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Abstract

Measurements of the pressure drop and the velocity distribution at three-dimensional roughnesses in a rectangular channel of variable channel width were performed with air. The friction factors of the extreme roughness ($p/h=2.5$, $g/e=1$) were found to be the highest which were measured up to now. The velocity distribution showed great differences to that observed at other roughnesses. The 'law of the wall' is not adequate to describe the velocity profile over these roughnesses, the velocity profiles are much flatter.

Geschwindigkeitsverteilung und Druckverlust an dreidimensionalen Rauigkeiten

Zusammenfassung

In einem rechteckförmigen Kanal mit variablem Plattenabstand wurden Messungen des Druckverlustes und der Geschwindigkeitsverteilung an dreidimensionalen Rauigkeiten mit Luft durchgeführt.

Die mit der extremen Rauigkeit ($p/h=2.5$, $g/e=1$) gefundenen Reibungskoeffizienten sind die höchsten bisher gemessenen. Die Geschwindigkeitsverteilung unterscheidet sich sehr von der an anderen Rauigkeiten beobachteten. So kann das Wandgesetz das Geschwindigkeitsprofil an dieser Rauigkeit nicht mehr richtig beschreiben, die Profile sind sehr viel flacher.

1. Introduction

Artificial roughness at the surface of fuel element rods is used in gas cooled reactors to improve the heat transfer capacity at a limited mass flow rate. The roughness which is presently used consists of repeated ribs or threads on the rod surface. With parts of the ribs cut away in an alternating pattern, we get a three-dimensional roughness, which was found to have an even higher heat transfer capacity /1,2/. The R-functions, which describe the friction characteristics of the roughness could, however, not be correlated in the same manner as it was done for two dimensional roughnesses /2/. In order to eliminate those inconsistencies measurements of the flow distribution at artificial three-dimensional roughnesses were undertaken. This report is the third in a series about measurements in a rectangular channel. The other two were on two-dimensional rectangular roughnesses and on round edged ribs /3,4/ (see also /5,6,7/).

2. Experimental Setup

Since the test rig and measuring methods were described in great detail in reference /3,5/ and /7/ only the main features and modifications shall be described here.

The measurements were performed with air near the open outlet of a vertical rectangular channel (Fig.1). The internal dimensions of the channel are 700 mm in the wide direction (z) and 60 mm minimum and 210 mm maximum in the y-direction with tolerances of ± 0.5 mm. The roughness elements which were made of aluminium were fixed to one of the wide walls, respectively to both wide walls. The axial pressure drop was measured by 13 pressure taps (0.2 mm i.d.) in the smooth wide wall over a length of 6500 mm.

The velocities were measured by means of a circular Pitot tube with an outer diameter of 0.6 mm. The corresponding static pressures were measured with a second tube, axially aligned to the flow direction, which has four holes at its circumference and an ellipsoid shaped head. The lateral distance between static tube and Pitot tube was set to 8 mm. The wall shear stress at the smooth wall was determined by the Preston method, using the same Pitot tube and a static pressure tap in the smooth wall.

The cross slide, which was used to position the probes with an accuracy of 0.01 mm at any position of the flow cross section, was installed 150 mm downstream of the channel outlet in order not to block the flow. The probe support with a diameter of 4 mm at its end extended approximately 300 mm into the channel, where the velocity measurements were taken. Velocity traverses parallel to the wide walls at the position of maximum velocity at the center, show a small zone in the center which seems to be unaffected by the short side walls (Fig.2). Here the velocity measurements were taken. There is, however, a strong variation of the velocity close to the short side walls, especially for the narrow channel widths. This overshooting of the velocity was not observed at two dimensional roughnesses. Because of difficulties in fixing the last row of elements closest to the side walls, this row was left out for this particular roughness (No.5), thus creating a narrow strip of smooth wall near the corners. This might be one reason for the peculiar velocity distribution.

3. Test parameters

Five different three-dimensional roughnesses were tested. The geometry is shown in figure 3 and the parameters are listed in table 1. The parameters were chosen in such a way, that the friction factors to be expected were high. In a simple experiment where the spacing p and g was systematically varied and the force upon a roughness rib was measured a relationship between friction factor and roughness geometry had been found /3,6/.

Figure 4 shows those results in terms of a friction factor based on the maximum velocity in the flow cross section for $e/h=4$. The parameters which were chosen for this investigation are marked (No.1,2 and 3) in figure 4. Roughness No.1 had, however, a ratio $e/h=6$, which had been found to yield higher friction factors than a ratio of $e/h=4$.

The roughness with the highest friction factor (No.2) was tested again with a smaller roughness height (No.5) and a somewhat different e/h -ratio of $e/h=5$. This roughness was also used to test a symmetrical rough channel, by placing roughnesses on both wide walls. Different from the rectangular roughness geometry is roughness No.4 with a rhombic geometry (see figure 3b). This shape was chosen because it was used in single pin tests in our laboratory /2/.

The relative roughness height was varied by changing the channel width L . Thus, the aspect ratio of the channel varied between 3.3:1 and 11.7:1.

The mean velocities were in the range between 8 and 30 m/s which resulted in a Reynolds number range of $0.8 \cdot 10^5 < Re < 5 \cdot 10^5$. The values of h^+ ranged between 500 and 3300. Most roughness-channel width combinations were tested at four different mass flow rates.

4. Evaluation

The time mean velocity u was calculated with the differential pressure between Pitot tube and static tube and the density of the humid air. The position of the Pitot tube close to the smooth wall was corrected according to Mac Millan /8/. A correction of the velocities for the effect of turbulence was not applied.

The shear stress at the smooth wall was determined by Preston tubes using the Patel /15/ calibration. The shear stress at the rough wall was determined by the knowledge of the axial pressure drop dp/dx and the shear stress at the smooth wall τ_s . Since there is a region in the center of the channel in which the influence of the short side walls on the flow is negligible, a force balance of steady flow yields

$$(\tau_s + \tau_r) dz dx = L dz dp. \quad (1)$$

Because of the discrete roughnesses the average over one pitch must be taken at least, thus the mean shear stress is

$$\tau_r = L \frac{\Delta p}{\Delta x} - \tau_s \quad (2)$$

The reference surface of this shear stress is the smooth wall between the ribs ($\epsilon=0$, see Fig.5). For another reference surface ($\epsilon \neq 0$) τ_r changes to

$$\tau_r = (L - \epsilon) \frac{\Delta p}{\Delta x} - \tau_s \quad (3)$$

The extension of the zones influenced by the smooth (\hat{y}_s) and rough (\hat{y}_r) wall respectively is given by the ratio of the wall shear stresses

$$\frac{\hat{y}_s}{\hat{y}_r} = \frac{\tau_s}{\tau_r} \quad (4)$$

With $\hat{y}_s = L - \hat{y}_r - \epsilon$, the length of the rough velocity profile is given by

$$\hat{y}_r = \frac{L - \epsilon}{1 + \tau_s / \tau_r} \quad (5)$$

It can be seen that the position of zero shear, however, is not dependent on ϵ if it is defined by \hat{y}_s . From equation (2) and (5) we get

$$\hat{y}_s = \frac{\tau_s}{\Delta p / \Delta x} \quad (6)$$

The average velocities in the two zones were determined by numerical integration of the measured values between the respective walls and the zero shear stress line. At the rough wall the integration started at the rib tip or at the root depending on the measuring position.

The friction factors are calculated by

$$f_{r,s} = 2/\bar{u}_{r,s}^{+2} \quad (7)$$

with

$$\bar{u}_{r,s}^{+} = \frac{\bar{u}_{r,s}}{(\tau_{r,s}/\rho)^{1/2}} \quad (8)$$

The Reynolds numbers are given by

$$Re_{r,s} = \frac{\bar{u}_{r,s}^4 \hat{y}_{r,s}}{\nu} \quad (9)$$

The experiments with single pins in annuli and with rod bundles are usually evaluated by the assumption of logarithmic velocity profiles near the smooth and rough wall with

$$u_s^+ = A_s \ln y_s^+ + B \quad (10)$$

and

$$u_r^+ = A_r \ln (y_r/h) + R. \quad (11)$$

If the results of the present experiment are to be used in this way the profile parameters A_s, A_r, B and R must be determined from integral quantities. As described before /3-7/, the following conditions must be met:

$$\bar{u}_s = \frac{u_{\tau s}}{\hat{y}_s} \int_0^{\hat{y}_s} (A_s \ln y^+ + B) dy \quad (12)$$

$$\bar{u}_r = \frac{u_{\tau r}}{\hat{y}_r} \int_0^{\hat{y}_r} (A_r \ln \frac{y}{h} + R) dy \quad (13)$$

$$u_{smax} = u_{rmax} \quad (14)$$

Integration of equation (12) and (13) for the flow in a plane channel yields

$$\bar{u}_s = u_{\tau s} \left[A_s \ln \left(\frac{\hat{y}_s u_{\tau s}}{v} \right) + B - A_s \right] \quad (15)$$

$$\bar{u}_r = u_{\tau r} \left[A_r \ln \left(\frac{\hat{y}_r}{h} \right) + R - A_r \right] \quad (16)$$

and equation (14) reads

$$\left[A_s \ln \left(\frac{\hat{y}_s u_{\tau s}}{v} \right) + B \right] u_{\tau s} = \left[A_r \ln \left(\frac{\hat{y}_r}{h} \right) + R \right] u_{\tau r} \quad (17)$$

with $\hat{y}_r = L - \epsilon - \hat{y}_s$

Of the five variables A_s, B, A_r, R and ϵ , two must be preset the other three are determined by equations (15), (16) and (17). A presetting of the parameters ϵ and B has turned out to be most suitable.

5. Results

5.1 Velocity profiles

Figures 6.1 - 6.27 show all measured velocity profiles reduced with the respective maximum velocity. The position of the Pitot tube relative to the roughness ribs was varied and is given by the x/h -value. Values of $x/h > 0.5$ represent measurements starting at the smooth wall between the ribs. The other measurements were taken above a rib. Most measurements were taken at the exact center of the channel and of a rib (z -position). The influence of the measuring position of the velocity distribution was especially large for narrow channel widths and the roughness with the highest friction factors (No.2). Here we get also different velocity distributions at the same x/h -position but different z -positions. This was verified by measuring the velocity at a fixed distance from the wall in z -direction (not shown). The level of turbulence at these measurements was extremely high, which made the measurements difficult because of the large fluctuations of the Pitot-readings. In figures 6.24-6.27 the velocity profiles of the channel with two rough walls are shown. The velocity distribution turned out to be exactly symmetrical, therefore only one half of the velocity profile was measured and is shown here twice. The velocity distribution at different relative roughness heights does not show a systematic variation. For some roughnesses it shows a tendency to become flatter with higher h/L -ratios, but this trend is reversed for other roughnesses. In order to check, whether there was an effect of the channel outlet, some measurements were performed 840 mm upstream instead of only 300 mm, but no difference was found.

5.2 Friction factors

The bulk friction factors were determined by

$$f_B = \frac{\Delta P}{\Delta x} \frac{(L-\epsilon)}{\rho \bar{u}^2} \quad (18)$$

with the bulk velocity \bar{u} from integration of the velocity profiles and the axial pressure drop $\Delta p/\Delta x$. From the 12 differential pressures measured along the channel an average value was formed without the first four measurements at the entrance and the last one at the exit because of deviations from a linear pressure gradient.

Figure 7 shows the friction factors determined with $\epsilon=0$, i.e. neglecting the volume of the roughness, as function of the relative roughness height. Plotted are the average values of all measurements at a certain channel width with the range of scatter. The large scatter stems mainly from the different bulk velocities due to different axial measuring positions. Since there is a lateral velocity variation due to the limited length of the roughness ribs, the integrated velocity at a certain traverse is different at different axial positions. This was not the case at two-dimensional roughnesses. As for other roughnesses the friction factor grows with increasing relative roughness height. But there is one distinctive difference to former measurements. The function $f=f(h/L)$ is not linear but has a maximum at a certain value of h/L for three roughnesses which were tested up to the highest h/L -ratio. This does not change if the friction factor of the rough zone f_r is plotted over the h/\hat{y}_r -ratio (Fig.8). Neither does another definition of the hydraulic diameter change this tendency. For the other extreme of $\epsilon, \epsilon=h$, the friction factors are lower but the maximum is even more pronounced (Fig.9). The maximum is reached at different h/\hat{y}_r -ratios for the different roughnesses; for the roughness with the highest friction factor it is reached at the lowest h/\hat{y}_r -ratio. At the highest h/\hat{y}_r -ratio all three

roughnesses have the same friction factor.

Because of this strange characteristics the roughness No.5 was tested with similar parameters as No.2 but with a smaller height. The agreement in the friction factors for the lower h/\hat{y}_r -ratios is very good, but this roughness does not show a maximum friction factor at a high h/\hat{y}_r -ratio. It could, however, not be tested up to the same h/\hat{y}_r -ratio.

Errors in measurements of pressures and velocities or in the integration of the velocity cannot account for this characteristics. It is also very unlikely that secondary flows can cause a difference in the friction factor by a factor of two, between the measured and the expected ones. Besides, secondary flows tend to increase the friction factors, rather than decrease them. The possible explanation, that the flow at the mid-plane is not unaffected by the short side walls, is not convincing since the aspect ratio at high h/L -ratios is also the highest. So, it can only be speculated about the reason for these results. It might be possible that due to the restricted flow cross section the dead water regions behind the ribs are changed in length, and a roughness which had a high friction factor at low relative roughness heights is no longer in the optimum range of the parameter field, but a different pitch would result in higher friction factors. No measurements at roughnesses, however, are known to the author, which yielded friction factors as high as the present ones. Measurements at single pins with three-dimensional roughness contained in different smooth tubes yielded friction factors which were increasing less than linear with increasing relative roughness height, but no maximum was found /9/. These friction factors were similar to those of roughness No.4.

The Reynolds number range of those measurements was large and the flow was in the fully rough regime at Reynolds numbers and h^+ -values which were reached in the present experiment. The

highest friction factors obtained by Nikuradse /10/ with sand grain roughness are well below the present ones and are even lower than the highest friction factors obtained with two dimensional roughnesses /3,6,7,11/.

The results at high h/L-ratios cannot be generalized, more measurement would be necessary. So, the further discussion should be restricted mainly to the results for low h/L-ratios, up to h/L=0.08, especially for roughness No.5, where agreement with No.2 is good.

This roughness was tested with one wall rough and with both walls rough, i.e. a symmetrical rough channel. The friction factors of the rough zone should be the same for both cases at the same relative roughness height h/\hat{y}_r . However, this friction factor depends on the definition of the hydraulic diameter. The shear stress at the rough wall in the channel with one rough wall is given by equation (3),

$$\tau_r = (L-\epsilon) \frac{\Delta p}{\Delta x} - \tau_s , \quad (3)$$

and for the channel with two rough walls it is

$$\tau_r = \left(\frac{L}{2}-\epsilon\right) \frac{\Delta p}{\Delta x} . \quad (19)$$

Since the friction factor is

$$f_r = \frac{2(\tau_r/\rho)^{1/2}}{\bar{u}_r} , \quad (20)$$

and \bar{u}_r is changed by the same amount in both cases if ϵ is changed, there must be one ϵ -value which produces the same friction factors for the asymmetrical and the symmetrical case.

Figure 10 shows the friction factor over the h/\hat{y}_r -ratio for different definitions of the hydraulic diameter.

Besides $\epsilon=h$ and $\epsilon=0$, there is a volumetric definition $\epsilon=hb(1+g/e)/p$ and the definition $\epsilon=-h$, which is physically absurd but which yields the same friction factors for both cases, the symmetrical and the asymmetrical flow. Experiments at two-dimensional roughnesses ($p/h=7.7$) had shown that the same friction factors are obtained with a volumetric definition ($\epsilon=hb/p$).

From other experiments it is known that the friction factor of the smooth zone, f_s , increases due to the opposite roughness. Therefore the friction factor f_s reduced by the theoretical friction factor f_o of a smooth pipe at the same Reynolds number

$$1/f_o^{1/2} = 4 \log(\text{Re}_s f_o^{1/2}) - 0.4 \quad (21)$$

was plotted over various parameters in figure 11-14.

In figure 11 the parameter is the ratio of the friction factors of the rough and the smooth zone. The relation from Warburton /12/ gives too low friction factor multipliers, which was already found for two-dimensional roughnesses /3,6,7/. A better relation would be

$$\frac{f_s}{f_{os}} = 1.025 + 0.008 \frac{f_r}{f_s} \quad (22)$$

The relation from Warburton and Pirie /13/, which uses the parameter f_r/f_{or} , with f_{or} being the friction factor of a smooth tube at the Reynolds number of the rough zone, gives too high friction factor multipliers. Here the data are better correlated by

$$\frac{f_s}{f_{os}} = 1.036 + 0.004 \frac{f_r}{f_{or}} \quad (23)$$

The correlation

$$\frac{f_s}{f_{OS}} = 1.056 + 0.062 \frac{h}{\hat{y}_s} \quad (24)$$

which was found for two-dimensional roughnesses /3,6,7/ correlates the present data as well with a tolerable amount of scatter. Plotting f_s/f_{OS} over $\log(h/\hat{y}_s)$ reduces the amount of scatter somewhat (Fig.14).

Since the correlation (24) using the roughness height and length of the smooth profile as parameter is the only one which is independent of the parameter ϵ and fits for both two-dimensional and three-dimensional roughnesses, it is regarded superior to the other two.

5.3 Logarithmic velocity profiles at the smooth wall

The mean velocity profiles near the smooth wall are plotted in figures 15.1-15.22 in non-dimensional form, together with a straight line representing the 'law of the smooth wall'

$$u^+ = A_s \ln y^+ + B \quad (10)$$

with $A_s = 2.5$ and $B = 5.5$

Especially for the roughness No.2 the measured velocity profiles depart from this line already far from the maximum velocity, which is due to the large turbulent energy transfer from the rough zone. The line of zero shear stress and the position of maximum velocity lie extremely far apart. For this roughness the distance from the smooth surface to the position of maximum velocity was approximately $2 \hat{y}_s$ for all channel widths.

The slope A_s of almost all measured profiles is lower than 2.5 while the parameter B is higher, both tendencies getting stronger for higher relative roughness heights.

The slope A_s calculated from integral quantities by equation (15) with $B=5.5$ is plotted in figure 16.

The correlation

$$A_s = 2.55 + 0.4 / \ln (0.1 \frac{h}{\hat{y}_s}) \quad (25)$$

which was found for two-dimensional square roughnesses, fits quite well.

For a fixed slope of $A_s=2.5$, the resulting B-values would be in the range of $3.8 < B < 5.8$, decreasing with increasing h/\hat{y}_s but with a relatively large amount of scatter (see B_{SA} in table 2).

5.4 The breakdown of the 'law of the rough wall'

Figures 17.1-17.27 show the non-dimensional profiles at the rough wall together with a line representing the 'law of the rough wall' with the slope $A_r=2.5$. The origin of the velocity profile was put at the root of the ribs ($\epsilon=0$). No effect of the Reynolds number could be detected, but there is a strong effect of the measuring position for most roughnesses.

The big discrepancy between the measured velocity profiles and the 'law of the wall' is evident. Except for some profiles for low relative roughness heights, it does not seem to be sensible to force a straight line through the points in order to define a slope A_r and the parameter R. All slopes would be less than 2.5 and some would be as low as 0.5. A displacement of the origin ϵ by several rib heights would be necessary in order to get a value of 2.5. There is no big difference between symmetric and asymmetric flow (Fig.17.21-17.27). Since the position of zero shear stress is far beyond the position of

maximum velocity a large region of the velocity field cannot lie on any straight line, and a description of the flow by a law of the wall would be inadequate. Nevertheless, for the calculation of the flow in an asymmetrical channel some description of the velocity field is necessary. For lack of a better one, the logarithmic description was used. The values of A_r and R were calculated by equation (16) and (17). Most values of A_r are between 0.5 and 1.5, but there is no systematic variation (see table 2). For roughness No.5 the average value is $A_r=0.52$.

From a plot $(2/f_r)^{1/2}$ versus $(\ln \frac{\hat{y}}{h} - 1)$ the parameters A_r and R can be obtained /14/. Figure 18 shows the result for roughness No.5. For symmetrical flow we get $A_r=1.8$ and for asymmetrical flow $A_r=2.9$.

6. Transformation of bulk data with different profile parameters

For the transformation of results from pressure drop measurements at rough rods in annular channels and the application of these transformed data on a bundle geometry it is important to know the errors which arise from using different velocity profile parameters.

Therefore the transformation was applied on the measurements with roughness No.5 (For a derivation of the transformation equations for parallel plates see /3/). Input data are the bulk friction factor and the bulk Reynolds number (table 2), a constant $B=5.5$ and the different values of A_s and A_r . For A_s , only $A_s=2.5$ and the correlation (25) was used. For A_r , the values 0.52, 1.8, 2.9 and 2.5 were used.

The calculated friction factors of the smooth and rough zone and the position of zero shear stress were compared to the experimentally determined results (1. line) and the difference in per cent is given in table 3. Also the calculated R -value is given. While the differences in the smooth friction factor are rather large, the error in the friction factor of the rough

zone is relatively small.

Even for the simple assumption of a constant A_S and $A_R=2.5$ the maximum error is only 4.75%. It is of interest, that a modification of A_S alone does not change the friction factor considerably. Except for the extreme value of $A_R=2.9$, where the errors in f_S and f_R are greatest, no constant R value is obtained.

7. Eddy viscosity

Calculations of turbulent flow are often based on semi-empirical turbulence models /16/, using the distribution of the eddy viscosity as an input parameter.

In analogy with the fluid viscosity ν an eddy viscosity ϵ in turbulent flow can be defined by

$$-\rho \overline{u'v'} = \rho \epsilon \frac{\partial u}{\partial y}, \quad (26)$$

with $\overline{u'v'}$ as the time-mean correlation of the velocity fluctuations in the main direction of the flow (u') and in the direction normal to the wall (v'), u is the time-mean fluid velocity and y is coordinate normal to the wall.

In contrast to the viscosity ν the eddy viscosity ϵ is not a constant in the flow field. One of the models for the distribution of the eddy viscosity in smooth tubes which can be expressed explicitly was given by Reichardt /17/:

$$\frac{\epsilon}{u_\tau R} = \frac{\kappa}{3} \left[1 - \left(\frac{R-y}{R} \right)^2 \right] \frac{1}{2} + \left(\frac{R-y}{R} \right)^2 \quad (27)$$

with $\kappa = 0.4$, R being the radius of the tube and u_τ the friction velocity.

If the turbulent shear stress is defined by equation (26) than the total shear stress is given by

$$\tau = \rho (\epsilon + \nu) \frac{\partial u}{\partial y} \quad (28)$$

The assumption of a linear variation of the shear stress in a plane channel

$$\frac{\tau}{\tau_w} = 1 - \frac{y}{\hat{Y}} \quad (29)$$

together with (28) leads to the relation

$$\epsilon^+ = \frac{\epsilon}{u_\tau \hat{Y}} = \left[\frac{1 - y/\hat{Y}}{\partial u^+ / \partial (y/\hat{Y})} - \frac{\nu}{\hat{Y} u_\tau} \right] \quad (30)$$

where ϵ^+ is the nondimensional eddy viscosity. In asymmetrical flow the friction velocities $u_{\tau S}$ and $u_{\tau R}$ and the lengths of the smooth (\hat{Y}_S) and rough (\hat{Y}_R) velocity profiles in the respective zones are used to evaluate ϵ_S^+ and ϵ_R^+ . The eddy viscosity was determined at discrete positions where velocity measurements were taken.

For roughness No.5 the non-dimensional eddy viscosity ϵ^+ is shown in figure 19 for the smooth zones and in figure 20 for the respective rough zones versus the non-dimensional distance $(y/\hat{Y})_{S,R}$ from the wall. For comparison with the smooth tube data the line according to equation (27) is added in the figures.

The eddy viscosity distributions of the smooth zones look similar. Compared to the smooth tube values they are higher in the center of the smooth zone by approximately 50% and decrease towards zero close to the zero shear stress line. In this region the scatter is large due to the low velocity gradient which approaches zero.

There are however differences close to the wall depending on the channel width. At large channel widths the data measured close to the smooth wall are coincident with the smooth tube results. At narrow channel widths which corresponds to a large relative roughness height the data very close to the wall are below the smooth tube results and show a steep increase.

Looking for possible errors in measurement or evaluation there are two facts which might cause errors. The outer diameter of the Pitot tube which was used was 0.6 mm. This is 1/6 of the length of the smooth zone in figure 20.4. The measuring position was defined by $y = y_c + 0.15 d_{\text{pitot}}$ according to Mac Millan, with y_c being the position of the Pitot tube center. The measurements close to the wall were taken in steps of 0.1 mm and were evaluated at the position half-way between two points taking their difference of u^+ and y for equation (30). Taking a close look at the data, it does not seem possible that these facts can account for the deviation of the eddy viscosity from the smooth tube results. The measurements in a rectangular water channel with one rough wall /11/ showed similar results. Higher values far from the wall were also reported by Rehme /18/, which were obtained in a subchannel of a rod bundle.

The distribution of the eddy viscosity in the rough zone shows a large deviation of that in the smooth zone. Although the scatter of the points is large, it is obvious that the eddy diffusivity is much higher for these three-dimensional roughnesses than for two-dimensional roughnesses as reported in /11/ and /19/. There are no points near $y/\hat{y}=1.0$ because the maximum velocity is reached already at lower y/\hat{y} values.

8. Conclusion

The investigation on three-dimensional roughnesses gave the following results:

1. Three-dimensional roughnesses produce very high friction factors, higher ones than ever measured before.
2. The friction factor does not raise linearly with increasing relative roughness height, but shows a maximum value.
3. There is no agreement between the friction factors in symmetrical and asymmetrical flow at the same relative roughness height, unless the hydraulic diameter is defined with a negative ϵ .
4. The non-coincidence of the position of zero shear stress and that of maximum velocity is more pronounced than for two-dimensional roughness.
5. The slope of the non-dimensional velocity profile at the smooth wall is lowered by the roughness opposite, which means higher friction factors. The correlation, which was found for two-dimensional roughnesses describes also the present results.
6. The 'law of the rough wall' does not hold for the description of the rough velocity profiles.
7. The transformation of bulk data with profile parameters as they were used up to now, result in friction factors up to 20% too low for the smooth zone and up to 5% too high for the rough zone for roughness heights smaller than 10% of the channel width.

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Nomenclature

A	slope of the logarithmic velocity profile
B	constant of the logarithmic velocity profile at smooth walls;
b	width of the roughness rib (m)
d_h	hydraulic diameter (m)
e	length (z-direction) of a rib (m)
f	friction factor = $2\tau/\rho u^2$
f_o	friction factor of a smooth tube
g	gap (z-direction) between two ribs (m)
h	height of roughness rib (m)
h^+	dimensionless height of roughness rib = $h u_\tau/\nu$
L	width of channel (m)
p	axial pitch of the repeated roughness rib (m)
p	pressure (Nm^{-2})
u	mean velocity (ms^{-1})
u_τ	friction velocity = $(\tau/\rho)^{1/2}$ (ms^{-1})
u^+	dimensionless velocity = u/u_τ
\bar{u}	average velocity in a section (ms^{-1})
R	parameter of the logarithmic velocity profile at rough walls

Re	Reynolds number = $\bar{u}d_h/\nu$
x	axial distance
y	distance normal to the wall
y^+	dimensionless distance from the wall = yu_τ/ν
\hat{y}	position of the zero shear stress line, length of respective zones
z	distance parallel to the wall normal to the flow

Greek symbols

ϵ	displacement of the origin of the velocity profile at rough walls (m)
ν	kinematic viscosity (m^2s^{-1})
ρ	density (kgm^{-3})
τ	shear stress (Nm^{-2})

Subscripts

max	maximum
r	at the rough wall or pertaining to the rough zone
s	at the smooth wall or pertaining to the smooth zone
vol	volumetric definition of origin of velocity profile

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No	p	h	e	p/h	g/e	h/L
1	35	10.4	60	3.37	0.5	0.050/0.077/0.122
2	25	10.4	40	2.40	1.0	0.050/0.065/0.077/0.095/0.122/0.173
3	35	10.4	40	3.37	1.25	0.049/0.065/0.095/0.122/0.172
4	32	10.2	60	3.14	0.783	0.049/0.076/0.120/0.170
5	15	6.1	30	2.46	1.0	0.029/0.045/0.072/0.100
5	same parameters, but both walls rough					

Table 1: Geometrical parameters of the roughnesses and channel

h/L	A _S	A _R	R	Δf _S (%)	Δf _r (%)	Δh/ \hat{Y}_r (%)
0.029	2.43	0.55	6.26	-	-	-
	2.43	0.52	6.33	0.18	-0.01	-0.09
	2.43	1.8	2.80	-7.29	1.61	3.20
	2.43	2.9	0.50	-11.05	3.77	5.43
	2.43	2.5	1.55	-9.41	2.66	4.39
	2.5	2.5	1.53	-12.21	3.02	3.81
0.045	2.38	0.74	4.83	-	-	-
	2.38	0.52	5.27	1.63	-0.13	-0.44
	2.38	1.8	2.35	-7.49	1.67	2.59
	2.38	2.9	0.44	-11.90	4.14	4.67
	2.38	2.5	1.31	-10.01	2.87	3.71
	2.5	2.5	1.28	-14.69	3.25	2.93
0.072	2.31	0.33	4.40	-	-	-
	2.31	0.52	4.11	-1.92	0.09	0.36
	2.31	1.8	1.83	-13.39	2.39	3.36
	2.31	2.9	0.35	-18.63	5.60	5.45
	2.31	2.5	1.02	-16.40	3.95	4.48
	2.5	2.5	0.99	-22.83	4.22	3.51
0.100	2.26	0.52	3.36	-	-	-
	2.26	0.52	3.36	0.0	0.0	0.0
	2.26	1.8	1.55	-14.41	2.82	2.96
	2.26	2.9	0.38	-20.61	6.83	5.07
	2.26	2.5	0.92	-17.99	4.75	4.09
	2.5	2.5	0.89	-25.71	4.82	3.04

Table 3: Comparison of friction factors evaluated with different parameters for roughness No.5

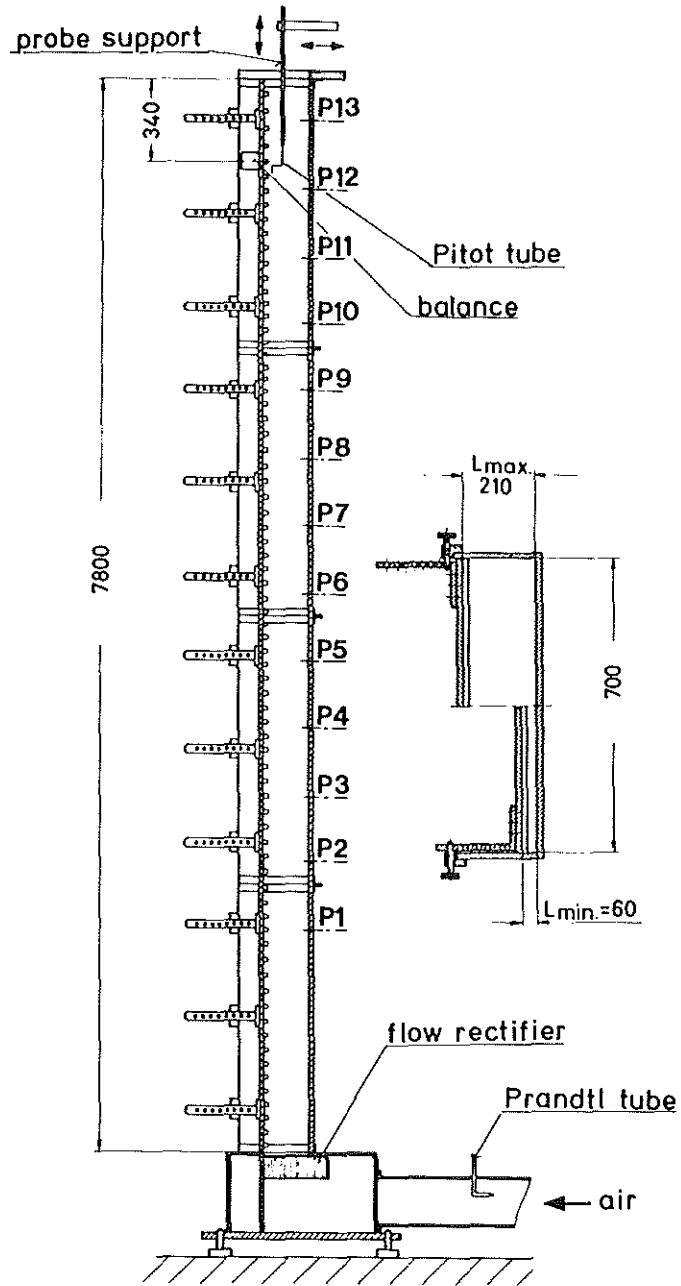


Fig. 1. The test section.

