IF((E(K)+AE(J)).GE.EH) GOTO 88	1110
	IE=NE-NEND	560	57 EU=E(K)+AE(J)	1120
			GO TO 44	
			9 IF((E(K+1)+AE(J)).LT.E(I)) GOTO 3	
			58 EU=E(I)	
			44 IF((E(K+1)+AE(J)).GE.EH) GOTC 61	

59	EO=E(K+1)+AE(J)	1130	IF(IF1.EQ.IAF) GOTO 185	1690
	GO TO 62	1140	LIE=IF1	1700
C 61	EO=EH	1150	KMB1=KMIN	1710
	DETERMINATION OF SIGMA AND THE WEIGHTING FUNCTION AT EU,EO	1160	GOTO 182	1720
62	KJ=KMAX(J)	1170	185 LIE=IAF	1730
	IAB=0	1180	77 KMB1=KMIN+IB1	1740
	IAF=0	1190	182 KMAF=KMIN+IAF	1750
	IF1=IAF+1	1200	DO 60 L=IB1,LIE	1760
	IB=IAB+1	1210	IF(IAB.NE.0) GOTO 183	1770
	DO 4 IG=1,KJ	1220	IF(IF1.EQ.IAF) GOTO 183	1780
	KM=KMIN+IG	1230	KML=KMIN+L-1	1790
	IF(ET(KM).GE.EU) GOTO 66	1240	GOTO 184	1800
65	IAB=IG	1250	183 KPL=KMIN+L	1810
	IAF=IG	1260	184 IF(IB1.LT.LIE) GOTO 63	1820
	IF(ET(KM+1).LT.EO) GOTO 4	1270	64 F1=SGP1*F11	1830
	IF(ET(KM+1).GE.EO) GOTO 67	1280	F2=SGP2*F12	1840
66	IF(IG.NE.1) GOTO 72	1290	DIFF=EO-EU	1850
	SGP1=SGIP(J,1)	1300	GO TO 45	1860
	CALL TRA(EU,FI1,NFE,EF,FI)	1310	63 IF(L.NE.IB1) GOTO 49	1870
	IAF=1	1320	48 F1=SGP1*F11	1880
	IF(ET(KM).LT.EO) GOTO 4	1330	DIFF=ET(KMB1+1)-EU	1890
164	SGP2=SGIP(J,1)	1340	GO TO 43	1900
	CALL TRA(EU,FI2,NFE,EF,FI)	1350	49 IF(L.NE.LIE) GOTO 46	1910
	GO TO 80	1360	47 F1=F2	1920
72	IF(IAB.NE.0) GOTO 70	1370	DIFF=EO-ET(KMAF)	1930
73	IF(ET(KM).GE.EO) GOTO 75	1380	F2=SGP2*F12	1940
76	IAF=IG	1390	GO TO 45	1950
	GO TO 4	1400	46 F1=F2	1960
70	IF(ET(KM).LT.EO) GOTO 74	1410	DIFF=ET(KML+1)-ET(KML)	1970
75	IF1=IAF+1	1420	43 IF(IAB.EQ.0) GOTO 181	1980
	KMF=KMIN+IF1	1430	F2=SGIP(J,L+1)*WERT(KML+1)	1990
	KMAF=KMIN+IAF	1440	GOTO 45	2000
	SGP2=SGIP(J,IAF)+(SGIP(J,IF1)-SGIP(J,IAF))*(EO-ET(KMAF))/(ET(KMF)-	1450	181 F2=SGIP(J,L)*WERT(KML+1)	2010
	1ET(KMAF))	1460	45 SUM=SUM+DIFF*(DBLE(F1)+DBLE(F2))/2.	2020
	CALL TRA(EU,FI2,NFE,EF,FI)	1470	60 CONTINUE	2030
	GO TO 80	1480	C SUMMATION LOOP FOR INTEGRAL CLOSED	2040
74	IAF=IG	1490	3 CONTINUE	2050
67	IB=IAB+1	1500	C LCOP OVER EXCITAIGN LEVELS J CLOSED	2060
	KMB=KMIN+IB	1510	88 IF(SU(I).EQ.0.0) GOTO 42	2070
	KMAB=KMIN+IAB	1520	51 QUER(K)=SUM/DBLE(SU(I))	2080
	SGP1=SGIP(J,IAB)+(SGIP(J,IB)-SGIP(J,IAB))*(EU-ET(KMAB))/(ET(KMB)-E	1530	GO TO 52	2090
	1T(KMAB))	1540	42 QUER(K)=DBLE(0.0)	2100
	CALL TRA(EU,FI1,NFE,EF,FI)	1550	52 GES=GES+SNGL(QUER(K))	2110
	GO TO 4	1560	2 CONTINUE	2120
4	CONTINUE	1570	C LCOP OVER ALL INSCATTERING GROUPS K CLOSEC	2130
C	CALCULATION OF THE INTEGRAL BETWEEN EU AND EO BY TRAPEZOIDAL RULE	1580	1002 IF(I-NAB)111,115,131	2140
80	IF(IAB.EQ.0) GOTO 79	1590	115 IF(EMAX .EQ.ET(LMAX)) GOTO 111	2150
78	KMB=KMIN+IB	1600	110 IF(EMAX .EQ.E(NAB+1)) GOTO 111	2160
	IF(ET(KMB).NE.EU) GOTO 82	1610	IF(EMAX .LT.E(NAB+1)) GOTO 133	2170
81	IB1=IB	1620	131 DI=0.	2180
	LIE=IAF	1630	GO TO 134	2190
	GO TO 77	1640	133 DI=1.0-GES/QUOT(NAB)	2200
82	IB1=IAB	1650	IF(DI.LE.0.0) GOTO 111	2210
	LIE=IAF	1660	C CALCULATION OF THE SCATTERING PROBABILITIES IN THE CONTINUUM RANGE	2220
	GO TO 77	1670	134 CALL XKON(NAB,I,DI, EMAX,PROB,A,NE,AG,NFE,EF,FI,XNUE,NX)	2230
79	IB1=1	1680	111 QUD=0.	2240

	CALL DOUB(8HSGI ,NAM3)	220	4 IF(FELD(5).LT.AE(J)) GOTO 121	780
C	READING OF ENERGIES,SIGMA TOT,LEVEL EXCITATION SIGMA FROM KEDAK	230	120 K=K+1	790
	KFEL(1)=NAM	240	IF(K.LE.ISG) GO TO 303	800
	KFEL(2)=NAM2	250	K=1	810
	KFEL(3)=NAM3	260	ISGF=ISGF+1	820
C	INELASTIC SCATTERING CROSS SECTIONS WITH ENERGY SCALE FROM KNDF	270	303 SGIP(J,K)=FELD(6)	830
C	ABOVE THRESHOLD FCR INELASTIC SCATTERING EMIN	280	121 CALL NDFNXT(N,NFEL,FELD,NDAT,NCO)	840
	NFEL(1)=3	290	IF(N.EQ.1) GOTO 20	850
10	CALL NDFLOC (N,NFEL,FELD,NDAT,NCO)	300	IF(ISGF.EQ.0) GO TO 304	860
	L=0	310	ISGP=(ISGF-1)*ISG+K+1	870
11	IF(FELD(5).NE.0.0) GOTO 36	320	RETURN	880
	L=1	330	304 IF(J.NE.1) GOTO 22	890
	SGIT(1)=FELD(5)	340	EMAX =ET(K)	900
	ET(1)=FELD(4)	350	C KMAX(J) - CONSTANT WHICH HAS TO BE ADDED TO EMIN (AND ALL HIGHER	910
	CALL NDFNXT(N,NFEL,FELD,NDAT,NCO)	360	C ENERGIES) TO OBTAIN THE FIRST ENERGY (AND HIGHER ENERGIES RESP.)	920
	IF(N.EQ.1) GOTO 11	370	C AT WHICH CROSS SECTION VALUES FOR THE LEVEL J ARE AVAILABLE ON	930
	GOTO 1000	380	C KEDAK	940
36	L=L+1	390	22 KMAX(J)=K	950
	IF(L.LE.NET) GO TO 300	400	IF(FELD(4).NE.AE(J)) GO TO 21	960
	L=1	410	IF(NAEF.EQ.0) GO TO 305	970
	NETF=NETF+1	420	NAEP=(NAEF-1)*NAE+J+1	980
300	ET(L)=FELD(4)	430	RETURN	990
	SGIT(L)=FELD(5)	440	305 JMAT=J	1000
	IF(FELD(4).GE.E(NE)) GOTO 12	450	DO 311 J=1,JMAT	1010
	CALL NDFNXT (N,NFEL,FELD,NDAT,NCO)	460	IF(KMAX(1).GE.KMAX(J)) GOTO 311	1020
	IF(N.EQ.1) GOTO 36	470	WRITE(NA,312) J	1030
12	LMA=L	480	312 FORMAT (' ***ERROR 5.02 : THE NUMBER OF ENERGIES STORED FOR THE	1040
	IF(NETF.EQ.0) GO TO 301	490	1 EXCITATION CROSS SECTIONS OF LEVEL J IS LARGER THAN THAT FOR THE	1050
	NETP=(NETF-1)*NET+L+1	500	2 FIRST (LOWEST) LEVEL.THIS IS AN ERROR ON KEDAK AND HAS TO BE	1060
	RETURN	510	3 CORRECTED. J = ',I2)	1070
301	EMIN =ET(1)	520	GOTO 1000	1080
C	READING OF INELASTIC EXCITATION CROSS SECTIONS AND DETERMINATION	530	311 CONTINUE	1090
C	OF THE NUMBER OF EXCITATION LEVELS JMAT AND THE MAXIMUM ENERGY	540	NE1=NE-1	1100
C	EMAX, UP TO WHICH LEVEL EXCITATION CROSS SECTIONS ARE AVAILABLE ON	550	180 FORMAT (' ***ERROR 5.03 : ONE OF THE BOUNDARIES OF THE ENERGY	1110
C	KEDAK (CORRESPONDS BOUNDARY BETWEEN CONTINUUM AND DISCRETE REGION)	560	1 RANGE IN WHICH INELASTIC SCATTERING IS POSSIBLE EMAX (- UPPER	1120
C	A KEDAK CONVENTION IS PRESUPPOSED HERE: THE TOTAL INELASTIC	570	2 ENERGY LIMIT OF THE DISCRETE REGION) OR EMIN (- LAST KEDAK-ENERGY	1130
C	SCATTERING CROSS SECTION AND THE INELASTIC EXCITATION CROSS	580	3 AT WHICH THE TOTAL INELASTIC SCATTERING CROSS SECTION IS STILL	1140
C	SECTIONS FOR THE DIFFERENT LEVELS HAVE TO BE STORED AT THE SAME	590	4 EQUAL TO ZERO) IS NOT CONTAINED IN ANY OF THE ENERGY GROUPS OF	1150
C	ENERGY POINTS WHERE DATA FOR THE EXCITATION FUNCTIONS ARE GIVEN	600	5 THE GIVEN GROUP STRUCTURE. THAT MEANS THE GIVEN STRUCTURE OF	1160
C	STARTING AT THE HIGHEST ENERGY BELOW THE LEVEL ENERGY OF THE	610	6 ENERGY GROUPS DOES NOT COVER THE WHOLE ENERGY RANGE.')	1170
C	EXCITATION LEVEL CONCERNED	620	C DETERMINATION OF THE ENERGY OF GROUP IA, IN WHICH EMIN IS LYING, AND	1180
	NFEL(1)=4	630	C OF THE ENERGY GROUP NAB, IN WHICH EMAX IS LYING.	1190
	FELD(4)=0.0	640	DO 84 L=1,NE1	1200
	CALL DOUB (8HSGIZ ,NAM3)	650	IF(EMIN .GE.E(L).AND.EMIN .LT.E(L+1)) GOTO 208	1210
	KFEL(3)=NAM3	660	IF(EMIN .GE.E(L+1)) GOTO 84	1220
	CALL NDFLOC (N,NFEL,FELD,NDAT,NCO)	670	IF(EMIN .GE.E(L)) GOTO 211	1230
	AE(1)=FELD(4)	680	IF(L.NE.1) GOTO 209	1240
	J=0	690	211 WRITE(NA,180)	1250
21	CALL NDFLOC (N,NFEL,FELD,NDAT,NCO)	700	GOTO 1000	1260
	J=J+1	710	208 IA=L	1270
	IF(J.LE.NAE) GO TO 302	720	209 IF(EMAX .GT.E(L).AND.EMAX .LE.E(L+1)) GOTO 210	1280
	J=1	730	IF(EMAX .GT.E(L+1)) GOTO 84	1290
	NAEF=NAEF+1	740	WRITE(NA,180)	1300
302	AE(J)=FELD(4)	750	GOTO 1000	1310
	K=0	760	84 CONTINUE	1320
20	IF(FELD(5).LT.ET(1)) GOTO 121	770	210 NAB=L	1330

	IF(ET(LMA).LT.E(NAB+1)) GOTO 315	1340			1900
	DO 313 L=1,LMA	1350		90 LA=0	1910
	IF(ET(L).GE.E(NAB+1)) GOTO 314	1360		91 SU(II)=0.	1920
313	CONTINUE	1370		SUMM=0.	1930
C	LMAX = (NUMBER OF KEDAK ENERGIES + 1) BETWEEN EMIN AND THE UPPER	1380		LMX1=LMAX-1	1940
C	ENERGY LIMIT E(NAB+1) OF GROUP NAB	1390		DO 40 LL=1,LMX1	1950
314	LMAX=L	1400		IF(LA.NE.0) GOTO 102	1960
	GOTO 316	1410		101 IF(E(II).LT.ET(LL)) GOTO 40	1970
315	LMAX=LMA	1420		118 IF(E(II).GE.ET(LL+1)) GOTO 40	1980
C	INTERPOLATION OF THE WEIGHTING FUNCTION AT THE KEDAK ENERGY POINTS	1430		117 LA=LL	1990
316	KEN1=KEN-1	1440		SG1=SGIT(LL)+(SGIT(LL+1)-SGIT(LL))*(E(II)-ET(LL))/(ET(LL+1)-ET(LL))	2000
	IF(LMAX.LE.IWE) GO TO 306	1450		1)	2010
	IWP=LMAX-IWE+1	1460		EA=E(II)	2020
	IWF=1	1470		CALL TRA(EA,FHI1,NFE,EF,FI)	2030
	RETURN	1480		102 IF(E(II+1).GT.ET(LL)) GOTO 122	2040
306	IF(KEN.EQ.0) GOTO 156	1490		GC TO 40	2050
15	DO 32 LF=1,LMAX	1500		C THE END GROUP IS THE CASE IN QUESTION	2060
	WERT(LF)=0.	1510		122 IF(II.NE.NAB) GOTO 107	2070
	DO 16 IF=1,KEN1	1520		123 NE1=NE-1	2080
	IF(ET(LF).LT.EF(IF)) GOTO 16	1530		IF(NAB.NE.NE1) GOTO 107	2090
	IF(ET(LF).GE.EF(IF+1)) GOTO 1116	1540		124 IF(E(NAB+1).LT.ET(LMAX)) GOTO 107	2100
38	WERT(LF)=FI(IF)+(ET(LF)-EF(IF))*(FI(IF+1)-FI(IF))/(EF(IF+1)-EF(IF))	1550		126 EE=ET(LMAX)	2110
	1)	1560		LE=LMAX	2120
	GC TO 32	1570		SG2=SGIT(LMAX)	2130
1116	IF(ET(LF).NE.EF(KEN)) GOTO 16	1580		FHI2=WERT(LMAX)	2140
	WERT(LF)=FI(KEN)	1590		GC TO 100	2150
	GOTO 32	1600		107 IF(E(II+1).GT.ET(LL+1)) GO TO 40	2160
16	CONTINUE	1610		127 LE=LL+1	2170
32	CONTINUE	1620		SG2=SGIT(LL)+(SGIT(LL+1)-SGIT(LL))*(E(II+1)-ET(LL))/(ET(LL+1)-ET(LL))	2180
	GC TO 39	1630		1)	2190
156	DO 18 KF=1,LMAX	1640		EE=E(II+1)	2200
	EK=ET(KF)	1650		CALL TRA(EE,FHI2,NFE,EF,FI)	2210
	WERT(KF)=PHI(EK)	1660		GC TO 100	2220
18	CONTINUE	1670		40 CONTINUE	2230
C	INTERPOLATION OF SIGMA TOTAL AND THE WEIGHTING FUNCTION AT THE	1680		GC TO 41	2240
C	ENERGY GROUP BOUNDARIES	1690		C CALCULATION OF THE INTEGRAL OVER THE WEIGHTING FUNCTION AND THE	2250
39	DO 50 II=IA,NAB	1700		C INTEGRAL OVER SIGMA TOTAL*WEIGHTING FUNCTION BY TRAPEZOIDAL RULE	2260
	NE1=NE-1	1710		100 LE1=LE-1	2270
C	CRITICAL CASES	1720		DO 30 IL=LA,LE1	2280
	IF(E(IA+1).GT.ET(1)) GOTO 7	1730		IF(LA.NE.LE1) GOTO 104	2290
8	IF(II.EQ.1A) GOTO 50	1740		3 IF(E(II).GE.ET(1)) GOTO 103	2300
9	IA1=IA+1	1750		44 FU1=SGIT(LA)*WERT(LA)	2310
	IF(II-IA1) 90,108,90	1760		FU2=SG2*FHI2	2320
7	IF(II-IA) 90,108,90	1770		DIF1=ET(LA)-EA	2330
108	IF(E(II).GE.ET(1)) GOTO 112	1780		DIFF=EE-ET(LA)	2340
109	LA=1	1790		SUP=DIF1*(WERT(LA)+FHI1)/2.+DIFF*(FHI2+WERT(LA))/2.	2350
	SG1=0.	1800		GC TO 28	2360
	EA=E(II)	1810		103 FU1=SG1*FHI1	2370
	CALL TRA(EA,FHI1,NFE,EF,FI)	1820		FU2=SG2*FHI2	2380
	GC TO 91	1830		DIFF=EE-EA	2390
112	IF(E(II).EQ.ET(1)) GOTO 116	1840		SUP=DIFF*(FHI2+FHI1)/2.	2400
	GC TO 1000	1850		GC TO 28	2410
116	LA=1	1860		104 IF(IL.NE.LA) GOTO 19	2420
	SG1=SGIT(1)	1870		5 IF(E(II).GE.ET(1)) GOTO 34	2430
	FHI1=WERT(1)	1880		6 FU1=SGIT(LA)*WERT(LA)	2440
	EA=E(II)	1890		DIF1=ET(LA)-EA	2450

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DIFF=ET(LA+1)-ET(LA)
SUP=DIFF*(WERT(LA)+FHI1)/2.+DIFF*(WERT(LA+1)+WERT(LA))/2.
GOTO 35
34 FU1=SG1*FHI1
DIFF=ET(LA+1)-EA
SUP=DIFF*(WERT(LA+1)+FHI1)/2.
GO TO 35
19 IF(IL.NE.LE1) GOTO 26
27 FU1=FU2
DIFF=EE-ET(LE1)
SUP=DIFF*(FHI2+WERT(LE1))/2.
FU2=SG2*FHI2
GO TO 28
26 FU1=FU2
DIFF=ET(IL+1)-ET(IL)
SUP=DIFF*(WERT(IL+1)+WERT(IL))/2.
35 FU2=SGIT(IL+1)*WERT(IL+1)
28 SUMM=SUMM+DIFF*(FU1+FU2)/2.
SU(II)=SU(II)+SUP
30 CONTINUE
41 IF(SU(II).EQ.0.) GOTO 31
29 QUOT(II)=SUMM/SU(II)
GO TO 50
31 QUOT(II)=0.
50 CONTINUE
RETURN
1000 STOP
END

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SUBROUTINE XKON(NZG,NUGR,DI, EMAX,W,A,NE,ENG,NF,E,F,XNUE,NX)
C XKON CALCULATES TRANSITION PROBABILITIES FOR INELASTIC SCATTERING
C IN THE 'CONTINUUM REGION'.
REAL*8 ELO,EUD,EXEM,EINS,SU1,SU2,EMD,WD,EXEL,EXD,
1 NAM, EMAX,W,SW
DIMENSION W(NX), ENG(NE),F(NF),E(NF)
COMMON NAM,ISTRUK,ISPA,NA,LIZ
C XNUE IS THE ADJUSTABLE PARAMETER FOR THE CALCULATION OF THE
C NUCLEAR TEMPERATURE AND IS AN INPUT QUANTITY.
CALL FSPIE
XN=XNUE*A
KE=0
KEN=NF
IF(NF.EQ.1) KEN=0
401 IF(DI)2,2,1
C DI>0. MEANS:IN ENERGY GROUP NUGR BOTH RESOLVED AND UNRESOLVED
C LEVELS HAVE TO BE CONSIDERED.EMAX(MM) IS THE HIGHEST ENERGY,
C FOR WHICH THE W ARE CALCULATED FROM RESOLVED ENERGY LEVELS.
1 D=DI
EANF=SNGL(EMAX)
202 EEND=ENG(NUGR+1)
C INDIRECT ASSUMPTION FOR DEFINITION OF EEND :
C HIGHEST ENERGY OF THE GROUP STRUCTURE IS LESS OR EQUAL THAN