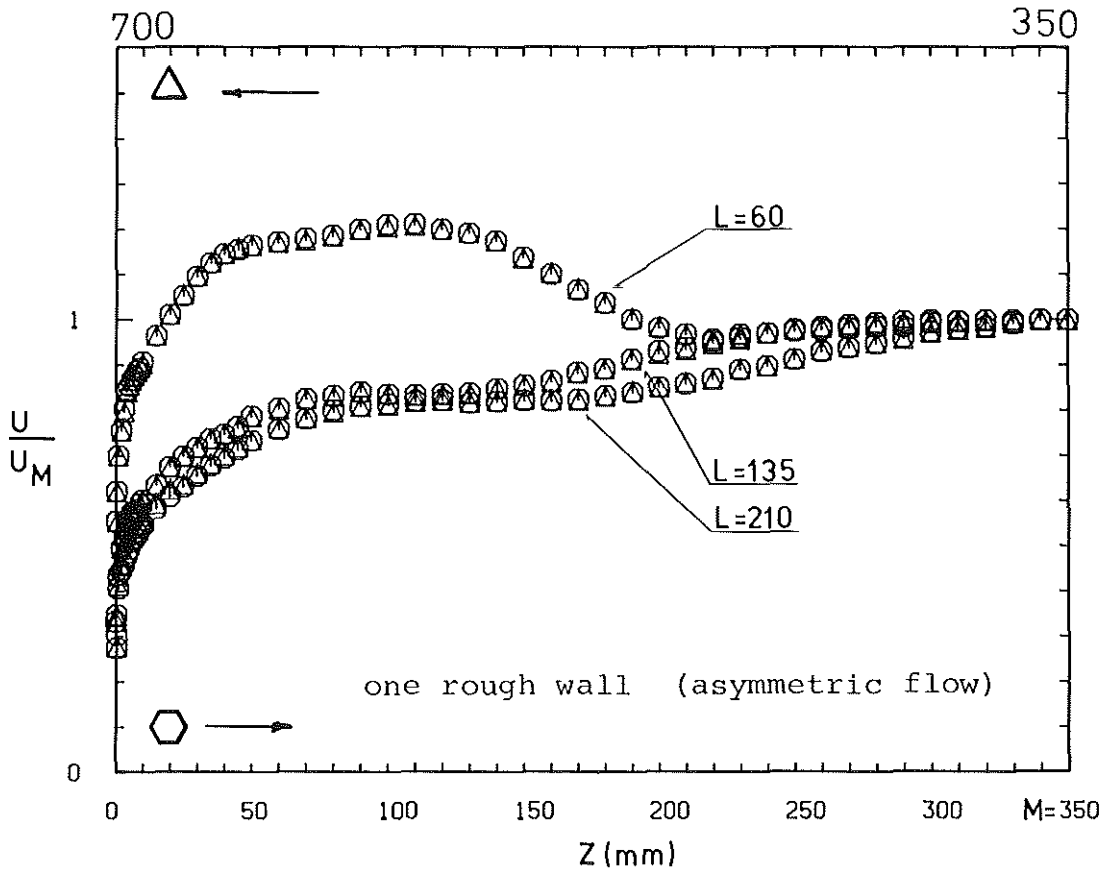
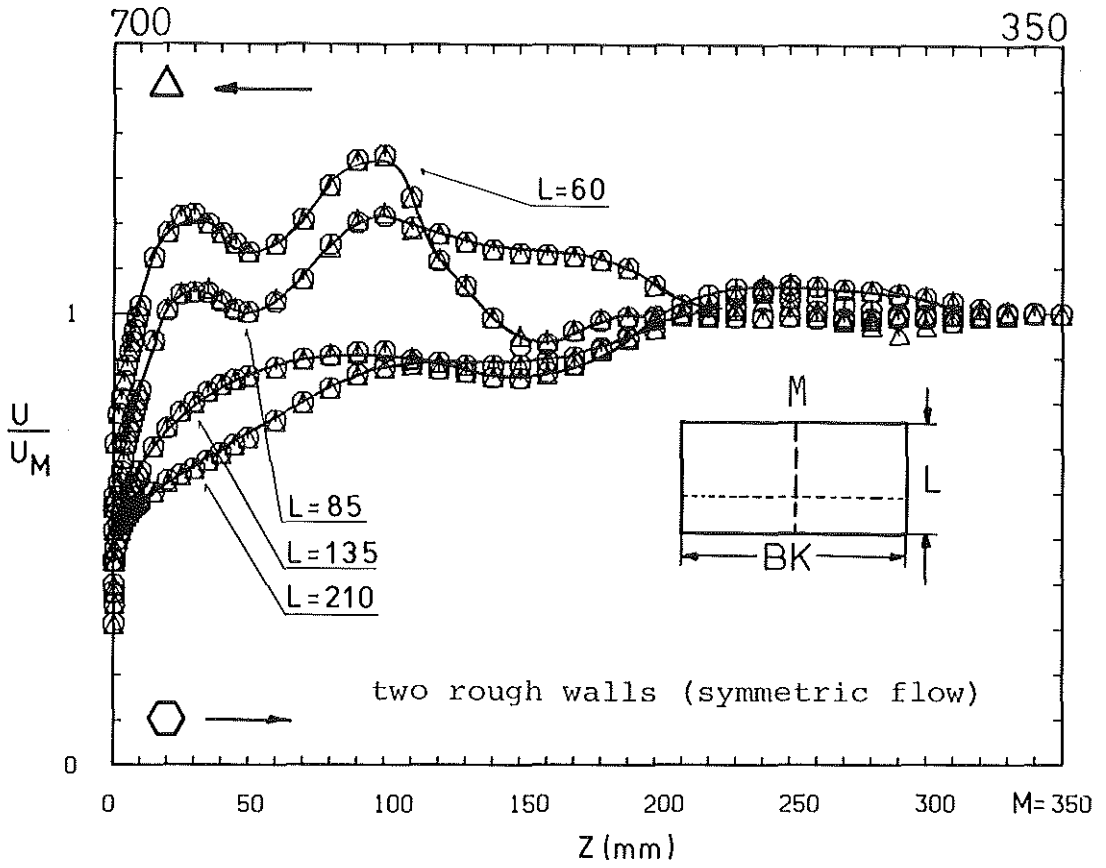


Fig.2: Velocity profiles parallel to the rough wall

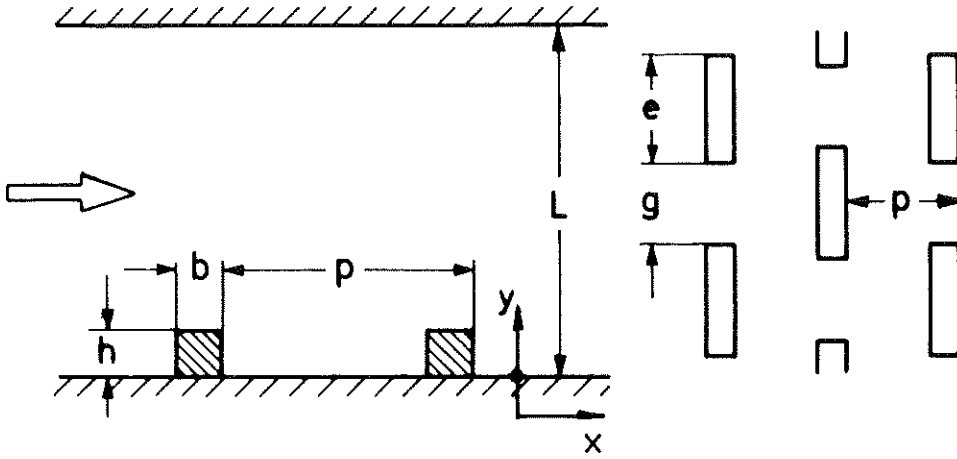


Fig. 3.a: Roughness geometry of No.1,2,3 and 5

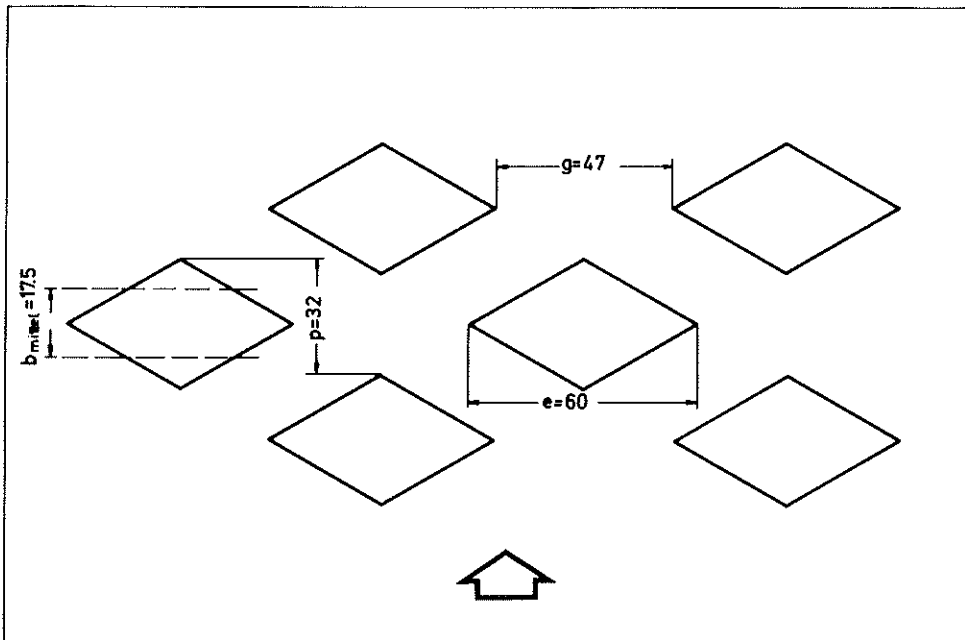


Fig.3.b: Roughness geometry of No.4

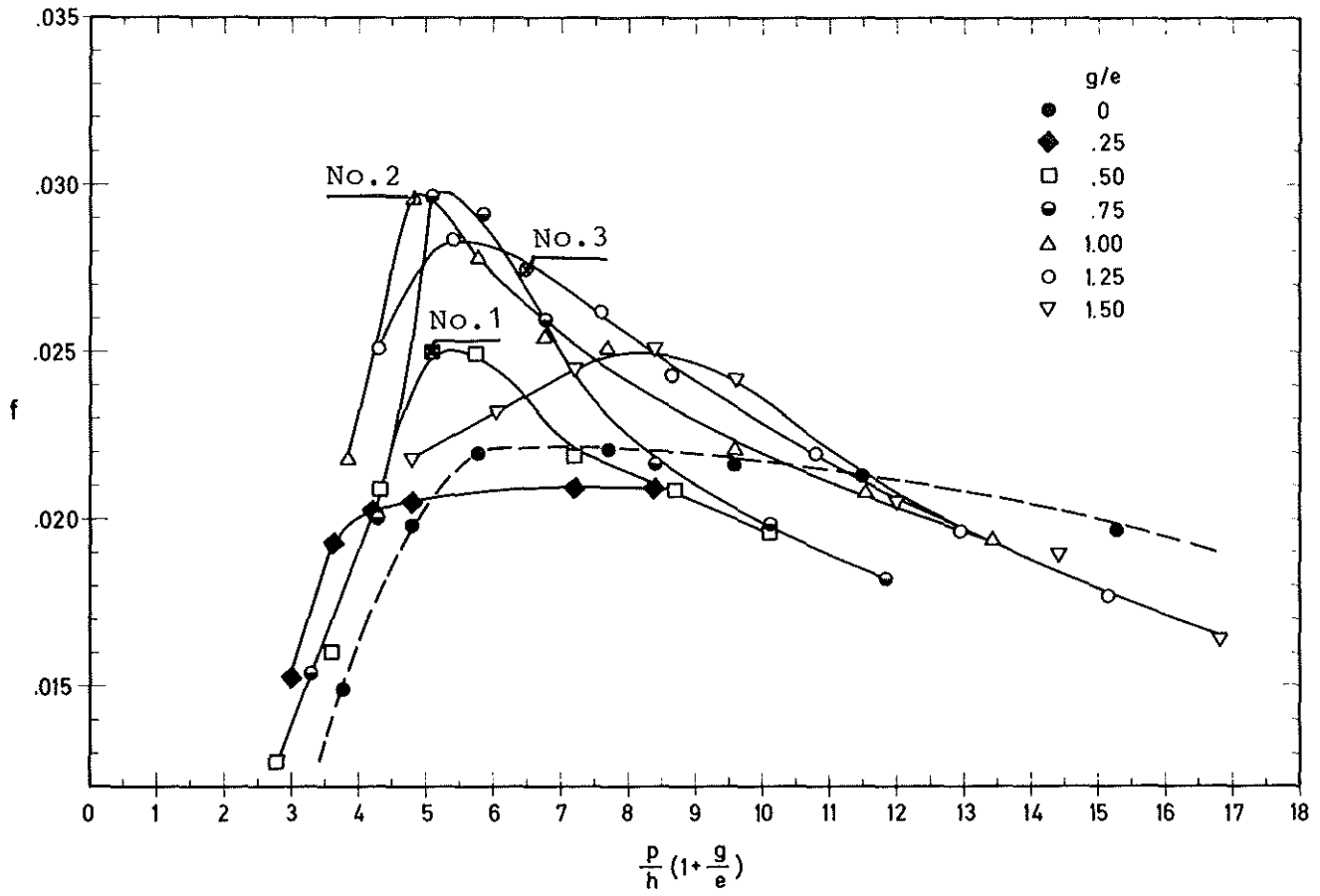


Fig.4: Friction factors as function of the roughness geometry ($e/h=4$)

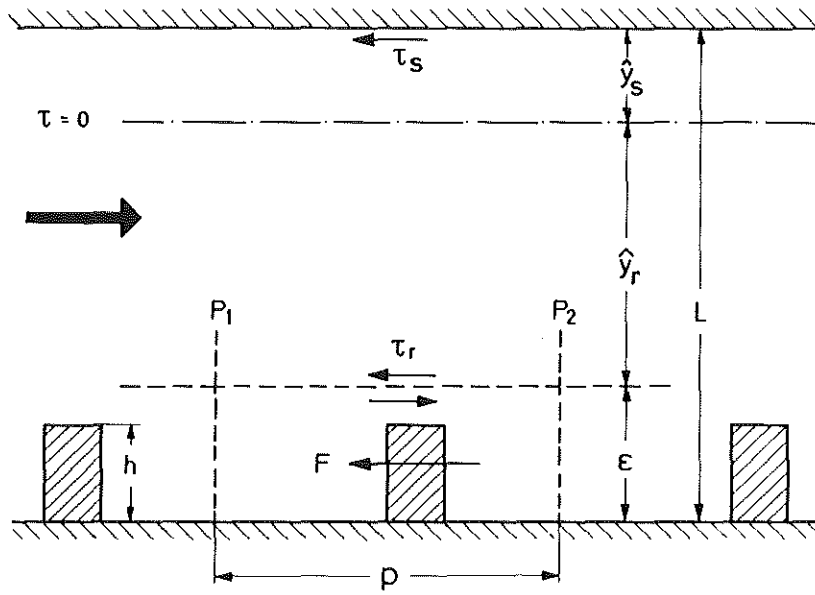


Fig.5: Scheme of flow cross section

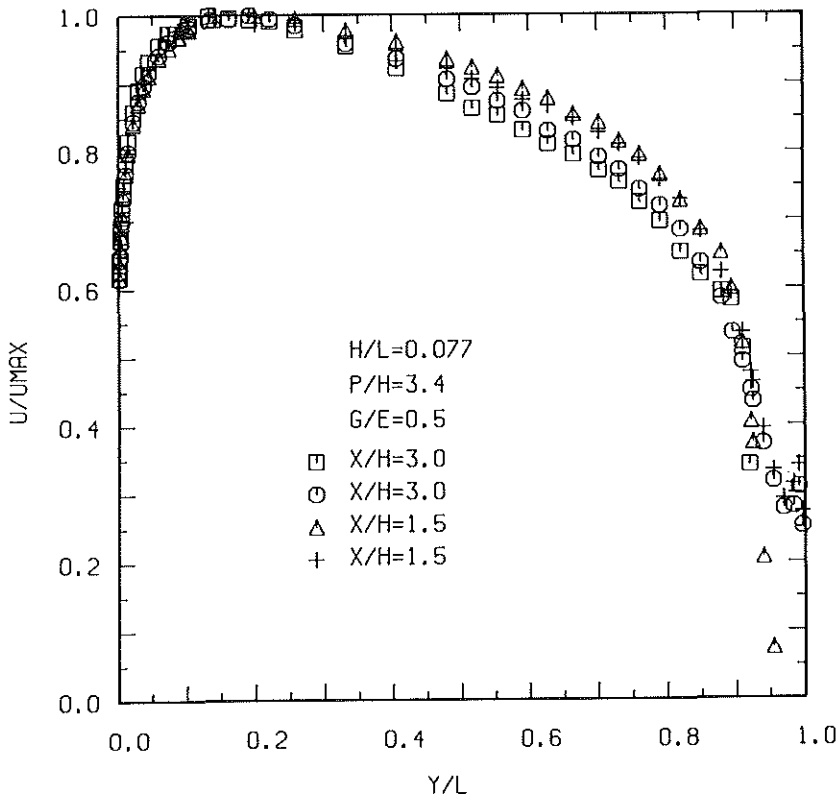
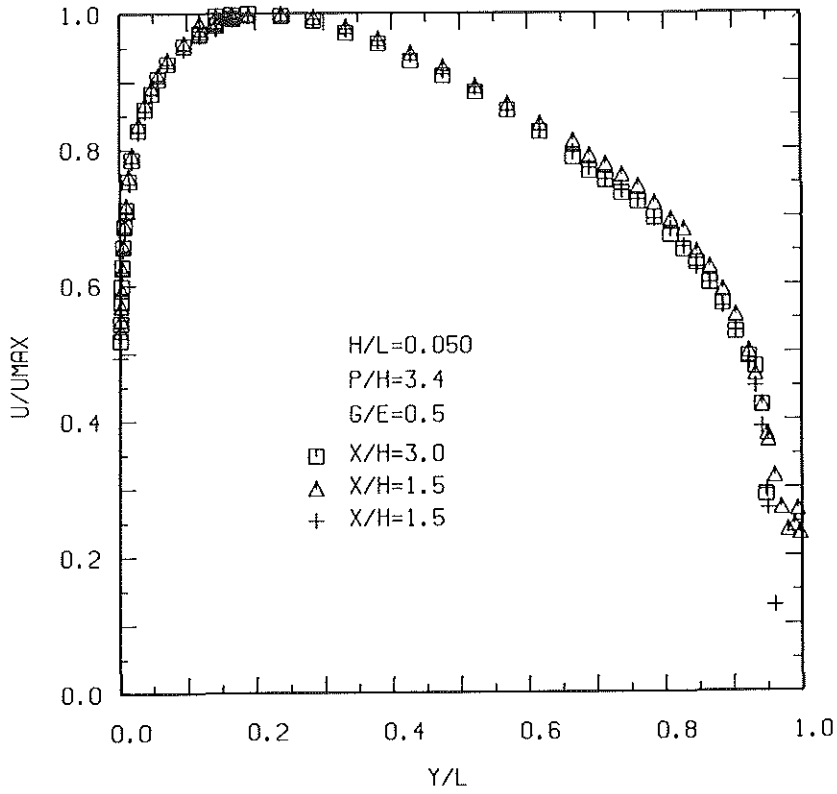


Fig.6.1-6.2: Velocity profiles (No.1)

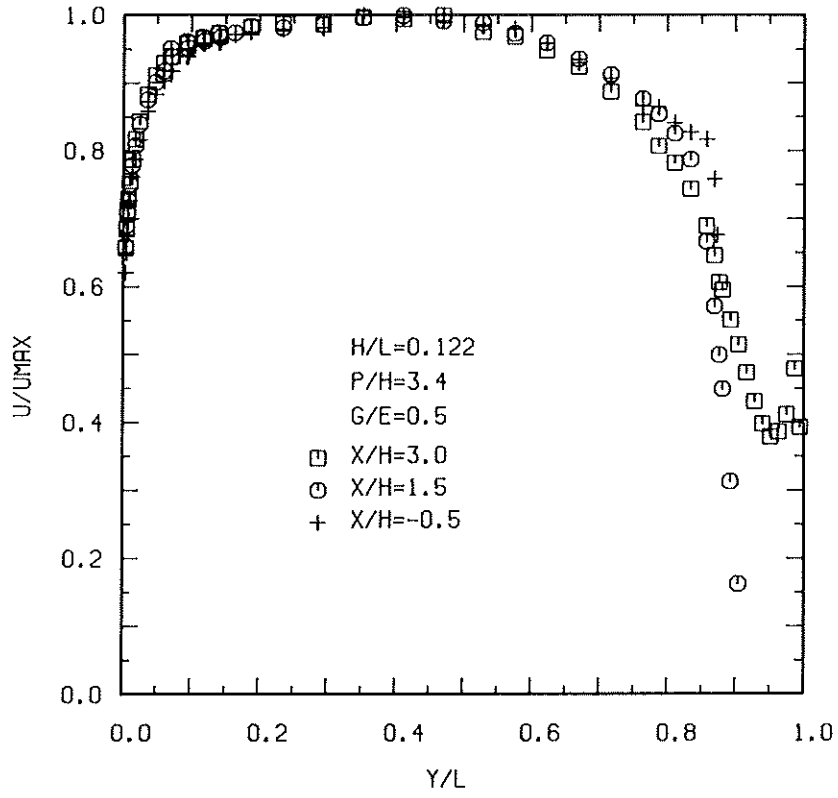


Fig.6.3: Velocity profiles (No.1)

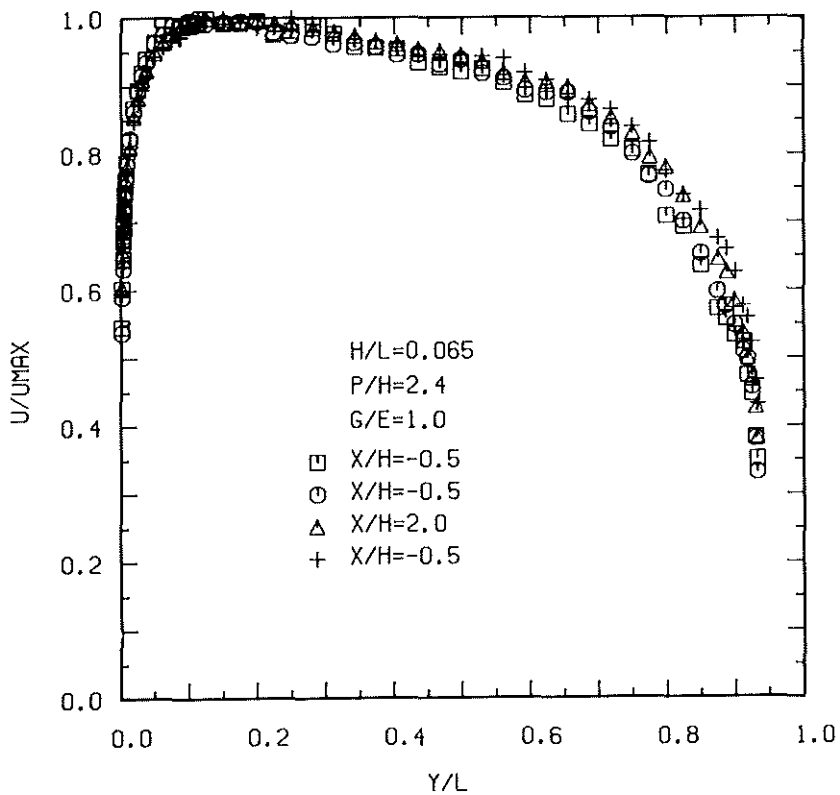
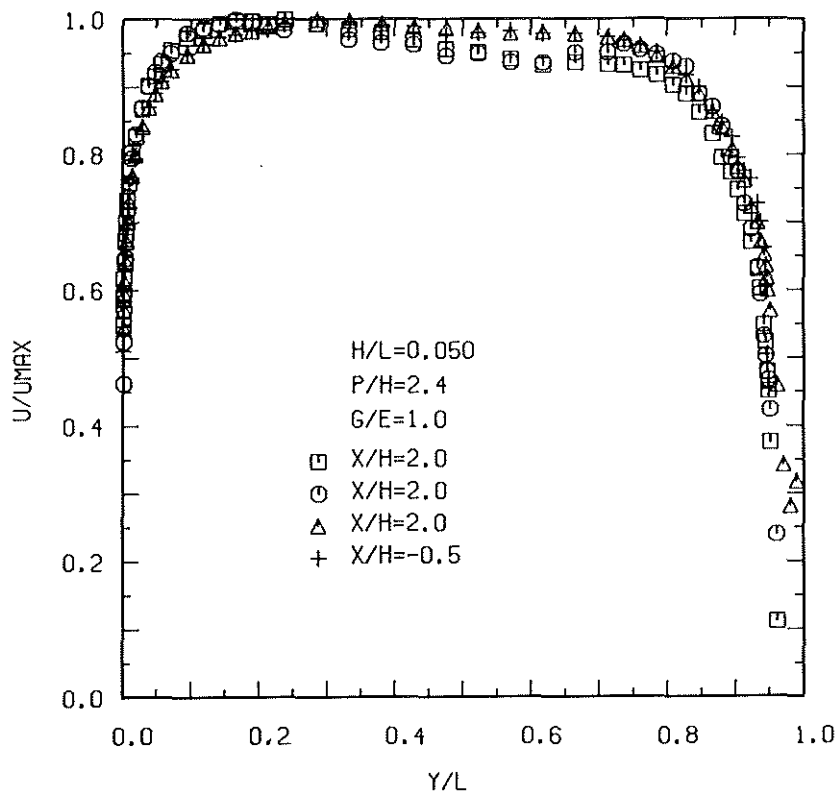


Fig.6.4-6.5: Velocity profiles (No.2)

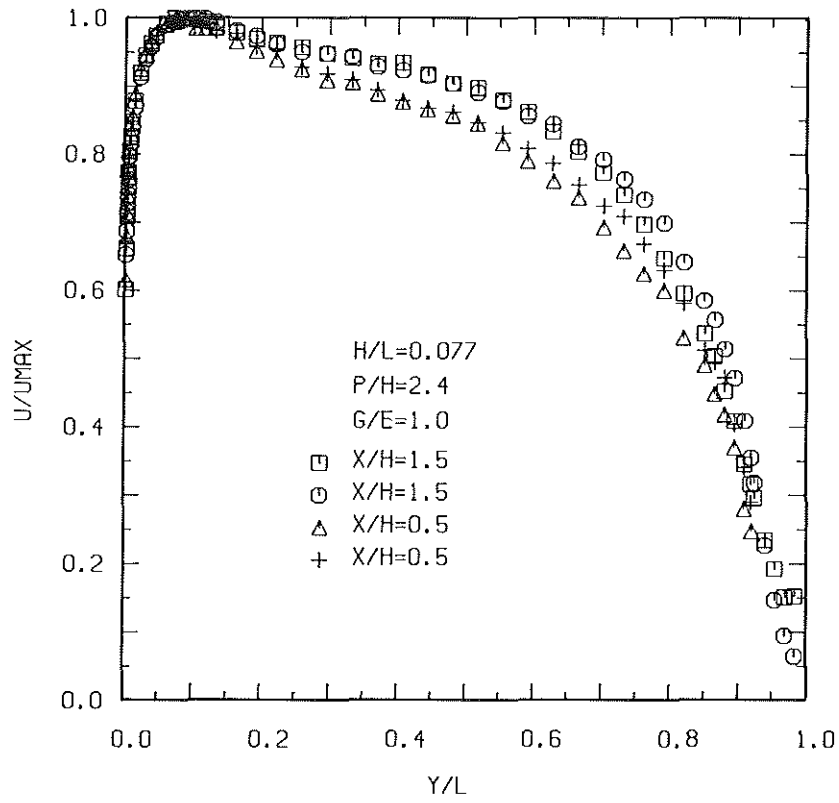


Fig.6.6: Velocity profiles (No.2)

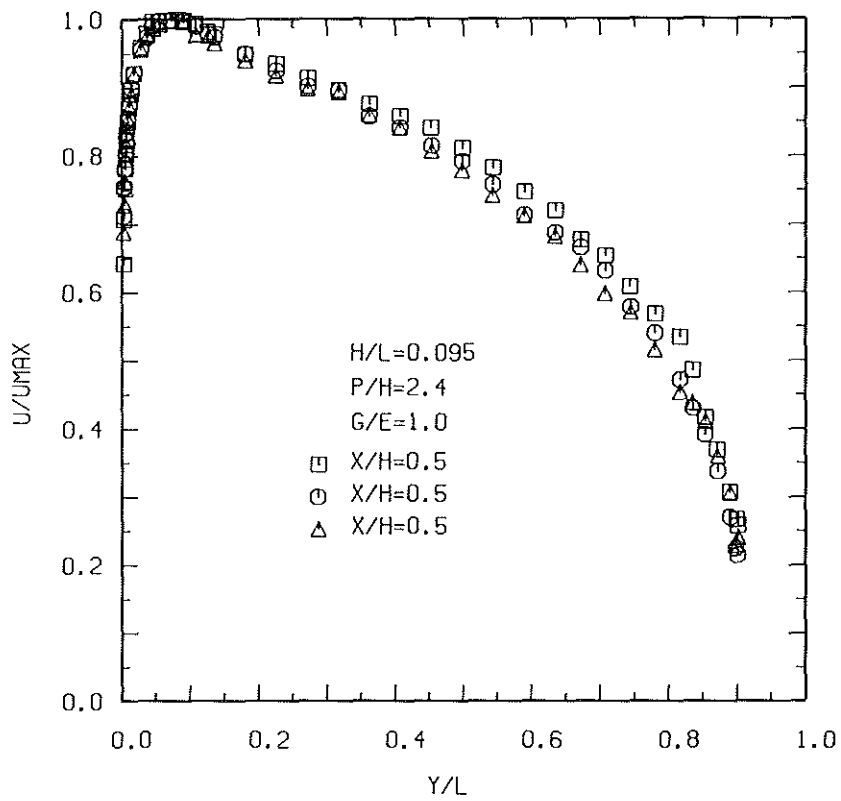
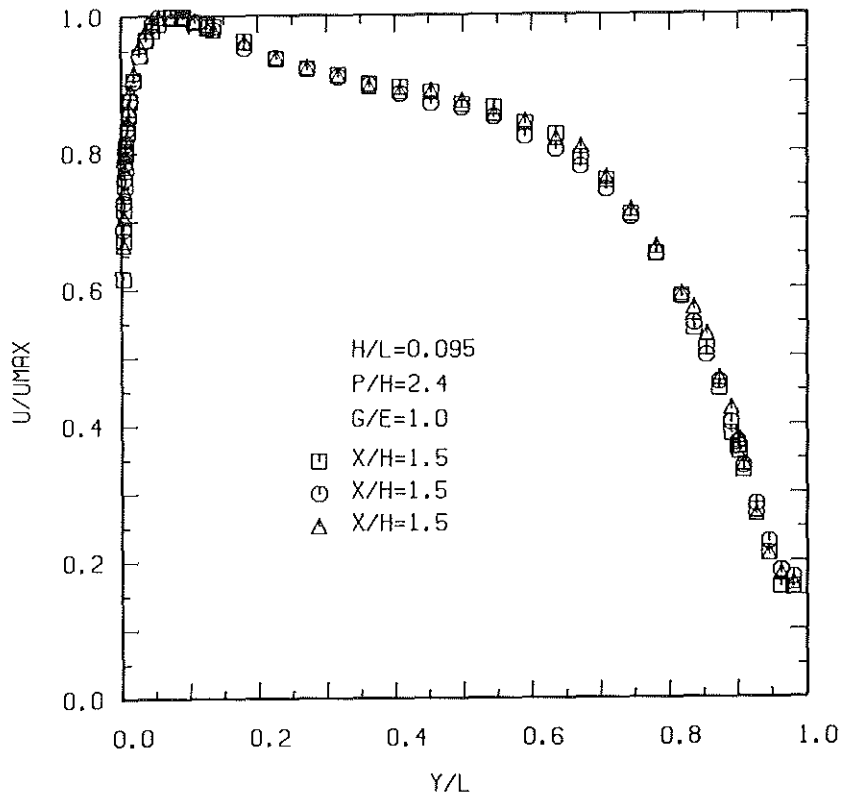


Fig.6.7-6.8: Velocity profiles (No.2)

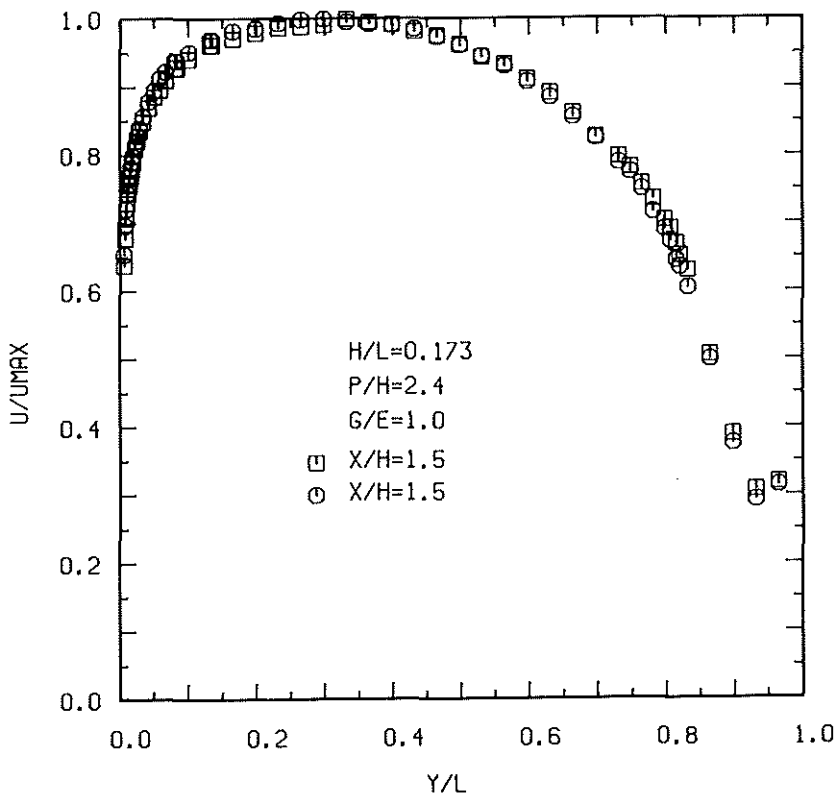
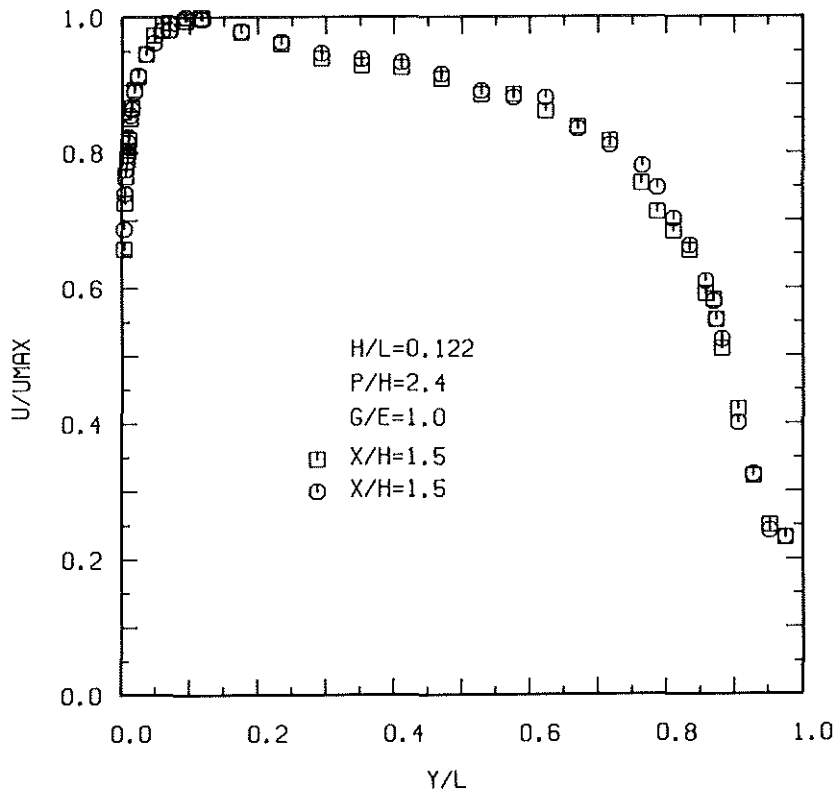


Fig.6.9-6.10: Velocity profiles (No.2)

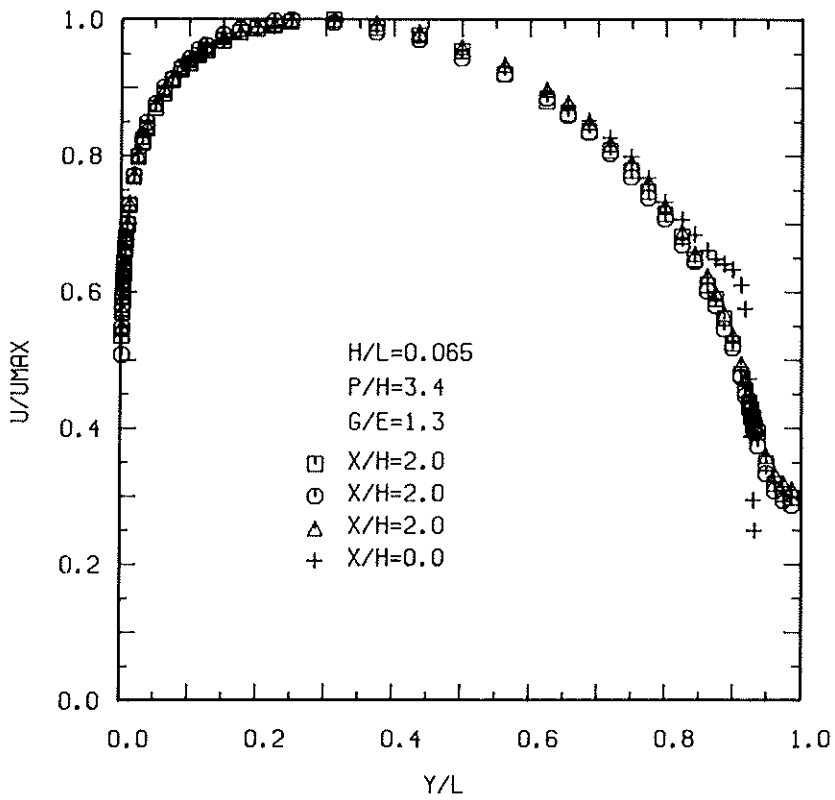
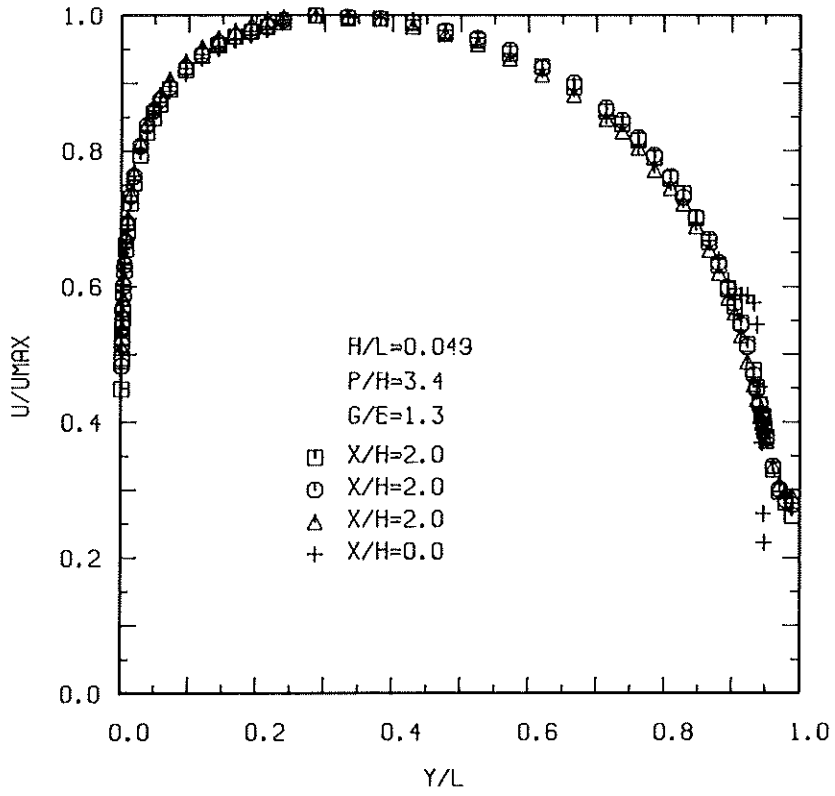


Fig.6.11-6.12: Velocity profiles (No.3)

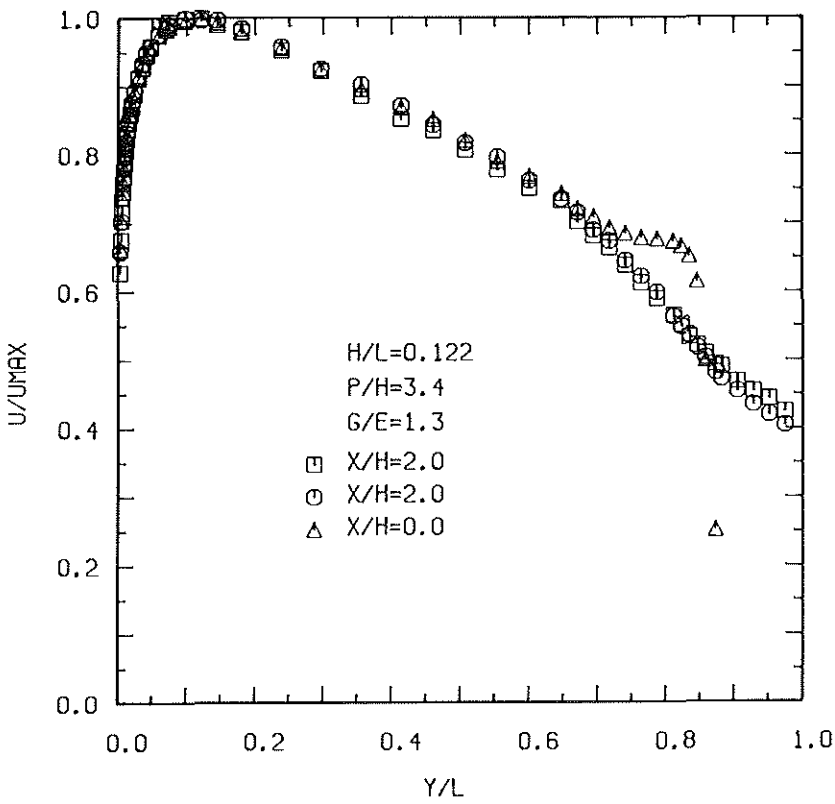
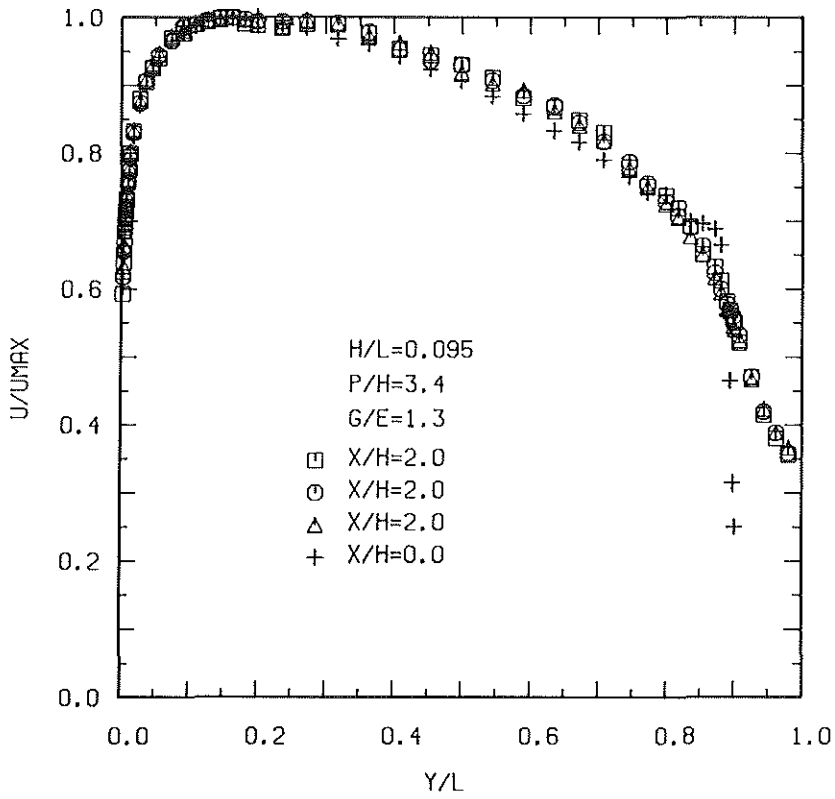


Fig.6.13-6.14: Velocity profiles (No.3)

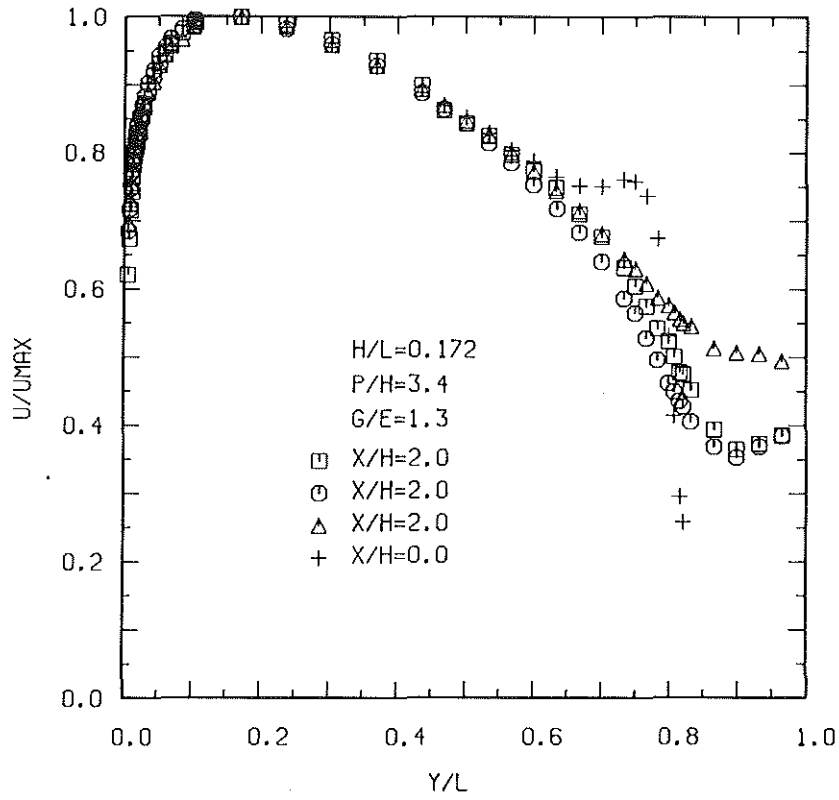


Fig.6.15: Velocity profiles (No.3)

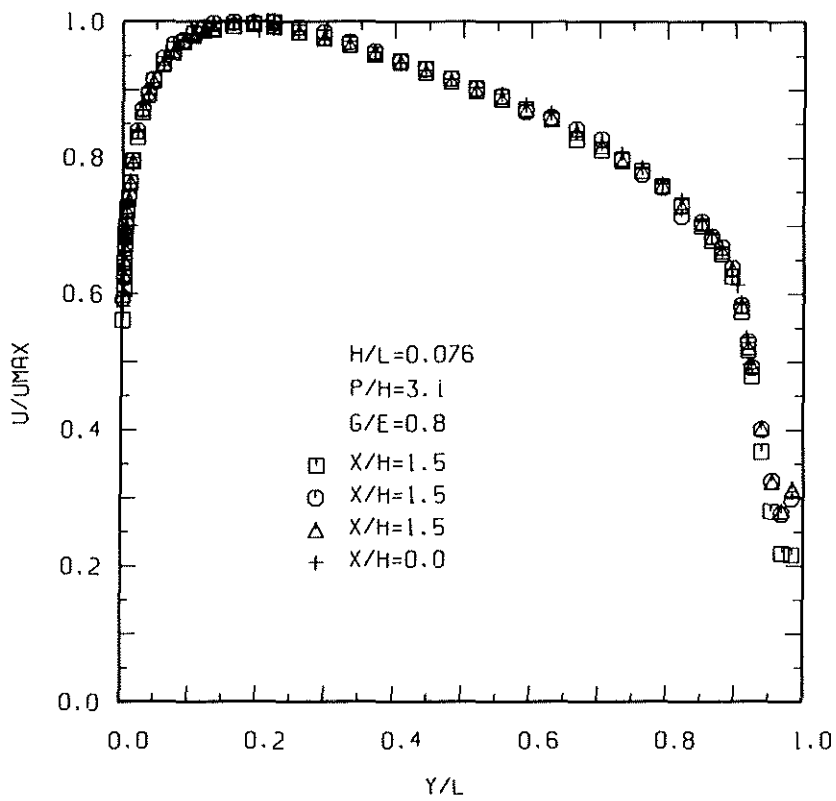
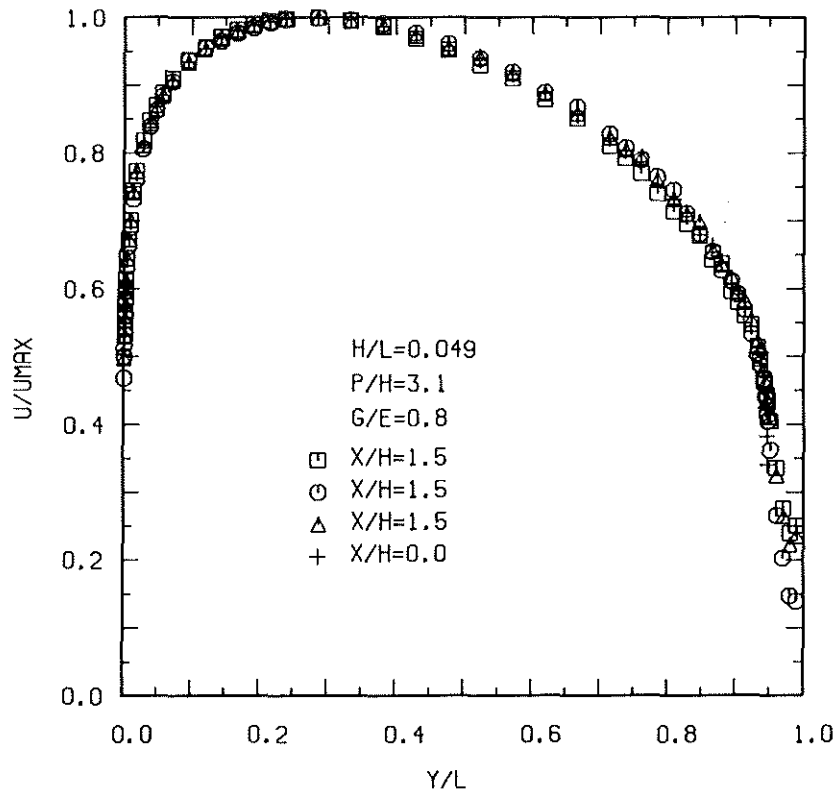


Fig.6.16-6.17: Velocity profiles (No.4)

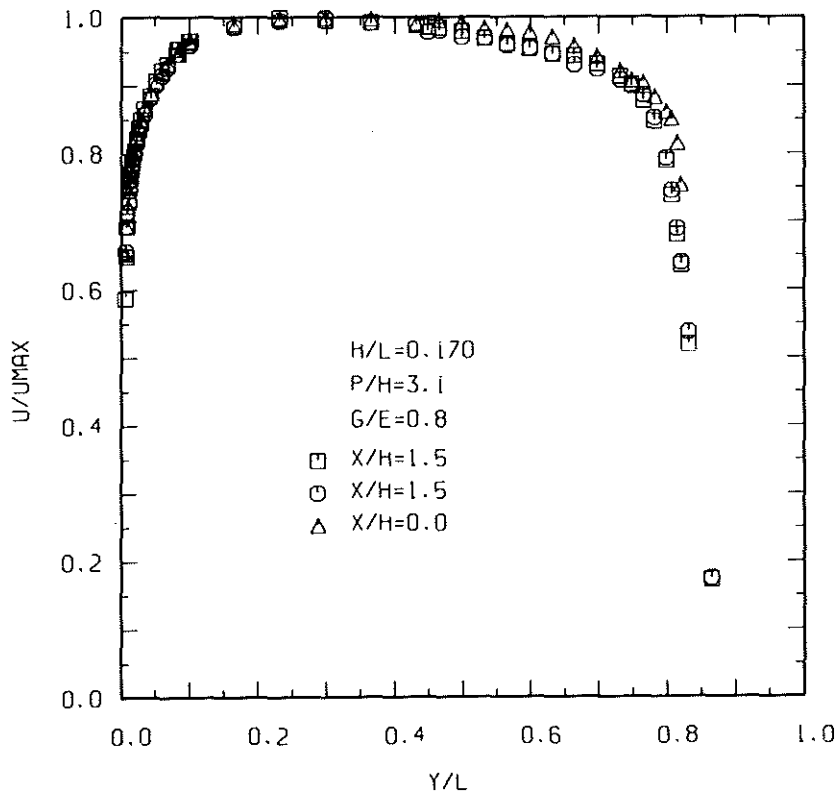
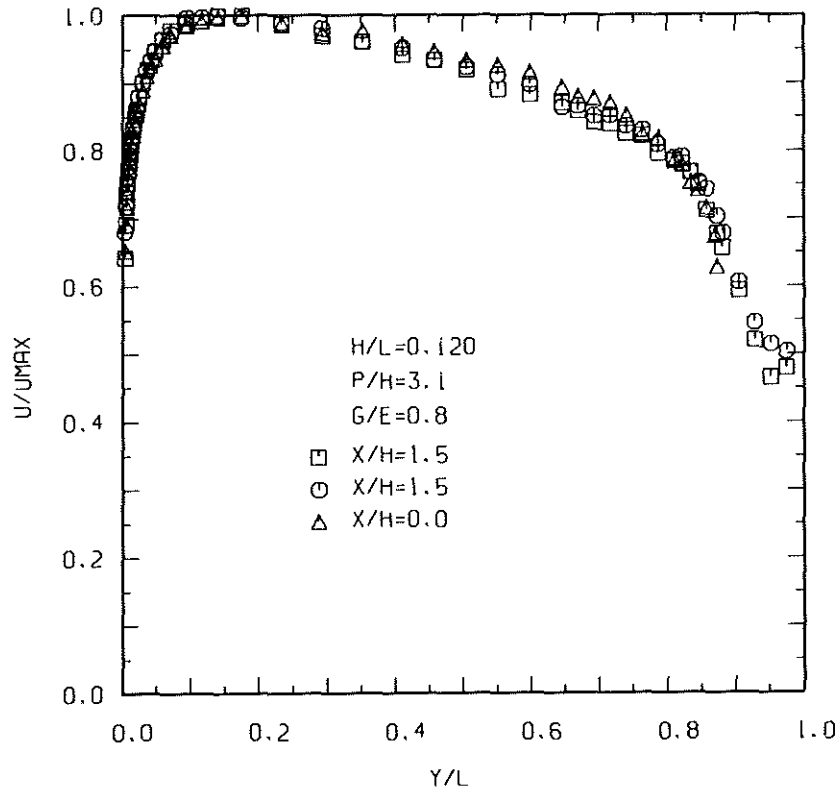


Fig.6.18-6.19: Velocity profiles (No.4)

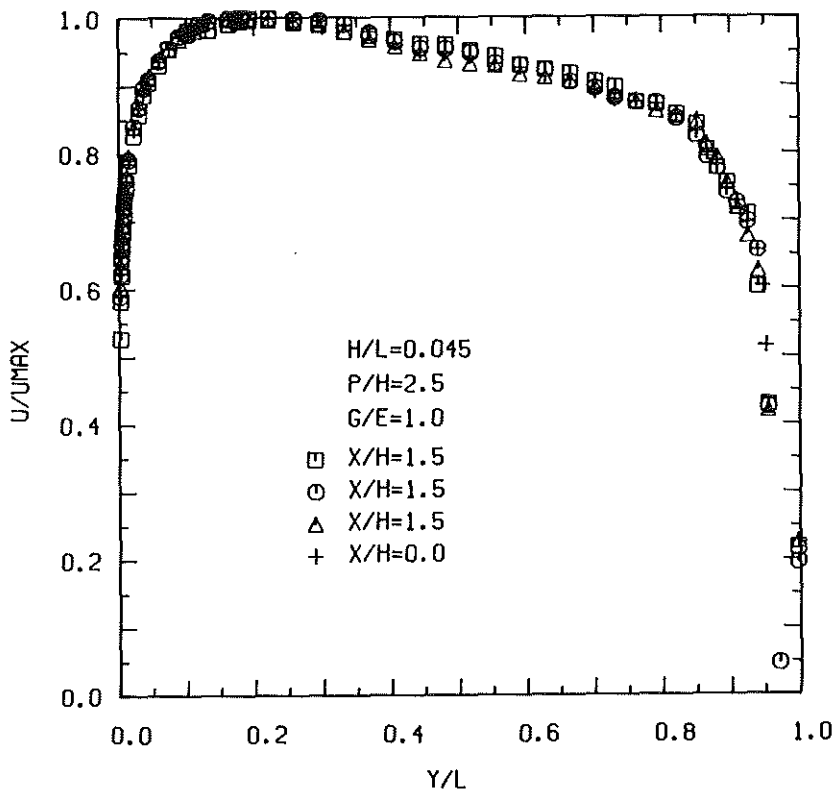
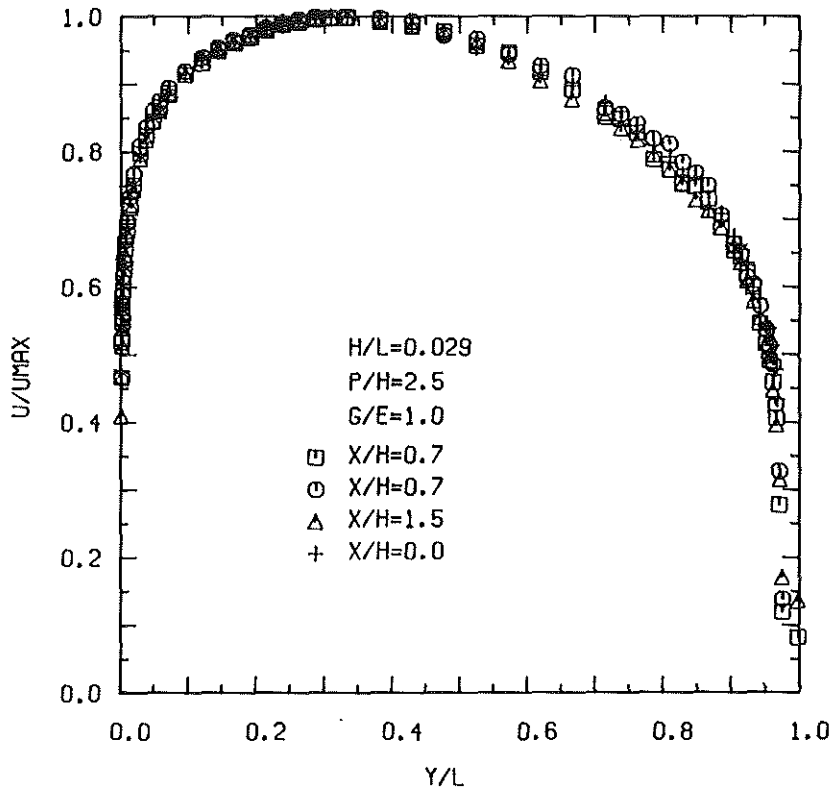