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2460 C THE UPPER ENERGY LIMIT ON KEDAK 240
2470 GOTO3 250
2480 2 D=1. 260
2490 EANF=ENG(NUGR) 270
2500 GO TO 202 280
2510 3 IF(KEN)4,4,5 290
2520 C FOLLOING IS THE CALCULATION OF THE AVERAGED NUCLEAR TEMPERATURE. 300
2530 C FOR KEN>0(STATEMENT NUMBER 5) THE AVERAGED NUCLEAR TEMPERATURE 310
2540 C IS CALCULATED USING AN ENERGY POINTWISE FLUX, SUBMITTED BY INPUT. 320
2550 5 DO 198 M=1,KEN 330
2560 IF(E(M)-EANF)198,500,197 340
2570 500 MANF=M 350
2580 GO TO 195 360
2590 197 MANF=M 370
2600 GO TO 194 380
2610 198 CONTINUE 390
2620 195 FANF=F(MANF) 400
2630 GO TO 200 410
2640 194 FANF=F(MANF-1)+((F(MANF)-F(MANF-1))*(EANF-E(MANF-1)))/(E(MANF)-
2650 1E(MANF-1)) 420
2660 200 MAN=MANF-1 430
2670 DO 196 M=MAN ,KEN 440
2680 IF(E(M)-EEND) 196,191,192 450
2690 191 MEND=M 460
2700 GC TO 190 470
2710 192 MEND=M 480
2720 GO TO 189 490
2730 196 CONTINUE 500
190 FEND=F(MEND) 510
GO TO 201 520
189 FEND=F(MEND-1)+ (F(MEND)-F(MEND-1))*(EEND-E(MEND-1))/(E(MEND)-
1E(MEND-1)) 530
201 JA=MANF 540
JE=MEND-2 550
IF(JE.LT.JA) GO TO 250 560
S0=SQRT(E(MANF)/XN)*F(MANF) 570
SINT=(SQRT(EANF/XN)*FANF+S0)/2.*(E(MANF)-EANF)+(SQRT(E(MEND-1)/XN)
1*F(MEND-1)+SQRT(EEND/XN)*FEND)/2.*(EEND-E(MEND-1)) 580
XINT=(FANF+F(MANF))/2.*(E(MANF)-EANF)+(F(MEND-1)+FEND)/2.*
1(EEND-E(MEND-1)) 590
DO 196 J=JA,JE 600
DELT=E(J+1)-E(J) 610
S1=S0 620
S0=SQRT(E(J+1)/XN)*F(J+1) 630
SINT=SINT+(S1+S0)/2.*DELT 640
186 XINT=XINT+(F(J)+F(J+1))/2.*DELT 650
THE =SINT/XINT*0.001 660
251 THE2=THE**2 670
GO TO 6 680
250 THE =SQRT((EEND+EANF)/(2.*XN))*0.001 690
GO TO 251 700
C IF KEN=0(STATEMENT NUMBER4),THE AVERAGE NUCLEAR TEMPERATURE 710
C IS WEIGHTED BY A FLUX,WHICH IS CALCULATED FROM A FUNCTION PHI(E). 720
4 DE=EEND-EANF 730
DED=DE/14. 740
EO=EANF+DED 750

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ET 13

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X1=PHI(EANF)
X0=PHI(E0)
S1=SQRT(EANF/XN)*X1
S0=SQRT(E0/XN)*X0
SINT=(S1+S0)/2.*DED
XINT=(X1+X0)/2.*DED
DC 99 KF=1,13
EO=E0+DED
X1=X0
S1=S0
X0=PHI(E0)
S0=SQRT(E0/XN)*X0
SINT=SINT+(S1+S0)/2.*DED
99 XINT=XINT+(X1+X0)/2.*DED
THE =SINT/XINT*0.001
THE2=THE**2
C FOLLOWING IS THE CALCULATION OF THE INELASTIC SCATTERING
C TRANSITION PROBABILITIES.
6 SW=0.
EU=ENG(1)/THE*0.000001
NU=NUGR-1
DO 7 K=1,NU
EL=EU
EU=ENG(K+1)/THE*0.000001
EM=EU-EL
ELD=DBLE(EL)
EUD=DBLE(EU)
EMD=DBLE(EM)
EINS=DBLE(1.0)
IF(EM-0.01)600,600,601
IF (E(H)/THETA-E(H+1)/THETA) LESS OR EQUAL 0.01,(STATEMENT
C NUMBER 600),EXD IS USED TO CALCULATE THE TRANSITION PROBABILITIES,
C OTHERWISE(STATEMENT NUMBER 601) THE DOUBLE PRECISION EXPONENTIAL
C FUNCTION IS SUFFICIENT.
600 WD=-ELD*EXD(-EMD)-EMD*EXD(-EMD)-EXD(-EMD)-EMD
EXEL=DEXP(-ELD)
WD=EXEL*WD
GO TO 602
601 EXEM=DEXP(-EMD)
SU1=EINS+ELD
SU2=EINS+EUD
SU2=SU2*EXEM
WD=SU1-SU2
602 EXEL=DEXP(-ELD)
WD=EXEL*WD
W(K)=WD*DBLE(THE2)
IF(W(K).LT.0.0) KE=K
7 CONTINUE
IF(KE.GT.0) GOTQ15
GO TO 16
15 DO 9 NG=1,KE
9 W(NG)=0.0
16 EL=EU
EU=EEND/THE*0.000001
EM=EU-EL
ELD=DBLE(EL)

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EUD=DBLE(EU)
EMD=DBLE(EM)
EINS=DBLE(1.0)
EXEM=DEXP(-EMD)
EXEL=DEXP(-ELD)
SU1=EINS+ELD
SU2=EINS+EUD
SU2=EXEM*SU2
WD=SU1-SU2
WD=EXEL*WD
W(NUGR)=WD*DBLE(THE2)
DO 501 K=1,NUGR
501 SW=SW+W(K)
DO 8 K=1,NUGR
8 W(K)=W(K)/SW*DBLE(D)
RETURN
END

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C DOUBLE PRECISION FUNCTION EXD(X)
EXD CALCULATES THE SERIES EXP(X)-1.
REAL*8 X,XN,DNFAK,XN1
N=1
1 XN=X**N/DNFAK(N)
IF(DABS(XN).LE.1.D-50) GO TO 2
N=N+1
GO TO 1
2 IF(N.EQ.1) GO TO 4
N1=N-1
EXD =XN
DC 3 K=1,N1
XN1=XN*DFLOAT(N-K+1)/X
EXD = EXD +XN1
3 XN=XN1
GO TO 5
4 EXD=XN
5 RETURN
END

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C DOUBLE PRECISION FUNCTION DNFAK(NN)
DNFAK CALCULATES N*(N-1)*(N-2)*.....*(N-(N-1)) FOR THE
C SERIES DEVELOPED IN EXD.
REAL*8 N
N=DFLOAT(NN)
DNFAK=N
1 IF(N.LE.(1.D0+1.D-5)) RETURN
N=N-1.D0
DNFAK=DNFAK*N
GO TO 1
END

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SUBROUTINE DOUB(IK,K)		4	EG(NX),R(NX),RSP(NX),	240
REAL*8 IK,K	10	5	ESP(NSPEK),SPEK(LSPEK),	250
K=IK	20	6	NST(6),LEG(6),MAZ(2),	260
RETURN	30	7	SGNC(ICOS,ISM),FEKO(ISM,NFCU),ELSIG(6,NECU,NX),	270
END	40	8	ELTOT(2,NX),	280
	50	9	MAT(KT),DAT(KT),	290
		A	ET(NTTT),ST(NTTT)	300
			COMMON STOFF,ISTRUK,ISPA,NOUT,KPR,IM,IL,KL	310
		C		320
SUBROUTINE TRA(EFI,FIS,NF,EF,FI)	10	C	DATA NST/81,81,161,161,321,321/	330
DIMENSION EF(NF),FI(NF)	20	C	NST(L)= ANZAHL DER AEQUIDISTANTEN STREUKOSINUSSTUEZPUNKTE DER	340
KEN=NF	30	C	SGNC IM L-SYSTEM FUER DAS (L-1)-TE MOMENT	350
IF(NF.EQ.1) KEN=0	40	C	DATA LEG/2,1,3,4,5,6/	360
KUA=0	50	C	LEG = REIHENFOLGE BEI DER MOMENTENBERECHNUNG	370
IF(KEN.NE.0) GO TO 2	60	C	WRITE (NOUT ,9000)	380
1 FIS=PHI(EFI)	70	9000	FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 6')	390
GO TO 3	80	C	WRITE (NOUT ,9001)	400
2 DO 105 KI=1,KEN	90	9001	FORMAT(' PROGRAMM ZUR BERECHNUNG ELASTISCHER STREUMATRIZEN'/)	410
IF(EF(KI).GE.EFI) GO TO 114	100	C	CALL FSPIE	420
113 KUA=KI	110	C		440
GO TO 105	120	C*****SETZEN INTERNER KONSTANTEN*****		450
114 KAB=KUA+1	130	C	ISM = ANZAHL DER ENERGIESTUEZPUNKTE PRO MAKROGRUPPE IM GROBNETZ	460
GO TO 106	140	C	ISD = MAXIMALE ANZAHL DER ENERGIESTUEZPUNKTE PRO MAKROGRUPPE	470
105 CONTINUE	150	C	IM FEINNETZ	480
106 FIS=FI(KUA)+(FI(KAB)-FI(KUA))*(EFI-EF(KUA))/(EF(KAB)-EF(KUA))	160	C	NK = ANZAHL VON GRUNDPUNKTEN IM UEBERSTREUBEREICH	490
3 RETURN	170	C	NR = MAXIMALZAHL VON GRUNDPUNKTEN IM NICHT-UEBERSTREUBEREICH	500
END	180	C	NS = MAXIMALE FEINGRUPPENZAHL PRO GROBGRUPPE	510
		C	PM = 1	520
		C	PM = ANTEIL VON MUEL BEI DER KORREKTUR DES 1. MOMENTES	530
SUBROUTINE FLUMMI(A,B,EA,EB,E,EN,SGN,ECO,SCO,V,W,F,AR,FU,ABN,INT,	10	C	NECU = MAXIMALF STREUBREITE	540
1 GR,FEKOE,FG,R,RSP,ESP,SPEK,SGNC,FEKO,ELSIG,	20	C	NMAX = MAXIMALER LEGENDRE-APPROXIMATIONSGRAD	550
2 ELTOT,ET,ST,MAT,DAT,NLA,NLE,ISEL,NMAX,NX,NE27,	30	C	NOUT = UNIT-NUMMER DER AUSGABE-DATEI	560
3 NSPEK,LSPEK,MAZ,NTK,KT,NTTT,NTTP,ICOS,ICOSP,	40	C	ICOS = ANZAHL DER AEQUIDISTANTEN STREUKOSINUSSTUEZPUNKTE DER VON	570
4 NECU,NECUP,ISM,ISMP,ISD,ISDP,ISCO,ISCP,ISEC,	50	C	KNDF GELESENEN DIFFERENTIELLEN STREUQUERSCHNITTE (SGNC)	580
5 ISECP,KIM,NS,NK,NR)	60	C	NF = 10	590
C	70	C	NF = SPEICHEREINHFIT FUER ENERGIEDIFF. LEGENDRE-KOEFF.	600
C	80	C		610
C*****	90	C*****INITIALISIERUNG INTERNER GROESSEN*****		620
C	100	C	KIM = 0	630
C	110	C	NEN = 0	640
C	120	C	ITA = 0	650
C	130	C	ISO=1	660
C	140	C	NSYS = 1	670
C	150	C	NSTIS = NST(6)	680
C	160	C	NUEB = NECU	690
C	170	C	CD = 2./(ICOS-1)	700
C	180	C	NECUP=0	710
C	190	C	ISDP=0	720
C	200	C	ISMP=0	730
C	210	C	ICOSP=0	740
C	220	C	ISCP=0	750
C	230	C	ISECP=0	760
C		C	NTTP=0	770
REAL*8 STOFF,MAT		C	KSPE=1	780
DIMENSION A(ISM),B(ISM),EA(ISM),EB(ISM),E(ISM),				790
1 EN(ISD),SGN(ISD),ECO(ISD),SCO(ISD),V(ISD),W(ISD),F(ISD),				
2 AR(ICOS),FU(ICOS),				
3 ABN(NE27),INT(NE27),GR(NE27),FEKOE(NE27),				

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C*****LESEN DES ATOMGEWICHTES***** 800
CALL LOOK0(XMAT,MASSE,COM,&1001) 810
ALFA = ((XMAT-1.)/(XMAT+1.))**2 820
C 830
C*****UMSORTIEREN DER GRUPPENGRENZEN***** 840
KSPEK=NSPEK 850
IF(NSPEK.GT.1) GO TO 305 860
KSPE=0 870
KSPEK=0 880
305 DO 302 I=1,NX 890
II=NX-I+1 900
302 ABN(I)=EG(II) 910
IF(ABN(NX)-1.E-3)301,303,304 920
301 WRITE(NOUT,350) 930
350 FORMAT(//' ***ERROR 6.8 : THE SMALLEST GROUP BOUNDARY MUST BE GREA 940
ITER THAN OR EQUAL TO 1.E-3') 950
GO TO 8000 960
304 NEGR=NE27 970
ABN(NEGR)=AMAX1(1.E-3,ALFA*ABN(NX)) 980
GO TO 4 990
303 NEGR=NX 1000
C 1010
C*****BESTIMMUNG DER GRUPPLETHARGIEN***** 1020
4 NGR = NEGR-1 1030
QA = ABN(1)/ABN(2) 1040
DO 5 I = 1,NGR 1050
RSP(I) = 0. 1060
5 R(I)=ALOG(ABN(I)/ABN(I+1)) 1070
C 1080
C*****LESEN DER SGNC-ENERGIEN UND BESTIMMUNG DER STREUBREITE***** 1090
IF(MASSE.EQ.1) GOTO 6 1100
CALL LOOK3(NEN,EB,ISO,NGR,ABN,ISM,ISMP,GR,NEGR,ISEL) 1110
IF(ISMP.GT.0) RETURN 1120
ES = EB(1) 1130
GOTO 7 1140
6 ALFA = 1.E-4 1150
ES = ABN(1) 1160
7 NUE=2 1170
DO 11 I=1,NGR 1180
DO 9 II=I,NEGR 1190
IF(ABN(II).LT.ABN(I)*ALFA) GO TO 10 1200
9 CONTINUE 1210
II=NEGR 1220
10 NUE=MAX0(NUE,II-I) 1230
11 CONTINUE 1240
IF(NUE.LE.NUEB) GOTO 8 1250
NECUP=NUE-NUEB 1260
RETURN 1270
8 NUEB = NUE 1280
ISCOP=NUEB*ISD-ISCO 1290
IF(ISCOP.LT.0) ISCO=0 1300
IF(ISCOP.GT.0) RETURN 1310
ISECP=ISD-ISEC 1320
IF(ISECP.LT.0) ISECP=0 1330
IF(ISECP.GT.0) RETURN 1340
C 1350

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C*****BESTIMMUNG DER MAKROGRUPPENEINTEILUNG***** 1360
12 KIM = MINO(NGR-IM,NUEB-1) 1370
IM = IM+KIM 1380
CALL MAKRO(ALFA,ISO,IL,IM,NGR,ABN,ISD,ISM,ISDP,NS,NK,NECU,NUEB,V, 1390
1 W,ELTOT,INTT,INT,NEGR,ISEL) 1400
IF(ISDP.GT.0) RETURN 1410
C 1420
13 NEK = NEN 1430
IF(NEN.EQ.0) GOTO 16 1440
IF(EB(NEN).GE.ABN(IL)) GOTO 16 1450
WRITE(NOUT,913) STOFF,EB(NEN) 1460
STOP 1470
16 CONTINUE 1480
C 1490
C 1500
C***** M A K R O G R U P P E N R E C H N U N G ***** 1510
C 1520
IR = 1 1530
20 NANF = INT(IR)+1 1540
IF(IR.EQ.1) NANF = NANF-1 1550
NEND = INT(IR+1) 1560
NEN = NEK 1570
ISOT = 1 1580
C 1590
C**** BESTIMMUNG DES EINSTREUINTERVALLES***** 1600
EANF = ABN(NEND+1) 1610
22 EEND=ABN(NANF) 1620
IF(MASSE.GT.1) GOTO 28 1630
NEN = 1 1640
EA(1) = ABN(IL) 1650
28 IF(MASSE.EQ.1) GOTO 62 1660
DO 60 NE = 1,NEN 1670
EA(NE) = EB(NE) 1680
62 CONTINUE 1690
C 1700
C*****MAKROSKOPISCHER TOTALER QUERSCHNITT FUER FEINSTRUKTURWICHTUNG***** 1710
NTT=1 1720
IF(NTK.EQ.1) CALL MIXSGT(KT,MAT,DAT,NTT,NTTP,NTT,ET,ST,EANF,EEND) 1730
IF(NTTP.GT.0) RETURN 1740
C 1750
C*****EINORDNEN DES AUSSTREUINTERVALLES UND LESEN DER SGNC***** 1760
IF(EA(1).GE.EEND) GOTO 88 1770
DO 64 NE = 1,NEN 1780
NI = NE 1790
IF(EA(NE).GE.EANF) GOTO 66 1800
64 CCNTINUE 1810
66 DO 68 NE = 1,NEN 1820
NO = NEN-NE+1 1830
IF(EA(NO).LE.EEND) GOTO 70 1840
68 CCNTINUE 1850
C 1860
70 IF(NI.EQ.1) GOTO 78 1870
NG=MINO(NO+2,NEN) 1880
NI=MAX0(NI-2,1) 1890
NEN = NO-NI+1 1900
DO 74 NE = NI,NO 1910

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2C ON KEDAK')
1002 FORMAT(// ' ***ERROR 6.7 : FOR ',A8,' SCATTERING MATRICES CANNOT B
1E CALCULATED FOR'/16X,'LACK OF ISOT1 (ATOMIC WEIGHT) ON KEDAK')
END

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12 S = SNGL(SN+SP)
RETURN
13 S = A(1)
IF(M.EQ.2) S = S+A(2)
RETURN
END

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SUBROUTINE SUM(M,A,S)
C *****SUMMATION*****
C REAL*8 SP,SN
DIMENSION A(M)
C IF(M.LE.2) GOTO 13
C *****SORTIEREN*****
N = 1
IA = 2
KK = 0
1 K = 0
DO 2 I = IA,M,2
IF(A(I-1).LE.A(I)) GOTO 2
B = A(I)
A(I) = A(I-1)
A(I-1) = B
K = 1
2 CONTINUE
N = -N
IA = IA-N
IF(K.EQ.1) GOTO 1
IF(KK.EQ.1) GOTO 3
KK = 1
GOTO 1
C *****SUMMIEREN*****
3 SP = 0.
SN = 0.
IF(A(1).GE.0.) GOTO 8
IF(A(M).LE.0.) GOTO 10
DO 4 I = 2,M
IF((A(I-1).LE.0.).AND.(A(I).GE.0.)) GOTO 5
4 CONTINUE
5 IA = I
IB = I-1
DO 6 I = IA,M
SP = SP+DBLE(A(I))
DO 7 I = 1,IB
7 SN = SN+DBLE(A(IB+1-I))
GOTO 12
8 DO 9 I = 1,M
9 SP = SP+DBLE(A(I))
GOTO 12
10 DO 11 I = 1,M
11 SN = SN+DBLE(A(M+1-I))

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SUBROUTINE IPOLA(M,A,B,N,X,Y,T)
C *****STETIG DIFFB. INTERPOLATION DURCH POLYNOME 4. GRADES*****
C DIMENSION A(M),B(M),T(M),X(N),Y(N)
C *****TANGENTEN*****
T(1)=(B(2)-B(1))/(A(2)-A(1))
T(M)=(B(M)-B(M-1))/(A(M)-A(M-1))
MN = M-1
DO 2 J = 2,MN
T(J) = 0.
IF((B(J+1).GT.B(J)).AND.(B(J).GT.B(J-1)))
1T(J) = (B(J+1)-B(J-1))/(A(J+1)-A(J-1))
IF((B(J+1).LT.B(J)).AND.(B(J).LT.B(J-1)))
1T(J) = (B(J+1)-B(J-1))/(A(J+1)-A(J-1))
2 CONTINUE
C *****INTERPOLATION*****
Y(1) = B(1)
Y(N) = B(M)
NM = N-1
J = 1
DO 26 I = 2,NM
18 IF(X(I)-A(J)) 20,22,24
20 S = X(I)-UM
Y(I) = ((A4*S+A3)*S+A2)*S+A1)*S+VM
GOTO 26
22 Y(I) = B(J)
GOTO 26
24 J = J+1
U = 0.5*(A(J)-A(J-1))
V = 0.5*(B(J)-B(J-1))
TP = 0.5*(T(J)+T(J-1))
TM = 0.5*(T(J)-T(J-1))
A1 = 0.5*(3*V/U-TP)
A2 = -0.5*TM/U
A3 = 0.5*(TP-V/U)/(U*U)
A4 = 0.5*TM/(U*U*U)
UM = 0.5*(A(J-1)+A(J))
VM = 0.5*(B(J-1)+B(J))
GOTO 18
26 CONTINUE
RETURN
END

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SUBROUTINE PUSUM(KA,A,KE,E,B)
C
C*****STUETZPUNKTMENGEN VEREINIGEN*****
C
DIMENSION A(KA),E(KE),B(KE)
C
DO 2 K = 1,KE
2 B(K) = E(K)
C
J=1
I=1
K=1
3 IF(K.GT.KA) GO TO 10
IF(I.GT.KE) GOTO 14
IF(A(K)-B(I)) 4,5,6
4 E(J) = A(K)
J = J+1
K = K+1
GOTO 3
5 E(J) = A(K)
J = J+1
I = I+1
K = K+1
GOTO 3
6 E(J) = B(I)
J = J+1
I = I+1
GOTO 3
10 IF(I.GT.KE) GO TO 18
DO 12 K=I,KE
E(J)=B(K)
12 J=J+1
GO TO 18
14 DO 16 I=K,KA
E(J)=A(I)
16 J=J+1
18 KE=J-1
RETURN
C
END