Fig.6.20-6.21: Velocity profiles (No.5)

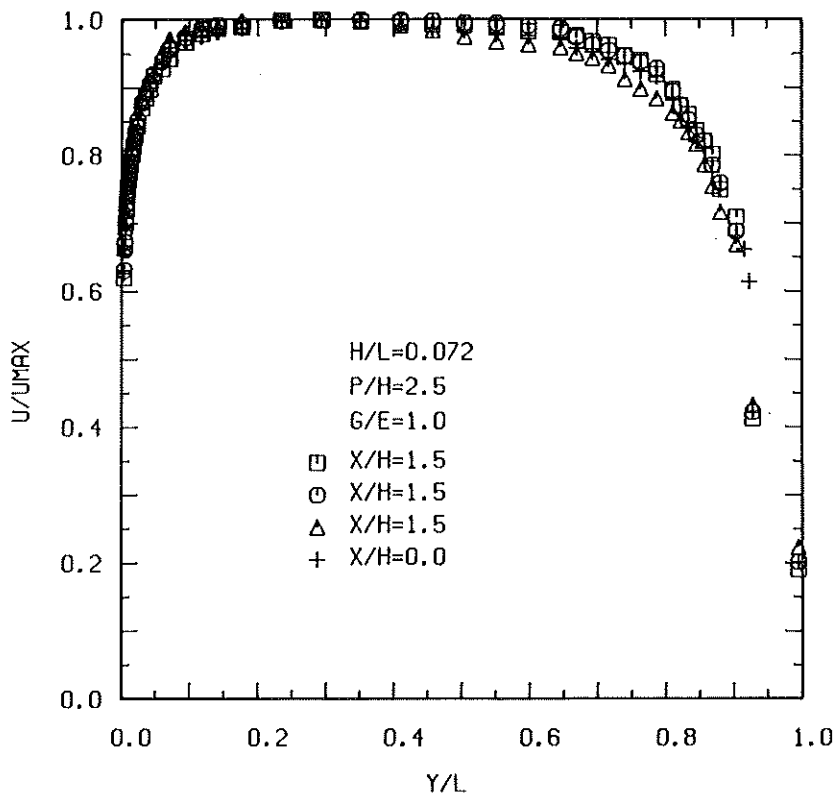
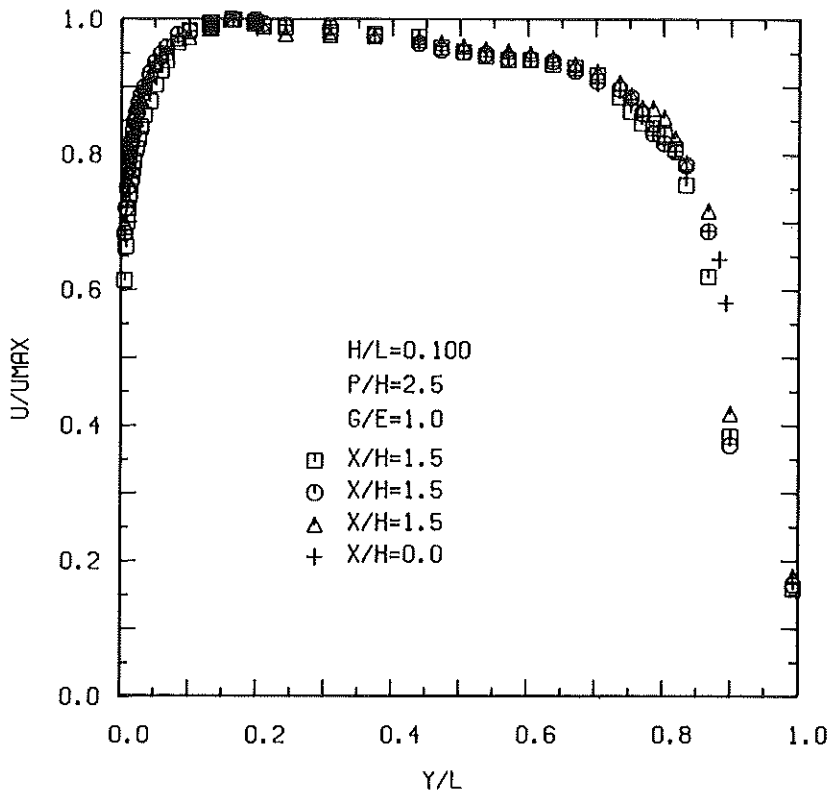


Fig.6.22-6.23: Velocity profiles (No.5)

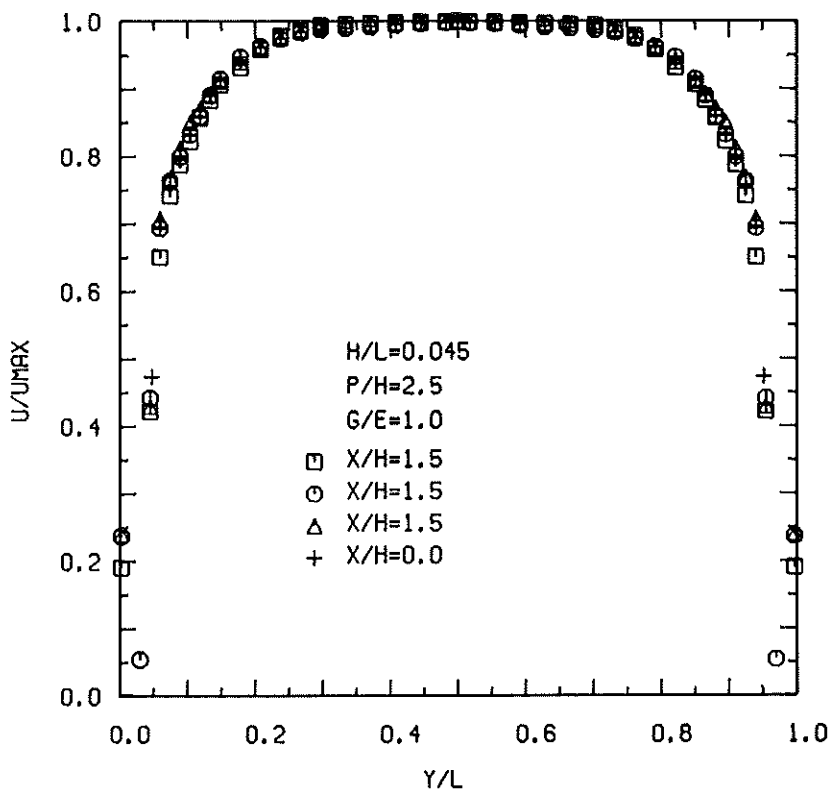
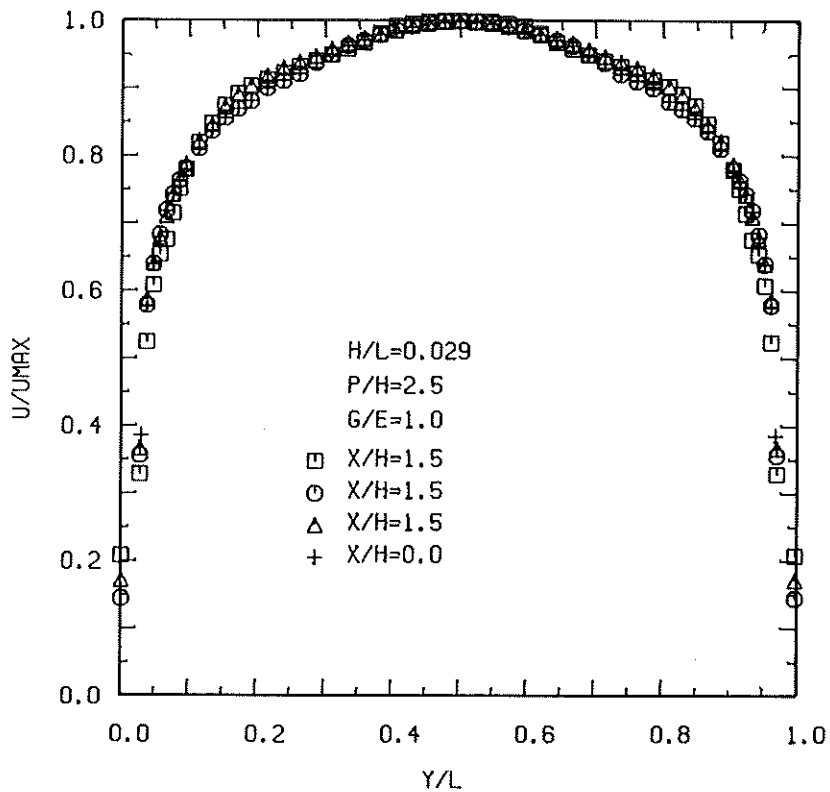


Fig.6.24-6.25: Velocity profiles (No.5 both walls rough)

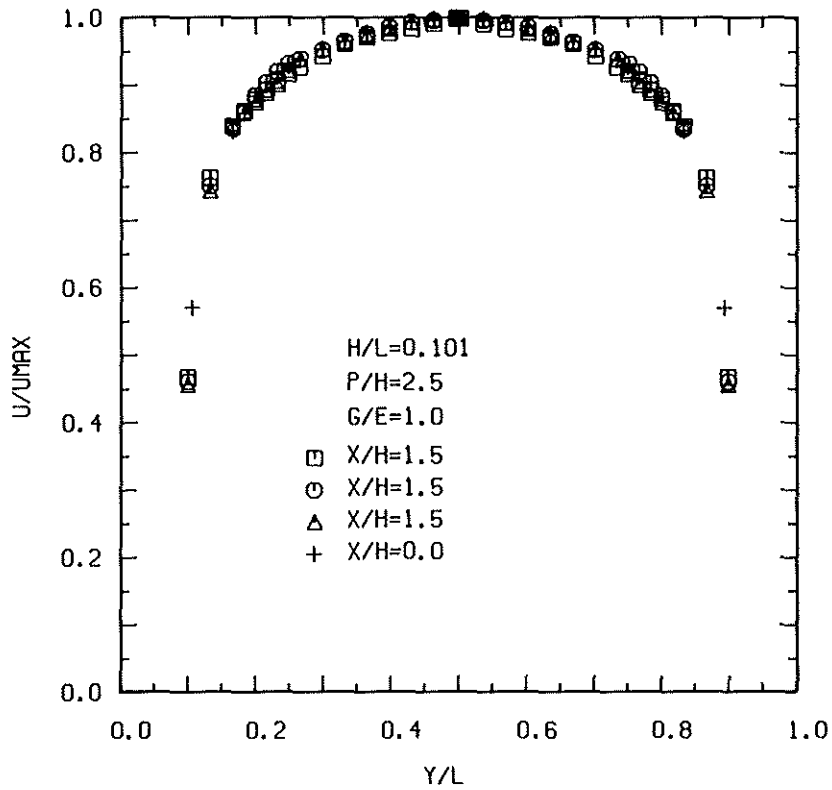
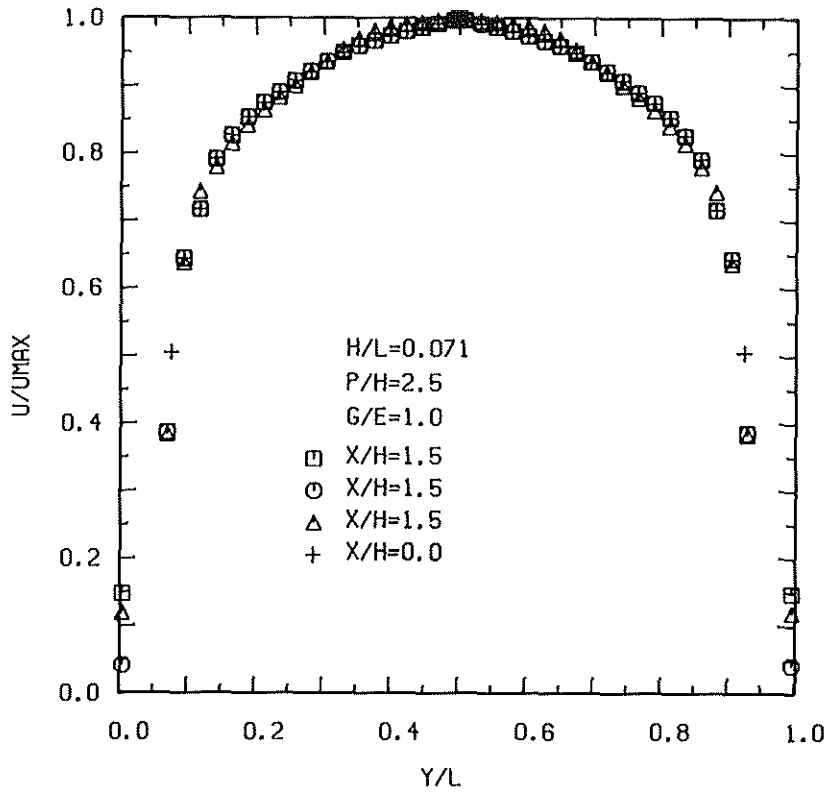


Fig.6.26-6.27: Velocity profiles (No.5 both walls rough)

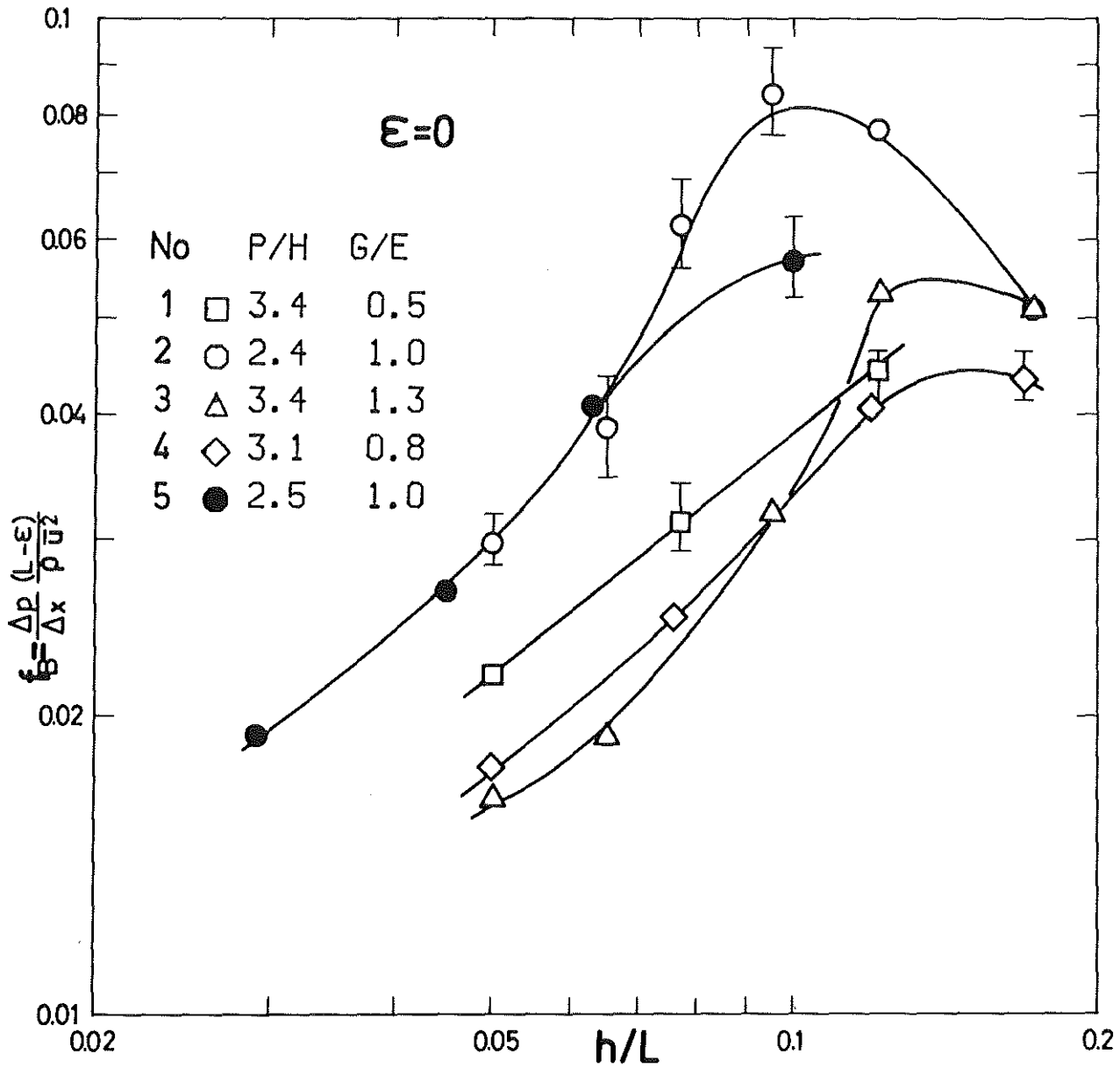


Fig.7: The total friction factor versus the relative roughness height ($\epsilon=0$)

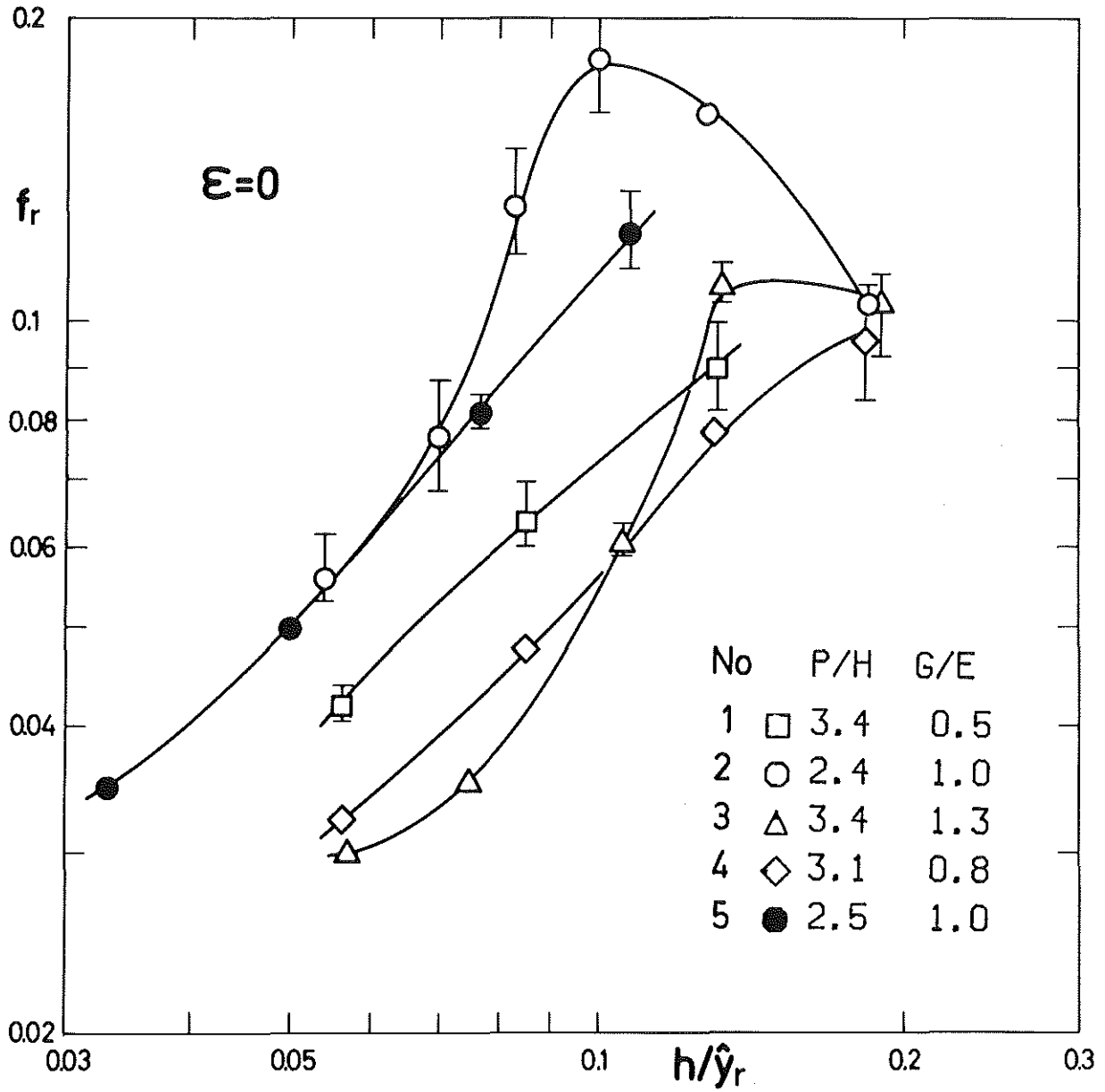


Fig.8: The friction factor of the rough zone versus the relative roughness height ($\epsilon=0$)

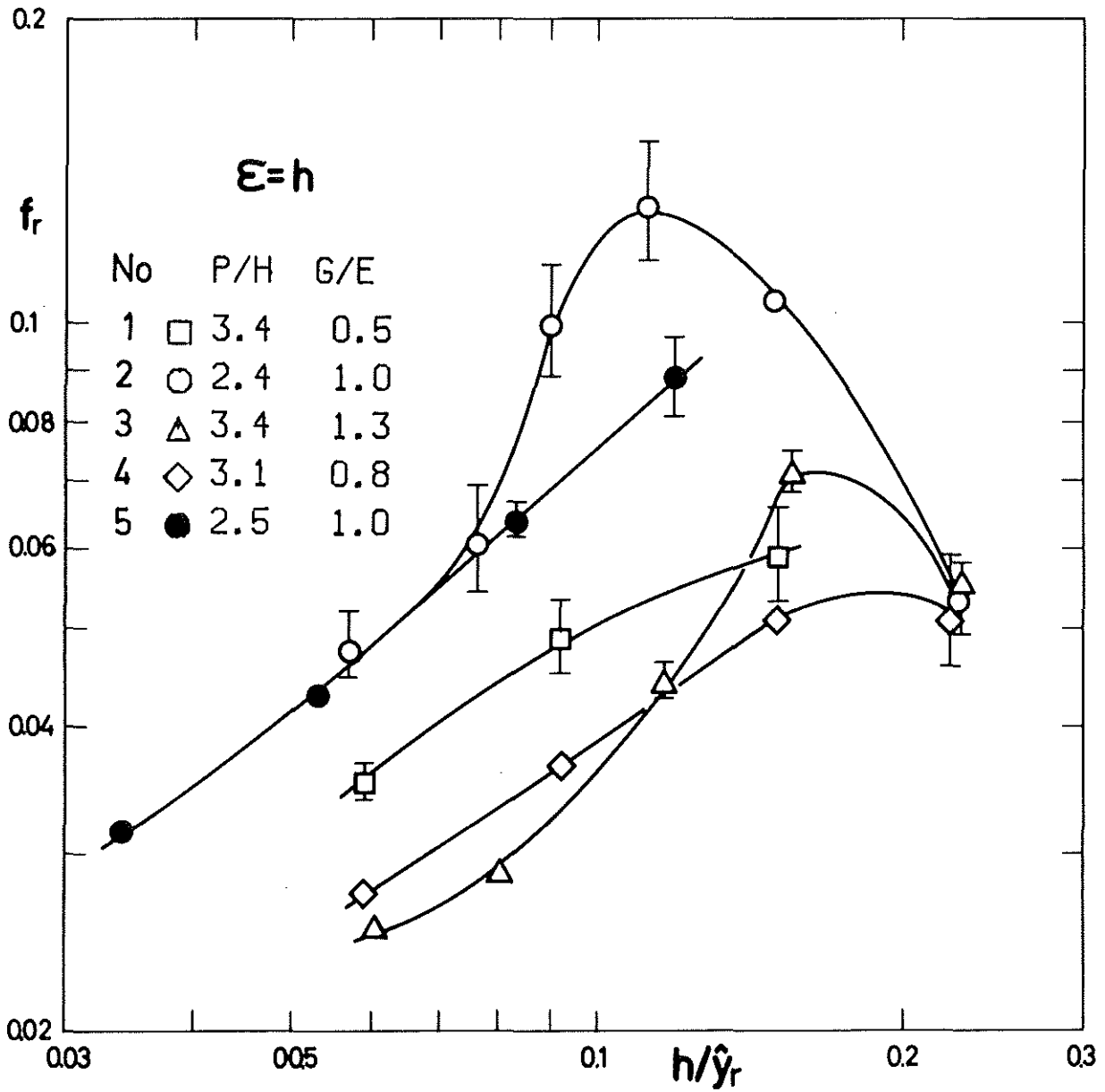


Fig.9: The friction factor of the rough zone versus the relative roughness height ($\epsilon=0$)

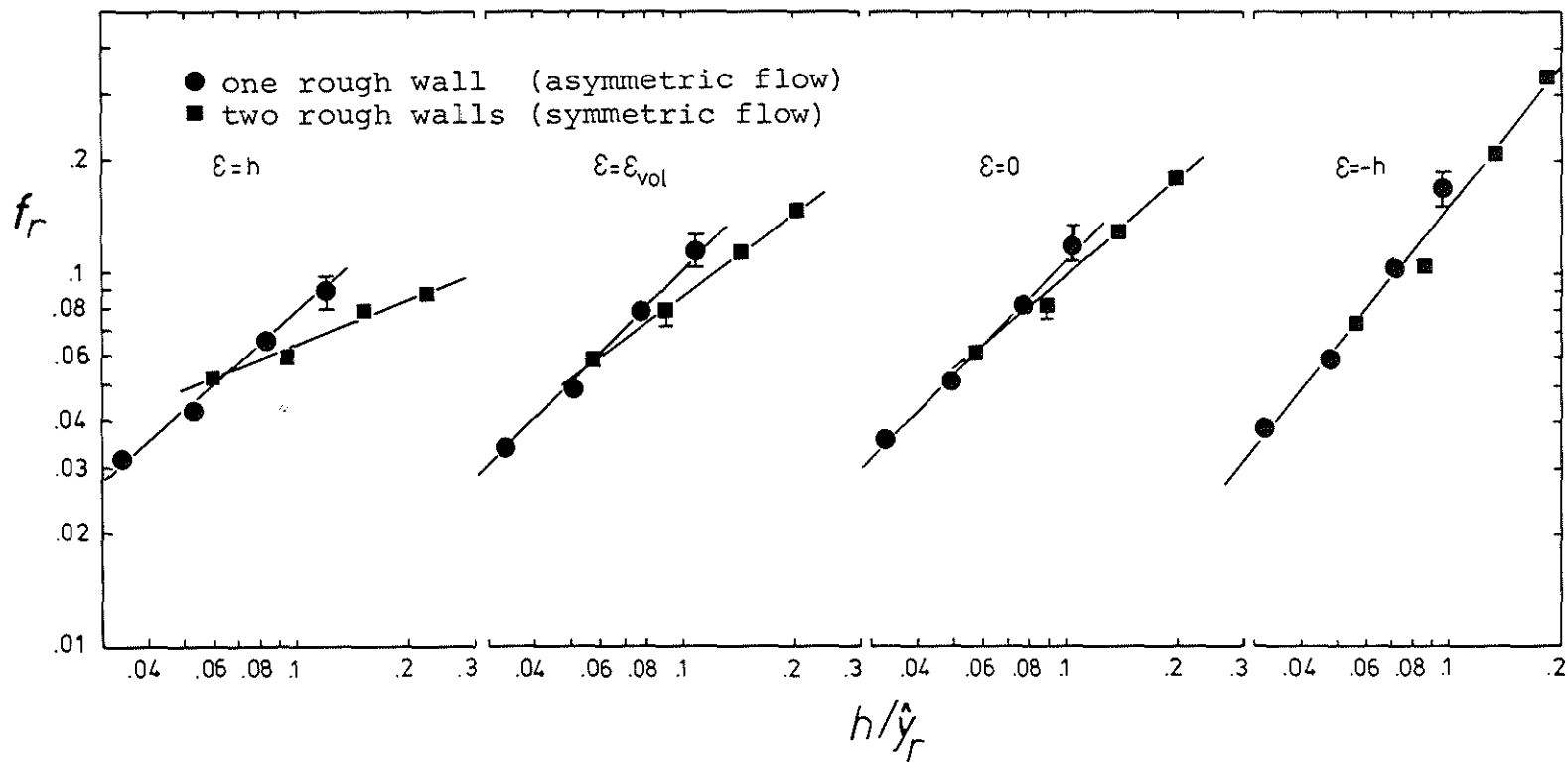


Fig.10: The friction factor of the rough zone in symmetrical and asymmetrical flow for different definitions of the hydraulic diameter.

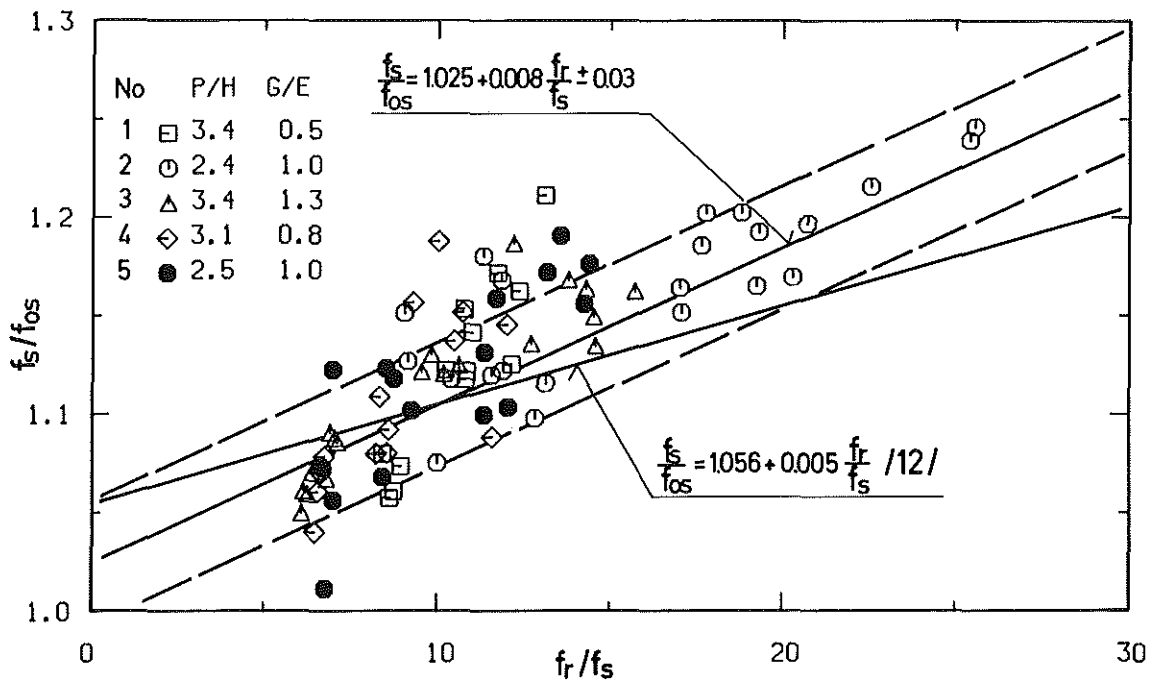


Fig.11: Variation of the friction factor of the smooth zone f_s with the friction factor of the rough zone f_r ($\epsilon = \epsilon_{vol}$)

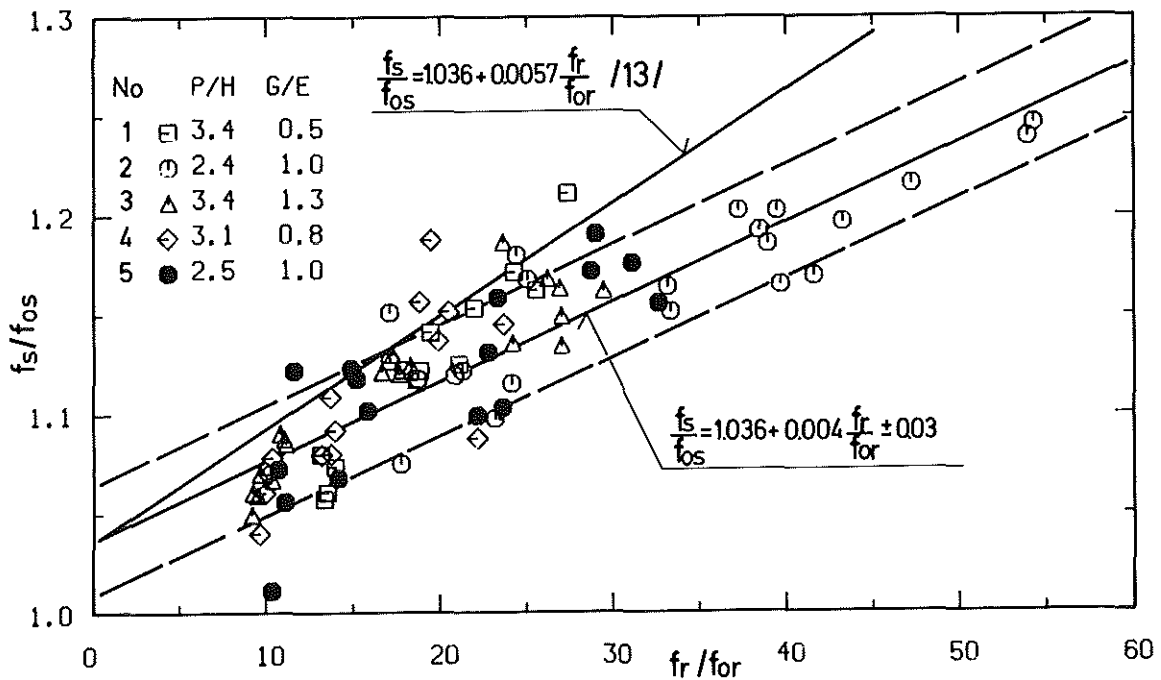


Fig.12: Variation of the friction factor of the smooth zone f_s with the friction factor of the rough zone f_r ($\epsilon = \epsilon_{vol}$)

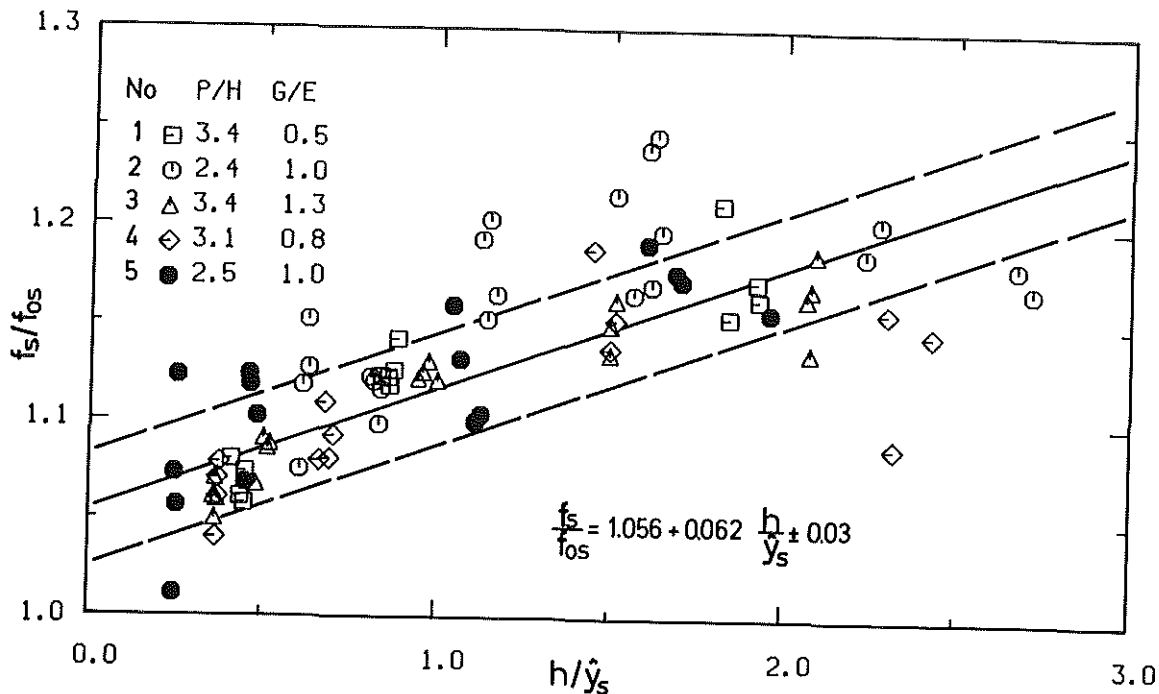


Fig.13: Variation of the friction factor of the smooth zone f_s with the relative length of the smooth zone.

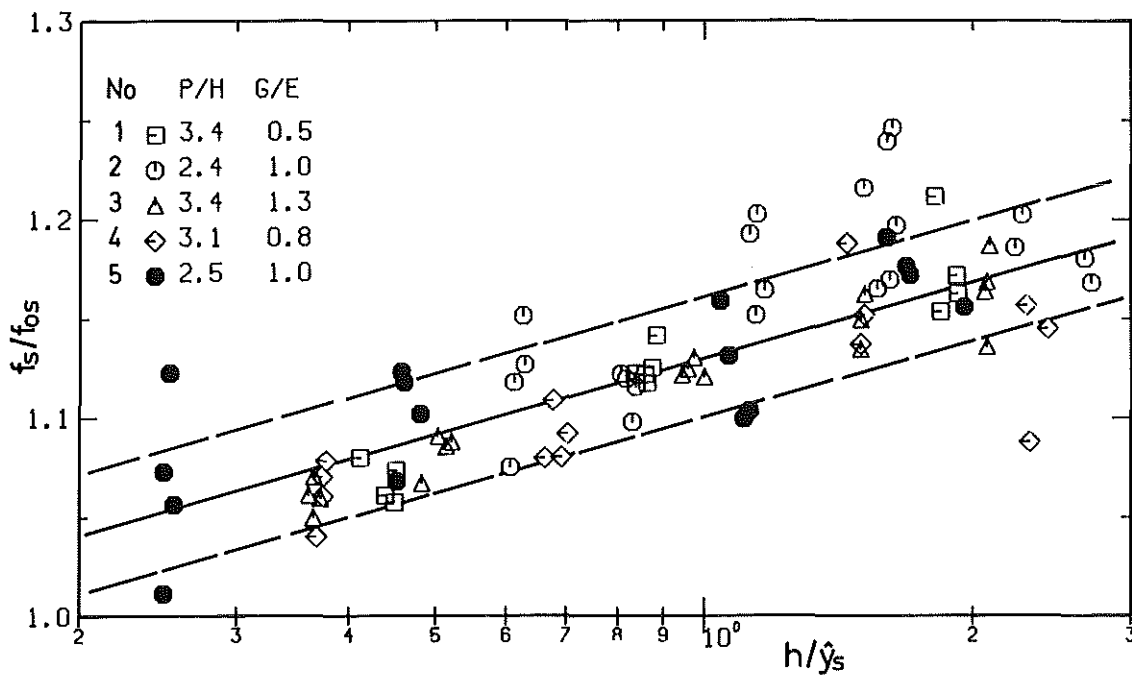


Fig.14: Variation of the friction factor of the smooth zone f_s with the relative length of the smooth zone.

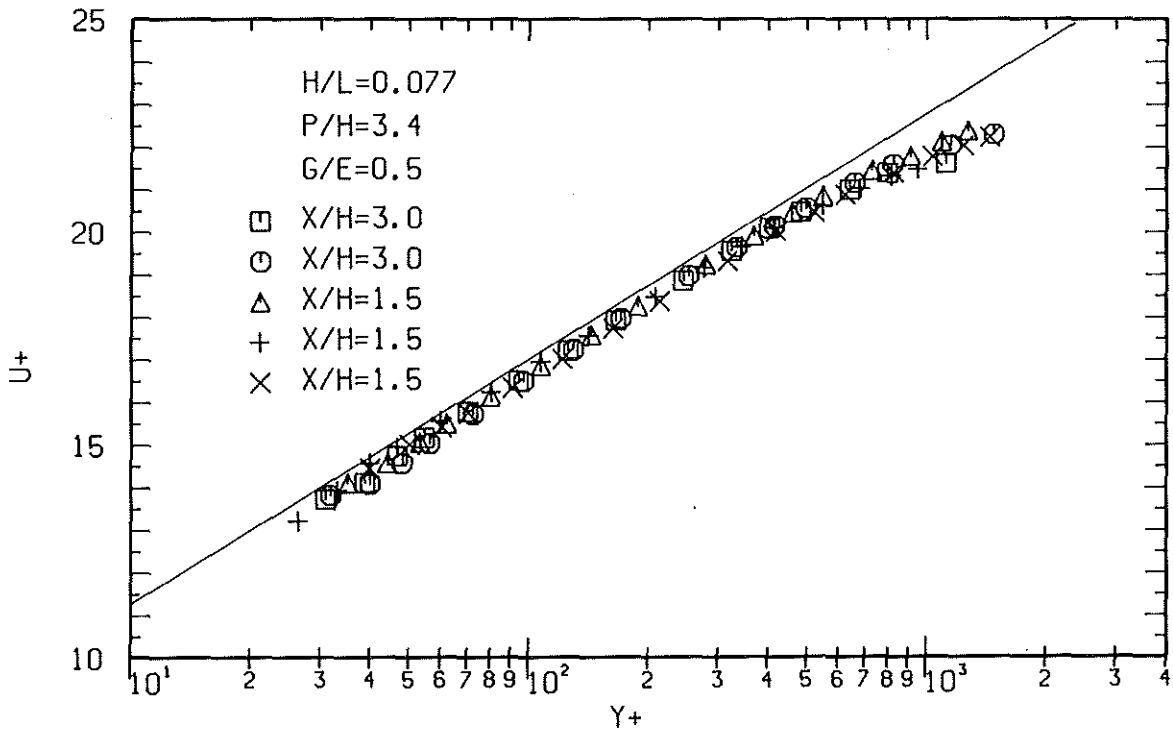
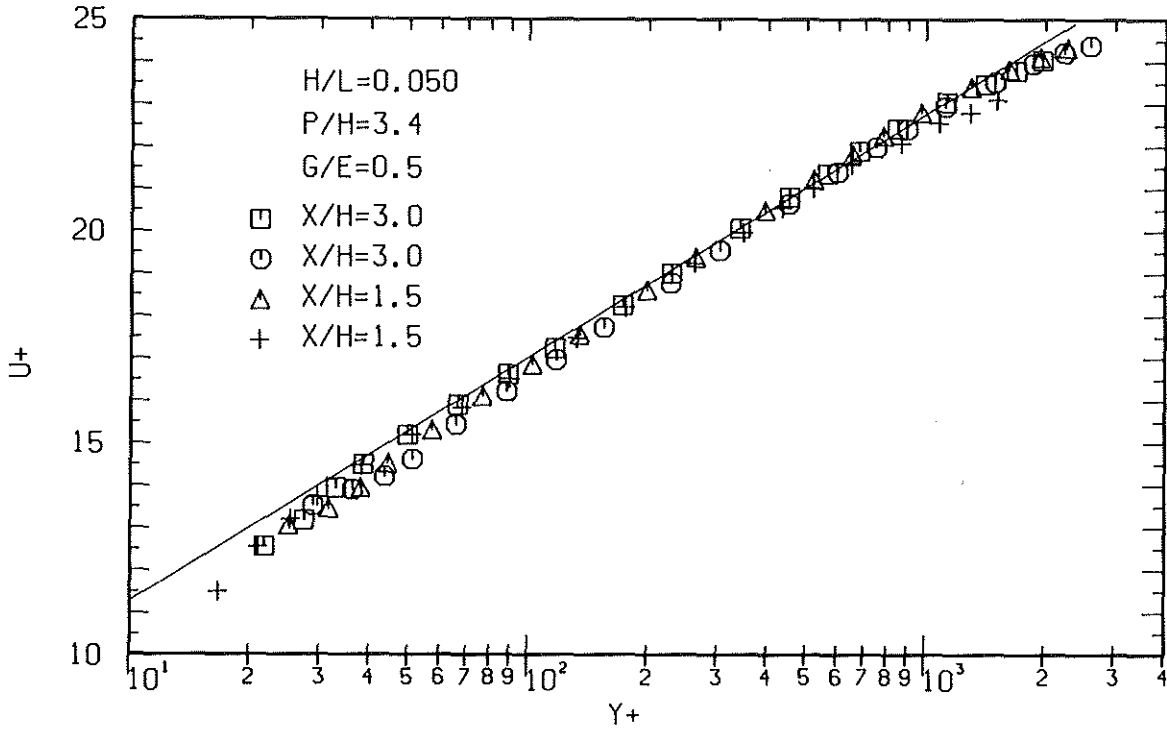


Fig.15.1-15.2: Non-dimensional velocity profiles at the smooth wall with roughness No.1 at opposite wall

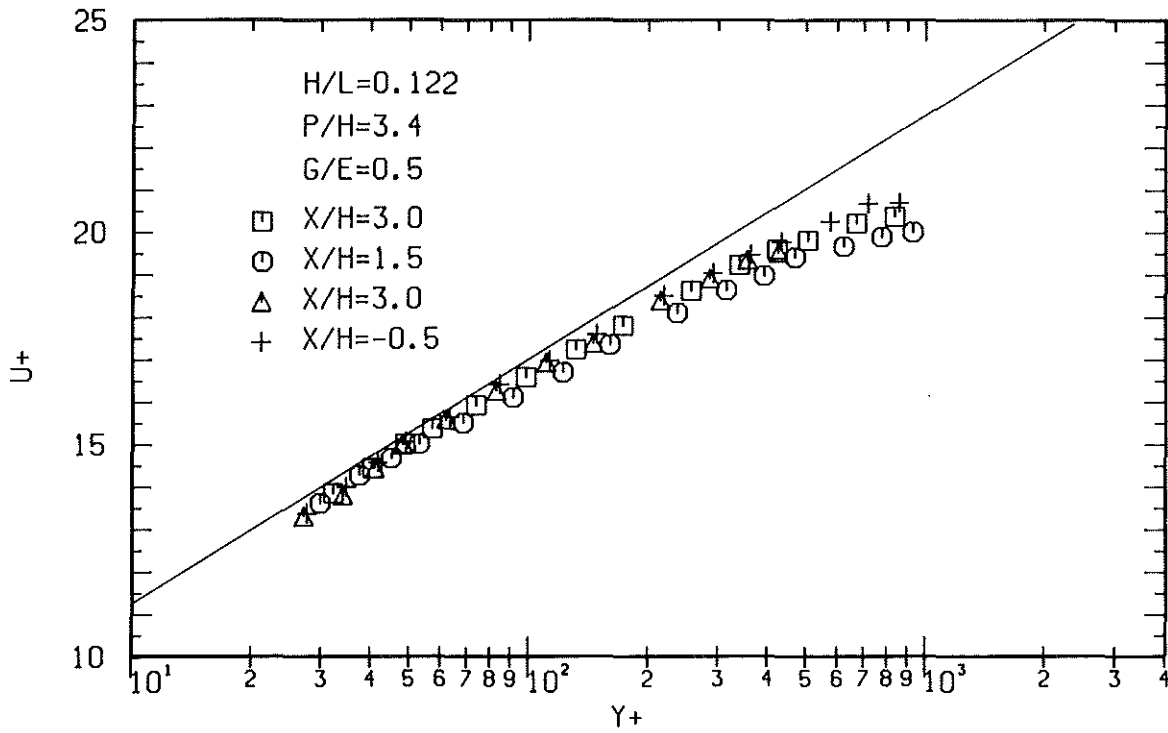


Fig.15.3: Non-dimensional velocity profiles at the smooth wall with roughness No.1 at opposite wall

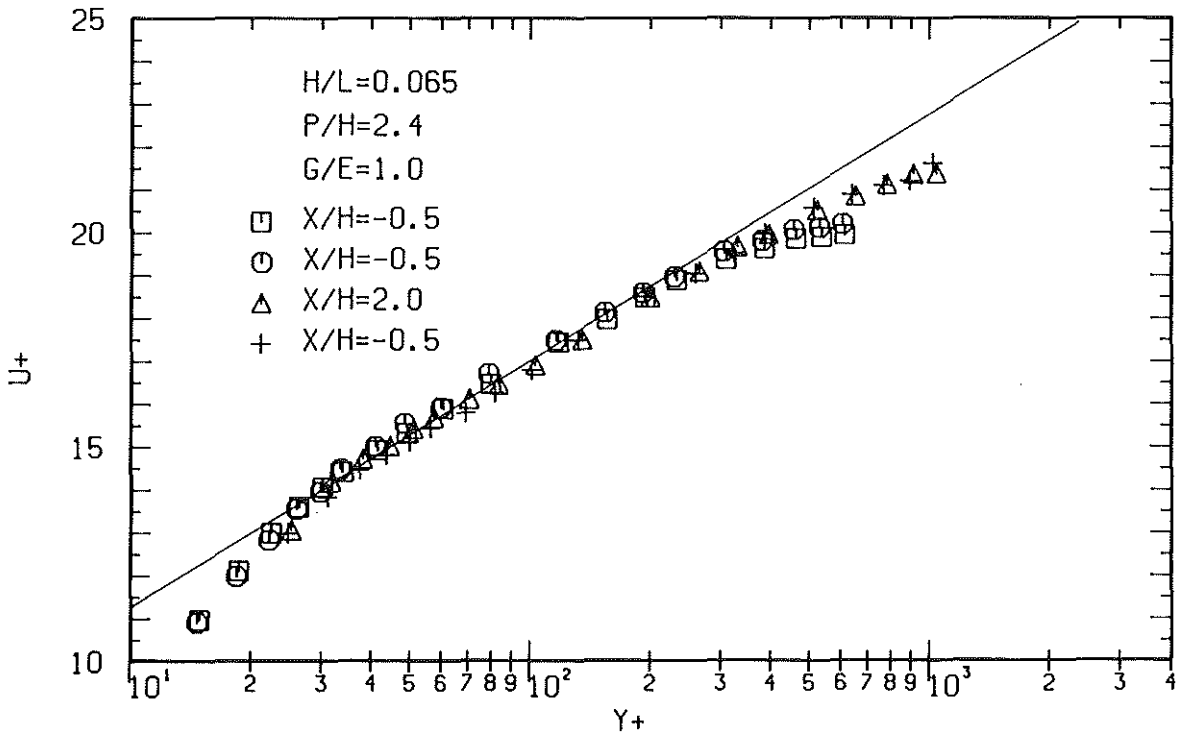
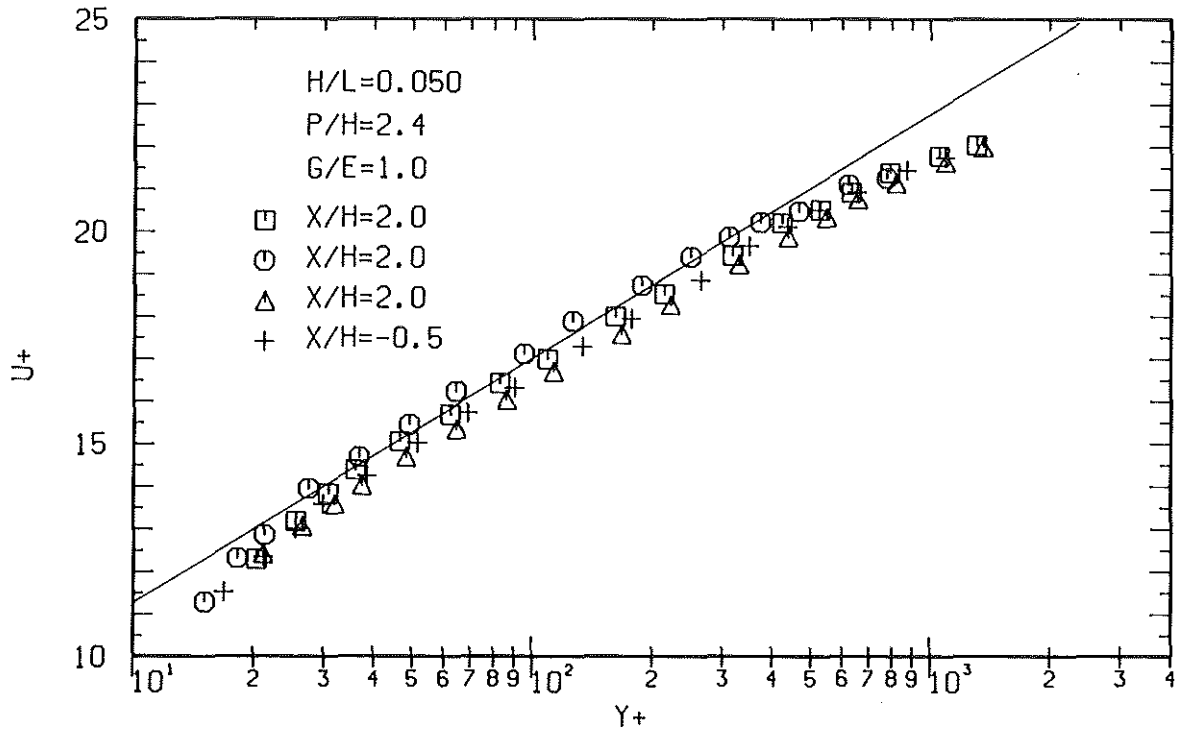


Fig.15.4-15.5: Non-dimensional velocity profiles at the smooth wall with roughness No.2 at opposite wall

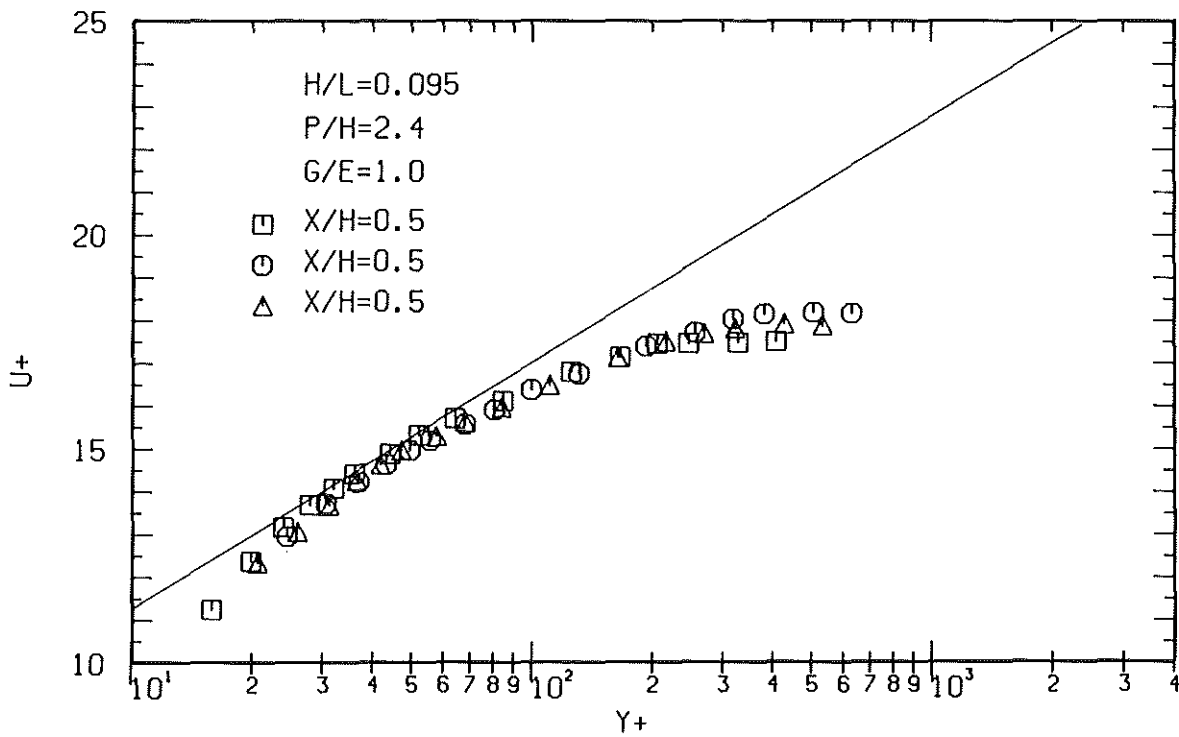
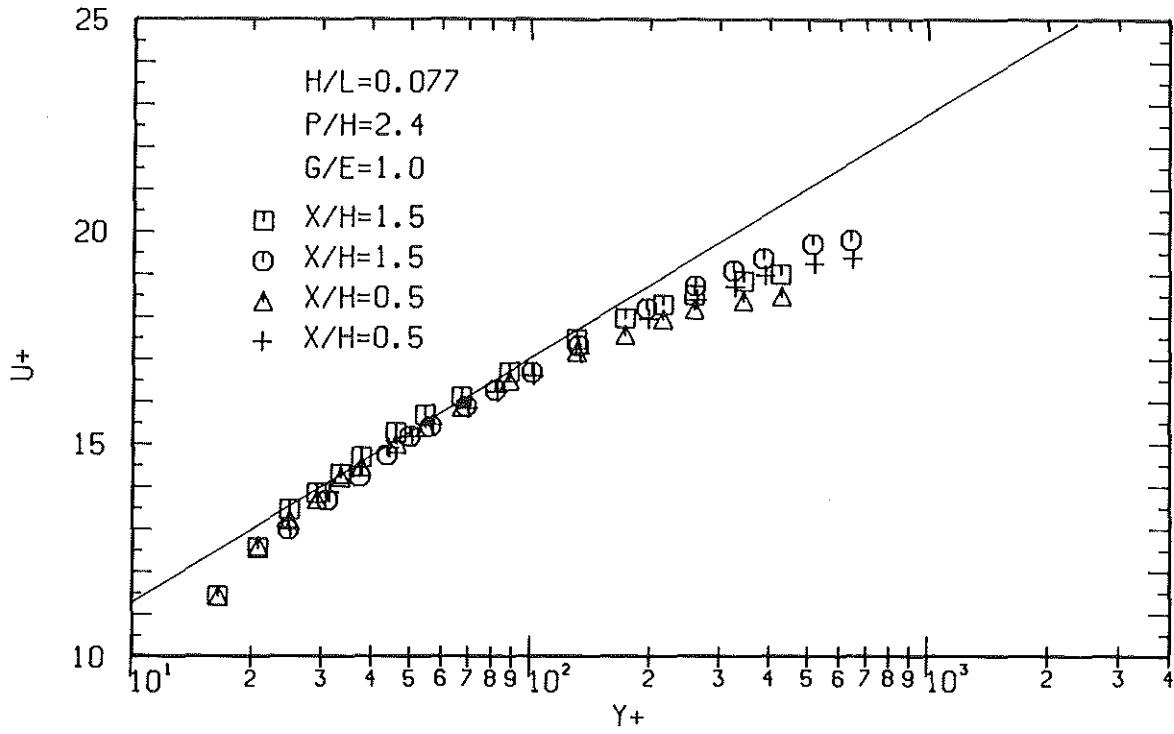


Fig.15.6-15.7: Non-dimensional velocity profiles at the smooth wall with roughness No.2 at opposite wall