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SUBROUTINE MAKRO(ALFA,ISO,IL,IM,NGR,ABN,ISD,ISM,ISDP,NS,NK,NECU,
1 NUEB,V,W,LST,INTT,INT,NEGR,ISEL)
C
C*****BESTIMMUNG EINER DEN SGN-STELLEN UND DER GRUPPENSTUETZPUNKTZAHL AN
C*****GEPASSTEN MAKROGRUPPENEINTEILUNG*****
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REAL*8 STOFF
DIMENSION ABN(NEGR),INT(NEGR),LST(NGR),V(ISD),W(ISD)
COMMON STOFF,ISTRUK,ISPA,NOUT
C
C*****ANZAHL DER SGN IN DER N-TEN GRUPPE*****
DO 10 N = IL,IM
10 LST(N) = 0
EA = ABN(IM+1)
EE = ABN(IL)
11 CALL LOOKI(ISD,K,V,W,EA,EE,0)
DO 12 N = IL,IM
A = ABN(N+1)
B = ABN(N)
DO 12 MM = 1,K
E = V(K-MM+1)
IF(A.LT.E.AND.E.LE.B) LST(N) = LST(N)+1
12 CONTINUE
IF(K.LT.ISD) GOTO 13
EA = V(K)
GOTO 11
13 IF(ISEL.LT.0) WRITE(NOUT,100)
ISCP = 0
NTO = 0
DO 14 N = IL,IM
NTO = NTO+LST(N)
IF(ISEL.LT.0) WRITE(NOUT,102) N,LST(N)
LSS = LST(N)
C
= SGN-STELLEN IN DER N-TEN GRUPPE
LSG = NK+5
C
= GRUNDPUNKTE PRO GRUPPE
LSM = (ISM-NS*(NK+5))/NS
C
= SGNC-ENERGIESTUETZSTELLEN PRO GRUPPE
14 ISDP = MAXO(ISDP,LSG+LSM+LSS-ISD)
IF(ISEL.LT.0) WRITE(NOUT,103) NTO
IF(ISDP.GT.0) RETURN
C
C*****MAKROGRUPPENEINTEILUNG*****
IT = ISO
ID = ISD-ISM
INT(1) = IL-1
IR = 2
3 II = INT(IR-1)+1
IS = 0
DO 4 I = II,IM
IS = IS+LST(I)
IF((IS.GE.ID).OR.(1+I-II.GT.NS)) GOTO 5
IF(I.NE.IT) GOTO 4
IF(I.EQ.IM.OR.I.EQ.II) GOTO 4
IT = 0
GOTO 5
4 CONTINUE
INT(IR) = IM
INT(1) = IL
INTT = IR
IE = IR-1
GOTO 7

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64 IA

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5 INT(IR) = MAXO(II,I-1)
IR = IR+1
GOTO 3
C
C*****SCHREIBEN DER MAKROGRUPPENEINTEILUNG*****
7 IF(ISEL.GT.0) RETURN
WRITE(NOUT,94)
DO 8 I = 1,IE
IP = INT(I)+1
IF(I.EQ.1) IP = IP-1
8 WRITE(NOUT,96) I,IP,INT(I+1)
C
RETURN
C
94 FORMAT(///11X,'M A K R O G R U P P E N'//)
96 FORMAT(1H0,2X,I2,'. MAKROGRUPPE',5X,I3,'. BIS ',I3,'. FEINGRUPPE')
100 FORMAT(//14X,'VERTEILUNG DER SGN'//)
102 FORMAT(1X,I3,'. FEINGRUPPE',12X,I4,' SGN')
103 FORMAT(1H0,6X,' INSGESAMT',12X,I4,' SGN')
C
END

SUBROUTINE MIXSGT(KT,MAT,DAT,NTT,NTP,NT,ET,ST,EA,EE)
C
C*****TOTALER QUERSCHNITT EINER MISCHUNG*****
REAL*8 MAT(KT),A(5),STOFF
DIMENSION DAT(KT),ET(NTT),ST(NTT),NW(4),NAD(2)
COMMON STOFF,ISTRUK,ISPA,NOUT
CATA NW(1),A(2),A(3) /3,'BEST ','SGT '/
C
DO 10 K = 1,KT
D = DAT(K)
A(1) = MAT(K)
NU = 0
CALL NDFLOC(KP,NW,A,NAD,KC)
IF(KP.NE.1) GOTO 12
E = A(4)
S = A(5)
IF(E-EA) 6,4,14
4 CALL ADD(K,NTT,NTP,NT,ET,ST,NU,D,E,S,EV,SV,EW,SW,EA,EE)
IF(NTP.GT.0) RETURN
6 EV = E
SV = S
CALL NDFNXT(KP,NW,A,NAD,KC)
IF(KP.EQ.1) GOTO 7
E = EE
S = SV
GOTO 8
7 E = A(4)
S = A(5)
IF(E-EA) 6,4,8

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8 CALL ADD(K,NTT,NTP,NT,ET,ST,NU,D,E,S,EV,SV,EW,SW,EA,EE)
IF(NTP.GT.0) RETURN
IF(E.LT.EE) GOTO 6
10 CONTINUE
RETURN
C
12 WRITE(NOUT,20) MAT(K)
STOP
14 WRITE(NOUT,22) MAT(K),EA
STOP
C
20 FORMAT(//' ***ERROR 6.5 : FOR ',A8,' IN THE FINE-STRUCTURE MIXTUR
1E NO SGT'/16X,'ARE ON KEDAK')
22 FORMAT(//' ***ERROR 6.6 : FOR ',A8,' IN THE FINE-STRUCTURE MIXTUR
1E NO SGT'/16X,'ARE ON KEDAK FOR ENERGIES BELOW',1PE10.3,' EV')
C
END

SUBROUTINE ADD(K,NTT,NTP,NT,ET,ST,NU,D,E,S,EV,SV,EW,SW,EA,EE)
C
DIMENSION ET(NTT),ST(NTT)
C
C*****ERSTES MATERIAL*****
IF(K.GT.1) GOTO 10
IF(NU.GT.0) GOTO 6
IF(E.GT.EA) GOTO 2
ET(1) = E
ST(1) = S*D
GOTO 4
2 T = (S-SV)/(E-EV)
ET(1) = EA
ST(1) = D*(S+T*(EA-E))
4 NU = 1
NT = 1
IF(E.LT.EE) RETURN
ET(2) = EE
ST(2) = D*(S+T*(EE-E))
NU = 2
NT = 2
RETURN
6 NU = NU+1
NT = NT+1
IF(NT.GT.NTT) GOTO 22
IF(E.LE.EE) GOTO 8
ET(NU) = EE
ST(NU) = ST(NU-1)
RETURN
8 ET(NU) = E
ST(NU) = S*D
RETURN
C
C*****WEITERE MATERIALIEN*****

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10 IF(NU.GT.0) GOTO 12
   EW = ET(1)
   SW = ST(1)
   IF(E.GT.EA) GOTO 12
   ST(1) = ST(1)+D*S
   EV = E
   SV = S
   NU = 1
   RETURN
12 T = (S-SV)/(E-EV)
   NP = NU+1
   DO 14 I = NP,NT
   IF(ET(I).GT.E) GOTO 16
   Z = SV+T*(ET(I)-EV)
   EW = ET(I)
   SW = ST(I)
14 ST(I) = ST(I)+D*Z
   RETURN
16 IF(E.EQ.ET(I-1)) GOTO 20
   Z = SW+(ST(I)-SW)/(ET(I)-EW)*(E-EW)
   IF(NT+1.EQ.NTT) GOTO 22
   DO 18 J = I,NT
   JJ = NT+I-J
   ET(JJ+1) = ET(JJ)
18 ST(JJ+1) = ST(JJ)
   NT = NT+1
   ET(I) = E
   ST(I) = Z+D*S
   NU = I
   EV = E
   SV = S
   RETURN
20 NU = I-1
   EV = E
   SV = S
   RETURN
C
22 NTP = 1000
   RETURN
   END

SUBROUTINE TRAFD(IC,NE,ISM,S,XM,AL,A)
C
C*****TRANSFORMATION DER SGNC INS L-SYSTEM*****
C
DIMENSION A(IC),S(IC,ISM)
C
XMM = XM-1
XMP = XM+1
QM = XM*XM
DO 2 I = 1,IC
V = SQRT(0.5*(1.+AL+(1.-AL)*A(I)))

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350 A(I) = 0.5*(XMP*V-XMM/V)
360 B = SQRT(A(I)*A(I)+QM-1.)
370 F = (A(I)+B)*(A(I)+B)/(XM*B)
380 DO 2 N = 1,NE
390 2 S(I,N) = S(I,N)*F
C
RETURN
END

SUBROUTINE INFORM(ALFA,NEGR,ABN,R,RSP,IL,IM,INTT,INT,NST,NSTIS,
1 IR,ISTT,K,NTK,NTT,NOUT)
C
C*****INFORMATIONEN*****
C
DIMENSION ABN(NEGR),R(NEGR),RSP(NEGR),INT(NEGR),NST(6),NA(6),
1 NB(6)
C
IF(IR.GT.1) GOTO 6
WRITE(NOUT,919)
WRITE(NOUT,935) ALFA
3 DO 5 I = 1,6
NA(I) = I-1
5 NB(I) = NSTIS
WRITE(NOUT,936) (NA(I),I = 1,6),(NB(I),I = 1,6),
1 (NST(I),I = 1,6)
6 IF(IR+1.LT.INTT) GOTO 7
WRITE(NOUT,932)
DO 4 I = IL,IM
4 WRITE(NOUT,933) I,ABN(I),ABN(I+1),R(I),RSP(I)
WRITE(NOUT,934)
7 I1 = INT(IR+1)+1
I2 = INT(IR)+1
IF(IR.EQ.1) I2 = I2-1
WRITE(NOUT,909)
EFIN = AMAX1(ALFA*ABN(I1),ABN(IM+1))
WRITE(NOUT,920) IR,IR,ABN(I1),ABN(I2)
WRITE(NOUT,921) IR,EFIN,ABN(I2)
WRITE(NOUT,922) IR,ISTT
WRITE(NOUT,925) K
IF(NTK.EQ.1) WRITE(NOUT,126) NTT
RETURN
C
909 FORMAT(1H0)
919 FORMAT(1H0//24X,'I N F O R M A T I O N E N'//)
920 FORMAT(1H0,I2,'. AUSSTREUINTERVALL = ',I2,'. MAKROGRUPPE', 5X,E10.
14,' EV', ' - ',E10.4,' EV')
921 FORMAT(1H0,I2,'. EINSTREUINTERVALL',23X,E10.4,' EV', ' - ',E10.4,
1' EV')
922 FORMAT(1H0,'ZAHL DER ENERGIESTUETZPUNKTE IM ',I2,'. AUSSTREUINTERV
1ALL',18X,I4)
925 FORMAT(1H0,'ZAHL DER GELESENEN SGN IM AUSSTREUINTERVALL',28X,I4)
932 FORMAT(///16X,'F E I N G R U P P E N E I N T E I L U N G'///)

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933 FORMAT(1X,I3,'.GRUPPE',1X,1PE10.3,' EV -',1PE10.3,' EV',1X,'DU = 440
1',E9.2,1X,'RSP = ',1PE11.4) 450
934 FORMAT(////18X,'DU = LETHARGIEDIFFERENZ'//18X,'RSP = GRUPPENINTEGR 460
1AL UEBERS SPEKTRUM') 470
935 FORMAT(1H0,'ALFA = ',F8.5) 480
936 FORMAT(1H0,'WINKELSTUETZSTELLEN'//12X,'MOMENT',5X,6(3X,I3)/12X,'IS 490
10TROPIS',3X,6(3X,I3)/12X,'ANISOTROPIE',1X,6(3X,I3)) 500
126 FORMAT(1H0,'ZAHL DER SGT FUER FEINSTRUKTURWICHTUNG IM AUSSTREUINTE 510
IRVALL',11X,I5) 520
END 530

SUBROUTINE PUNK(NGR,NEGR,ABN,NA,NE,NEN,ALFA,NK,NR,ISM,ISMP,ISN,E) 10
C 20
C*****GRUNDPUNKTE INTERVALLWEISE LOGARITHMISCH AEQUIDISTANT***** 30
C 40
REAL*8 D,ED,EX,EP 50
DIMENSION ABN(NEGR),E(ISM) 60
C 70
ISMP = MAX(0,(NE-NA+1)*(NK+NR)+NEN-ISM) 80
IF(ISMP.GT.0) RETURN 90
EX = DBLE(1./(NK-1)) 100
E(1) = ABN(NE+1) 110
J = 2 120
DO 10 NN = NA,NE 130
N = NA+NE-NN 140
ED = DBLE(E(J-1)) 150
AB = ABN(N) 160
A = AMINI(AB,ABN(N+1)/ALFA) 170
4 D = DBLE(A/ABN(N+1)) 180
D = D**EX 190
6 EP = ED*D 200
E(J) = SNGL(EP) 210
IF(E(J).GT.A) GOTO 8 220
ED = EP 230
J = J+1 240
GOTO 6 250
8 E(J-1) = A 260
IF(A.GE.AB) GOTO 10 270
ED = DBLE((AB/A)**(1./NR)) 280
IF(ED.GE.D) D = ED 290
ED = DBLE(A) 300
11 EP = ED*D 310
E(J) = SNGL(EP) 320
IF(E(J).GT.AB) GOTO 9 330
ED = EP 340
J = J+1 350
GOTO 11 360
9 E(J-1) = AB 370
10 CONTINUE 380
ISN = J-1 390
RETURN 400
END 410

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SUBROUTINE LOOKO(XMAT,MASSE,COM,*) 10
C 20
C*****LESEN DES ATOMGEWICHTS ***** 30
C 40
REAL*8 STOFF,F(6),ST 50
INTEGER*2 U/'H '/,V 60
DIMENSION NWORT(4),NADAT(2) 70
COMMON STOFF,ISTRUK,ISPA,NOUT 80
EQUIVALENCE(ST,V) 90
DATA F(2),F(3) /'BEST ','ISOT1'/ 100
C 110
NWORT(1) = 3 120
F(1) = STOFF 130
CALL NDFLOC(KP,NWORT,F,NADAT,KC) 140
IF(KP.NE.1) GO TO 1 150
F(4) = F(4)/1.008665 160
XMAT = F(4) 170
COM = 2/(3*F(4)) 180
MASSE = XMAT 190
IF(XMAT-MASSE.GT.0.5) MASSE = MASSE+1 200
RETURN 210
C 220
1 ST=STOFF 230
IF(U.NE.V) RETURN 1 240
XMAT=1. 250
MASSE=1 260
CCM=0.6666667 270
RETURN 280
END 290

SUBROUTINE LOOK1(ISD,K,E,SGN,EO,EE,NT) 10
C 20
C*****LESEN DER SGN UND MUEL VON KEDAK***** 30
C 40
REAL*8 STOFF,A(5),B(2) 50
DIMENSION E(ISD),SGN(ISD),NWORT(4),NADAT(2) 60
COMMON STOFF,ISTRUK,ISPA,NOUT 70
FG(XA,XB,XC,YA,YC) = YA+(YC-YA)/(XC-XA)*(XB-XA) 80
DATA A(2),B /'BEST ','SGN ','MUEL '/ 90
C 100
A(3) = B(NT+1) 110
1 NWORT(1) = 3 120
A(1) = STOFF 130
CALL NDFLOC(KP,NWORT,A,NADAT,KC) 140
IF(KP.NE.1) STOP 150
K = 1 160
E(1) = A(4) 170

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      SGN(1) = A(5)
      IF(E(1)-EO) 5,4,3
3  WRITE(NOUT,40) STOFF,EO,A(3)
      STOP
4  K = K+1
5  CALL NDFNXT(KP,NWORT,A,NADAT,KC)
      IF(KP.NE.1) GOTO 14
      IF(K.GT.1) GOTO 11
      IF(A(4)-EO) 8,9,10
8  E(1) = A(4)
      SGN(1) = A(5)
      GOTO 5
9  E(1) = A(4)
      SGN(1) = A(5)
      GOTO 4
10 K = 2
11 E(K) = A(4)
      IF(K.EQ.ISD) RETURN
      SGN(K) = A(5)
      IF(E(K)-EE) 4,16,12
12 Y = FG(E(K-1),EE,E(K),SGN(K-1),SGN(K))
      E(K) = EE
      SGN(K) = Y
      GOTO 16
14 E(K) = EE
      SGN(K) = SGN(K-1)
      WRITE(NOUT,42) A(3),STOFF,E(K-1),A(3),E(K-1)
16 IF(E(1).EQ.EO) RETURN
      Y = FG(E(1),EO,E(2),SGN(1),SGN(2))
      E(1) = EO
      SGN(1) = Y
      RETURN
C
40 FORMAT(//' ***ERROR 6.2 : FOR ',A8,' SCATTERING MATRICES CANNOT B
1E CALCULATED'/16X'FOR ENERGIES BELOW',1PE10.3,' EV FOR LACK OF ',
2A8,' ON KEDAK')
42 FORMAT(//' ***WARNING 6.1 : THE ',A8,' FOR ',A8,' AT ENERGIES AB
LOVE',1PE10.3,' EV'/18X,'ARE SET EQUAL TO THE ',A8,' AT',E10.3,' E
2V'//)
C
      END
      SUBROUTINE LOOK2(NEN,EA,ICOS,ICOSP,AR,ISM,SGNC)
C
C*****LESEN DER SGNC VOM KERNDATENBAND*****
C
      REAL*8 STOFF,F(6)
      DIMENSION EA(ISM),NWORT(4),NADAT(2),SGNC(ICOS,ISM),AR(ICOS)
      COMMON STOFF,ISTRUK,ISPA,NOUT
      DATA F(2),F(3) /'BEST ','SGNC '/
C
      ICOSP = 0

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2  NWORT(1) = 4
      F(1) = STOFF
      DO 4 NE = 1,NEN
      F(4) = EA(NE)
      CALL NDFLOC(KP,NWORT,F,NADAT,KC)
      IF(KP.NE.1) STOP
      IP = 1
3  AR(IP) = F(5)
      SGNC(IP,NE) = F(6)
      IP = IP+1
      CALL NDFNXT(KP,NWORT,F,NADAT,KC)
      IF(KP.EQ.1) GOTO 3
      IF(IP-1-ICOS) 8,4,10
4  CONTINUE
      RETURN
8  WRITE(NOUT,44) STOFF,ICOS,EA(NE)
      STCP
10 ICOSP = IP-1-ICOS
      RETURN
C
44 FORMAT(//' ***ERROR 6.3 : FOR ',A8,' THE SGNC COSINE MESH HAS LES
1S'/16X,'THAN ',I3,' MESHPOINTS AT',1PE10.3,' EV')
C
      END
      SUBROUTINE LOOK3(NE,EA,ISO,NGR,ABN,ISM,ISMP,LST,NEGR,ISEL)
C
C*****LESEN DER SGNC-ENERGIEN VOM KERNDATENBAND*****
C
      REAL*8 STOFF,F(6)
      DIMENSION EA(ISM),ABN(NEGR),LST(NEGR),NWORT(4),NADAT(2)
      COMMON STOFF,ISTRUK,ISPA,NOUT
      DATA F(2),F(3) /'BEST ','SGNC '/
C
2  NWORT(1) = 4
      F(1) = STOFF
      F(4) = 0.
      NE = 1
4  CALL NDFLOC(KP,NWORT,F,NADAT,KC)
      IF(KP.EQ.1) GO TO 5
      CALL NDFLOC(KP,NWORT,F,NADAT,KC)
      IF(KP.EQ.1) GO TO 5
      IF(NE-1)13,13,6
5  EA(NE)=F(4)
      F(4) = F(4)+F(4)*(1.E-06)
      NE = NE+1
      IF(NE.LE.ISM) GOTO 4
      ISMP=50
      RETURN
6  WRITE(NOUT,62)
      IF(EA(1).GE.ABN(1)) GO TO 12
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C*****ANZAHL DER SGNC IN DER N-TEN GRUPPE*****
NE=NE-1
NTC = 0
DO 10 N = 1,NGR
LST(N) = 0
DO 8 M = 1,NE
IF((ABN(N+1).LE.EA(M)).AND.(EA(M).LT.ABN(N))) LST(N) = LST(N)+1
8 CONTINUE
10 NTO = NTO+LST(N)
IF(ISEL.GT.0) GO TO 11
WRITE(NOUT,52)
DO 7 N=1,NGR
7 WRITE(NOUT,54) N,LST(N)
WRITE(NOUT,56) NTO
C
C*****ANISOTROPIE-SCHWELLE*****
11 DO 14 N = 1,NGR
NN = NGR-N+1
IF(LST(NN).NE.0) GOTO 16
14 CONTINUE
16 ISO = NN+1
RETURN
C
12 IF(ISEL.GT.0) GO TO 15
WRITE(NOUT,52)
WRITE(NOUT,58)
15 EA(1) = ABN(1)
ISO=1
RETURN
C
13 WRITE(NOUT,60) STOFF
STOP
C
52 FORMAT(1H0//14X,'VERTEILUNG DER SGNC'//)
54 FORMAT(1X,I3,' FEINGRUPPE',12X,I4,' SGNC')
56 FORMAT(1H0,6X,' INSGESAMT',12X,I4,' SGNC')
58 FORMAT(5X,'ENERGIEBEREICH LIEGT IM BEREICH ISOTROPER STREUUNG')
60 FORMAT(//' ***ERROR 6.4 : FOR ',A8,' SCATTERING MATRICES CANNOT B
IE CALCULATED'/16X'FOR LACK OF SGNC ON KEDAK')
62 FORMAT(//' ***MESSAGE 6.1 : WARNING NDF. 2 MAY BE IGNORED')
C
END
SUBROUTINE LECAL(MASSE,XMAT,ALFA,ISOT,LEG,NLE,NANF,NEND,NUEB,IL,
1 IM,NEGR,ABN,ICOS,AR,ISN,E,SGNC,NST,NSTIS,ISM,ISD,
2 NECU,ITA,NF,GR,FEKOE,FU,EW,A,H,V,W,F,FEKO)
C
C*****BERECHNUNG DER ENERGIEDIFFERENTIELLEN LEGENDRE-KOEFFIZIENTEN*****
DIMENSION NST(6),LEG(6),AR(ICOS),FU(ICOS),ABN(NEGR),GR(NEGR),
1 FEKOE(NEGR),A(ISD),H(ISD),E(ISM),FW(ISM),V(ISD),W(ISD),
2 F(ISD),SGNC(ICOS,ISM),FEKO(ISM,NFCU)

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C
EP = 0.8
XP = XMAT+1
XM = XMAT-1
IF(MASSE.EQ.1) XM = 0.
NM = NUEB-1
IA = NANF
IC = MINO(NEGR,NEND+NUEB)
IB = IC-1
IF(ITA.EQ.0.AND.ISOT.EQ.0)
1CALL LEGINT(NLE,NSTIS,XMAT,ITA)
REWIND NF
C
C*****ENERGIEDIFF. LEGENDRE-KOEFF. (L-1)-TER ORDNUNG*****
N = 0
L = 1
GOTO 2
1 N=N+1
L=LEG(N)
2 IF(ISOT.EQ.0) GOTO 3
NSB = NSTIS
NSA = NST(L)
CALL LEGPOL(NSA,H,NLE)
C
3 DO 28 J = 1,ISN
C
C*****GRENZEN FUER WINKEL INTEGRATION*****
GRM = 0.
EL = AMAX1(ABN(IM+1),ALFA*E(J))
DO 8 I = IA,IC
IF(EL.LT.ABN(I)) GOTO 6
GR(I) = -1.
GOTO 8
6 IF(E(J).GT.ABN(I)) GOTO 7
GR(I) = 1.
GOTO 8
7 G = SQRT(ABN(I)/E(J))
GR(I) = 0.5*(XP*G-XM/G)
GRA = ABS(GR(I))
IF((GRA.LT.1.).AND.(GRM.LT.GRA)) GRM = GRA
8 CCNTINUE
IF(ISOT.EQ.0) GOTO 11
C
C*****MOMENTE BEI ANISOTROPIE IM SP-SYSTEM*****
IF(GRM.LT.EP) NSB = NST(L)
IF(NSA.EQ.NSB) GOTO 9
NSA = NSB
CALL LEGPOL(NSA,H,NLE)
9 DO 10 IP = 1,ICOS
10 FU(IP) = SGNC(IP,J)
CALL IPOLA(ICOS,AR,FU,NSA,H,A,F)
CALL LEGANS(L,NSA,A,IA,IB,NEGR,GR,FEKOE,V,W)
NSB = NSTIS
GOTO 12
C
C*****MOMENTE BEI ISOTROPIE IM SP-SYSTEM*****

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11 CALL LEGIST(L,NSTIS,IA,IB,NEGR,GR,FEKOE) 660
C 670
C*****VORBEREITUNG ZUM ABSPEICHERN***** 680
12 II = 1 690
DO 22 I = IA,IB 700
IF(ABN(I+1).GT.E(J)) GOTO 22 710
FEKO(J,II) = 6.283185*FEKOE(I) 720
II = II+1 730
IF(II.GE.NUEB) GOTO 27 740
22 CONTINUE 750
DO 26 I = II,NM 760
26 FEKO(J,I) = 0. 770
27 FEKO(J,NUEB) = 6.283185*FEKOE(NEGR) 780
C 790
28 CONTINUE 800
C 810
C*****NORMIERUNG FUER ZWISCHENPUNKTE***** 820
IF(N.GT.0) GOTO 30 830
DO 29 J = 1,ISN 840
29 EW(J) = FEKO(J,NUEB) 850
GOTO 1 860
30 DO 31 J = 1,ISN 870
DO 31 I = 1,NUEB 880
31 FEKO(J,I) = FEKO(J,I)/EW(J) 890
C 900
C*****SPEICHERN DER ENERGIEDIFF. LEGENDRE-KOEFF. (L-1)-TER ORDNUNG***** 910
DO 32 I = 1,NUEB 920
32 WRITE(NF) (FEKO(J,I),J = 1,ISN) 930
IF(N.LT.NLE+1) GOTO 1 940
C 950
REWIND NF 960
C 970
RETURN 980
END 990

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SUBROUTINE LEGPOL(NST,A,N,NSTIS) 10
C 20
C*****BERECHNUNG VON LEGENDRE-POLYNOMEN***** 30
C 40
REAL*8 X(321),F(6,321),D 50
DIMENSION A(NSTIS) 60
COMMON /INTEG/ X,F,D 70
C 80
NSM = NST-1 90
NA = -100000 100
ND = 320/NSM*625 110
DO 6 I = 1,NST 120
X(I) = DFLOAT(NA)*1.D-05 130
A(I) = SINGL(X(I)) 140
6 NA = NA+ND 150
D = ND*1.E-05 160
C 170

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IF(N.GT.0) GOTO 2 180
DO 1 I = 1,NST 190
1 F(1,I) = 1. 200
RETURN 210
C 220
2 IF(N.GT.1) GOTO 4 230
DO 3 I = 1,NST 240
F(1,I) = 1. 250
3 F(2,I) = X(I) 260
RETURN 270
C 280
4 NM = N-1 290
DO 5 I = 1,NST 300
F(1,I) = 1. 310
F(2,I) = X(I) 320
DO 5 J = 1,NM 330
5 F(J+2,I) = ((2*J+1)*X(I)*F(J+1,I)-J*F(J,I))/(J+1) 340
C 350
RETURN 360
END 370

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SUBROUTINE LEGINT(N,NAK,XMAT,ITA) 10
C 20
C*****BERECHNUNG VON LEGENDRE-INTEGRALEN***** 30
C 40
REAL*8 H(321),GRAL(6,321),F(6),A,B,C,D,DE,G,P,F0,F1,F2,F3,F4,F5, 50
1 X,Y 60
COMMON /INTEG/ H,GRAL,D 70
C 80
ITA = 1 90
NST = 321 100
C 110
NA = -100000 120
DO 1 K = 1,NST 130
H(K) = DFLOAT(NA)*1.D-05 140
1 NA = NA+625 150
D = 0.00625 160
C 170
IF(XMAT.GT.1.5) GOTO 4 180
C*****WASSERSTOFF***** 190
DO 2 K = 1,160 200
DO 2 L = 1,6 210
2 GRAL(L,K) = 0. 220
DC 3 K = 161,321 230
A = H(K)*H(K) 240
B = A*H(K) 250
GRAL(1,K) = 2*A 260
GRAL(2,K) = 4*B/3 270
GRAL(3,K) = 0.5*A*(3*A-2) 280
GRAL(4,K) = 2*B*(A-1) 290
GRAL(5,K) = A*(35*(A-1)*(A-2./7)-1)/12 300
GRAL(6,K) = 4.5*B*(A-1)*(A-5./9) 310

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3 CCNTINUE
  GOTO 12
C
C*****NICHT-WASSERSTOFFE*****
4 A = DBLE(XMAT)
  DO 5 L = 1,6
5 GRAL(L,1) = 0
  IF(XMAT.GT.30.) GOTO 8
  C = 0.5*DSQRT((A+1)*(A-1))
  F(1) = C
  DO 6 L = 2,5
6 F(L) = F(L-1)*C
  P = A-1./A
  Y = DSQRT((A-1)/(A+1))
C
  DO 7 K = 2,NST
  X = (H(K)+DSQRT((H(K)+1)*(H(K)-1)+A*A))/(2*C)
  DE = X-Y
  F0 = 0.5*G(2,X,Y)
  F1 = F(1)*(G(3,X,Y)/3-1)
  F2 = F(2)*(0.25*G(4,X,Y)-G(2,X,Y)+DLOG(X/Y)/DE)
  F3 = F(3)*(0.2*G(5,X,Y)-G(3,X,Y)+3-1./(X*Y))
  F4 = F(4)*(G(6,X,Y)/6-G(4,X,Y)+3*G(2,X,Y)-4*DLOG(X/Y)/DE
  +0.5/(X*Y)*G(2,1./X,1./Y))
1 F5 = F(5)*(G(7,X,Y)/7-G(5,X,Y)+10*G(3,X,Y)/3-10+5./(X*Y)
1 -G(3,1./X,1./Y)/(3*X*Y))
  GRAL(1,K) = P*F0*DE
  GRAL(2,K) = P*F1*DE
  GRAL(3,K) = P*0.5*(3*F2-F0)*DE
  GRAL(4,K) = P*0.5*(5*F3-3*F1)*DE
  GRAL(5,K) = P*0.125*(35*F4-30*F2+3*F0)*DE
  GRAL(6,K) = P*0.125*(63*F5-70*F3+15*F1)*DE
7 CCNTINUE
  GOTO 11
C
8 B = 1./(A*A)
  C = B*B
  DO 10 K = 2,NST
  X = H(K)
  Y = -1.
  DO 9 L = 1,6
9 F(L) = -0.125*(5*G(L+4,X,Y)/(L+4)-6*G(L+2,X,Y)/(L+2)+G(L,X,Y)/L)*C
1 +0.5*(3*G(L+2,X,Y)/(L+2)-G(L,X,Y)/L)*B
2 +2*G(L+1,X,Y)/(A*(L+1))+G(L,X,Y)/L
  DE = X-Y
  GRAL(1,K) = F(1)*DE
  GRAL(2,K) = F(2)*DE
  GRAL(3,K) = 0.5*(3*F(3)-F(1))*DE
  GRAL(4,K) = 0.5*(5*F(4)-3*F(2))*DE
  GRAL(5,K) = 0.125*(35*F(5)-30*F(3)+3*F(1))*DE
  GRAL(6,K) = 0.125*(63*F(6)-70*F(4)+15*F(2))*DE
10 CCNTINUE
C
11 GRAL(1,NST) = 2.
  GRAL(2,NST) = 4/(3*A)
  GRAL(4,NST) = 0.

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GRAL(6,NST) = 0.
C
C*****REDUKTION VON NST AUF NAK *****
12 NP = N+1
13 IF(NAK.EQ.NST) RETURN
  K1 = (NST-1)/(NAK-1)
  KK = K1+1
  DO 15 K = 2,NAK
  H(K) = H(KK)
  DO 14 L = 1,NP
14 GRAL(L,K) = GRAL(L,KK)
15 KK = KK+K1
  D = 2./(NAK-1)
C
  RETURN
  END
C
FUNCTION G(N,X,Y)
C
C BERECHNUNG VON (X**N-Y**N)/(X-Y) DURCH ENTWICKLUNG NACH POTENZEN
C VON D = X-Y
C
REAL*8 G,X,Y,D
D = X-Y
GOTO(1,2,3,4,5,6,7,8,9,10),N
1 G = 1.
  RETURN
2 G = D+2*Y
  RETURN
3 G = D*D+3*Y*D+3*Y*Y
  RETURN
4 G = D*D*D+4*Y*D*D+6*Y*Y*D+4*Y*Y*Y
  RETURN
5 G = D**4+5*Y*D**3+10*Y*Y*D**2+10*Y*Y*Y*D+5*Y**4
  RETURN
6 G = D**5+6*Y*D**4+15*Y*Y*D**3+20*Y*Y*Y*D**2+15*Y**4*D+6*Y**5
  RETURN
7 G = D**6+7*Y*D**5+21*Y*Y*D**4+35*Y*Y*Y*D**3+35*Y**4*D*D
1 +21*Y**5*D+7*Y**6
  RETURN
8 G = D**7+8*Y*D**6+28*Y*Y*D**5+56*Y*Y*Y*D**4+70*Y**4*D*D*D
1 +56*Y**5*D*D+28*Y**6*D+8*Y**7
  RETURN
9 G = D**8+9*D**7*Y+36*D**6*Y*Y+84*D**5*Y*Y*Y+126*D**4*Y**4
1 +126*D**3*D*Y**5+84*D**2*Y**6+36*D*Y**7+9*Y**8
  RETURN
10 G = D**9+10*D**8*Y+45*D**7*Y*Y+120*D**6*Y*Y*Y+210*D**5*Y**4
1 +252*D**4*Y**5+210*D**3*D*Y**6+120*D**2*D*Y**7+45*D*Y**8+10*Y**9
  RETURN
  END

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SUBROUTINE LEGIST(L,NST,IA,IB,NEGR,GR,E)
C
C*****BERECHNUNG VON PARTIELLEN LEGENDRE-KOEFFIZIENTEN BEI ISOTROPIE IM
C*****SP-SYSTEM*****
C
REAL*8 X(321),GRAL(6,321),D,A,B
DIMENSION GR(NEGR),E(NEGR)
COMMON /INTEG/ X,GRAL,D
C
DO 7 I=IA,IB
GO = 1.
GU = -1.
IKO = NST
IKU = 1
IF(GR(I+1).GE.1.) GOTO 5
IF(GR(I).LE.-1.) GOTO 5
IF(GR(I).GE.1.) GOTO 2
GO = GR(I)
XIKO = 1.+0.5*(NST-1)*(GO+1.)
IKO = XIKO
IF(FLOAT(IKO).LT.XIKO) IKO = IKO+1
2 IF(GR(I+1).LE.-1.) GOTO 3
GU = GR(I+1)
XIKU = 1.+0.5*(NST-1)*(GU+1.)
IKU = XIKU
IF(FLOAT(IKU).LT.XIKU) IKU = IKU+1
3 A = GRAL(L,IKO)
IF(GO.LT.X( IKO)) A = A+(GO-X( IKO))/D*(GRAL(L,IKO)-GRAL(L,IKO-1))
B = GRAL(L,IKU)
IF(GU.LT.X( IKU)) B = B+(GU-X( IKU))/D*(GRAL(L,IKU)-GRAL(L,IKU-1))
4 E(I) = (A-B)*0.7957747E-01
GOTO 7
5 E(I) = 0.
7 CONTINUE
E(NEGR ) = GRAL(L,NST)*0.7957747E-01
C
RETURN
C
END

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SUBROUTINE LEGANS(L,NST,F,IA,IB,NEGR,GR,F,POLY,A)
C
C*****BERECHNUNG VON PARTIELLEN LEGENDRE-KOEFFIZIENTEN BEI ANISOTROPIE
C*****IM SP-SYSTEM*****
C
REAL*8 X(321),POL(6,321),D,B
DIMENSION F(NST),GR(NEGR),E(NEGR),POLY(NST),A(NST)
COMMON /INTEG/ X,POL,D
C
C*****BESTIMMUNG DER ZU INTEGRIERENDEN FUNKTION*****
DO 2 J = 1,NST
B = DBLE(F(J))

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2 POLY(J) = B*POL(L,J)
C
C*****SUCHEN DER GRENZINDICES MIT KORREKTURINTERVALLEN*****
ES = 0
I=IA-1
3 I = I+1
IF((GR(I+1).GE.1.).OR.(GR(I).LE.-1.)) GOTO 18
IF(GR(I).LT.1.) GOTO 4
II = I
GOTO 18
4 GO = GR(I)
GU=AMAX1(-1.,GR(I+1))
XIKO = 1.+0.5*(NST-1)*(1.+GO)-0.001
NO = XIKO
XIKU = 1.+0.5*(NST-1)*(1.+GU)+0.001
NU = XIKU
IF(FLOAT(NO).LT.XIKO) NO = NO+1
NUO = NO-NU
IF((NUO/2*2).NE.NUO) NO = N/2-1
DU = GU-X(NU)
DO = GO-X(NO)
IF(ABS(DU).LT.1.E-20) DU = 0
IF(ABS(DO).LT.1.E-20) DO = 0
C
C*****SIMPSON-INTEGRATION FUER UEBERSTREUUNG UND TOTALES MOMENT*****
5 NO1 = NO-2
NO2 = NO-1
NU1 = NU+2
NU2 = NU+1
IF(NO1.GE.NU1) GOTO 7
E(I) = (POLY(NU)+POLY(NO))*D/2.
IF(NO2.LT.NU2) GOTO 19
DO 6 J = NU2,NO2
6 E(I) = E(I)+D*POLY(J)
GOTO 19
7 DO 12 J = NU2,NO2,2
JJ = (J-NU2+2)/2
12 A(JJ) = POLY(J)
CALL SUM(JJ,A,SA)
DO 15 J = NU1,NO1,2
JJ = (J-NU1+2)/2
15 A(JJ) = POLY(J)
CALL SUM(JJ,A,SO)
SE = POLY(NO)+POLY(NU)
E(I) = (4*SA+2*SO+SE)*D/3
C
C*****RANDKORREKTUR DES SIMPSON-INTEGRALS*****
IF(I.EQ.NEGR) GOTO 20
TU = (POLY(NU2)-POLY(NU))/D
TO = (POLY(NO)-POLY(NO2))/D
E(I) = (POLY(NO)+0.5*DO*TO)*DO-(POLY(NU)+0.5*DU*TO)*DU+E(I)
GOTO 19
18 E(I) = 0.
19 ES = ES+E(I)
IF(I.LT.IB) GOTO 3
C

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C*****INTEGRATIONSBEREICH FUER DAS TOTALE MOMENT***** 690
  NU = 1 700
  NO = NST 710
  I=NEGR 720
  GOTO 5 730
C 740
C*****G,G-STREUUNG ***** 750
  20 E(II) = E(NEGR)-ES 760
  RETURN 770
C 780
  END 790

  SUBROUTINE GRUPIN(MASSE,XMAT,ALFA,COM,ISOT,PM,LEG,NLE,NANF,NEND, 10
  1 NUEB,IL,IM,NGR,NEGR,ABN,ISN,EH,K,H,F,M,G,EN, 20
  2 NTT,ET,ST,MAZ,KSPE,KSPEK,ESP,SPEK,NSPEK,LSPEK, 30
  3 ISM,ISD,NECU,ISCO,ISEC,NF,E,EW,GR,WA,U,V,W,ISTT, 40
  4 RSP,ELSIG) 50
C 60
C*****GEWICHTETE MULTIGRUPPEN-ENERGIEINTEGRATION***** 70
C 80
  DIMENSION LEG(6),ABN(NEGR),RSP(NEGR),ESP(NSPEK),SPEK(LSPEK), 90
  1 EH(ISM),EW(ISM),H(ISD),F(ISD),G(ISD),EN(ISD),V(ISD), 100
  2 W(ISD),E(ISD),WA(ISCO),U(ISEC),ET(NTT),ST(NTT),GR(NGR), 110
  3 MAZ(2),ELSIG(6,NECU,NGR) 120
C 130
  NP = NLE+1 140
  XP = XMAT+1 150
  XM = XMAT-1 160
  IF(MASSE.EQ.1) XM = 0. 170
  NM = NUEB-1 180
C 190
C*****VEREINIGUNG DER ENERGIESTUETZPUNKTMENGEN***** 200
  DO 2 J = 1,ISN 210
  2 E(J) = EH(J) 220
  ISTT = ISN 230
  CALL PUSUM(K,H,ISTT,E,V) 240
C 250
C*****INTERPOLATION DER TOTALEN ELASTISCHEN QUERSCHNITTE***** 260
  CALL IPOLIN(K,H,F,ISTT,E,W,U) 270
C 280
C*****INTERPOLATION DER MITTLEREN STREUKOSINUS***** 290
  IF(ISOT.EQ.0) GOTO 4 300
  CALL IPOLIN(M,G,EN,ISTT,E,V,U) 310
  GOTO 8 320
  4 DO 6 J = 1,ISTT 330
  6 V(J) = COM 340
C 350
  8 DO 66 N = 1,NP 360
  L = LEG(N) 370
  LISOT = L*ISOT 380
C*****LESEN UND INTERPOLATION DER NORMIERTEN ENERGIEDIFF. LEGENDRE-KOEFF 390
  DO 11 I = 1,NUEB 400

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  JI = (I-1)*ISTT 410
  READ(NF) (EW(J),J = 1,ISN) 420
  IF(ISOT.EQ.0.AND.I.EQ.NUEB) GOTO 9 430
  CALL IPOLIN(ISN,EH,EW,ISTT,E,WA(1+JI),U) 440
  GOTO 11 450
  9 XY = EW(ISN) 460
  DO 10 J = 1,ISTT 470
  10 WA(J+JI) = XY 480
  11 CONTINUE 490
C 500
C*****MUEL-ANGEPASSTE ENERGIEDIFF. LEGENDRE-KOEFF. ***** 510
  DO 44 J = 1,ISTT 520
  XY = 0. 530
  JP = J+NM*ISTT 540
  IF(LISOT.EQ.2) V(J) = PM*(V(J)-WA(JP)) 550
  DO 12 NN = NANF,NEND 560
  IF((ABN(NN+1).LE.E(J)).AND.(E(J).LE.ABN(NN))) GOTO 14 570
  12 CCNT INUE 580
  NN = NEND 590
  14 EL = AMAX1(ABN(IM+1),ALFA*E(J)) 600
  DO 42 I = 1,NUEB 610
  JI = J+(I-1)*ISTT 620
  NI = MINO(IM,NN+I-1) 630
  P = 0. 640
  IF(ISOT.EQ.0) GOTO 36 650
  GU = -1. 660
  IF(EL.GE.ABN(NI+1)) GOTO 16 670
  GU = SQRT(ABN(NI+1)/E(J)) 680
  GU = 0.5*(XP*GU-XM/GU) 690
  16 GO = 1. 700
  IF(I.EQ.1) GOTO 18 710
  LI = MINO(IM+1,NN+I-1) 720
  GO = -1. 730
  IF(EL.GE.ABN(LI)) GOTO 18 740
  GO = SQRT(ABN(LI)/E(J)) 750
  GO = 0.5*(XP*GO-XM/GO) 760
  18 IF(GU.LT.GO) GOTO 20 770
  P = 0. 780
  GOTO 36 790
  20 IF(GO-GU.LT.2.) GOTO 22 800
  P = 0. 810
  IF(L.EQ.2) P = 1. 820
  GOTO 36 830
  22 O2 = GO*GO 840
  U2 = GU*GU 850
  O3 = O2*GO 860
  U3 = U2*GU 870
  GOTO(24,26,28,30,32,34),L 880
  24 P = 0.75*(O2-U2) 890
  GOTO 36 900
  26 P = 0.5*(O3-U3) 910
  GOTO 36 920
  28 P = 0.1875*(O2*(3*O2-2)-U2*(3*U2-2)) 930
  GOTO 36 940
  30 P = 0.75*(O3*(O2-1)-U3*(U2-1)) 950
  GOTO 36 960