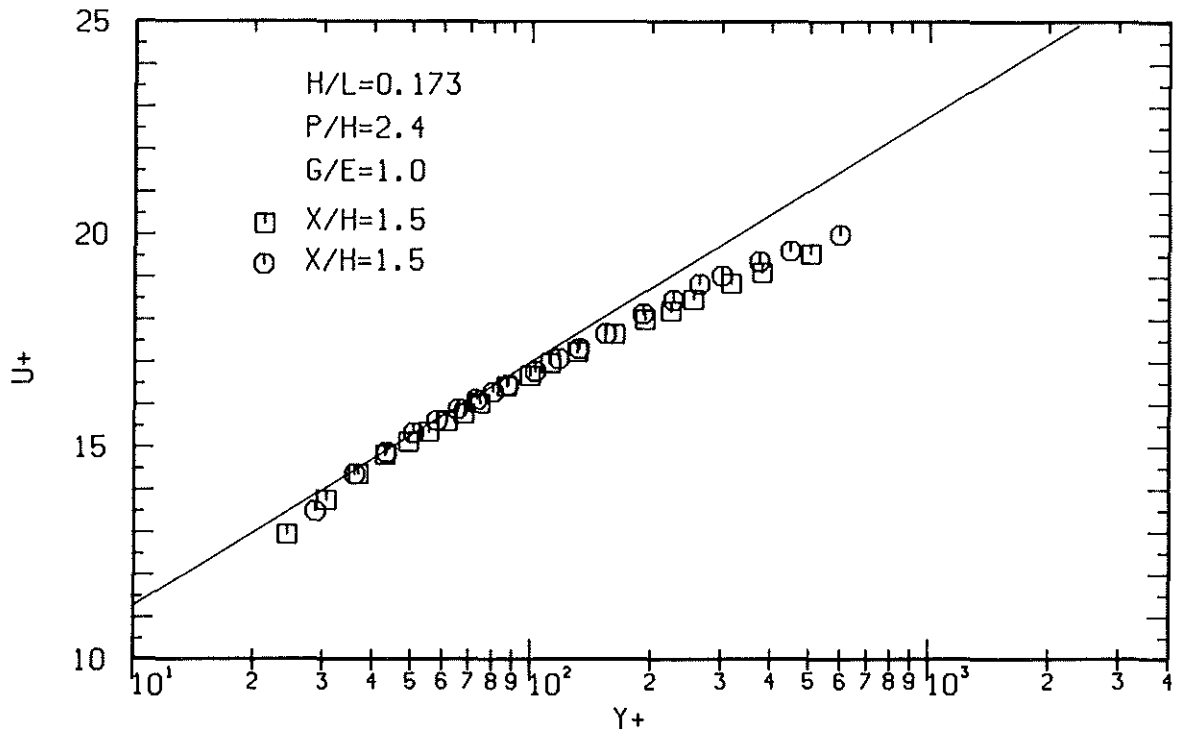
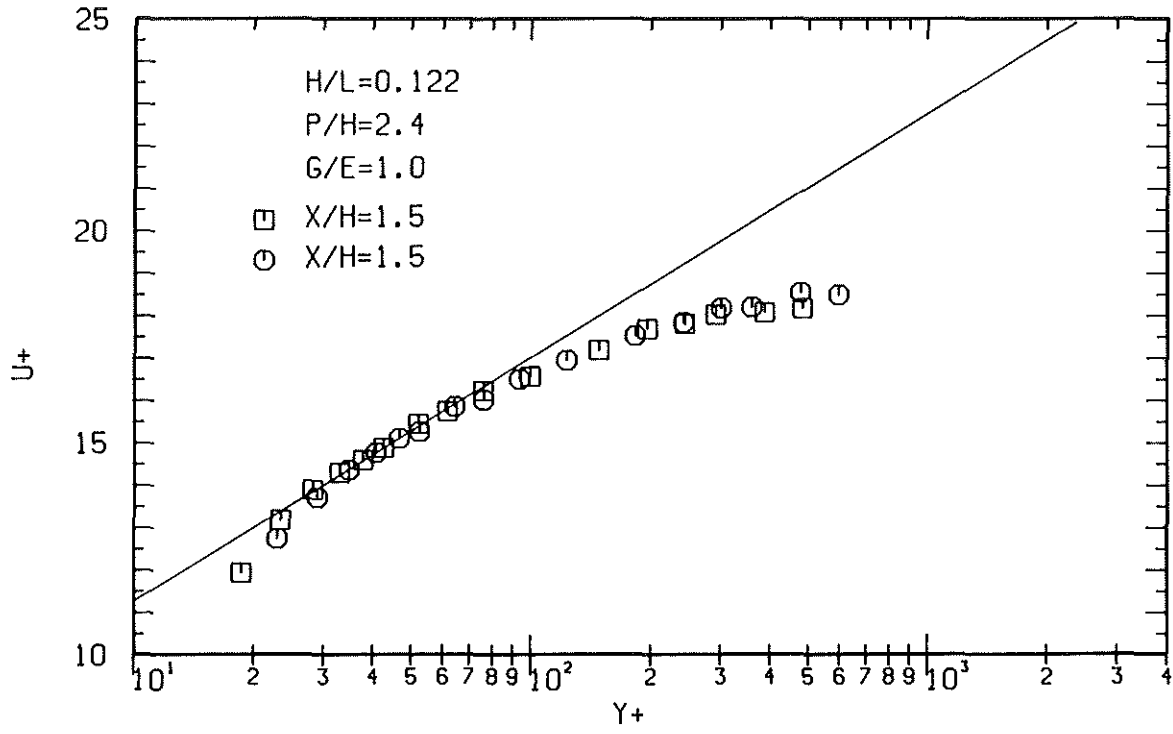


Fig.15.8-15.9: Non-dimensional velocity profiles at the smooth wall with roughness No.2 at opposite wall

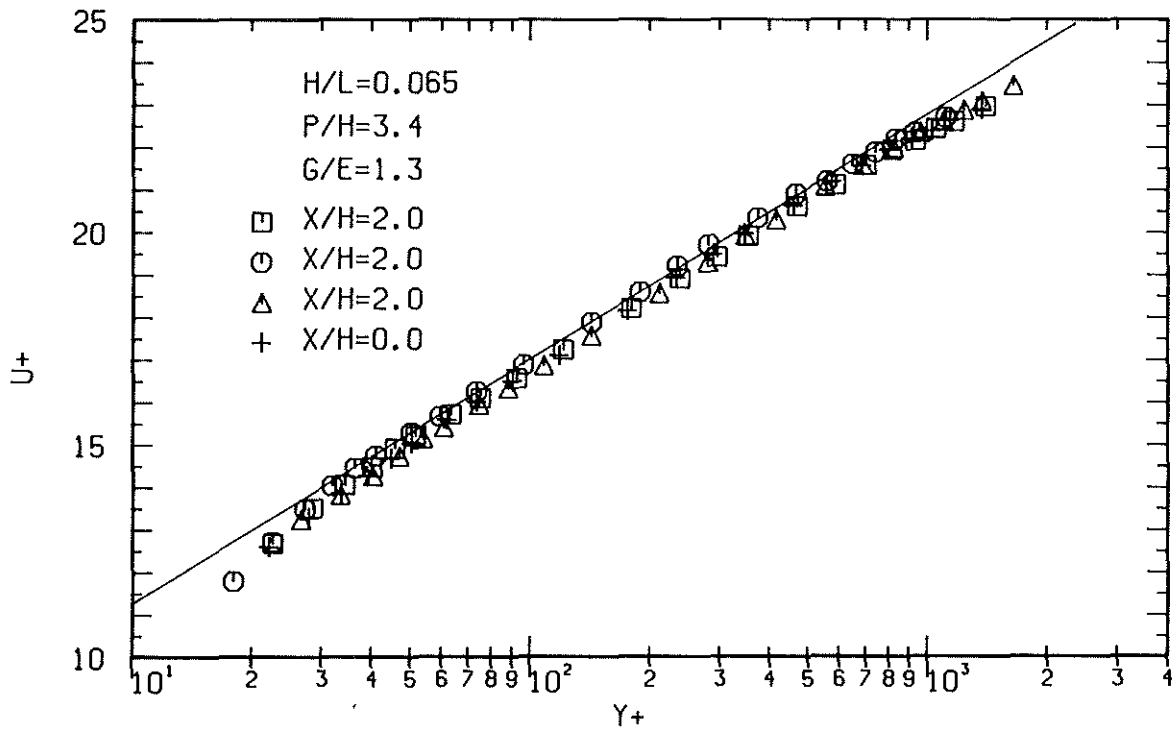
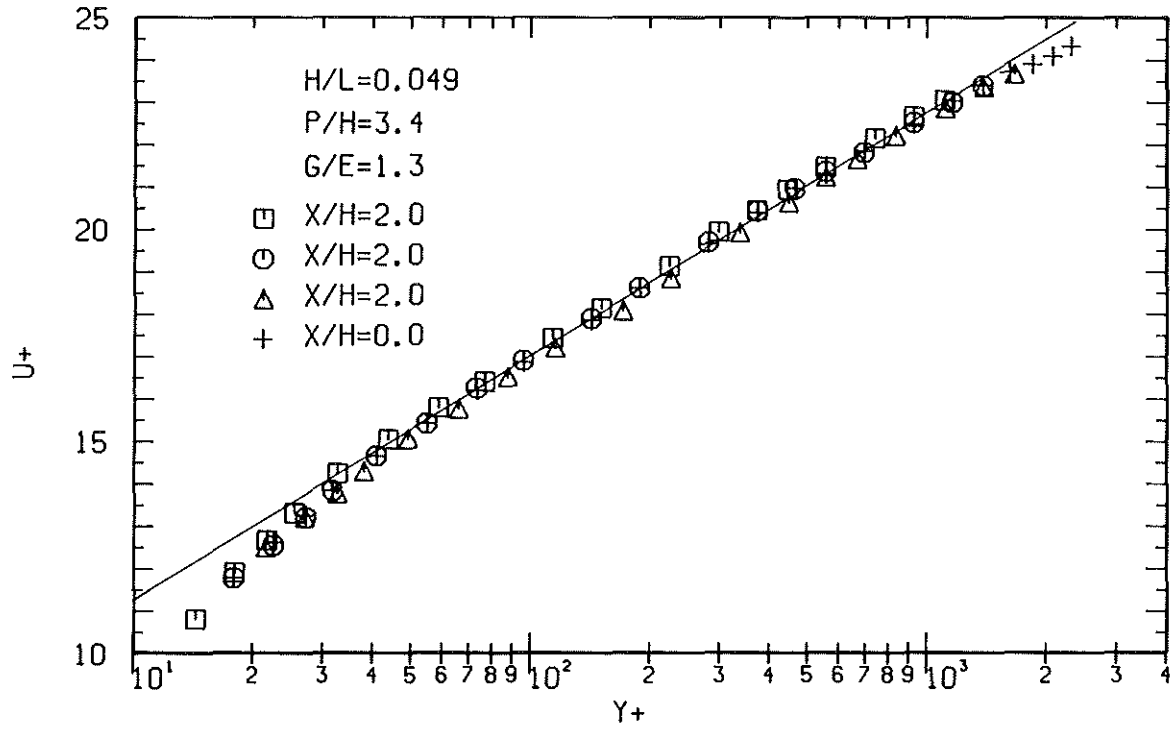


Fig.15.10-15.11: Non-dimensional velocity profiles at the smooth wall with roughness No.3 at opposite wall

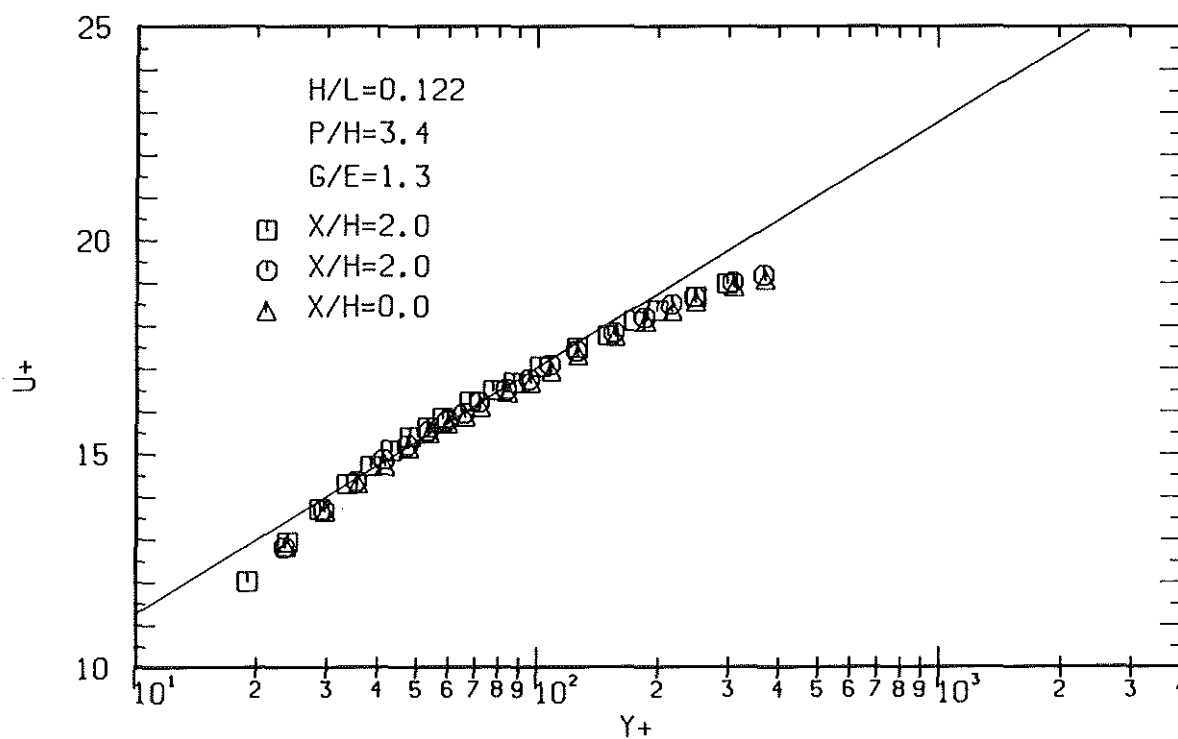
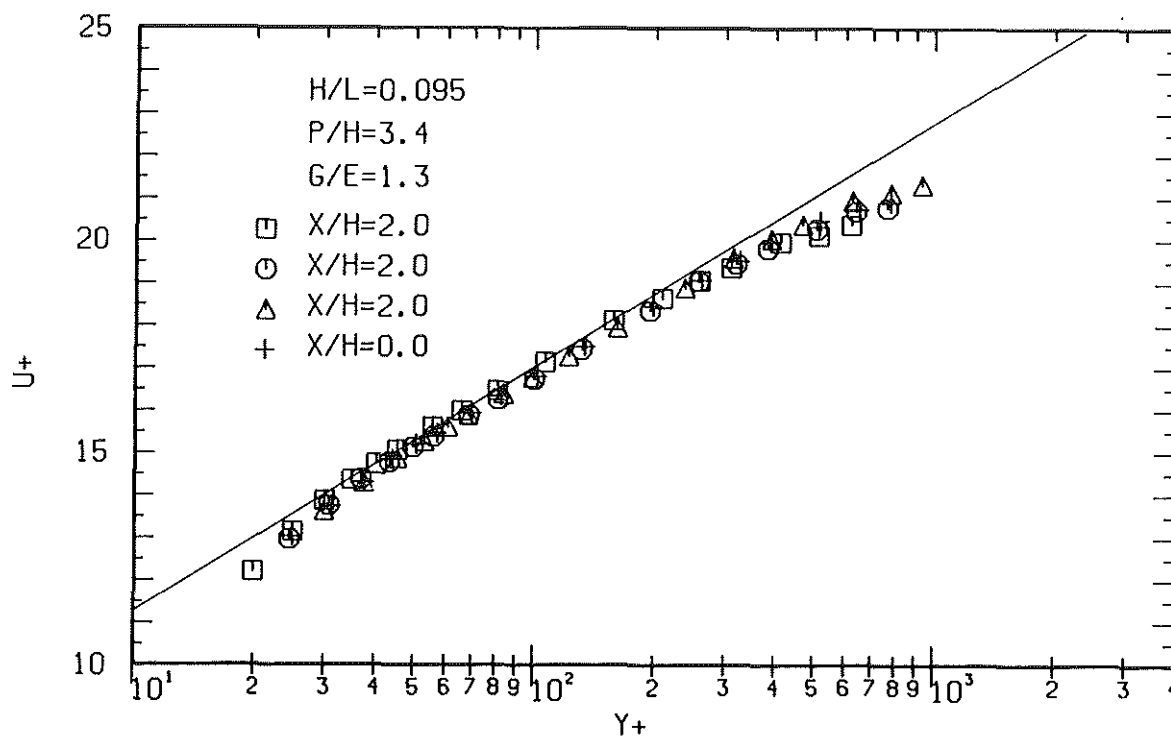


Fig.15.12-15.13: Non-dimensional velocity profiles at the smooth wall with roughness No.3 at opposite wall

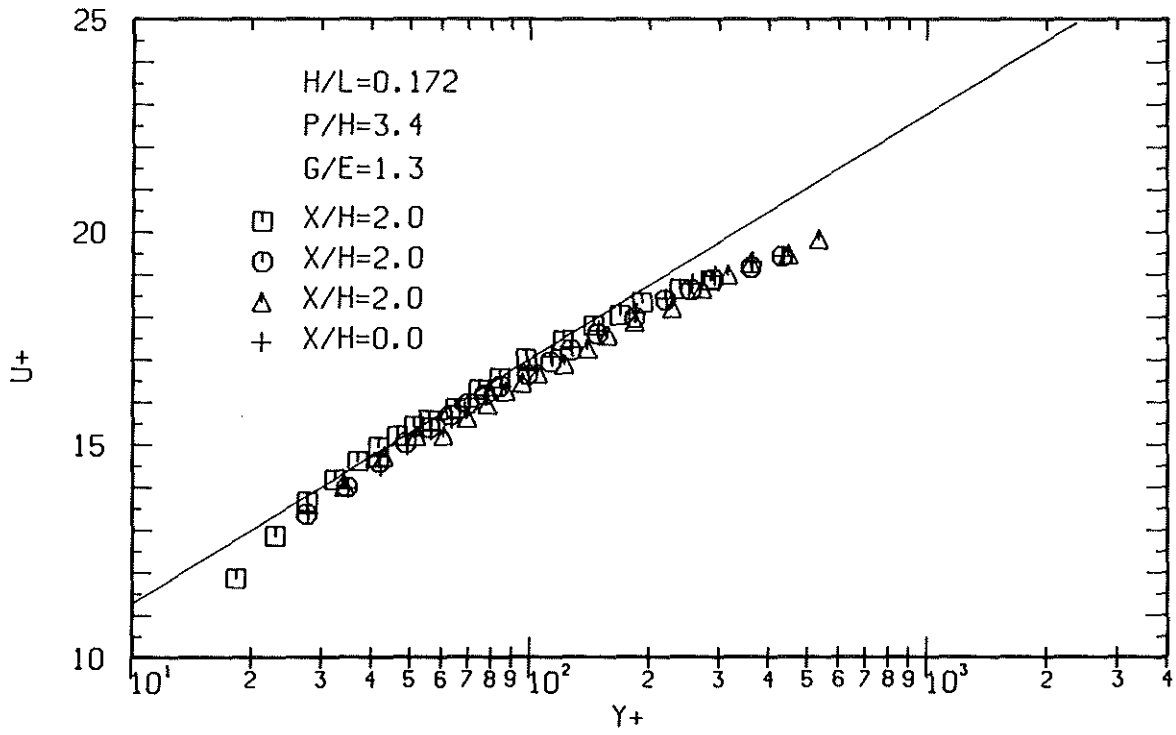


Fig.15.14: Non-dimensional velocity profiles at the smooth wall with roughness No.3 at opposite wall

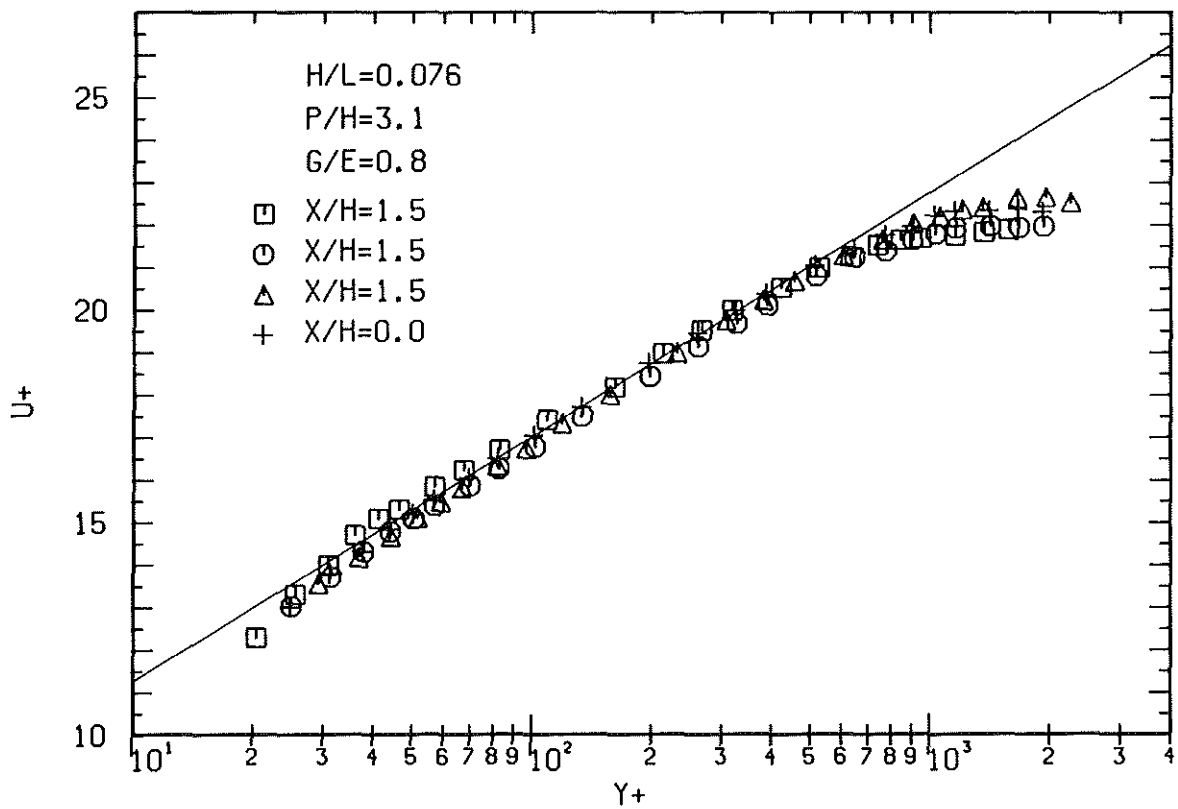
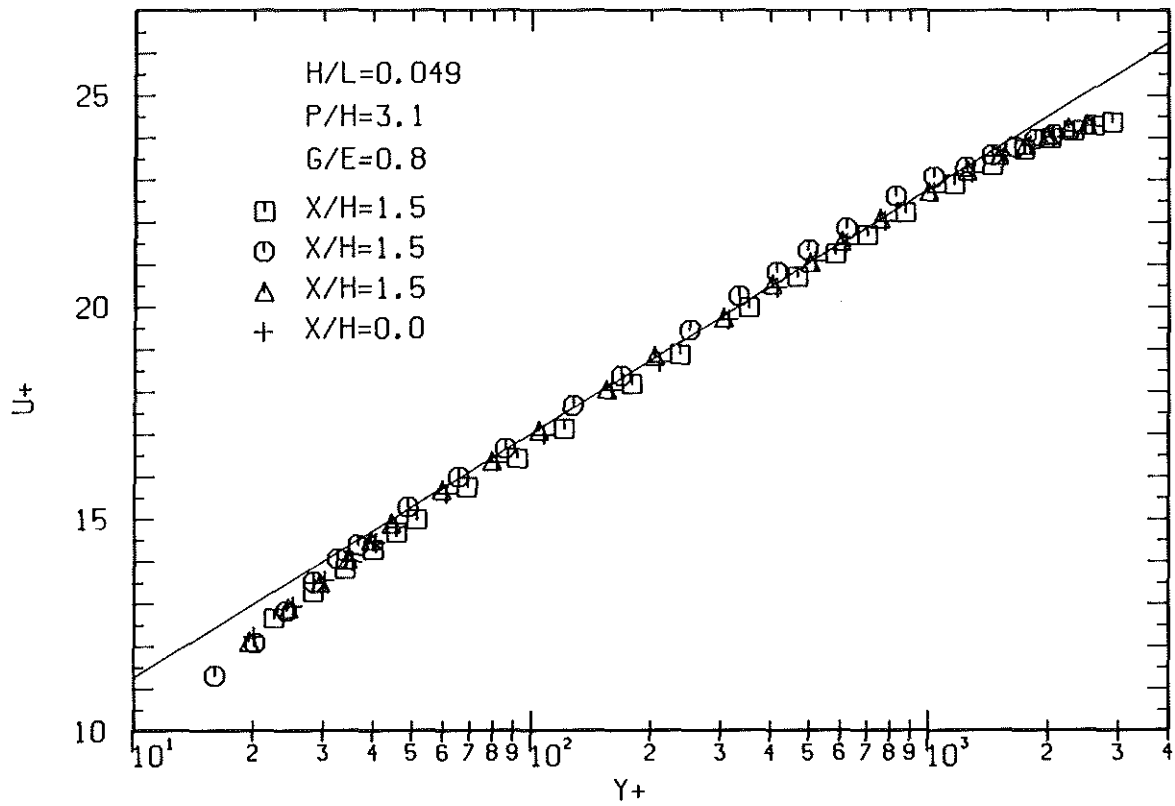


Fig.15.15-15.16: Non-dimensional velocity profiles at the smooth wall with roughness No.4 at opposite wall

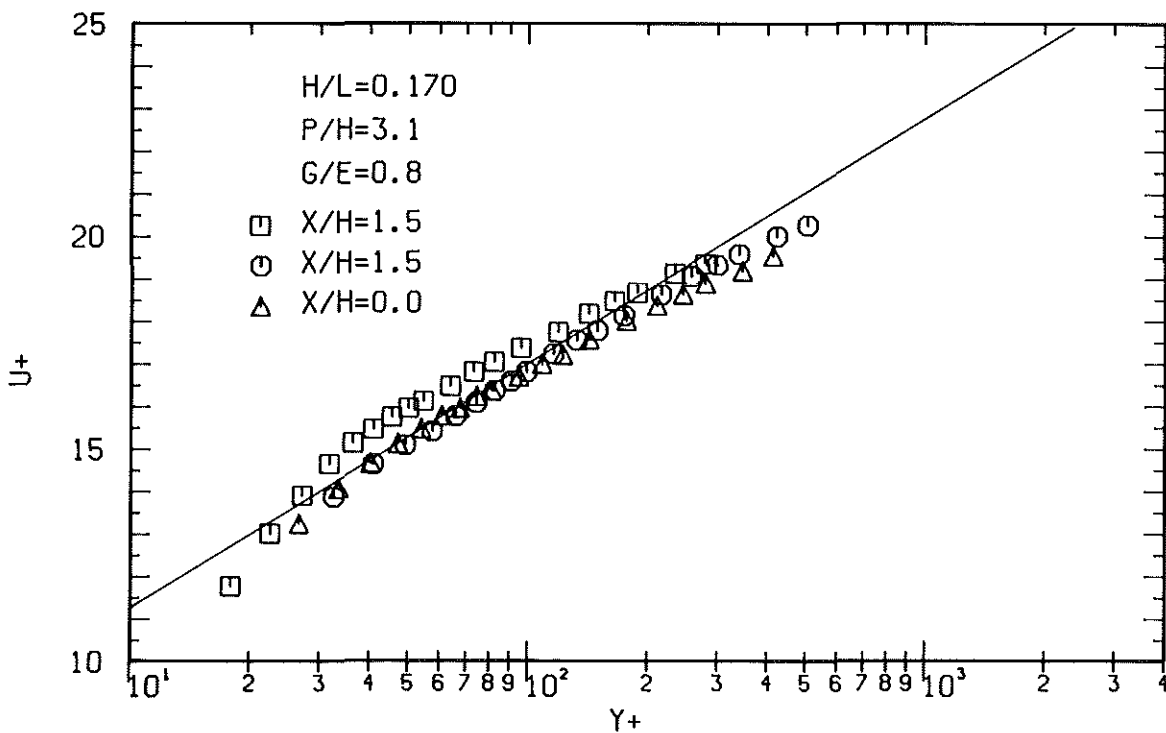
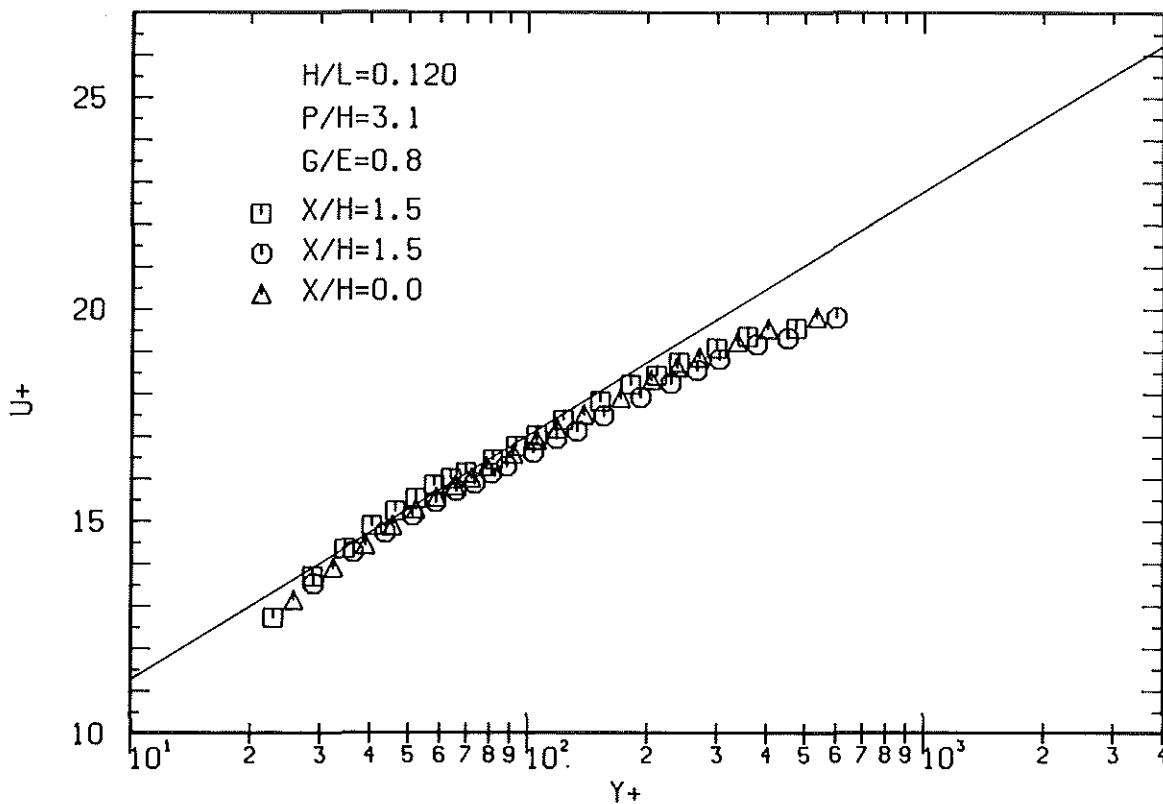


Fig.15.17-15.18: Non-dimensional velocity profiles at the smooth wall with roughness No.4 at opposite wall

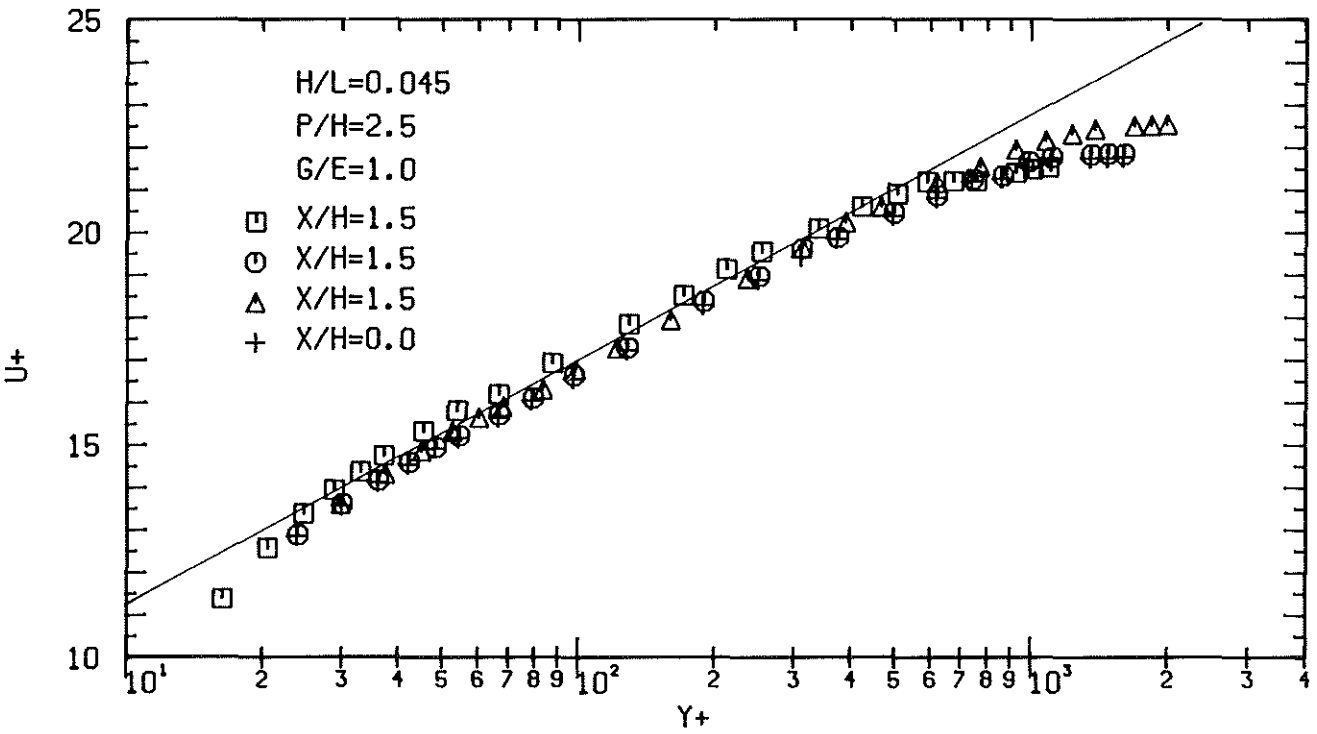
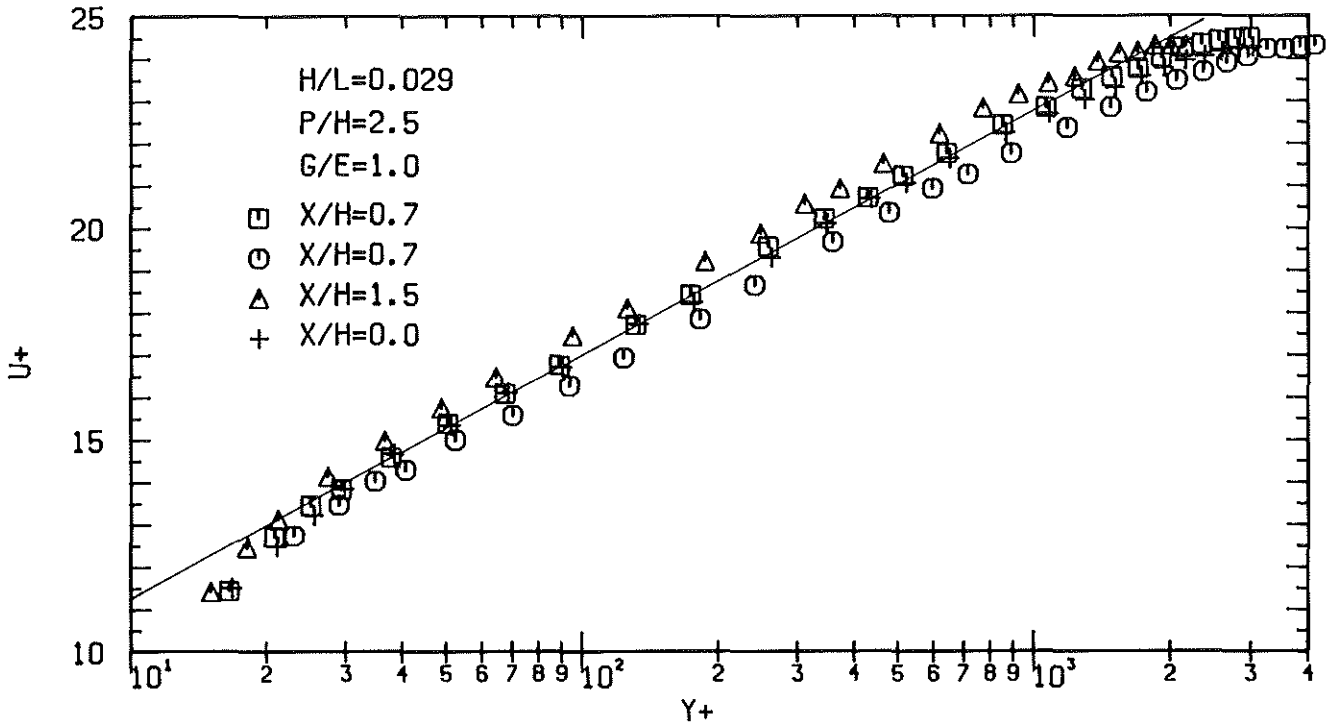


Fig.15.19-15.20: Non-dimensional velocity profiles at the smooth wall with roughness No.5 at opposite wall

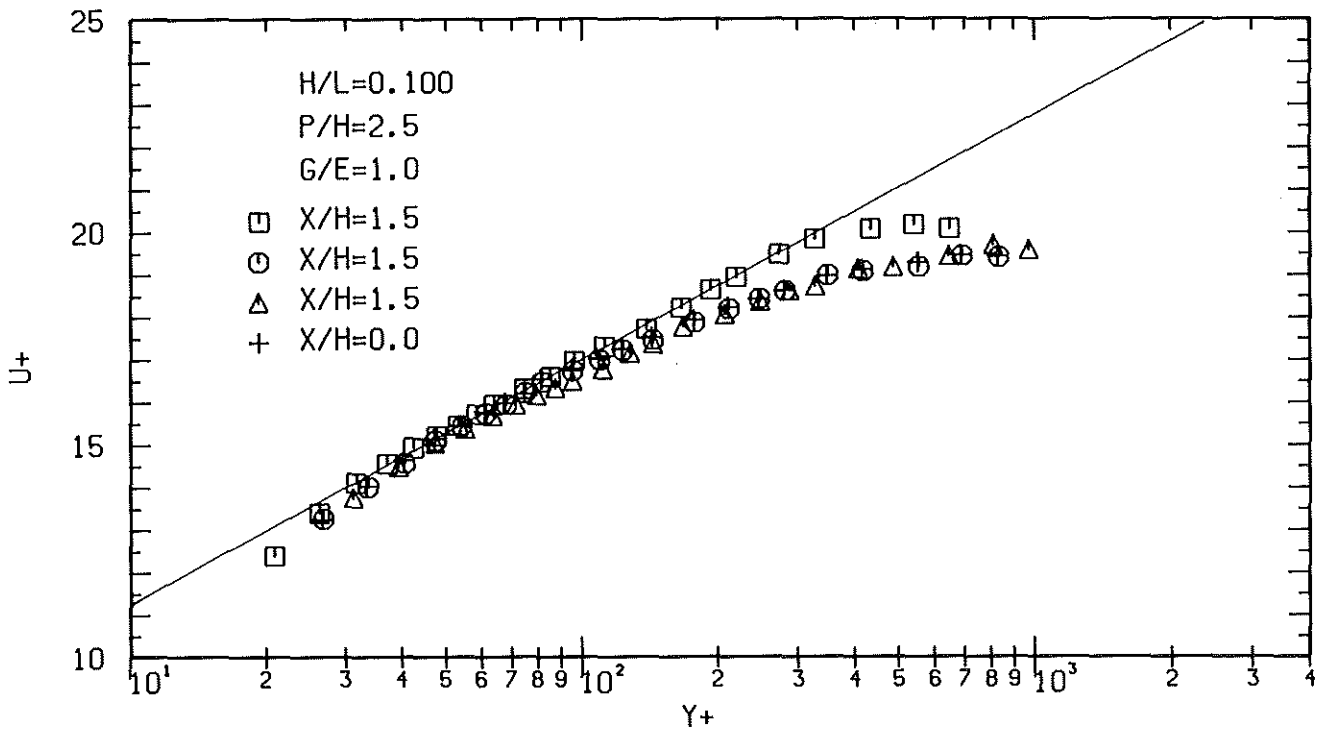
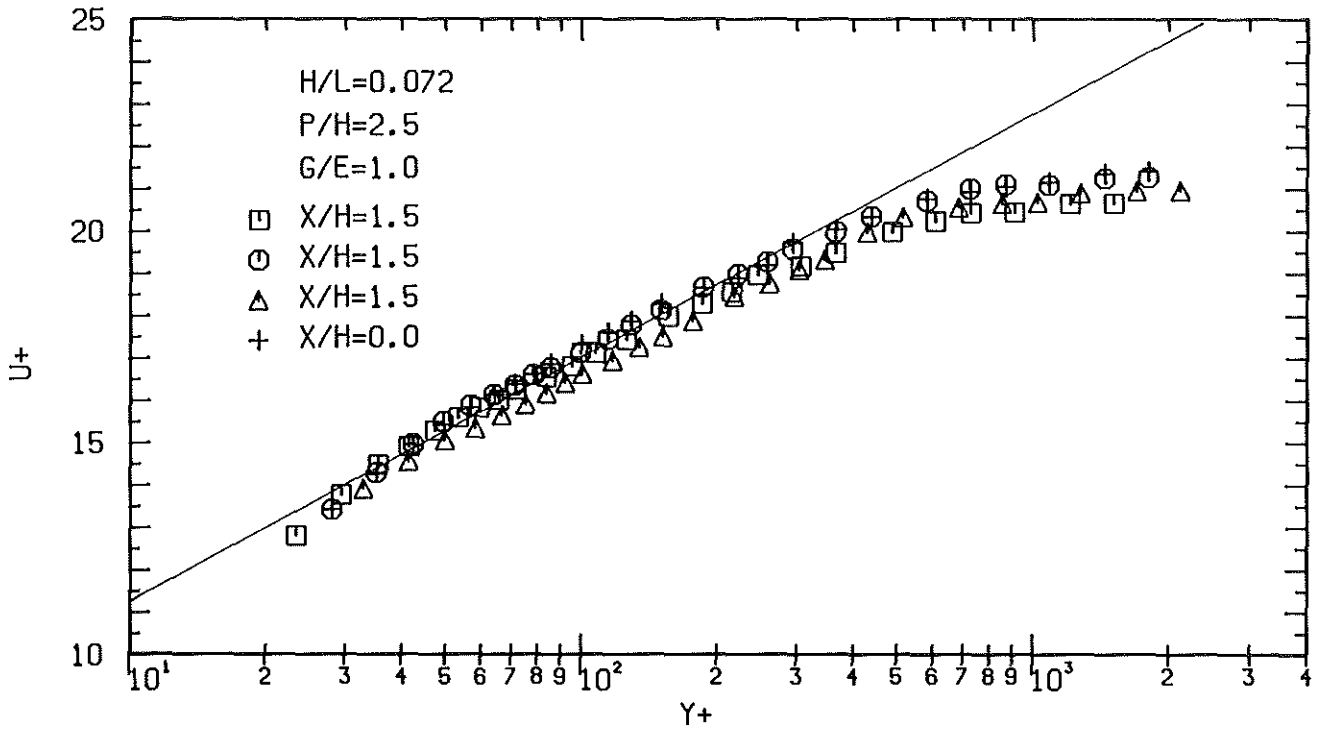


Fig.15.21-15.22: Non-dimensional velocity profiles at the smooth wall with roughness No.5 at opposite wall

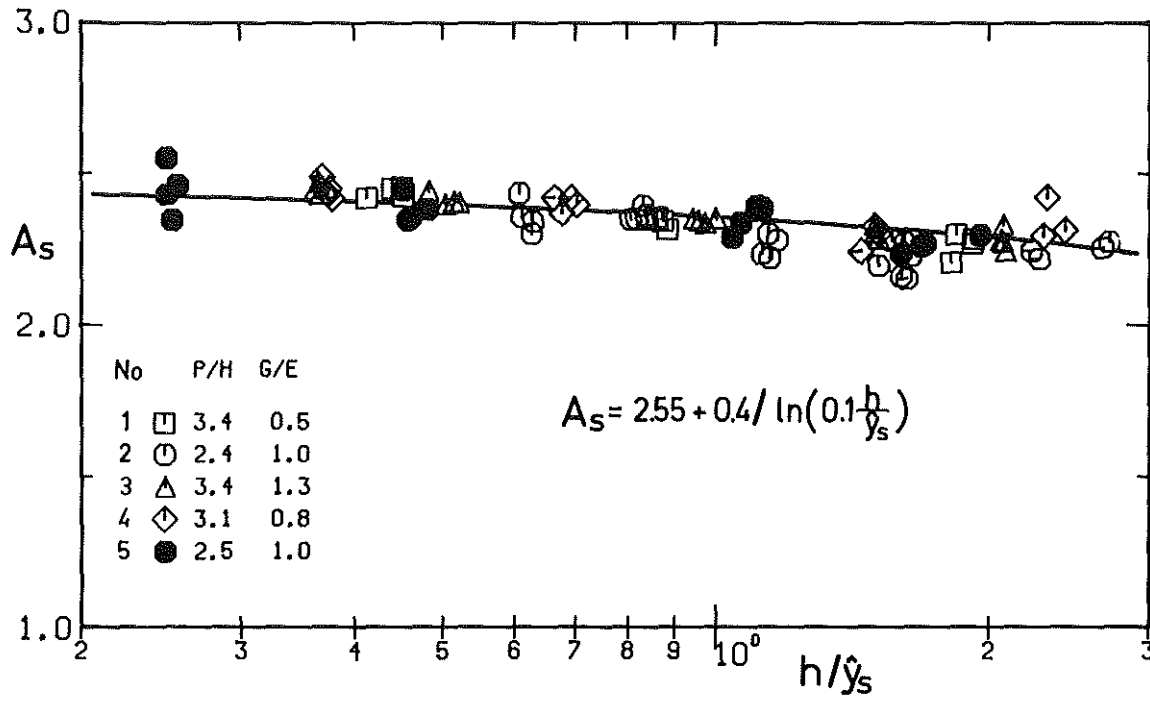


Fig.16: The profile parameter A_s determined by equation (15) with $B=5.5$

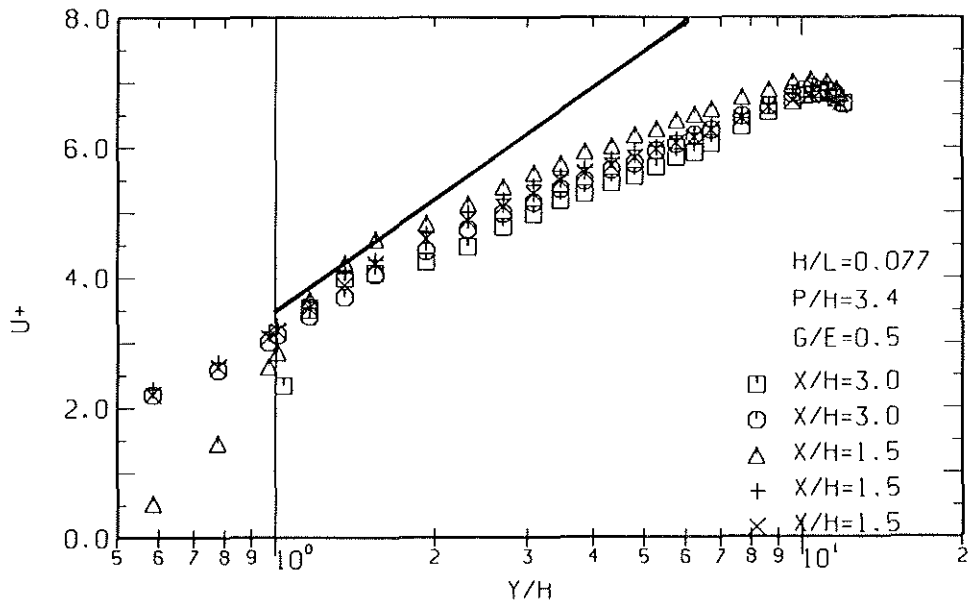
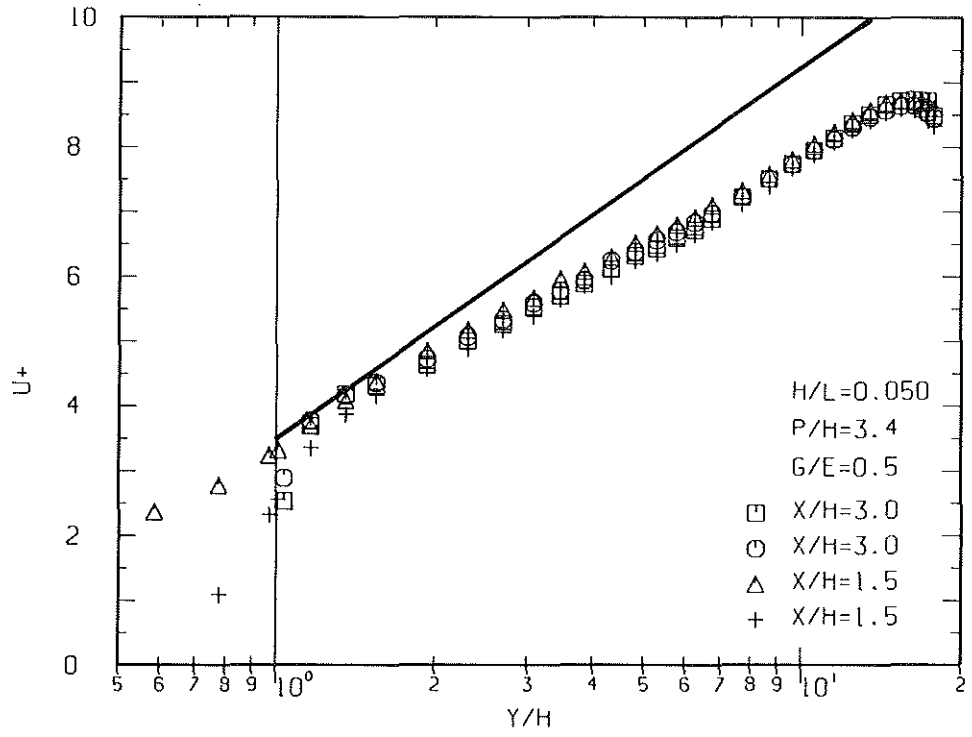


Fig.17.1-17.2: Non-dimensional velocity profiles at the rough wall (No.1)

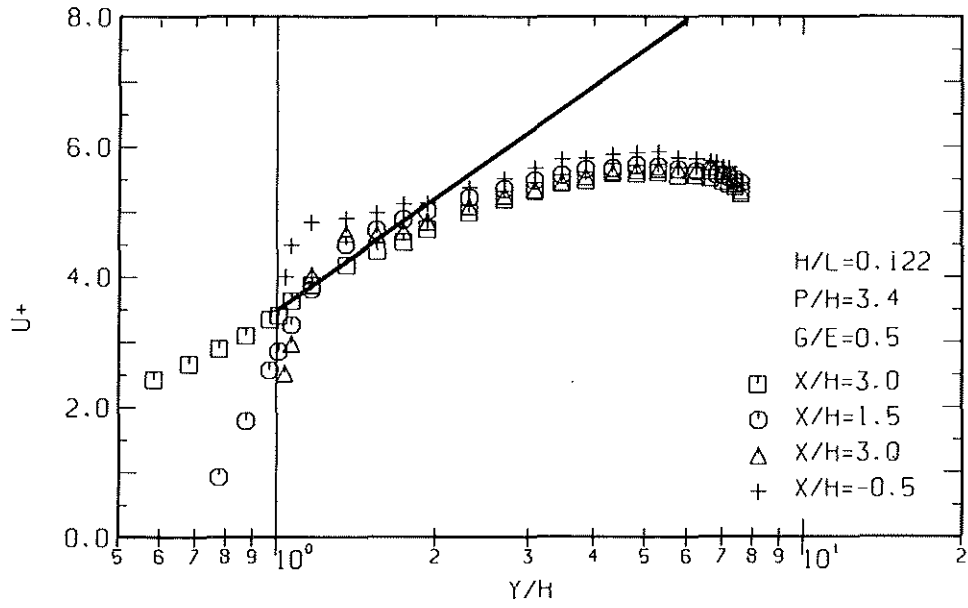


Fig.17.3: Non-dimensional velocity profiles at the rough wall (No.1)

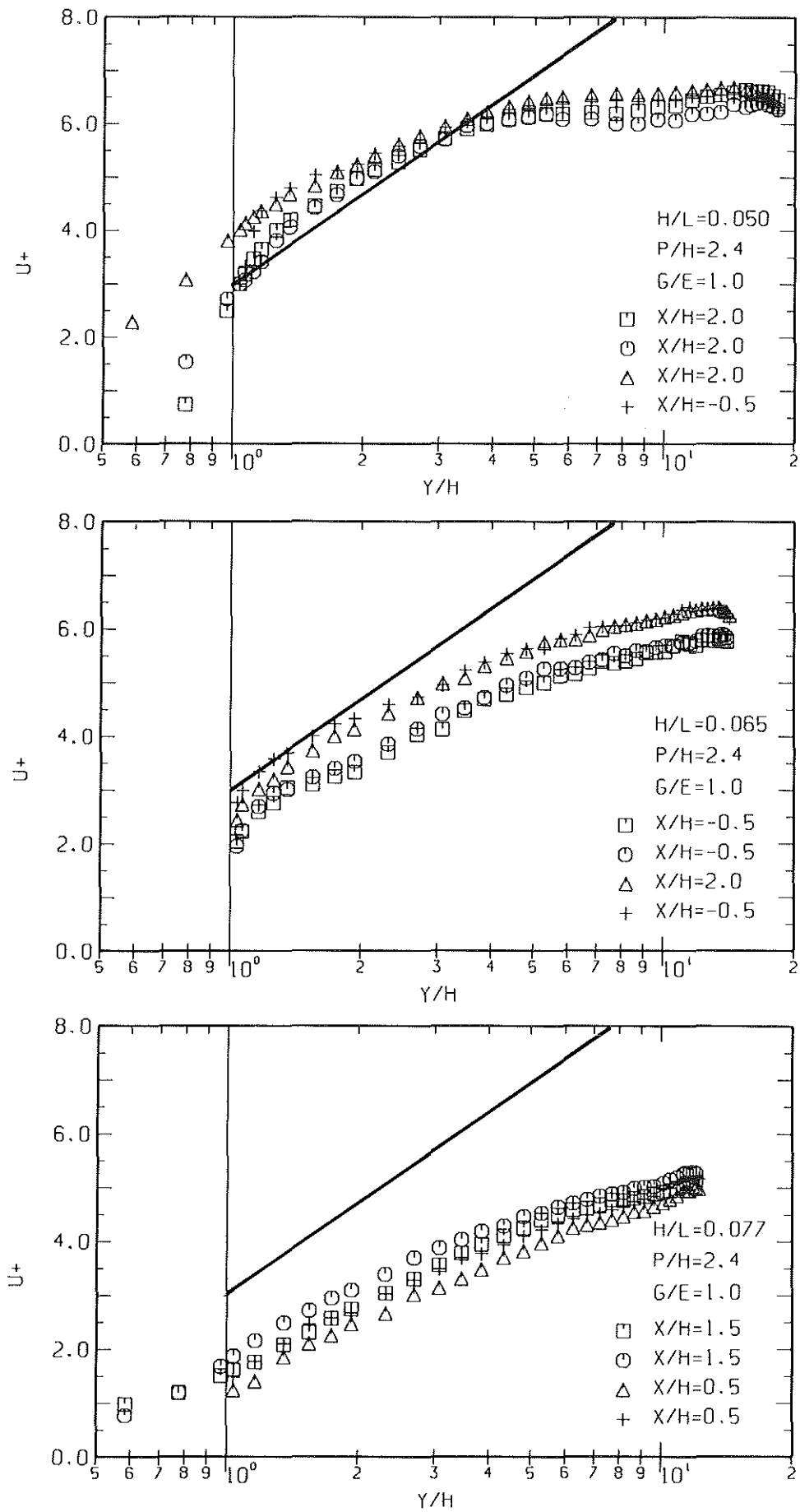


Fig.17.4-17.6: Non-dimensional velocity profiles at the rough wall (No.2)

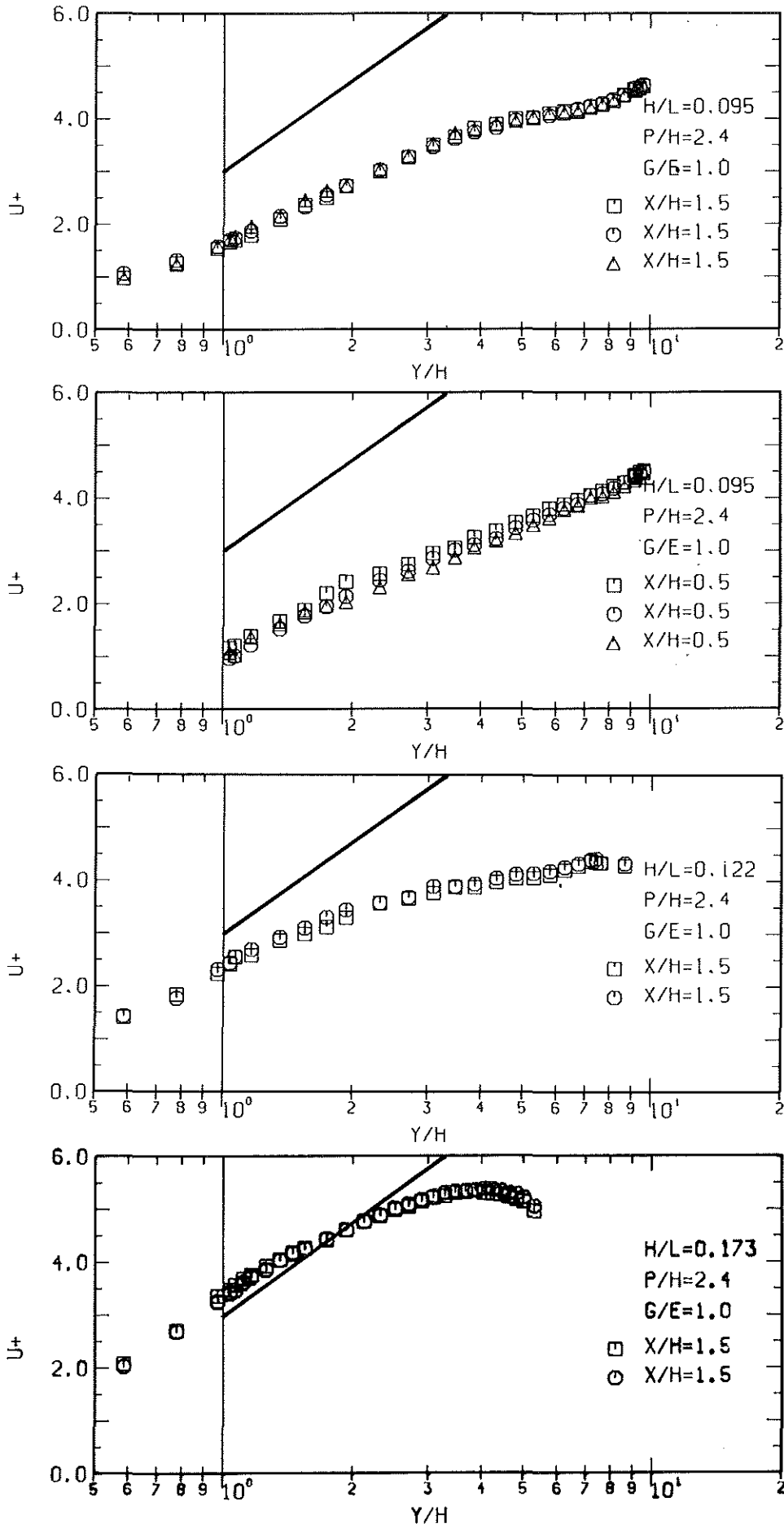


Fig.17.7-17.10: Non-dimensional velocity profiles at the rough wall (No.2)