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32 P = 0.03125*(02*((02-1)*(35*02-10)-1)-02*((02-1)*(35*02-10)-1))  970
   GOTO 36 980
34 P = 0.1875*(03*(02-1)*(9*02-5)-03*(02-1)*(9*02-5)) 990
36 CONTINUE 1000
   IF(EL.LT.ABN(NI+1).AND.I.LT.NUEB) GOTO 40 1010
   WA(JI) = WA(JP)-XY 1020
   IF(I.EQ.NUEB) GOTO 42 1030
   IA = I+1 1040
   DO 38 II = IA,NUEB 1050
38 WA(J+(II-1)*ISTT) = 0. 1060
   WA(JI) = (WA(JI)+P*V(J))*W(J) 1070
   GOTO 44 1080
40 XY = XY+WA(JI) 1090
42 WA(JI) = (WA(JI)+P*V(J))*W(J) 1100
44 CONTINUE 1110
C 1120
C*****AUSSORTIEREN ZUR GEWICHTETEN INTEGRATION***** 1130
   DO 64 NR = 1,NUEB 1140
   NI = (NR-1)*ISTT 1150
   JA = 1 1160
   DO 58 NT = NANF,NEND 1170
   NN = NANF+NEND-NT 1180
C NN-TE GRUPPE = AUSSTREUGRUPPE 1190
C (NN+NR-1)-TE GRUPPE = EINSTREUGRUPPE 1200
   AM = ABN(NN) 1210
   DO 46 J = JA,ISTT 1220
   IF(E(J).GT.AM) GOTO 48 1230
   JI = J-JA+1 1240
   EN(JI) = E(J) 1250
46 F(JI) = WA(J+NI) 1260
   J = ISTT+1 1270
48 IKK = J-JA 1280
   JA = J-1 1290
   IF(NN+NR-1.GT.IM) GOTO 58 1300
   IF(IKK.GE.3) GOTO 52 1310
   VAL = 0. 1320
   GOTO 56 1330
52 EN(IKK) = AM 1340
   T = (F(IKK-2)-F(IKK-1))/(EN(IKK-2)-EN(IKK-1)) 1350
   F(IKK) = F(IKK-1)+T*(AM-EN(IKK-1)) 1360
C 1370
C*****AUSWAHL DES BENOETIGTEN TEILSPEKTRUMS UND INTERPOLATION***** 1380
   CALL SPRAL(KSPE,KSPEK,MAZ,ESP,SPEK,NTT,ET,ST,L,IKK,EN,G,NSPEK, 1390
   1 LSPEK) 1400
C 1410
C*****ERZEUGUNG DER NORMIERUNGSINTEGRALE***** 1420
   IF(NR.GT.1) GOTO 54 1430
   CALL TRAPEZ(IKK,EN,G,F,GR(NN),U,0) 1440
   IF(L.EQ.1) RSP(NN) = GR(NN) 1450
C 1460
C*****SPEKTRUMSGEWICHTETE ENERGIEINTEGRATION***** 1470
54 CALL TRAPEZ(IKK,EN,F,G,VAL,U,1) 1480
   VAL = VAL/GR(NN) 1490
56 ELSIG(L,NR,NN) = VAL 1500
58 CONTINUE 1510
64 CCNTINUE 1520

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66 CCNTINUE 1530
C 1540
   RETURN 1550
   END 1560

SUBROUTINE IPOLIN(M,A,B,N,X,Y,HR) 10
C 20
C*****LINEARE INTERPOLATION***** 30
C DIMENSION A(M),B(M),X(N),Y(N),HR(M) 40
C 50
C NM = N-1 60
C 70
C 80
C*****KOEFFIZIENTEN HR ***** 90
   MM = M-1 100
   DO 6 I = 1,MM 110
   6 HR(I) = (B(I+1)-B(I))/(A(I+1)-A(I)) 120
C 130
C*****INTERPOLATION***** 140
   Y(1) = B(1) 150
   Y(N) = B(M) 160
   I = 1 170
   DO 14 J = 2,NM 180
   8 IF(A(I+1)-X(J)) 9,10,11 190
   9 I = I+1 200
   GOTO 8 210
10 Y(J) = B(I+1) 220
   GOTO 14 230
11 Y(J) = B(I)+HR(I)*(X(J)-A(I)) 240
14 CONTINUE 250
C 260
   RETURN 270
   END 280

SUBROUTINE SPRAL(KS,NS,MZ,E,S,NTT,FT,ST,L,IKK,EN,G,NP,NL) 10
C 20
C*****BEREITSTELLUNG DES WICHTUNGSSPEKTRUMS***** 30
C 40
C DIMENSION E(NP),S(NL),EN(IKK),G(IKK),ET(NTT),ST(NTT),MZ(2) 50
C 60
C F0(A,B,X,Y) = (1/X+1/Y)*(B-A) 70
C F1(A,B,X,Y) = (1/(X*X)+1/(Y*Y))*(B-A) 80
C F2(A,B,X,Y) = (1/(X*X*X)+1/(Y*Y*Y))*(B-A) 90
C F3(A,B,X,Y) = (1/X**4+1/Y**4)*(B-A) 100
C F4(A,B,X,Y) = (1/X**5+1/Y**5)*(B-A) 110
C F5(A,B,X,Y) = (1/X**6+1/Y**6)*(B-A) 120
C 130
C IF(KS.NE.0) GO TO 2 140

```

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C*****SPEKTRUM AUS DER FUNKTION PHI(E) BZW. PHIL(L,E) *****
 40 IF(L.GT.1.AND.MZ(1).EQ.1) GOTO 42
    DO 41 I = 1,IKK
 41 G(I) = PHI(EN(I))
    GOTO 7
 42 DO 43 I = 1,IKK
 43 G(I) = PHIL(L,EN(I))
    GOTO 7
C
C*****SPEKTRUM AUS PUNKTWEISER EINGABE*****
 2 NN = 2
  DO 6 I = 1,IKK
  DO 3 N = NN,NS
  M = N
  IF(MZ(1).NE.0) M = M+(L-1)*NS
  IF(E(N)-EN(I)) 3,4,5
 3 CONTINUE
 4 G(I) = S(M)
  NN = N
  GOTO 6
 5 G(I) = S(M-1)+(S(M)-S(M-1))/(E(N)-E(N-1))*(EN(I)-E(N-1))
  IF(N.GT.2) NN = N-1
 6 CONTINUE
C
C*****UNGEWICHTETE INTERVALLBREITEN*****
 7 IF(NTT.GT.1) GOTO 9
  IK = IKK-1
  DO 8 I = 1,IK
 8 EN(I) = EN(I+1)-EN(I)
  RETURN
C
C*****FEINSTRUKTURGEWICHTETE INTERVALLBREITEN*****
 9 N = 1
  STE = -1
  IF(EN(IKK).LE.ET(2)) STE = ST(2)
  IF(EN(1).GE.ET(NTT-1)) STE = ST(NTT-1)
  STA = STE
  IF(STE.LT.0) GOTO 10
  N = 0
  M = 0
10 NN = 1
  IK = IKK-1
  DO 38 I = 1,IK
  IF(N.EQ.0) GOTO 17
  DO 11 N = NN,NTT
  IF(ET(N)-EN(I)) 11,11,12
11 CONTINUE
  N = NTT
12 DO 14 M = N,NTT
  IF(ET(M)-EN(I+1)) 14,16,16
14 CONTINUE
  M = NTT
16 STA = ST(N)+(EN(I)-ET(N))*(ST(N-1)-ST(N))/(ET(N-1)-ET(N))
  STE = ST(M)+(EN(I+1)-ET(M))*(ST(M-1)-ST(M))/(ET(M-1)-ET(M))
17 IF(MZ(2).EQ.0) GOTO 18
  GOTO (18,21,24,27,30,33),L

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```

C***** 0. MOMENT
18 IF(M.GT.N) GOTO 19
  W = F0(EN(I),EN(I+1),STA,STE)
  GOTO 36
19 W = F0(EN(I),ET(N),STA,ST(N))+F0(ET(M-1),EN(I+1),STE,ST(M-1))
  IF(N.EQ.M-1) GOTO 36
  MM = M-2
  DO 20 J = N,MM
20 W = W+F0(ET(J),ET(J+1),ST(J+1),ST(J))
  GOTO 36
C***** 1. MOMENT
21 IF(M.GT.N) GOTO 22
  W = F1(EN(I),EN(I+1),STA,STE)
  GOTO 36
22 W = F1(EN(I),ET(N),STA,ST(N))+F1(ET(M-1),EN(I+1),STE,ST(M-1))
  IF(N.EQ.M-1) GOTO 36
  MM = M-2
  DO 23 J = N,MM
23 W = W+F1(ET(J),ET(J+1),ST(J+1),ST(J))
  GOTO 36
C***** 2. MOMENT
24 IF(M.GT.N) GOTO 25
  W = F2(EN(I),EN(I+1),STA,STE)
  GOTO 36
25 W = F2(EN(I),ET(N),STA,ST(N))+F2(ET(M-1),EN(I+1),STE,ST(M-1))
  IF(N.EQ.M-1) GOTO 36
  MM = M-2
  DO 26 J = N,MM
26 W = W+F2(ET(J),ET(J+1),ST(J+1),ST(J))
  GOTO 36
C***** 3. MOMENT
27 IF(M.GT.N) GOTO 28
  W = F3(EN(I),EN(I+1),STA,STE)
  GOTO 36
28 W = F3(EN(I),ET(N),STA,ST(N))+F3(ET(M-1),EN(I+1),STE,ST(M-1))
  IF(N.EQ.M-1) GOTO 36
  MM = M-2
  DO 29 J = N,MM
29 W = W+F3(ET(J),ET(J+1),ST(J+1),ST(J))
  GOTO 36
C***** 4. MOMENT
30 IF(M.GT.N) GOTO 31
  W = F4(EN(I),EN(I+1),STA,STE)
  GOTO 36
31 W = F4(EN(I),ET(N),STA,ST(N))+F4(ET(M-1),EN(I+1),STE,ST(M-1))
  IF(N.EQ.M-1) GOTO 36
  MM = M-2
  DO 32 J = N,MM
32 W = W+F4(ET(J),ET(J+1),ST(J+1),ST(J))
  GOTO 36
C***** 5. MOMENT
33 IF(M.GT.N) GOTO 34
  W = F5(EN(I),EN(I+1),STA,STE)
  GOTO 36
34 W = F5(EN(I),ET(N),STA,ST(N))+F5(ET(M-1),EN(I+1),STE,ST(M-1))
  IF(N.EQ.M-1) GOTO 36

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```

MM = M-2
DO 35 J = N,MM
35 W = W+F5(ET(J),ET(J+1),ST(J+1),ST(J))
36 EN(I) = 0.5*W
38 NN = M-1
C
RETURN
END

SUBROUTINE TRAPEZ(M,E,F,G,W,H,L)
C
C*****INTEGRATION NACH DER TRAPEZREGEL*****
C
DIMENSION E(M),F(M),G(M),H(M)
C
L = 0 : GEWICHTSFUNKTION KONSTANT = 1
L = 1 : GEWICHTSFUNKTION NICHT KONSTANT
C
IF(L.EQ.0) GOTO 8
DO 2 J = 2,M
2 H(J-1) = (F(J-1)*G(J-1)+F(J)*G(J))*E(J-1)
GOTO 12
8 DO 10 J = 2,M
10 H(J-1) = (F(J)+F(J-1))*E(J-1)
12 MM = M-1
CALL SUM(MM,H,W)
W = 0.5*W
C
RETURN
END

FUNCTION PHIL(L,E)
PHIL=1./E**L
RETURN
END

```

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1270
1280
1290
1300
1310
1320
1330
1340

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100
110
120
130
140
150
160
170
180
190
200
210

10
20
30
40

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SUBROUTINE MUKON(ELSIG,ELTOT,NLE,NECU,NUEB,IL,IM,NGR)
C
C*****MULTIGRUPPENKONSTANTEN*****
C
DIMENSION ELSIG(6,NECU,NGR),ELTOT(2,NGR)
C
NP = NLE+1
C*****TOTALER STREUQUERSCHNITT*****

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10
20
30
40
50
60
70
80

```

```

DO 14 NN = IL,IM
STO = 0.
DO 12 NS = 1,NUEB
NT = NS+NN-1
STO = STO+ELSIG(1,NS,NN)
IF(NT.EQ.IM) GOTO 14
12 CONTINUE
14 ELTOT(1,NN) = STO
C
C*****NORMIERUNG AUF DEN TOTALEN STREUQUERSCHNITT*****
22 DO 26 L = 1,NP
DO 26 NN=IL,IM
DO 24 I = 1,NUEB
NT=I+NN-1
ELSIG(L,I,NN) = ELSIG(L,I,NN)/ELTOT(1,NN)
IF(NT.EQ.IM) GO TO 26
24 CONTINUE
26 CONTINUE
C
C*****MITTLERER STREUKOSINUS*****
15 DO 20 NN = IL,IM
STO = 0.
DO 18 NS = 1,NUEB
NT = NS+NN-1
STO = STO+ELSIG(2,NS,NN)
IF(NT.EQ.IM) GOTO 20
18 CONTINUE
20 ELTOT(2,NN) = STO
C
C*****ABSCHNEIDEN (NULLSETZEN) VON GROESSEN KLEINER 1.E-6 *****
DO 40 L = 1,NP
DO 40 NN=IL,IM
DO 38 I = 1,NUEB
NT=I+NN-1
ELSIG(L,I,NN) = AINT(ELSIG(L,I,NN)*1.E+5)*1.E-5
IF(NT.EQ.IM) GO TO 40
38 CONTINUE
40 CONTINUE
C
RETURN
END

```

```

SUBROUTINE PRINT(ELSIG,ELTOT,NLA,NLE,NECU,NUEB,ISEL,NGR,KSPEK,
1 MAZ,NTK,KIM)
C
C*****SUBROUTINE ZUM AUSDRUCKEN DER STREUMATRIX,DER STREUKOSINUSMATRIX**
C
REAL*8 STOFF,MMM
DIMENSION ELSIG(6,NECU,NGR),ELTOT(2,NGR),MAZ(2)
COMMON STOFF,ISTRUK,ISPA,NOUT,KPR,NEND,NANF
C
NP = NLE+1

```

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100
110
120
130
140
150
160
170
180
190
200
210
220
230
240
250
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300
310
320
330
340
350
360
370
380
390
400
410
420
430
440
450
460
470
480
490

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40
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60
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80
90
100

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C	110	I=3	670
C*****	120	DO 31 NN=I1, I2	680
AUSDRUCKEN DER	130	WRITE (KPR) I, NN, ELTOT(1, NN), ELTOT(2, NN)	690
STREUMATRIZEN*****	140	31 CONTINUE	700
DO 3 L = 1, NP	150	DO 33 L = 1, NP	710
LM = L-1	160	LM=L-1	720
WRITE(NOUT,21) LM, LM, STOFF	170	DO 33 NN=I1, I2	730
IF(IABS(ISEL).EQ.1) WRITE(NOUT,28)	180	LL=MINO(NUEB, I2+KIM-NN+1)	740
I1 = NANF	190	I=LL+2	750
1 I2=I1+MINO(5, NEND-I1)	200	WRITE (KPR) I, LM, NN, (ELSIG(L, K, NN), K=1, LL)	760
WRITE(NOUT,24) (I, I = I1, I2)	210	33 CONTINUE	770
DO 2 I = 1, NUEB	220	C	780
I3=MINO(I2, NEND+KIM-I+1)	230	RETURN	790
IF(I3.LT.I1) GO TO 2	240	C	800
WRITE(NOUT,22) (ELSIG(L, I, NN), NN=I1, I3)	250	20 FORMAT(1H0)	810
2 CONTINUE	260	21 FORMAT(1H0//29X, 'SGNC', I1, 8X, 'ELASTISCHE STREUMATRIX ', I1, '. ORDN	820
WRITE(NOUT,20)	270	1UNG FUER ', A8//)	830
IF(I2.GE.NEND) GOTO 3	280	22 FORMAT(23X, 6F12.5)	840
I1 = I2+1	290	24 FORMAT(23X, 6(2X, I3, '.GRUPPE'))	850
GOTO 1	300	25 FORMAT(23X, 1P6E12.3)	860
3 CCNTINUE	310	28 FORMAT(56X, 'FUER SGN(E) = 1'/)	870
WRITE(NOUT,52)	320	27 FORMAT(1H0//31X, 'TOTALE ELASTISCHE STREUQUERSCHNITTE SGN UND STREU	880
C	330	1KOSINUS MUEL FUER ', A8//)	890
C*****	340	30 FORMAT(////42X, 'M A K R O W I C H T U N G '/')	900
INFORMATIONEN UEBER DIE	350	32 FORMAT(42X, 'ALLE MOMENTE WIE DAS 0. MOMENT MIT DEM'/	910
WICHTUNG*****	360	1 42X, 'EINGELESENEN PUNKTSPEKTRUM F(O,E)')	920
WRITE(NOUT,30)	370	34 FORMAT(42X, 'DAS L-TE MOMENT MIT DEM (L+1)-TEN EINGE-'/	930
IF(KSPEK.GT.0.AND.MAZ(1).EQ.0)WRITE(NOUT,32)	380	1 42X, 'LESENEN PUNKTSPEKTRUM F(L,E), L = 0,1,..', I2)	940
IF(KSPEK.GT.0.AND.MAZ(1).EQ.1)WRITE(NOUT,34) NLE	390	36 FORMAT(42X, 'ALLE MOMENTE WIE DAS 0. MOMENT MIT'/	950
IF(KSPEK.EQ.0.AND.MAZ(1).EQ.0)WRITE(NOUT,36)	400	1 42X, 'F(O,E) (STANDARD F(O,E) = 1/E)')	960
IF(KSPEK.EQ.0.AND.MAZ(1).EQ.1)WRITE(NOUT,38) NLE	410	38 FORMAT(42X, 'DAS L-TE MOMENT MIT F(L,E), L = 0,1,..', I2/	970
WRITE(NOUT,40)	420	1 42X, '(STANDARD F(L,E) = 1/E**((L+1)))')	980
IF(NTK.EQ.1.AND.MAZ(2).EQ.0)WRITE(NOUT,42)	430	40 FORMAT(////42X, 'M I K R O W I C H T U N G (FEINSTRUKTUR)'/)	990
IF(NTK.EQ.1.AND.MAZ(2).EQ.1)WRITE(NOUT,44) NLE	440	42 FORMAT(42X, 'ALLE MOMENTE WIE DAS 0. MOMENT MIT FS(O,E)=1/SGT(E)'/	1000
IF(NTK.EQ.0) WRITE(NOUT,46)	450	1 42X, 'SGT(E)=TOTALER QUERSCHNITT DER EINGELESENEN MISCHUNG')/	1010
C	460	44 FORMAT(42X, 'DAS L-TE MOMENT MIT FS(L,E), L = 0,1,..', I2/	1020
C*****	470	1 42X, '(STANDARD FS(L,E)=(1/SGT(E))**((L+1)))'/	1030
AUSDRUCKEN DER	480	2 42X, 'SGT(E)=TOTALER QUERSCHNITT DER EINGELESENEN MISCHUNG')/	1040
TOTALEN QUERSCHNITTE UND	490	46 FORMAT(42X, 'ALLE MOMENTE MIT FS(O,E) = 1'/	1050
STREUKOSINUS*****	500	1 42X, '(KEINE FEINSTRUKTURWICHTUNG)')	1060
WRITE(NOUT,27) STOFF	510	52 FORMAT(//29X, 'ERLAEUTERUNG'/29X, 'L-TE ORDNUNG, G-TE GRUPPE, I-TE Z	1070
IF(IABS(ISEL).EQ.1) WRITE(NOUT,28)	520	1EILE: MATRIXELEMENT L-TER ORDNUNG'/29X, 'FUER STREUUNG AUS DER G-TE	1080
I1 = NANF	530	2N GRUPPE IN DIE (G+I-1)-TE GRUPPE, BEZOGEN'/29X, 'AUF DEN TOTALEN E	1090
9 I2=I1+MINO(5, NEND-I1)	540	3LASTISCHEN QUERSCHNITT (TOTALES 0.MOMENT) DER'/29X, 'G-TEN GRUPPE')/	1100
WRITE(NOUT,24) (I, I = I1, I2)	550	54 FORMAT(//29X, 'ERLAEUTERUNG'/29X, 'G-TE GRUPPE, 1. ZEILE: TOTALER EL	1110
DO 10 I = 1, 2	560	1ASTISCHER QUERSCHNITT DER G-TEN'/29X, 'GRUPPE'/29X, 'G-TE GRUPPE, 2.	1120
10 WRITE(NOUT,25) (ELTOT(I, NN), NN = I1, I2)	570	2 ZEILE: MITTLERER STREUKOSINUS DER G-TEN GRUPPE')/	1130
WRITE(NOUT,20)	580	END	1140
IF(I2.GE.NEND) GOTO 11	590		
I1 = I2+1	600		
GOTO 9	610		
11 WRITE(NOUT,54)	620		
C	630		
C*****	640		
RESERVIEREN DER	650		
RESULTATE*****	660		
CALL DOPW(8HFLUM , MMM)			
I=0			
WRITE (KPR) I, MMM			
I1=NANF			
I2=NEND			
I=4			
L=NLE-NLA+1			
NN=I2-I1+1			
WRITE (KPR) I, STOFF, NN, L			

	SUBROUTINE SPALT(NE,ENG,X)	10	WRITE (NOUTP,12) MAT,E,K,J	570
C		20	12 FORMAT(1H0,A9,' EINFALLSENERGIE =',E16.8,' VON GRUPPE',I3,' BIS'	580
C	BERECHNUNG DES SPALTSPEKTRUMS	30	1, I3)	590
	REAL*8 MAT,FEST(10),NFEST(10),MMM	40	WRITE(NOUTP,14) I	600
	REAL*8 E,B,C,T,EF,EW,EG,EG1,XG,XG1,YG,YG1,XG2,YG2,XG12,YG12	50	14 FORMAT(' CHI ',I1)	610
	DIMENSION ENG(NE),NADAT(2),NSATZ(4),X(NE)	60	I=I+1	620
	COMMON MAT,ISTRUK,ISPA,NOUTP,JA,NANF,NEND,KL	70	WRITE (NOUTP,13) (X(IJ),IJ=L,IGRUP)	630
	EQUIVALENCE(FEST(1),NFEST(1))	80	13 FORMAT(1H ,7E16.8/(1X,7E16.8))	640
9000	WRITE (NOUTP,9000)	90	CALL NDFNXT(NSUCH,NSATZ,FEST,NUDAT,NC)	650
	FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 7')	100	IF(NSUCH) 5,5,7	660
	WRITE (NOUTP,9001)	110	5 KL=KL+1	670
9001	FORMAT(' PROGRAMM ZUR BERECHNUNG DES SPALTSPEKTRUMS')	120	RETURN	680
	N=1	130	END	690
	I=0	140		
	L=NE-NANF	150		
	IGRUP=NE-NEND	160	C	BESTIMMUNG DER 1/V-WERTE IM KFK-SATZ
	CALL DOPW (8HBEST ,NFEST(2))	170		10
	CALL DOPW (8HCHICR ,NFEST(3))	180	SUBROUTINE EDV (MM,A,NEF,XS,G,E,V,Y,NES,ES,F)	20
	CALL DOPW (8HSPALT ,MMM)	190	REAL*8 MAT,KZF	30
	NSATZ(1)=3	200	DIMENSION A(MM),XS(NEF),G(NEF),E(MM),V(MM),Y(MM)	40
	NFEST(1)=MAT	210	1,ES(NES),F(NES)	50
	CALL NDFLOC (NSUCH,NSATZ,FEST,NUDAT,NC)	220	COMMON MAT,ISTRUK,ISPA,NOUTP,JA,NANF,NEND,KL	60
	IF(NSUCH)6,6,7	230	WRITE (NOUTP,9000)	70
6	WRITE (NOUTP,8) FEST(1)	240	9000 FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 8')	80
8	FORMAT(' FUER DAS MATERIAL',A9,' GIBT ES KEIN SPALTSPEKTRUM')	250	WRITE (NOUTP,9001)	90
	GO TO 5	260	9001 FORMAT(' PROGRAMM ZUR BERECNUNG DES 1/V- GRUPPENMITTELWERTES')	100
7	E=FEST(4)	270	DO 5 I=1,NEF	110
	B=FEST(6)	280	ES(I)=XS(I)	120
	C=FEST(7)	290	5 F(I)=G(I)	130
	T=1./B	300	C=7.22929E-7	140
	EF= C/(4.*B*B)	310	NF=NEF	150
	EW=DSQRT(EF/T)	320	IF(NEF.EQ.1) NF=0	160
	DO 9 K=L,IGRUP	330	CALL DOPW (8HS1/V ,KZE)	170
	EG=DSQRT(ENG(K)/T)	340	DO 914 I=1,MM	180
	EG1=DSQRT(ENG(K+1)/T)	350	L=MM+1-I	190
	XG=EG1-EW	360	914 E(L)=A(I)	200
	XG1=EG-EW	370	IS=0	210
	YG=FG1+EW	380	J=MM	220
	YG1=EG+EW	390	IF(NF)25,3,25	230
	XG2=XG*XG	400	3 ES(1)=E(J)	240
	YG2=YG*YG	410	M=1	250
	XG12=XG1*XG1	420	L=MM-1	260
	YG12=YG1*YG1	430	DO 14 I=1,L	270
	X(K)=0.2820947917738/FW*(DEXP(-XG12)-DEXP(-XG2)-DEXP(-YG12)	440	K=MM-I	280
	1+DEXP(-YG2))-0.5*(DERF(XG1)-DERF(XG)+DERF(YG1)-DERF(YG))	450	ER1=E(K)	290
9	CONTINUE	460	ER2=E(K+1)	300
	IF(N.EQ.2) GO TO 10	470	DO 12 K=1,9	310
	N=2	480	M=M+1	320
	IJ=0	490	AK=K	330
	WRITE (JA) IJ,MMM	500	AK=AK*0.1	340
10	IJ=6	510	12 ES(M)=ER2+AK*(ER1-ER2)	350
	K=NE-L	520	IF(I-L)16,14,16	360
	J=NE-IGRUP	530	16 M=M+1	370
	WRITE (JA) IJ,MAT,E,K,J	540	14 ES(M)=ER1	380
	NNN=IGRUP+1-L	550		
	WRITE(JA) NNN,(X(IJ),IJ=L,IGRUP)	560		

DO 13 I=1,M	390	SZ2=ES(IS-1)-E(J)	950
13 F(I)=PHI(ES(I))	400	SZ=SZ1*SZ2	960
25 IS=IS+1	410	SUMZ=SUMZ+SZ	970
IF(E(J)-ES(IS))10,22,25	420	SN1=Y(J)+F(IS-1)	980
10 IF(IS-1)11,11,30	430	SN=SN1*SZ2	990
11 NFEHL=1	440	SUMN=SUMN+SN	1000
101 WRITE (NOUTP,100) NFEHL	450	501 IF(LAR)1,55,1	1010
100 FORMAT(1H0,6HFEHLER13)	460	55 IS=IS-1	1020
GO TO 999	470	1 V(J)=C*SUMZ/SUMN	1030
22 Y(J)=F(IS)	480	V(MM)=0.8862269*C/SQRT(0.0253)	1040
GO TO 40	490	WRITE (NOUTP,300) (I,V(I),I=1,MM)	1050
30 XM=(F(IS)-F(IS-1))/(ES(IS)-ES(IS-1))	500	300 FORMAT(1H0,6HGRUPPE,7X,3H1/V/(15,E17.8))	1060
Y(J)=XM*(E(J)-ES(IS-1))+F(IS-1)	510	916 L=0	1070
IS=IS-1	520	WRITE (JA)L,KZE	1080
40 KS=MM-1	530	DO 502 I=1,MM	1090
DO 1 I=1,KS	540	KP=MM-I+1	1100
J=MM-I	550	502 Y(I)=V(KP)	1110
4 KZ=0	560	WRITE (JA) MM,(Y(I),I=1,MM)	1120
IJ1=IS	570	KL=KL+1	1130
77 IF(E(J)-ES(IS))68,71,70	580	RETURN	1140
70 KZ=KZ+1	590	999 STOP	1150
80 IS=IS+1	600	END	1160
GO TO 77	610		
68 XM=(F(IS)-F(IS-1))/(ES(IS)-ES(IS-1))	620		
Y(J)=XM*(E(J)-ES(IS-1))+F(IS-1)	630		
LAR=0	640		
GO TO 550	650	SUBROUTINE REMO(NG,ABN,NSP,F,E,ERR,NLA,NLE,NGRE,NFG,NFI,NJM,NUJM,	10
71 Y(J)=F(IS)	660	ISEL,IWORK,WORK,L)	20
LAR=1	670	DIMENSION IWORK(1),WORK(1),KDAT(2)	30
550 SUMZ=0.	680	REAL*8 FNEV(3)	40
SUMN=0.	690	REAL*8 NHID	50
S1=Y(J+1)/SQRT(E(J+1))	700	REAL*8 MATN	60
S2=F(IJ1+1)/SQRT(ES(IJ1+1))	710	REAL*8 REM	70
SZ1=S1+S2	720	DIMENSION ABN(1),F(1),E(1),NFG(1),NFI(1)	80
SZ2=E(J+1)-ES(IJ1+1)	730	COMMON MATN,IQQ(2),NOUT,LI,M1,M2,KLL	90
SZ=SZ1*SZ2	740	EQUIVALENCE(N7,N8)	100
SUMZ=SUMZ+SZ	750	EQUIVALENCE(N3,N4)	110
SN1=F(IJ1+1)+Y(J+1)	760	DATA NHID/'H H1'/	120
SN=SN1*SZ2	770	DATA FNEV/'SGNC ','SGT ','SGN '/	130
SUMN=SUMN+SN	780	DATA REM/'REMO '/	140
IF(KZ-2)501,500,600	790	KLL=KLL+1	150
600 KZ=IJ1+KZ-2	800	WRITE (NOUT ,90C0)	160
IJ1=IJ1+1	810	9000 FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 9')	170
DO 2 N=IJ1,KZ	820	WRITE (NOUT ,9001)	180
S1=F(N)/SQRT(ES(N))	830	9001 FORMAT(' PROGRAMM ZUR BERECHNUNG DER ELASTISCHEN STREUMATRIZEN FUE	190
S2=F(N+1)/SQRT(ES(N+1))	840	1R DIE REMO-KORREKTUR'/)	200
SZ1=S1+S2	850	IWORK(24)=0	210
SZ2=ES(N)-ES(N+1)	860	IWORK(33)=35	220
SZ=SZ1*SZ2	870	MO=IWORK(33)+2+2*NJM	230
SUMZ=SUMZ+SZ	880	MM=MO+NG+1	240
SN1=F(N)+F(N+1)	890	IF(MM-(MM/2)*2.EQ.0)MM=MM+1	250
SN=SN1*SZ2	900	IWORK(34)=MM	260
2 SUMN=SUMN+SN	910	IWORK(8)=ISEL	270
500 S1=F(IS-1)/SQRT(ES(IS-1))	920	IWORK(14)=NG	280
S2=Y(J)/SQRT(E(J))	930	IWORK(15)=NSP	290
SZ1=S1+S2	940		

IWORK(16)=NLA	300	CALL DATNUM(MATN,FNEV(J+1),NDAT,NW,WORK(M0+M2-1),WORK(M1+M0),	860
IWORK(17)=NLE	310	1&10,&10,&10)	870
IWORK(19)=NJM+1	320	IF(KCAT(J).LT.NDAT)KDAT(J)=NDAT	880
IWORK(20)=NUJM	330	GOTO 6	890
WORK(27)=ERR	340	10 WRITE(NOUT,71)MATN,FNEV(J+1)	900
WORK(32)=-0.1	350	71 FORMAT(' ***ERROR 9.5 : FOR ',A8,' THERE IS NO ',A8)	910
DO 30 I=1,NG	360	6 CONTINUE	920
I1=NG-I+1	370	IF(KCAT(2).LT.KDAT(1)) KDAT(2)=KDAT(1)	930
30 WORK(M0+I1-1)=ABN(I)	380	N5=N4+KDAT(2)	940
WORK(M0+NG)=0.	390	N6=N5+KDAT(2)	950
CALL MASSIN(WORK,WORK(M0))	400	KK=KDAT(2)	960
IMAX=IWORK(18)	410	7 IMAX=IWORK(18)	970
NO=IWORK(34)+(NLF+1)*IMAX*2	420	N7=N6+KK	980
IWORK(1)=NO	430	N9=N8+ICOS	990
IF(MATN.EQ.NHID)GOTO 3	440	N10=N9+IMAX*NUJM*(NLE+1)	1000
CALL NIVNUM(MATN,FNEV(1),NIV,WORK(M0+M2-1),WORK(M0+M1),FNW,	450	N11=IMAX*(NLE+1)+N10	1010
1&3,&3,&3)	460	N11=N11+1+IMAX	1020
NIV=NIV+2	470	IWORK(6)=N5	1030
GOTO 4	480	IWORK(7)=N6	1040
3 NIV=0	490	IWORK(9)=N8	1050
ICOS=0	500	IWORK(10)=N9	1060
N2=NO	510	IWORK(11)=N10	1070
N3=NO	520	IWORK(12)=N11	1080
N4=NO	530	LM=N11+KDAT(2)	1090
IS=M2-1	540	IF(L.LT.LM)GOTO 20	1100
4 IWORK(26)=NIV	550	NZM=1	1110
N1=NO+NIV	560	DO 9 I=M2,M1	1120
IWORK(2)=N1	570	NZ=NFG(I)*NFI(I)	1130
IF(NIV.EQ.0)GOTO 5	580	NZM=MAX0(NZM,NZ)	1140
NW=1	590	9 CONTINUE	1150
CALL DATNUM(MATN,FNEV(1),ICOS,NW,1.0,-1.0,&3,&3,&3)	600	IZV=IMAX*(NLE+1)	1160
LM=N1+ICOS	610	IZV=MAX0(4,IZV)	1170
IF(LM.GT.L)GOTO 20	620	IWORK(13)=N11+NZM*IZV	1180
IF(FNW.GE.WORK(M0+M2-1))GOTO 3	630	JJ=2+IMAX*NZM	1190
IS=ICSDP(FNW,WORK(M0))	640	JJ=MAX0(JJ,KCAT(2))	1200
IF(FNW.EQ.WORK(M0+IS-1))IS=IS-1	650	LM=IWORK(13)+JJ	1210
N2=N1+ICOS	660	20 CONTINUE	1220
N3=N2+ICOS*NIV	670	IF(L.LT.LM)WRITE(NOUT,70)LM	1230
L1=N3+NIV	680	70 FORMAT(' ***ERROR 9.1 : REQUIRED WORKING FIELD',I8,' WORDS')	1240
L2=L1+ICOS	690	IF(L.LT.LM)RETURN	1250
LM=L2+ICOS*NIV	700	NQ=0	1260
IF(M1.LT.IS)IS=M1	710	WRITE(LI)NQ,REM	1270
IF(LM.GT.L)GOTO 20	720	NQ=4	1280
CALL KEDLEC(MATN,FNEV(1),NIV,WORK(NO),IWORK(N3),WORK(M0+M2-1),	730	NM1=NLF-NLA+1	1290
IWORK(M0+IS),	740	WRITE(LI)NQ,MATN,NG,NM1	1300
IWORK(N1),WORK(L1),WORK(L2),WORK(N2),1.0,-1.0,&3,&3,&3)	750	CALL REMP(WORK(M0),NFG,NFI,F,E,WORK,IWORK,IS)	1310
IWORK(26)=NIV	760	RETURN	1320
5 CCNTINUE	770	END	1330
IWORK(21)=ICOS	780		
IWORK(3)=N2	790		
IWORK(5)=N4	800		
KCAT(1)=1	810		
KCAT(2)=1	820		
LM=N4	830		
DO 6 I=M2,M1	840		
DO 6 J=1,2	850		

<pre> SUBROUTINE 1 KEDDAT(NAMIZ,NAMTYP,NW,NW1,FNW,NDAT,FA,FF,EP,EM,*,*,*) REAL*8 FM(6),NAMIZ,NAMTYP,NGNC,NGIZ DIMENSION EA(1),IEA(1),FA(1),FF(1),FAW(1),FFW(1),NWORT(4),NDT(2) DATA FM(2)/*BEST *//,KERST/0/,NGNC/'SGNC *//,NGIZ/'SGIZ */ KK=1 IF(NAMTYP.EQ.NGNC.OR.NAMTYP.EQ.NGIZ)KK=2 GOTO 1 ENTRY DATNUM(NAMIZ,NAMTYP,NDAT,NW,EP,EM,*,*,*) KK=3 IF(NAMTYP.EQ.NGNC.OR.NAMTYP.EQ.NGIZ)KK=4 GOTO 1 ENTRY KEDLEC(NAMIZ,NAMTYP,NW1,EA,IEA,EPW,EMW,FA,FF,FAW,FFW,EP,EM, 1*,*,*) KK=5 NW=0 GOTO 1 ENTRY NIVNUM(NAMIZ,NAMTYP,NW1,EPW,EMW,EFR,*,*,*) NW=0 KK=6 1 IF(KERST.NE.0)GOTO 20 L=1 KERST=1 20 FM(1)=NAMIZ FM(3)=NAMTYP GOTO (2,3,2,3,3,3),KK 2 CONTINUE NWORT(1)=3 LL=0 CALL NDFLOC(NT,NWORT,FM,N1,N2) IF(NT.EQ.0)RETURN 1 22 NDAT=1 IF(FM(4+LL).LE.EM)GOTO 21 IF(KK.EQ.3)GOTO 40 FA(1)=EM FF(1)=FM(5+LL) 40 CONTINUE NDAT=2 21 CONTINUE IF(KK.EQ.3)GOTO 8 FA(NDAT)=FM(4+LL) FF(NDAT)=FM(5+LL) 8 NDAT=NDAT+1 CALL NDFNXT(NT,NWORT,FM,N1,N2) IF(NT.EQ.0)GOTO 9 IF(FM(4+LL).LE.EM)NDAT=1 IF(KK.EQ.3)GOTO 41 FA(NDAT)=FM(4+LL) FF(NDAT)=FM(5+LL) 41 CONTINUE IF(FM(4+LL).GE.EP)GOTO 30 GOTO 8 30 IF(LL.EQ.0)RETURN IF(KK.EQ.1.OR.KK.EQ.3)RETURN GOTO 10 9 CONTINUE </pre>	<pre> 10 20 30 40 50 60 70 80 90 100 110 120 130 140 150 160 170 180 190 200 210 220 230 240 250 260 270 280 290 300 310 320 330 340 350 360 370 380 390 400 410 420 430 440 450 460 470 480 490 500 510 520 530 540 550 560 </pre>	<pre> IF(KK.EQ.3)GOTO 42 FA(NDAT)=EP FF(NDAT)=FF(NDAT-1) 42 CONTINUE GOTO 30 3 CONTINUE NWORT(1)=4 LL=1 NW1=0 FM(4)=1.0 23 CALL NDFLOC(NT,NWORT,FM,N1,N2) CALL NDFLOC(NT,NWORT,FM,N1,N2) IF(NT.EQ.0.AND.NW1.EQ.0)RETURN 1 IF(NT.EQ.0)GOTO 18 IF(KK.EQ.6.AND.NW1.EQ.0)EFR=FM(4) NW1=NW1+1 IF(NW.EQ.0.AND.FM(4).GE.EMW)GOTO 24 IF(NW.EQ.NW1)GOTO 24 FM(4)=(1+1.E-6)*FM(4) GOTO 23 24 CONTINUE IF(KK.GE.7)GOTO 26 IF(KK.GE.5)GOTO 25 FNW=FM(4) KK=KK-1 GOTO 22 25 NW=NW1-1 IF(EMW.EQ.FM(4))NW=NW1 LZ=0 IF(NW.EQ.0)LZ=1 IF(NW.EQ.0)NW=1 NPN=0 KK=KK+2 GOTO 3 26 FNW=FM(4) IF(KK.EQ.7)GOTO 22 IF(FM(4).GE.EPW)RETURN FM(4)=(1+1.E-6)*FM(4) NW=NW+1 GOTO 23 10 IF(LZ.NE.0)GOTO 88 DO 90 I=1,NDAT FAW(I)=FA(I) 90 FFW(I)=FF(I) IEA(1)=NDAT EA(1)=EMW LZ=2 NPN=NDAT 88 CONTINUE DO 91 I=1,NDAT FAW(I+NPN)=FA(I) 91 FFW(I+NPN)=FF(I) EA(LZ)=FNW IEA(LZ)=NDAT NPN=NPN+NDAT LZ=LZ+1 </pre>	<pre> 570 580 590 600 610 620 630 640 650 660 670 680 690 700 710 720 730 740 750 760 770 780 790 800 810 820 830 840 850 860 870 880 890 900 910 920 930 940 950 960 970 980 990 1000 1010 1020 1030 1040 1050 1060 1070 1080 1090 1100 1110 1120 </pre>
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      IF(FNW.GE.EPW)GOTO 89
      NW=NW+1
      FM(4)=(1+1.E-6)*FNW
      GOTO 23
18  FNW=EPW
      IF(KK.EQ.7)GOTO 88
      NW1=NW1+1
      RETURN
89  NW1=LZ-1
      RETURN
      END

```

```

FUNCTION BCM(EI,EO,WORK)
DIMENSION WORK(1)
AM=WORK(28)
Q=WORK(32)
IF(EO.EQ.0.0)GOTO 1
IF(Q.LT.0)Q=0.0
U=0.5*((AM+1)*EO-(AM-1)*EI+AM*Q)/SQRT(EI*EO)
IF(U.LT.-1)GOTO 1
IF(U.GT.1)U=1.0
Z=U*U-1
F=1.0/AM
IF(Q.NE.0.0)F=F*SQRT(EI/(EI-(AM+1)/AM*Q))
BCM=F*(Z+U*SQRT(Z+AM*AM-AM*(AM+1)/EI*Q))
RETURN
1  BCM=-1
RETURN
END

```

```

FUNCTION ICSOP(E,ABN)
DIMENSION ABN(1)
J=0
1  J=J+1
IF(E.LE.ABN(J).AND.ABN(J).GT.0.)GOTO 1
ICSOP=J-1
RETURN
END

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DOUBLE PRECISION FUNCTION PTL(N,XX)
REAL*8 XX
GOTO (1,2,3,4,5,6),N
1  PTL=1.0
RETURN

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2  PTL=XX
RETURN
3  PTL=0.5*(3*(XX**2)-1.0)
RETURN
4  PTL=0.5*(5*(XX**3)-3*XX)
RETURN
5  PTL=0.125*(35*(XX**4)-30*(XX**2)+3)
RETURN
6  PTL=0.125*(63*(XX**5)-70*(XX**3)+15*XX)
RETURN
END

```

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SUBROUTINE ZWIN(KL,EP,EM,ES,J2,J1)
DIMENSION ES(KL)
J1=KL
J2=1
KLI=KL-1
DO 1 I=1,KLI
IF(EP.GE.ES(I).AND.EP.LT.ES(I+1))GOTO 2
GOTO 1
2  J1=I+1
GOTO 11
1  CONTINUE
11 CONTINUE
DO 12 I=2,KL
IF(EM.LE.ES(I).AND.EM.GT.ES(I-1))GOTO 13
GOTO 12
13  J2=I-1
GOTO 14
12  CONTINUE
14  CONTINUE
RETURN
END

```

```

SUBROUTINE MASSIN(WORK,ABN)
REAL*8 MATN
REAL*8 HID
DIMENSION WORK(1),ABN(1),A(6),AC(13)
DIMENSION FF(6),FA(6)
COMMON MATN,IS(2),NOUT,LI,IR(2)
DATA A/0.57735,0.774596,0.33998,0.86114,0.53847,0.90618/
DATA HID/'H H1'/
IF(MATN.EQ.HID)GOTO 40
CALL KEDDAT(MATN,'ISOT1 ',0,NW1,FNW,NDAT,FA,FF,300.0,1.0,
1&40,&40,&40)
AME=FA(2)
GOTO 41
40  CONTINUE

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41	AME=1.008665	150	FUNCTION FXINT(EP,EM,SGN,ES,FLUX,EFLUX,NSP,NDAT)	10
	CONTINUE	160	DIMENSION FLUX(1),EFLUX(1)	20
	AME=AME/1.008665	170	DIMENSION SGN(1),ES(1)	30
	DO 10 I=1,6	180	FINTP(XA,XB,XC,YA,YC)=YA+(YC-YA)/(XC-XA)*(XB-XA)	40
	XA=A(I)**2-1	190	FG(XA,XB,XC,YA,YC)=YA+(YC-YA)/ALOG(XC/XA)*ALOG(XB/XA)	50
	AC(I+6)=(XA-A(I))*SQRT(XA+AME**2)/AME	200	CALL ZWIN(NDAT,EP,EM,ES,J2,J1)	60
10	AC(I)=(XA+A(I))*SQRT(XA+AME**2)/AME	210	Y=0	70
	AC(13)=-1/AME	220	SGN1=SGN(J1)	80
	XM=(AME+1)**2/2/AME	230	SGN2=SGN(J2)	90
	QM=2.0E8	240	ES1=ES(J1)	100
	IF(XM.NE.2)QM=XM/(XM-2)	250	ES2=ES(J2)	110
	QMX=ALOG(QM)	260	SGN(J1)=FG(ES(J1),EP,ES(J1-1),SGN(J1),SGN(J1-1))	120
	WORK(28)=AME	270	SGN(J2)=FG(ES(J2),EM,ES(J2+1),SGN(J2),SGN(J2+1))	130
	WORK(29)=QM	280	ES(J1)=EP	140
	WORK(30)=XM	290	ES(J2)=EM	150
	WORK(31)=QMX	300	J11=J1-1	160
	CALL IROMES(AC)	310	IF(NSP.EQ.1)GOTO 1	170
	I1=IR(1)	320	CALL ZWIN(NSP,EP,EM,EFLUX,JF2,JF1)	180
	I2=IR(2)	330	FF1=FINTP(EFLUX(JF2),EM,EFLUX(JF2+1),FLUX(JF2),FLUX(JF2+1))	190
	IMAX=2	340	GOTO 2	200
	DO 11 I=I2,I1	350	1 FF1=PHI(EM)	210
	E=ABN(I+1)/QM	360	2 CONTINUE	220
	IM=ICSOP(E,ABN(I))	370	DO 16 I=J2,J11	230
	IF(IM.GT.IMAX)IMAX=IM	380	DX=ES(I+1)-ES(I)	240
11	CONTINUE	390	IF(NSP.EQ.1)GOTO 3	250
	CALL IWIN(18,IMAX,WORK)	400	CALL ZWIN(NSP,ES(I+1),ES(I+1),EFLUX,ISG,ISF)	260
	RETURN	410	FF2=FINTP(EFLUX(ISG),ES(I+1),EFLUX(ISG+1),FLUX(ISG),FLUX(ISG+1))	270
	END	420	GOTO 4	280
			CONTINUE	290
			FF2=PHI(ES(I+1))	300
			4 Y=Y+DX*(SGN(I)*FF1+SGN(I+1)*FF2)	310
			FF1=FF2	320
			16 CONTINUE	330
			FXINT=Y/2	340
			ES(J1)=ES1	350
			ES(J2)=ES2	360
			SGN(J1)=SGN1	370
			SGN(J2)=SGN2	380
			RETURN	390
			END	400
	SUBROUTINE INTEN(E,SG,SGNC,EA,ICOS)	10		
	DIMENSION SG(1),SGNC(ICOS,1),EA(1)	20		
	NE=0	30		
3	NE=NE+1	40		
	IF(E.GT.EA(NE)+EA(NE)*1.0E-4)GOTO 3	50		
	IF(NE.EQ.1)GOTO 5	60		
	A=ALOG(EA(NE)/EA(NE-1))	70		
	B=ALOG(E/EA(NE-1))	80		
	DO 4 I=1,ICOS	90		
4	SG(I)=(SGNC(I,NE)-SGNC(I,NE-1))/A+B+SGNC(I,NE-1)	100		
	RETURN	110		
5	DO 6 I=1,ICOS	120		
	SG(I)=0.5	130		
6	CONTINUE	140		
	RETURN	150		
	END	160		
			SUBROUTINE SMORN(PI,N,IM,WORK)	10
			DIMENSION PI(N,IM),WORK(1)	20
			Q=WORK(32)	30
			JM=1	40
			IF(Q.GT.0)JM=IWO(26,WORK)	50
			AS=0	60
			DO 1 I=JM,IM	70
1	AS=PI(1,I)+AS	80		
	IF(AS.EQ.0)RETURN	90		
	DO 2 I=JM,IM	100		
	DO 2 J=1,N	110		

2	PI(J,I)=PI(J,I)/AS	120	SUBROUTINE REMP(ABN,IGR,NFI,FLUX,EFLUX,WORK,IWORK,IS)	10
	RETURN	130	DIMENSION ABN(1),IGR(1),NFI(1),FLUX(1),EFLUX(1),WORK(1),IWORK(1)	20
	END	140	REAL*8 MATN	30
			REAL*8 FNEV(4)	40
	SUBROUTINE SEARCH(NFG,NFI,WORK,IWORK,FLUX,EFLUX,ABN)	10	COMMON MATN, IIZ(2),NOUT,LI,IL,IM	50
	DIMENSION WORK(1),IWORK(1),FLUX(1),EFLUX(1),ABN(1),SGS(4)	20	DATA FNEV/'SGNC ','SGT ','SGN ','MUEL '/	60
	REAL*8 MATN	30	NO=IWORK(1)	70
	COMMON MATN,MAY(2),NOUT,LI	40	N2=IWORK(3)	80
	DATA SGS/'SGT ','SGN ','MUEL', 'FLUX'/	50	N4=IWORK(5)	90
	NDAT=IWORK(25)	60	N5=IWORK(6)	100
	N9=IWORK(13)	70	N6=IWORK(7)	110
	DO 3 I=1,NDAT	80	NE=IWORK(26)	120
3	WORK(N9+I-1)=1.0	90	NLA=IWORK(16)	130
	N4=IWORK(7)	100	NLE=IWORK(17)	140
	NZ=NFG*NFI	110	N11=IWORK(12)	150
	N11=IWORK(12)	120	IMAX=IWORK(18)	160
	IW=N11+3*NZ	130	AM=WORK(28)	170
	IG=IWORK(4)	140	DO 10 IGI=IM,IL	180
	NSP=IWORK(15)	150	IG=IM+IL-IGI	190
	DU=ALOG(ABN(1)/ABN(2))/NFG	160	IWORK(4)=IG	200
	E1=ABN(2)	170	IF(IG-IS)20,20,21	210
	DO 20 I=1,NFG	180	20 CONTINUE	220
	E2=E1*EXP(DU)	190	IWORK(23)=1	230
	DE=ALOG(E2/E1)/NFI	200	CALL ZWIN(NE,ABN(IG),ABN(IG+1),WORK(NO),JW2,JW1)	240
	FF=FXINT(E2,E1,WORK(N9),WORK(N4),FLUX,EFLUX,NSP,NDAT)	210	CALL NORM(JW2,JW1,WORK(N2),WORK)	250
	EM=E1	220	GOTO 22	260
	DO 21 J=1,NFI	230	21 IWORK(23)=0	270
	EP=EM*EXP(DE)	240	22 CONTINUE	280
	WORK(IW)=FXINT(EP,EM,WORK(N9),WORK(N4),FLUX,EFLUX,NSP,NDAT)/FF	250	IWORK(22)=N11	290
	IW=IW+1	260	JN=2	300
21	EM=EP	270	CALL KEDDAT(MATN,FNEV(JN),NW,NW1,FNW,NDAT,WORK(N6),WORK(N4),	310
20	E1=E2	280	1ABN(IG),ABN(IG+1),&2,&3,&4)	320
	IWORK(22)=N11	290	IWORK(25)=NDAT	330
	NQ=4*NZ+3	300	CALL SINT(ABN(IG),IGR(IG),NFI(IG),FLUX,EFLUX,WORK,N4,N6)	340
	J=4*NZ	310	JN=3	350
	WRITE(LI)NQ,IG,NFG,NFI,(WORK(N11+I-1),I=1,J)	320	CALL KEDDAT(MATN,FNEV(JN),NW,NW1,FNW,NDAT,WORK(N6),WORK(N4),	360
	WRITE(NOUT,70)IG,NZ,MATN	330	1ABN(IG),ABN(IG+1),&2,&3,&4)	370
	DO 1 M=1,4	340	JN=4	380
	WRITE(NOUT,71)SGS(M)	350	CALL KEDDAT(MATN,FNEV(JN),NW,NW1,FNW,NDAT,WORK(N6),WORK(N5),	390
	IW=N11+(M-1)*NZ-1	360	1ABN(IG),ABN(IG+1),&5,&3,&4)	400
1	WRITE(NOUT,72)(WORK(IW+I),I=1,NZ)	370	7 CONTINUE	410
70	FORMAT(/'_GROUP=' ,I3,' NUMBER OF FINE INTERVALS=' ,I5,	380	IWORK(25)=NDAT	420
	1' MATERIAL',2X,A8)	390	N4=-N4	430
71	FORMAT(' THE VALUES OF ',A4)	400	CALL SINT(ABN(IG),IGR(IG),NFI(IG),FLUX,EFLUX,WORK,N4,N6)	440
72	FORMAT(' ',1P6E12.4)	410	JN=4	450
	RETURN	420	CALL KEDDAT(MATN,FNEV(JN),NW,NW1,FNW,NDAT,WORK(N6),WORK(N5),	460
	END	430	1ABN(IG),ABN(IG+1),&6,&3,&4)	470
			8 CONTINUE	480
			CALL SEARCH(IGR(IG),NFI(IG),WORK,IWORK,FLUX,EFLUX,ABN(IG))	490
			CALL ISOFAL(ABN(IG),IGR(IG),NFI(IG),FLUX,EFLUX,WORK)	500
			DO 16 IO=NLA,NLE	510
			16 CALL SUCHM(IGR(IG),NFI(IG),WORK,IWORK,IO,ABN)	520
			GOTO 10	530
			5 ASSIGN 7 TO NNN	540
			GOTO 9	550
			6 ASSIGN 8 TO NNN	560

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9   W=2.0/3.0/AM
DO 15 J=1,NDAT
15  WORK(N5+J-1)=W
GOTO NNN,(7,8)
10  CONTINUE
RETURN
2   WRITE(NOUT,80)MATN
RETURN
3   WRITE(NOUT,90)MATN,FNEV(JN)
RETURN
4   WRITE(NOUT,91)MATN,FNEV(JN)
RETURN
80  FORMAT(' ***ERROR 9.2 : MATERIAL ',A8,' IS NOT FOUND')
90  FORMAT(' ***ERROR 9.3 MATERIAL ',A8,' TYP ',A8,
1' ARE NOT FOUND')
91  FORMAT(' ***ERROR 9.4 MATERIAL ',A8,' TYP ',A8,
1' ARE NOT FOUND IN LIST')
END

SUBROUTINE NORM(J2,J1,SC,WORK)
REAL*8 ANINT
DIMENSION SC(1),WORK(1)
NN=IWO(33,WORK)
NJM=IWO(19,WORK)
NMU=IWO(17,WORK)
KJJ=1
N5=IWO(2,WORK)
ICOS=IWO(21,WORK)
BL=WORK(N5)
BU=WORK(N5+ICOS-1)
NJ=0
CALL IWIN(24,NJ,WORK)
DO 1 I=J2,J1
IW=(I-1)*ICOS
ASG=ANINT(BU,BL,1,SC(IW+1),WORK(N5),0.0,ICOS,WORK,WORK(NN),NJM)
NJ=IWO(24,WORK)
IF(NJ.GE.NJM-1)KJJ=2
CALL IWIN(23,KJJ,WORK)
NJ=0
CALL IWIN(24,NJ,WORK)
DO 2 J=1,ICOS
2   SC(IW+J)=SC(IW+J)/ASG
1   CONTINUE
RETURN
END

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570 SUBROUTINE SINT(ABN,NG,NI,FLUX,EFLUX,WORK,N,N4)
580 DIMENSION ABN(1),FLUX(1),EFLUX(1),WORK(1)
590 N2=IWO(5,WORK)
600 N3=IWO(6,WORK)
610 N7=IWO(22,WORK)
620 N9=IWO(13,WORK)
630 NSP=IWO(15,WORK)
640 NDAT=IWO(25,WORK)
650 IG=1
660 L=N
670 N=IABS(N)
680 N2=IABS(N2)
690 NZ=NI*NG
700 DO 3 I=1,NDAT
710 IF(L.LT.0)WORK(N3+I-1)=WORK(N3+I-1)*WORK(N2+I-1)
3   WORK(N9+I-1)=1.0
DU=ALOG(ABN(IG)/ABN(IG+1))/NG
E1=ABN(IG+1)
IW=N7
DO 1 IP=1,NG
E2=E1*EXP(DU)
DDU=ALOG(E2/E1)/NI
EM=E1
DO 2 IR=1,NI
EP=EM*EXP(DDU)
WORK(IW)=FXINT(EP,EM,WORK(N),WORK(N4),FLUX,EFLUX,NSP,NDAT)
FF=FXINT(EP,EM,WORK(N9),WORK(N4),FLUX,EFLUX,NSP,NDAT)
IF(L.GT.0)GOTO 4
WORK(IW+NZ)=FXINT(EP,EM,WORK(N3),WORK(N4),FLUX,EFLUX,NSP,NDAT)
1/WORK(IW)
4   WORK(IW)=WORK(IW)/FF
EM=EP
2   IW=IW+1
E1=E2
1   CONTINUE
IF(L.LT.0)IW=IW+NZ
CALL IWIN(22,IW,WORK)
RETURN
END

SUBROUTINE ISOFAL(ABN,IGR2,NFI1,FLUX,EFLUX,WORK)
REAL*8 STOFF
REAL*8 DU,DDU
DIMENSION ABN(1),FLUX(1),EFLUX(1),WORK(1)
IG=1
N9=IWO(10,WORK)
QQ=WORK(32)
IC=IG
IGI=IWO(4,WORK)
IF(QQ.GT.0)CALL EGRENZ(ABN(1),ABN,IC,IV,WORK)
IFI=IGI+IG-1
CALL IWIN(4,IFI,WORK)

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VI 70


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NM1=IWO(17,WORK)+1
NJM=IWO(19,WORK)
NUJM=IWO(20,WORK)
AME=WORK(28)
IMAX=IWO(18,WORK)
IX=IWO(11,WORK)
N7=IWO(22,WORK)
DU=ABN(IG)/ABN(IG+1)
DU=DLOG(DU)/IGR2
2  IG1=IG+1
E1=ABN(IG+1)
FNW=0.
IF(QQ.GT.0.)FNW=QQ*AME/(AME-1.)
NU=NM1*IMAX
DO 15 IP=1,IGR2
E2=E1*DEXP(DU)
DDU=DU/NF11
EM=E1
DO 5 IR=1,NF11
EP=EM*DEXP(DDU)
DO 9I=1,NU
9  WORK(N7+I-1)=0.0
IF(EP.LE.FNW)GOTO 7
IF(QQ.LE.0)GOTO 6
IF(EM.LE.FNW)EM=FNW
6  CALL WAHRS(EP,EM,IMAX,WORK(N7),ABN(IC),NM1,NUJM,FLUX,EFLUX,
1WORK(N9),WORK,WORK(IX))
CALL SMORN(WORK(N7),NM1,IMAX,WORK)
7  CONTINUE
N7=N7+NM1*IMAX
5  EM=EP
15 E1=E2
CALL IWIN(22,N7,WORK)
RETURN
END

SUBROUTINE WAHRS(EP,EM,IMAX,PI,ABN,NM1,NUJM,FLUX,EFLUX,TE,WORK,EZ)
DIMENSION FLUX(1),EFLUX(1),WORK(1)
DIMENSION TE(NM1,IMAX,NUJM)
DIMENSION ABN(IMAX),PI(NM1,IMAX)
DIMENSION EZ(1)
COMMON MATN(4),NOUT
QM=WORK(29)
NN=IWO(34,WORK)
IG=1
AM=WORK(28)
Q=WORK(32)
KL=IWO(25,WORK)
N2=IWO(5,WORK)
KJ=IWO(23,WORK)
N3=IWO(6,WORK)
N4=IWO(7,WORK)

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N10=IWO(11,WORK)
N10=N10+ 1+IMAX
EZ(1)=EM
IJ=1
JM=1
E=EM/QM
IMAX1=ICSOP(E,ABN)
IF(E.EQ.ABN(IMAX1))IMAX1=IMAX1-1
IF(0.GT.0)CALL EGRENZ(EM,ABN,IK,IMAX1,WORK)
IMAX1=IMAX1+1
IG=IWO(4,WORK)
NG=IWO(14,WORK)
IT=NG-IG
IT=MINO(IMAX,IT)
IF(0.GT.0)GOTO 45
DO 42 I=1,IT
EP1=QM*ABN(IT-I+2)
IF(EP1.GE.EP.OR.EP1.LE.EM)GOTO 42
IJ=IJ+1
EZ(IJ)=EP1
42 CONTINUE
45 CONTINUE
DO 44 LJ=1,IJ
EZ(IJ+1)=EP
IMAX1=IMAX1-1
CALL ZWIN(KL,EZ(LJ+1),EZ(LJ),WORK(N4),J2,J1)
J11=J1-1
EM1=EZ(LJ)
DO 44 IK=J2,J11
EP1=WORK(N4+IK)
IF(1K.F0.J11)EP1=EZ(LJ+1)
CALL LMI(EP1,EM1,WORK(N4+IK),WORK(N4+IK-1),WORK(N2+IK),
1WORK(N2+IK-1),WORK(N3+IK),WORK(N3+IK-1),IMAX1,ABN,TE,NM1,NUJ,NUJM,
2FLUX,EFLUX,WORK,WORK(N10),IMAX,WORK(NN))
IF(0.GT.0)JM=IWO(26,WORK)
IF(IMAX1.GT.1.OR.KJ.GT.0)GOTO 3
IF(NM1.LT.3)GOTO 3
TE(3,JM,NUJ)=TE(1,JM,NUJ)/AM/AM/5
IF(NM1.LT.5)GOTO 3
TE(5,JM,NUJ)=-TE(1,JM,NUJ)/AM/AM/AM/AM/63
3  CONTINUE
NJ=IWO(24,WORK)
IF(NJ.GT.0)WRITE(NOUT,70)EP1,EM1
70  FORMAT(' ***WARNING 9.2 : FEW MESH POINTS IN ANGLE INTEGRATION',
11P2E12.4)
NJ=0
CALL IWIN(24,NJ,WORK)
EM1=EP1
DO 43 J=JM,IMAX1
DO 43 K=1,NM1
43  PI(K,J)=PI(K,J)+TE(K,J,NUJ)
44  CONTINUE
RETURN
END

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SUBROUTINE LMI (EP,EM,E1,E2,SG1,SG2,AMU1,AMU2, IMAX,ABN,TE,
INM1,NUJ,NUJM,FLUX,EFLUX,WORK,PER,IMAM,PL)
REAL*8 PLK,PL(NM1,IMAM)
REAL*8 ANINT
DIMENSION ABN(IMAX),WORK(1)
DIMENSION EFLUX(1),FLUX(1)
DIMENSION PER(NM1,IMAM),TE(NM1,IMAM,NUJM)
COMMON MATN(4),NOUT
EQUIVALENCE(JM,JMI)
FG(XA,XB,XC,YA,YC)=YA+(YC-YA)/ALOG(XC/XA)*ALOG(XB/XA)
NJM=IWO(19,WORK)
AM=WORK(28)
QM=WORK(29)
KJ=IWO(23,WORK)
XM=WORK(30)
Q=WORK(32)
NSP=IWO(15,WORK)
ERR=WORK(27)
NLA=IWO(16,WORK)
NLE=IWO(17,WORK)
JM=1
IF(Q.GT.0)FNW=Q/AM/(AM+1.)
IF(Q.GT.0)CALL EGRENZ(EP,ABN,JM,II,WORK)
ICOS=IWO(21,WORK)
NNE=IWO(26,WORK)
N1=IWO(3,WORK)
N5=IWO(2,WORK)
NR=IWO(33,WORK)
N8=IWO(9,WORK)
NG=IWO(14,WORK)
N6=IWO(4,WORK)
NL1=1
NL2=NLE+1
IF(NSP.GT.1)CALL ZWIN(NSP,EP,EM,EFLUX,JF2,JF1)
IF(Q.GT.0)GOTO 9
NO=IWO(1,WORK)
IF(KJ.EQ.0)GOTO 9
CALL ZWIN(NNE,EP,EM,WORK(NO),JW2,JW1)
JW=JW2+NO-1
IW=(JW2-1)*ICOS+N1
9 CONTINUE
DO 31 I=1,NM1
DO 31 J=1,IMAX
PER(I,J)=1.0
PL(I,J)=0
31 CONTINUE
DU =EP-EM
M1=1
NUJ=0
M2=1
ANS=1.0
SE=1
ISEL=IWO(8,WORK)
JZZ=0
41 NU1=2**NUJ+1
DX=DU/(NU1-1)

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10 DO 21 I=M1,NU1,M2
20 E=EM+(I-1)*DX
30 IF(Q.GT.0)WORK(8)=E
40 IF(ISEL.EQ.1)SE=FG(E1,E,E2,SG1,SG2)
50 ASE=FG(E1,E,E2,AMU1,AMU2)
60 AZ=(ASE-2.0/3.0/AM)*3/(1.0-3.0/5.0/(AM**2))
70 IF(NSP.GT.1)GOTO 2
80 FF=PHI(E)
90 GOTO 3
2 DO 22 J=JF2,JF1
100 ISG=J-1
110 IF(EFLUX(J).GT.F)GOTO 23
120 CONTINUE
130 22 CONTINUE
140 23 CONTINUE
150 FF=(FLUX(ISG+1)-FLUX(ISG))/(EFLUX(ISG+1)-EFLUX(ISG))*(E-
160 1EFLUX(ISG))+FLUX(ISG)
170 3 CONTINUE
180 ANS=0.0
190 BU=1.0
200 J=JM-1
210 IF(KJ.EQ.0)GOTO 62
220 CALL INTEN(E,WORK(N8),WORK(IW),WORK(JW),ICOS)
230 BL=-1.0
240 ANS=ANINT(BU,BL,2,WORK(N8),WORK(N5),AZ,ICOS,WORK,WORK(NR),NJM)
250 ANS=ASE-ANS
62 J=J+1
270 BL=-1.0
280 IF(J.GE.NG-N6+1)GOTO 30
290 IF(IMAX.EQ.1)GOTO 30
300 BL=BCM(F,ABN(J+1),WORK)
310 IF(BL.LE.-1.0)BL=-1.0
320 30 CONTINUE
330 DO 8 K=NL1,NL2
340 IF(PER(K,J).LT.ERR)GOTO 8
350 IF(BL.EQ.BU)GOTO 4
360 IF(BL.LT.BU)GOTO 6
370 JZZ=1
380 BL1=BL
390 BU1=BU
400 BL=BU
410 BU=BL1
420 6 CONTINUE
430 IF(BL.GE.1)GOTO 4
440 PLK=ANINT(BU,BL,K,WORK(N8),WORK(N5),AZ,ICOS,WORK,WORK(NR),NJM)
450 PLK=PLK+AKOR(BU,BL,K,AM)*ANS
460 GOTO 5
470 4 PLK=0.0
480 5 CONTINUE
490 IF(JZZ.EQ.0)GOTO 66
500 JZZ=0
510 BL=BL1
520 BU=BU1
530 66 CONTINUE
540 PLK=PLK*FF
550 IF(I.EQ.1)PL(K,J)=0.5*PLK*DX*SE
560 IF(I.EQ.NU1)PL(K,J)=PL(K,J)+0.5*PLK*DX*SE
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VI 72

	IF(I.EQ.1.OR.I.EQ.NU1)GOTO 8	1130		1690
	PL(K,J)=PL(K,J)+PLK*DX*SE	1140	10 CONTINUE	
8	CONTINUE	1150	70 FORMAT(10X,'SCATTERING MATRIX_',2I4, 3X,1PE12.4, ' ERROR ',	1700
	BU=BL	1160	11PE12.4)	1710
	IF(BU.EQ.-1.0)GOTO 20	1170	71 FORMAT(' ***WARNING 9.3 : FEW ENERGY_MESH POINTS BETWEEN '	1720
	GOTO 62	1180	11P2E12.4,' EV')	1730
20	CONTINUE	1190	CALL IWIN(8,ISEL,WORK)	1740
21	CONTINUE	1200	RETURN	1750
	IF(NU1.GT.2)GOTO 27	1210	END	1760
	DO 50 J=JM,IMAX	1220		
	DO 50 K=NL1,NL2	1230		
	TE(K,J,1)=PL(K,J)	1240		
50	PL(K,J)=0	1250	FUNCTION IWO(L,IWORK)	10
	NUJ=NUJ+1	1260	DIMENSION IWORK(1)	20
	M1=2	1270	IWO=IWORK(L)	30
	M2=2	1280	RETURN	40
	IF(IMAX.GT.JM.OR.KJ.GT.0)GOTO 41	1290	END	50
	IF(NL2.LT.3)GOTO 41	1300		
	DO 25 J=3,NL2	1310		
	PER(J,JM)=0.	1320		
	IF(J.NE.5)TE(J,JM,1)=0.	1330		
25	CONTINUE	1340	SUBROUTINE IWIN(L,N,IWORK)	10
	GOTO 41	1350	DIMENSION IWORK(1)	20
27	LX=0	1360	IWORK(L)=N	30
	DO 26 J=JM,IMAX	1370	RETURN	40
	DO 7 K=NL1,NL2	1380	END	50
	IF(PER(K,J).LT.ERR)TE(K,J,NUJ+1)=TE(K,J,NUJ)	1390		
	IF(PER(K,J).LT.ERR)GOTO 7	1400		
	AJ=0.5*TE(K,J,1)+PL(K,J)	1410		
	AL=TE(K,J,NUJ)	1420		
	ZC=1	1430	SUBROUTINE SUCHM(NFG,NFI,IWORK,WORK,MOM,ABN)	10
	DO 52 IB=1,NUJ	1440	DIMENSION ABN(1)	20
	ZC=4*ZC	1450	DIMENSION WORK(1),IWORK(1)	30
	AP=(ZC*AJ-TE(K,J,IB))/(ZC-1)	1460	REAL*8 MATN	40
	TE(K,J,IB)=AJ	1470	COMMON MATN,MAY(2),NOUT,LI ,IL,IMI	50
52	AJ=AP	1480	IG=IWORK(4)	60
	TE(K,J,NUJ+1)=AP	1490	Q=WORK(32)	70
	PER(K,J)=ABS(AP)	1500	MK=2	80
	IF(ABS(AP).GT.1.E-10)PER(K,J)=ABS((AP-AL)/AP)	1510	IK=IG	90
	IF(PER(K,J).GT.ERR)LX=1	1520	IF(Q.LT.0)GOTO 15	100
	PL(K,J)=0	1530	CALL EGRENZ(ABN(IG),ABN,IK,IV,WORK)	110
7	CONTINUE	1540	AM=WORK(28)	120
26	CONTINUE	1550	FNW=Q*AM/(AM-1.)	130
	NUJ=NUJ+1	1560	EE=AMAX1(FNW,ABN(IG+1))	140
	IF(NUJ.GE.NUJM)GOTO 100	1570	CALL EGRENZ(EE,ABN,I,IV,WORK)	150
	IF(LX.EQ.1)GOTO 41	1580	MK=4	160
	IF(Q.GT.0)CALL IWIN(26,JMI,WORK)	1590	15 CONTINUE	170
	CALL IWIN(8,ISEL,WORK)	1600	N9=IWORK(13)	180
	RETURN	1610	N11=IWORK(12)	190
100	CONTINUE	1620	NLE=IWORK(17)	200
	IF(Q.GT.0)CALL IWIN(26,JMI,WORK)	1630	IM=IWORK(18)	210
	WRITE(NOUT,71)EP,EM	1640	IW=N11	220
	DO 10 J=JM,IMAX	1650	NZ=NFG*NFI	230
	DO 10 K=1,NM1	1660		
	IF(PER(K,J).LE.ERR)GOTO 10	1670		
	WRITE(NOUT,70)K,J,TE(K,J,NUJ),PER(K,J)	1680		

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DO 11 J=1,NZ
DO 12 K=1,IM
L=(K-1)*NZ+MK+N9
LL=IW+(K-1)*(NLE+1)+MOM
WORK(L+J)=WORK(LL)
12 CONTINUE
11 IW=IW+(NLE+1)*IM
IWORK(N9)=IM*NZ+MK
L=IWORK(N9)+N9
IWORK(N9+1)=MOM
IWORK(N9+2)=IG
IF(Q.LT.0)GOTO 5
IWORK(N9+3)=IK
IWORK(N9+4)=IV
5 CONTINUE
WRITE(LL)(WORK(I),I=N9,L)
NGR=IWORK(14)
WRITE(NOUT,70)MOM,MATN,IG
DO 2 I=1,IM
IZ=IK+I-1
IF(IZ.GT.NGR)RETURN
WRITE(NOUT,71)IG,IZ
IW=N9+MK+(I-1)*NZ
2 WRITE(NOUT,72)(WORK(IW+J),J=1,NZ)
70 FORMAT(/', ELASTIC SCATTERING MATRIX SGNC',I1,' FOR ',A8,
1' GROUP=',I4)
71 FORMAT(' FROM GROUP ',I4,' TO GROUP ',I4)
72 FORMAT(' ',1P6E12.4)
RETURN
END

SUBROUTINE AMESH(B1,B2,Z,J2,NX,NL)
DIMENSION Z(6),ZA(6),ZF(6)
DIMENSION A(13),AC(1)
DATA NW/10/
GOTO (1,2,3,4,5,6),NL
1 NX=1
J1=0
J2=1
Z(1)=B1
RETURN
2 N=1
Z(1)=A(13)
GOTO 7
3 N=2
Z(1)=A(7)
Z(2)=A(1)
GOTO 7
4 N=3
Z(1)=A(8)
Z(2)=A(13)
Z(3)=A(2)

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GOTO 7
5 N=4
Z(1)=A(10)
Z(2)=A(9)
Z(3)=A(3)
Z(4)=A(4)
GOTO 7
6 Z(1)=A(12)
Z(2)=A(11)
Z(3)=A(13)
Z(4)=A(5)
Z(5)=A(6)
N=5
7 CONTINUE
IF(AC(13).GT.-1+1.0/NW)GOTO 11
M=N
N=N/2
IF(N.EQ.0)GOTO 1
DO 12 I=1,N
Z(I)=Z(M-N+I)
11 CONTINUE
J1=N
J2=1
ZA(J2)=Z(J2)-(Z(J2)+1)/NW
ZF(J1)=Z(J1)+(1-Z(J1))/NW
IF(J1.EQ.1)GOTO 9
J1=J1-1
DO 8 I=J2,J1
SW=Z(I+1)-Z(I)
SW=SW/NW
ZA(I+1)=Z(I+1)-SW
ZF(I)=Z(I)+SW
8 CONTINUE
9 CONTINUE
DO 10 I=1,N
IF(ZA(I).LE.B2)J2=J2+1
IF(ZF(I).GE.B1)J1=J1-1
10 CONTINUE
NX=J1-J2+2
Z(J1+1)=B1
RETURN
ENTRY IRDMES(AC)
DO 20 I=1,13
A(I)=AC(I)
RETURN
END

DOUBLE PRECISION FUNCTION ANINT(BU,BL,K,SG,XL,AZ,ICOS,WORK,T,NJM)
REAL*8 YA,YB
REAL *8 MATN
REAL*8 T(NJM)
REAL*8 ANG,XX,GAM

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VI 74

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	REAL *8	XA,XB,XA1,YA1,AJ,AI,AK,AL,AP,AXX	60	DM=(XB-XA)/(NP-1)	620
	REAL*8	PTL	70	AJ=0.5*AI	630
	DIMENSION	SG(1),XL(1),Z(7),WORK(1)	80	DO 40 IZ=2,NP,2	640
	COMMON	MAT(4),NOUT	90	XA=XA+DM	650
	EQUIVALENCE	(MATN,MAT(1))	100	YA=ANG(XA)	660
C	DATA	NHD/'H H1'/	110	YA=PTL(K,YA)	670
	ANG(XX)	=(XX+GAM)/DSQRT(GAM**2+2*XX*GAM+1.0)	120	YA=YA*WINK(XL,SG,KJ,ICOS,AZ,XA)	680
	ERR	=WORK(27)	130	XA=XA+DM	690
	AM	=WORK(28)	140	40 AJ=AJ+YA*DM	700
	GAM	=1/AM	150	AK=AJ	710
	Q	=WORK(32)	160	AL=T(NJ)	720
	IF(Q.LT.0)	GOTO 56	170	ZC=1	730
	EI	=WORK(8)	180	DO 51 IB=1,NJ	740
	GAM	=GAM*SQRT(EI/(EI-(AM+1)/AM*Q))	190	ZC=ZC*4	750
56	CONTINUE		200	AP=(ZC*AK-T(IB))/(ZC-1)	760
	IF(AM.LT.1.1)	CALL HIDR(BU,BL,K,AXX,&50)	210	T(IB)=AK	770
	IF(AM.LT.1.1)	GOTO 50	220	51 AK=AP	780
	XM	=WORK(30)	230	T(NJ+1)=AP	790
	KJ	=IWD(23,WORK)	240	EPS=(AP-AL)/AP	800
	CALL	AMESH(BU,BL,Z,L2,NX,K)	250	EPS=10.0*EPS	810
	NU	=1	260	IF(NJ.GE.NJM-1)	820
	XA	=BL	270	IF(ABS(EPS).GT.ERR)	830
	YA	=ANG(XA)	280	GOTO 29	840
	YA	=PTL(K,YA)	290	28 CONTINUE	850
	YA	=YA*WINK(XL,SG,KJ,ICOS,AZ,XA)	300	CALL IWIN(24,NJ,WORK)	860
	AXX	=0	310	29 CONTINUE	870
	IL	=1	320	AXX=AXX+T(NJ+1)	880
18	CONTINUE		330	XA=XB	890
	XA1	=XA	340	YA=YB	900
	XB	=Z(IL+L2-1)	350	IL=IL+1	910
	IF(KJ.NE.2)	GOTO 20	360	IF(IL.LE.NX)	920
	DO 21	I=1,ICOS	370	50 CONTINUE	930
	J	=1	380	ANINT=AXX	940
	IF(XL(I).GT.XA+ERR.AND.XL(I).LT.XB-ERR)	GOTO 22	390	RETURN	950
21	CONTINUE		400	END	960
	GOTO	20	410		
22	XB	=XL(J)	420		
	IL	=IL-1	430		
20	CONTINUE		440		
	YB	=ANG(XB)	450	FUNCTION WINK(XL,SG,KJ,ICOS,AZ,AM)	10
	IF(YB.EQ.0.AND.(K-2*(K/2)).EQ.0)	GOTO 57	460	DIMENSION XL(1),SG(1)	20
	YB	=PTL(K,YB)	470	COMMON M(4),NOUT	30
	GOTO	58	480	IF(KJ)100,101,100	40
57	YB	=0.0	490	101 WINK=0.5*(1+AZ*AM)	50
58	CONTINUE		500	RETURN	60
	YB	=YB*WINK(XL,SG,KJ,ICOS,AZ,XB)	510	100 DO 1 I=2,ICOS	70
	NJ	=0	520	IF(AM-XL(I))2,2,1	80
	AJ	=0.5*(YA+YB)*(XB-XA)	530	2 WINK=(SG(I)-SG(I-1))/(XL(I)-XL(I-1))*(AM-XL(I-1))+SG(I-1)	90
	XA1	=XA	540	RETURN	100
	YA1	=YA	550	1 CONTINUE	110
	T(1)	=AJ	560	WRITE(NOUT,70)AM	120
10	NJ	=NJ+1	570	WINK=0.5	130
	XA	=XA1	580	70 FORMAT(' ***WARNING 9.1 : THE ANGLE ',E11.4,' IS OUTSIDE')	140
	YA	=YA1	590	RETURN	150
	AI	=AJ	600	END	160
	NP	=2**NJ+1	610		

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SUBROUTINE HIDR(BU,BL,K,A,*)
REAL*8 A,X1,X2,B1,B2
B1=BU
B2=BL
X2=DSQRT(0.5*(1.0+B1))
X1=DSQRT(0.5*(1.0+B2))
GOTO (1,2,3,4,5,6),K
1 A=X2**2-X1**2
RETURN 1
2 A=2.0/3.0*(X2**3-X1**3)
RETURN 1
3 A=3.0/4.0*(X2**4-X1**4)-0.5*(X2**2-X1**2)
RETURN 1
4 A=0.5*(X2**5-X2**3-X1**5+X1**3)
RETURN 1
5 A=0.125*(35.0/6.0*(X2**6-X1**6)-7.5*(X2**4-X1**4)+1.5*(X2**2-
1X1**2))
RETURN 1
6 A=0.125*(9*(X2**7-X1**7)-14*(X2**5-X1**5)+5.*(X2**3-X1**3))
RETURN 1
END

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FUNCTION AKOR(B1,BL,K,AM)
ANG(X)=(X+GAM)/SQRT(GAM**2+2*X*GAM+1.)
Q0(X)=0.75*X**X
Q1(X)=0.5*(X**3)
Q2(X)=3./16.*X**X*(3*X*X-2)
Q3(X)=0.75*(X**3)*(X*X-1)
Q4(X)=1./32.*((X*X-1.)*(35.*X*X-10)-1)*X**X
Q5(X)=3./16.*(X**3)*(X*X-1.)*(9*X*X-5)
GAM=1./AM
IF(GAM.LT.0.99)GOTO 7
A1=SQRT((B1+1.)/2)
A2=SQRT((BL+1.)/2.)
GOTO 8
7 CONTINUE
A1=ANG(B1)
A2=ANG(BL)
8 CONTINUE
GOTO(1,2,3,4,5,6),K
1 AKOR=Q0(A1)-Q0(A2)
RETURN
2 AKOR=Q1(A1)-Q1(A2)
RETURN
3 AKOR=Q2(A1)-Q2(A2)
RETURN
4 AKOR=Q3(A1)-Q3(A2)

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RETURN
5 AKOR=Q4(A1)-Q4(A2)
RETURN
6 AKOR=Q5(A1)-Q5(A2)
RETURN
END

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SUBROUTINE EGRENZ(E,ABN,IK,IV,WORK)
RETURN
END

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C BERECHNUNG DES QUERSCHNITTS IN DER THERMISCHEN GRUPPE
C
SUBROUTINE THERM(NE,NTY,TYP)
REAL*8 MAT,TYP(NTY),FEST(5),A,B,C,D,P,F,G,H,Z,X
INTEGER*2 IHC(2),IMP(2)
DIMENSION N(4),E(2),S(2)
COMMON MAT,ISTRUK,ISPA,NOUT,LIZ,NANF,NEND,KL
EQUIVALENCE(IHC(1),HM)
DATA IMP,'PU','U'/'
WRITE(NOUT,9000)
9000 FORMAT(1H0/1H0/' PROGRAMM KENNZIFFER 10')
WRITE(NOUT,9001)
9001 FORMAT(' PROGRAMM ZUR BERECHNUNG DER THERMISCHEN QUERSCHNITTE'/)
N(1)=3
CALL DOPW(8HBEST,FEST(2))
CALL DOPW(8HTHERM,A)
CALL DOPW(8HSGF,B)
CALL DOPW(8HSGA,C)
CALL DOPW(8HSGC,D)
CALL DOPW(8HALPHA,P)
CALL DOPW(8HETA,F)
CALL DOPW(8HNUE,G)
CALL DOPW(8HSGG,H)
CALL DOPW(8HMUEL,Z)
HM=MAT CALL DOPW(8HSGN,X)
INR=0
SC=0.
LL=0
IF(IHC(1).EQ.IMP(1)) GO TO 802
IF(IHC(1).NE.IMP(2)) GO TO 822
802 LAR=1
DO 803 JJ=1,NTY
IF(TYP(JJ).EQ.C) GO TO 805
IF(TYP(JJ).NE.B) GO TO 803
805 LL=LL+1
IF(LL-2)803,801,803
803 CONTINUE

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      IF(LL-2)807,801,807
807 WRITE(NOUT,808) MAT
808 FORMAT(1H0/' ***WARNING 10. 1 : THE GROUP CROSS SECTION SGC CAN NO
IT BE CALCULATED FOR ',A6/' BECAUSE THE REACTION TYPES SGF AND SGA
ZARE NOT SPECIFIED IN THE INPUT')
      GO TO 801
822 LAR=0
      DO 821 JJ=1,NTY
      IF(TYP (JJ).EQ.C) GO TO 801
821 CONTINUE
      GO TO 807
801 I=0
      WRITE (LIZ)I,A
      FEST(1)=MAT
      DO 1 I=1,NTY
      KSIK=0
      IF(TYP(I).NE.P.AND.TYP(I).NE.F) GO TO 871
      IF(LAR.EQ.C) GO TO 875
      KSIK=1
      FEST(3)=H
      GO TO 872
875 WRITE(NOUT,876) MAT,TYP(I)
876 FORMAT(1H0/' ***WARNING 10. 2 : THE GROUP CROSS SECTION FOR ',2A8,
1' CAN NOT BE CALCULATED BECAUSE THE VALUE OF SGF IS ZERO')
      GO TO 1
874 KSIK=2
      FEST(3)=B
      GO TO 872
879 KSIK=3
      FEST(3)=G
      GO TO 872
871 FEST(3)=TYP(I)
872 CALL NDFLOC (J,N,FEST,K,K)
      IF(J)2,2,3
      2 WRITE (NOUT,4) TYP(I),MAT
      4 FORMAT(1H0/' ***WARNING 10. 3 : THE CROSS SECTION TYPE ',A9,' IS N
10T AVAILABLE IN THE KEDAK LIBRARY FOR MATERIAL ',A9)
      T=0.
      GO TO 890
      3 E(1)=FEST(4)
      S(1)=FEST(5)
      IF(E(1)-0.0253)24,5,6
      6 WRITE (NOUT,7) TYP(I),MAT
      7 FORMAT(1H0/' ***WARNING 10. 4 : IN THE KEDAK LIBRARY IS NO ENERGY
1POINT LESS THAN OR EQUAL TO 0.0253/' AVAILABLE FOR THE CROSS SECT
2ION TYPE ',A9,' FOR MATERIAL ',A9)
      GO TO 1
      24 CALL NDFNXT(J,N,FEST,K,K)
      IF(J)2,2,8
      8 E(2)=FEST(4)
      S(2)=FEST(5)
      IF(E(2)-0.0253)24,10,11
10 E(1)=E(2)
      S(1)=S(2)
      5 IF(TYP(I).EQ.G.OR.TYP(I).EQ.Z) GO TO 30
      T=S(1)*0.8862269

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      GO TO 9
30 T=S(1)
      GO TO 9
11 SE=S(1)+(0.0253-E(1))*((S(2)-S(1))/(E(2)-E(1)))
      IF(TYP(I).EQ.G.OR.TYP(I).EQ.Z) GO TO 31
      T=SE*0.8862269
      GO TO 9
31 T=SE
      9 IF(KSIK.EQ.3) GO TO 881
      IF(KSIK.EQ.2) GO TO 877
      IF(KSIK.EQ.0) GO TO 890
      TT=T
      GO TO 874
877 TT=TT/T
      IF(TYP(I).EQ.F) GO TO 879
      T=TT
      GO TO 890
881 T=T/(1.+TT)
890 WRITE (NOUT,12) MAT,TYP(I),NE
      12 FORMAT(1H /1H ,A9,A9,I10)
      WRITE (NOUT,13) T
      13 FORMAT(E16.8)
      K=5
      WRITE(LIZ) K,MAT,TYP(I),NE
      K=1
      WRITE(LIZ) K,T
      IF(TYP(I).EQ.C) GO TO 810
      IF(TYP(I).EQ.B) GO TO 812
      GO TO 1
810 SC=SC+T
      IF(INR.EQ.1) GO TO 818
      INR=1
      GO TO 1
812 SC=SC-T
      IF(INR.EQ.1) GO TO 818
      INR=1
      GO TO 1
818 WRITE(NOUT,12) MAT,D,NE
      WRITE(NOUT,13) SC
      K=5
      WRITE(LIZ)K,MAT,D,NE
      K=1
      WRITE(LIZ) K,SC
      1 CONTINUE
      KL=KL+1
      RETURN
      END

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.OR.TYP(I).EQ.X

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