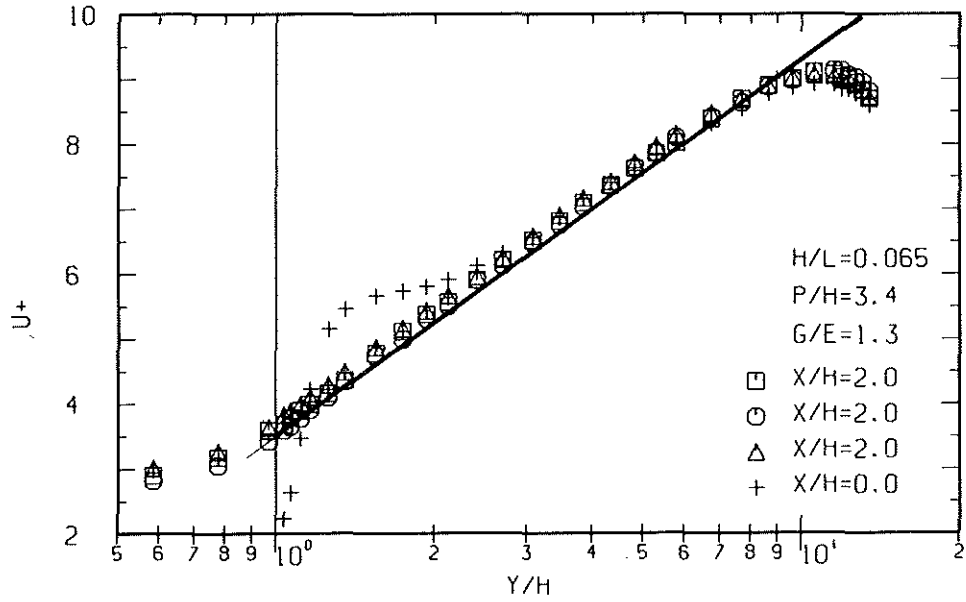
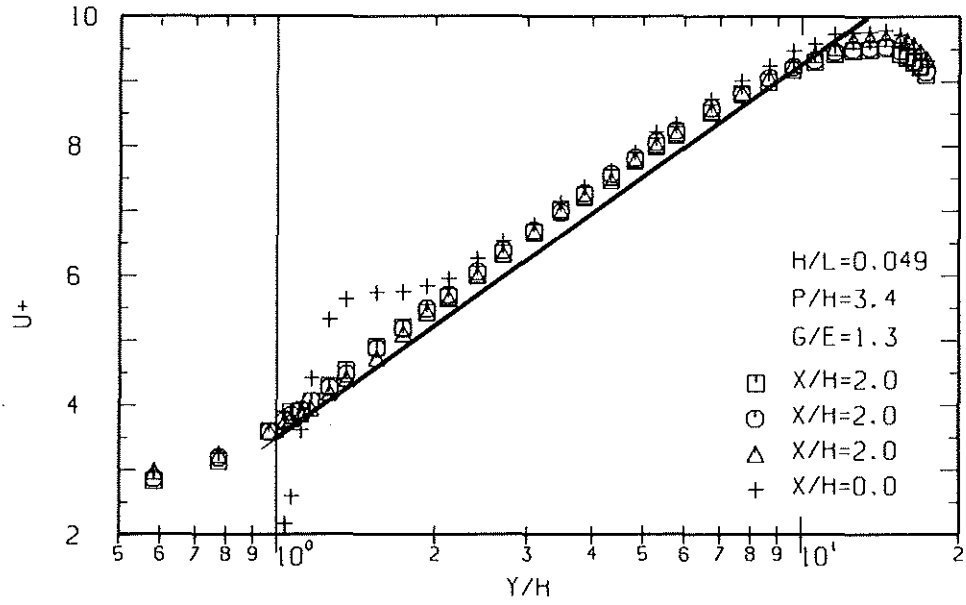


Fig.17.11-17.12: Non-dimensional velocity profiles at the rough wall (No.3)

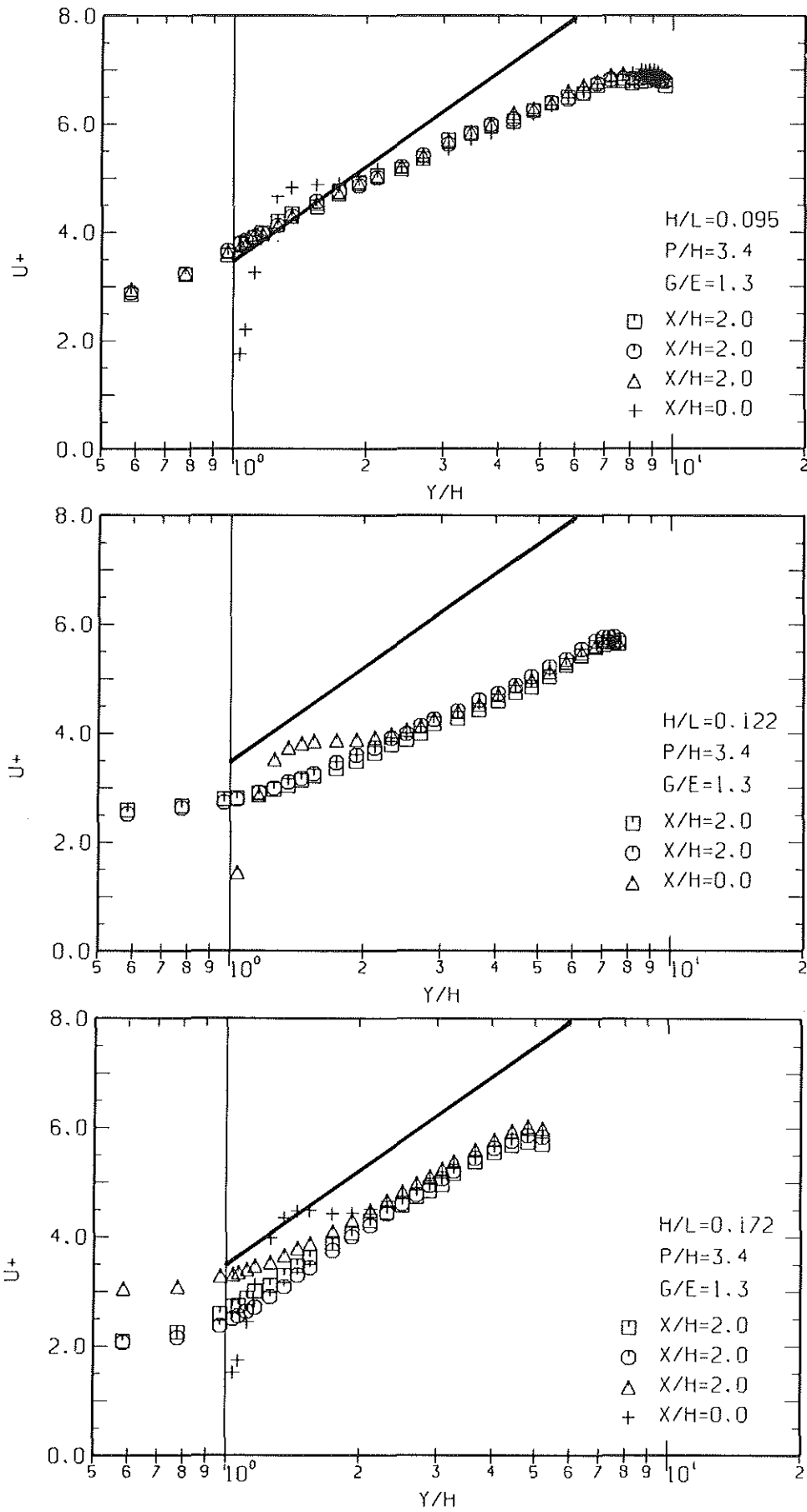


Fig.17.13-17.15: Non-dimensional velocity profiles at the rough wall (No.3)

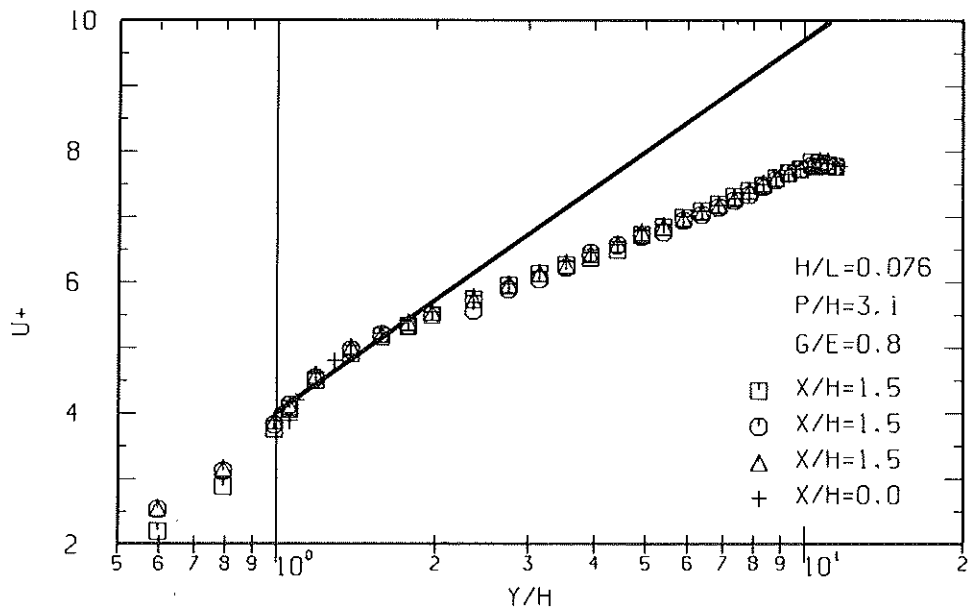
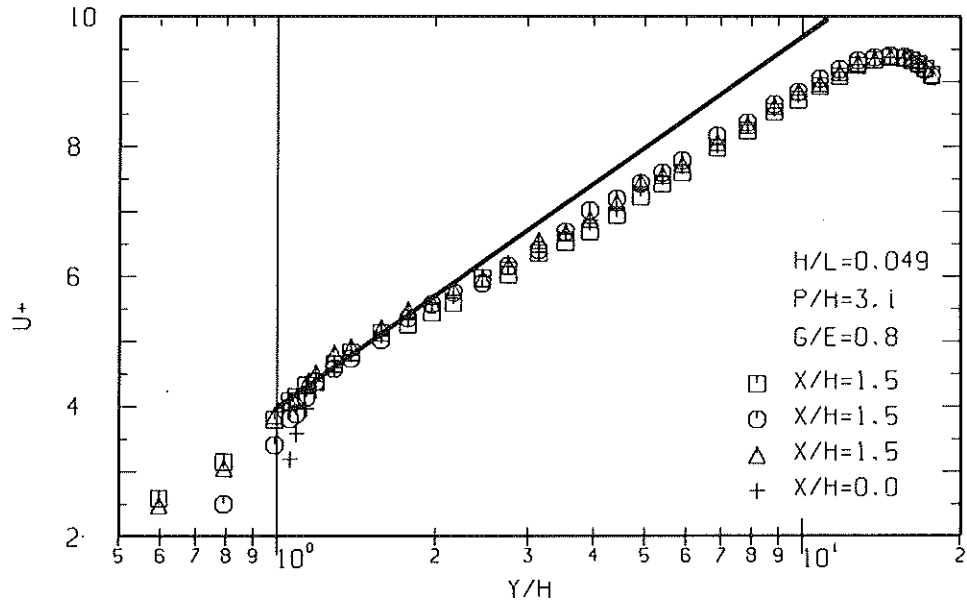


Fig.17.16-17.17: Non-dimensional velocity profiles at the rough wall (No.4)

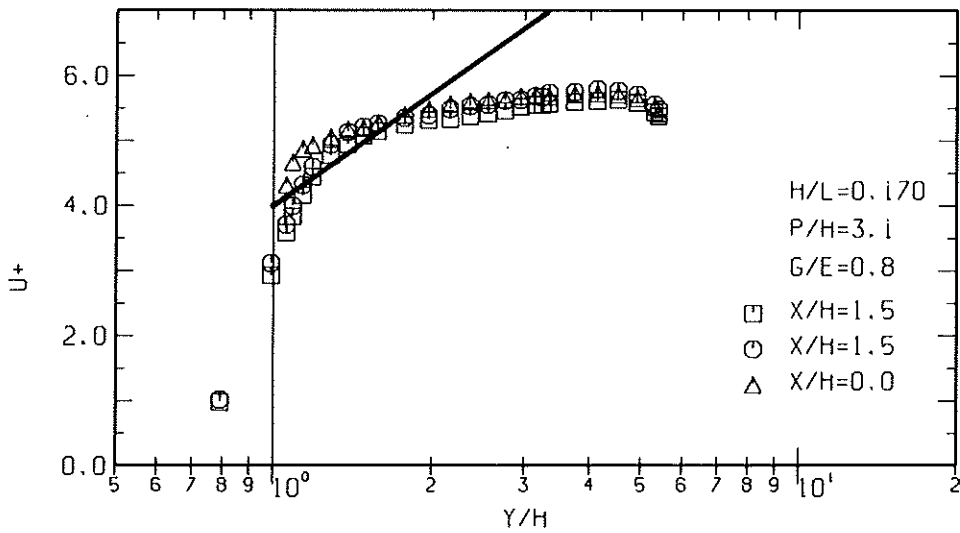
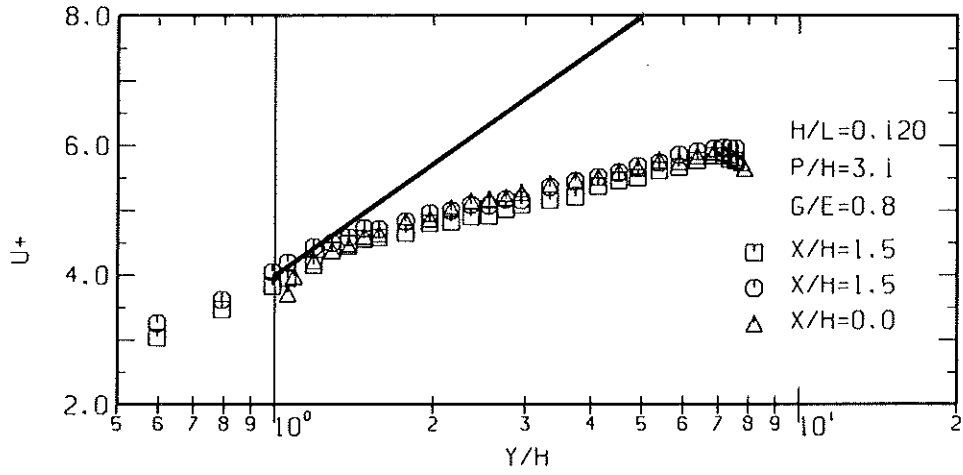


Fig.17.18-17.19: Non-dimensional velocity profiles at the rough wall (No.4)

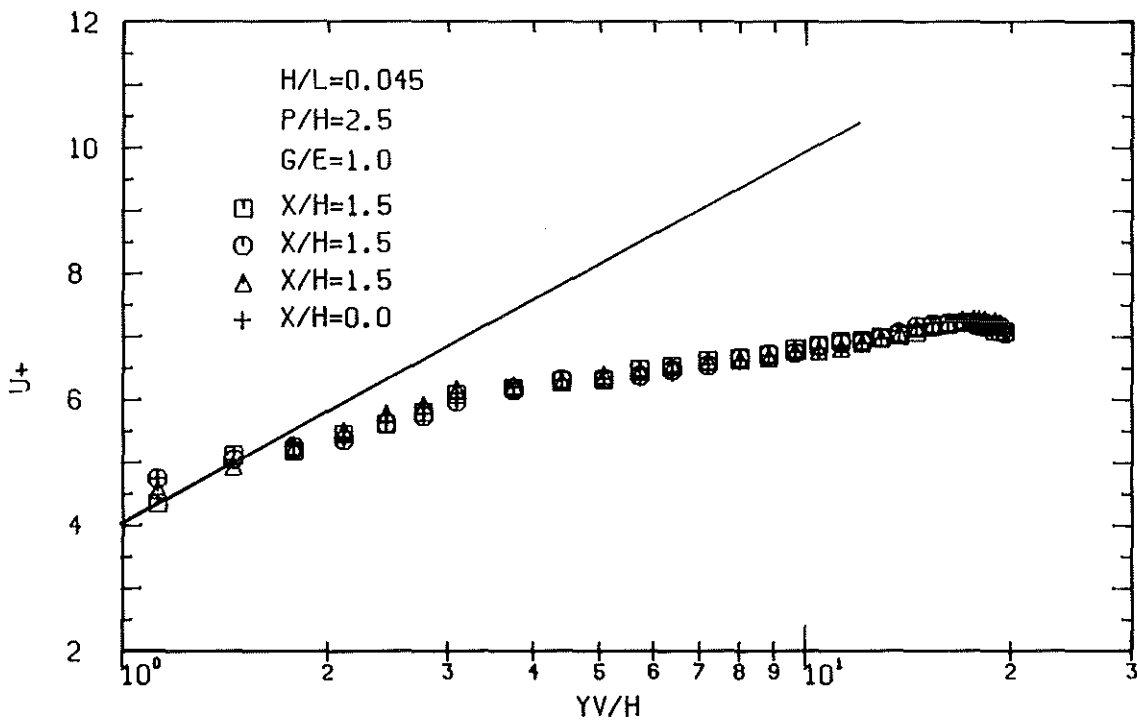
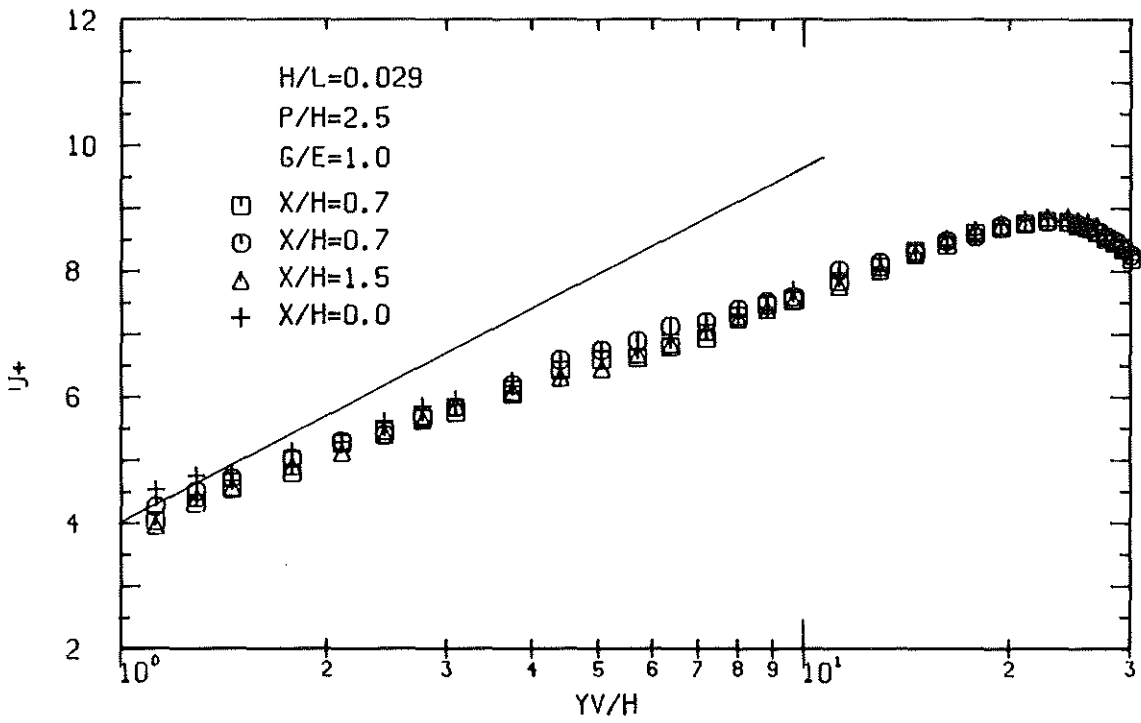


Fig.17.20-17.21: Non-dimensional velocity profiles at the rough wall (No.5)

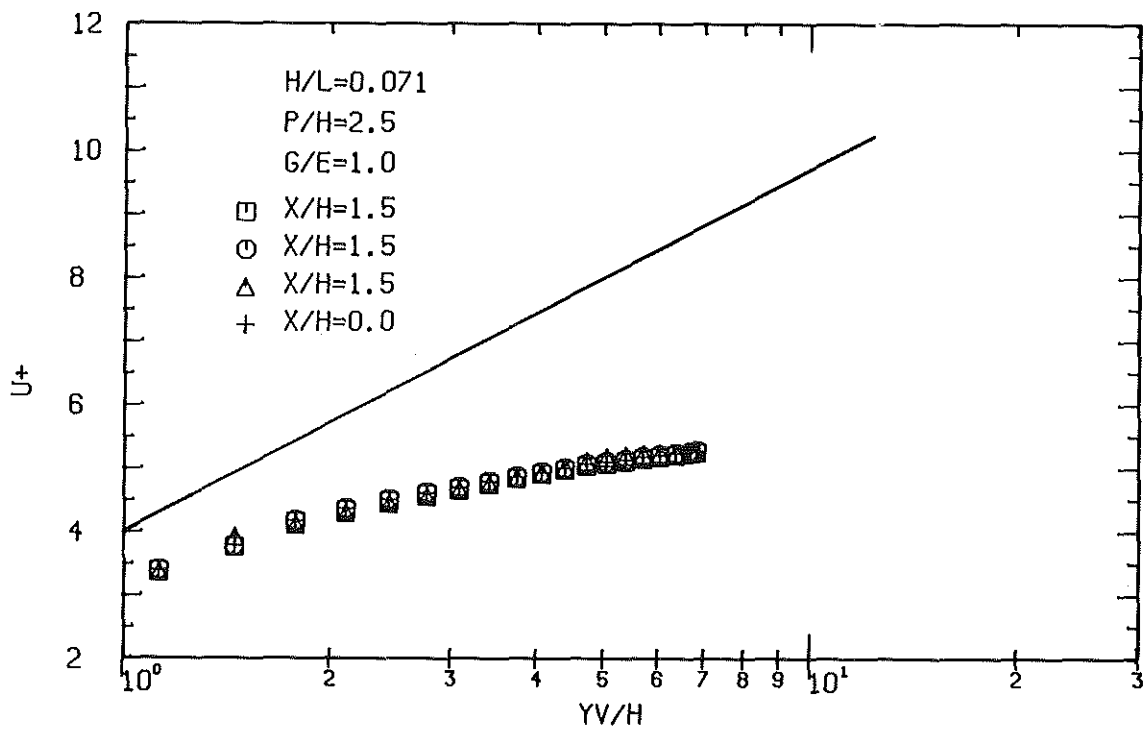
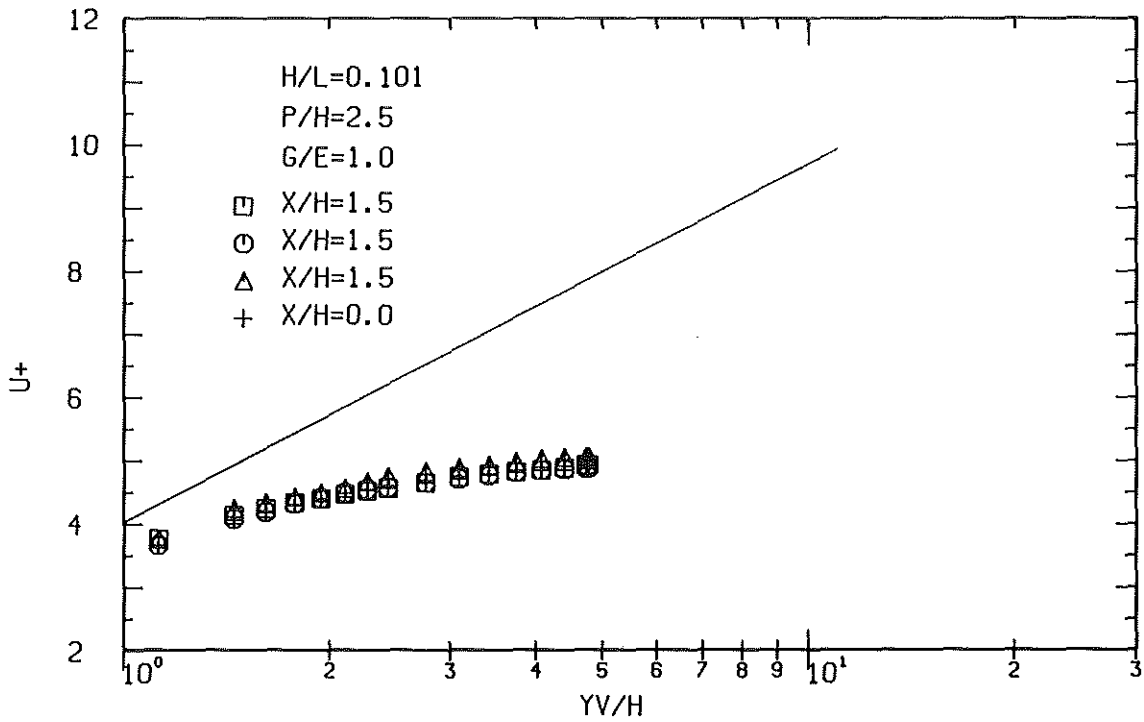


Fig.17.22-17.23: Non-dimensional velocity profiles at the rough wall (No.5)

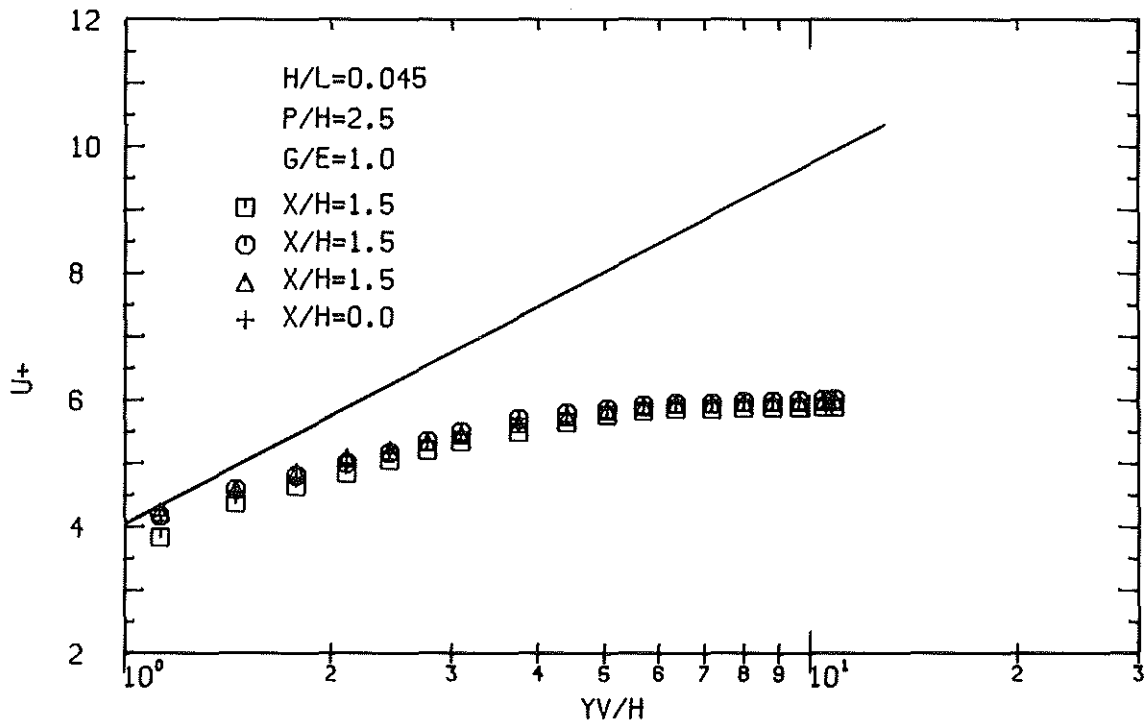
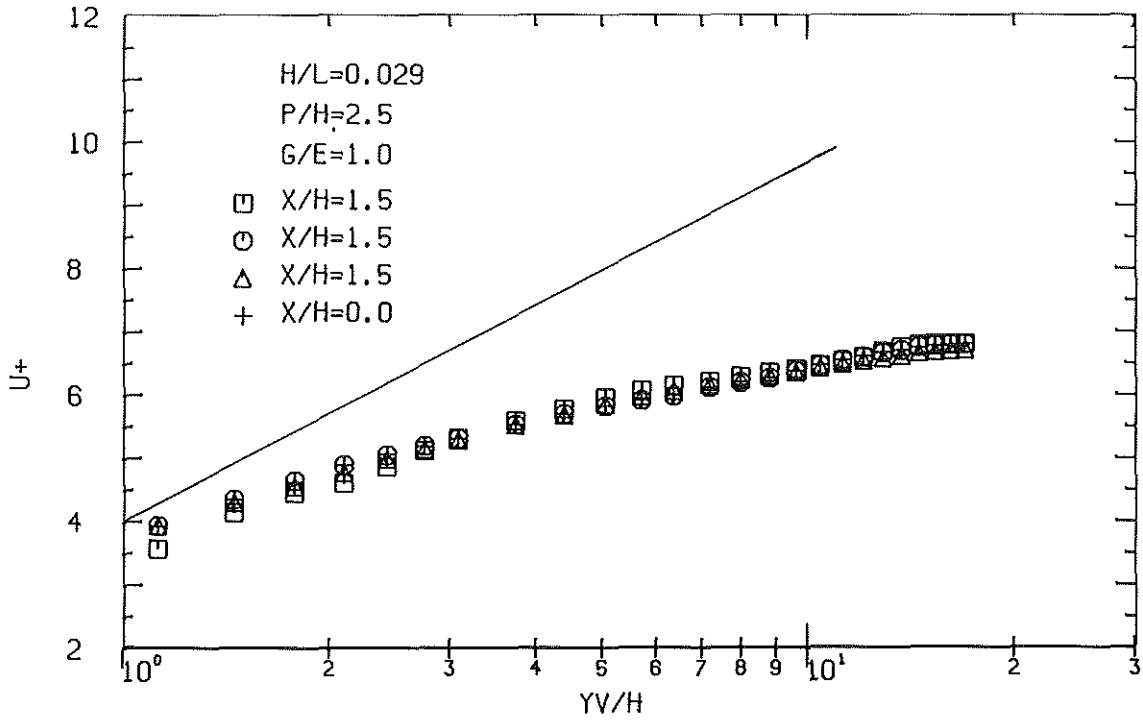


Fig.17.24-17.25: Non-dimensional velocity profiles at the rough wall (No.5, both walls rough)

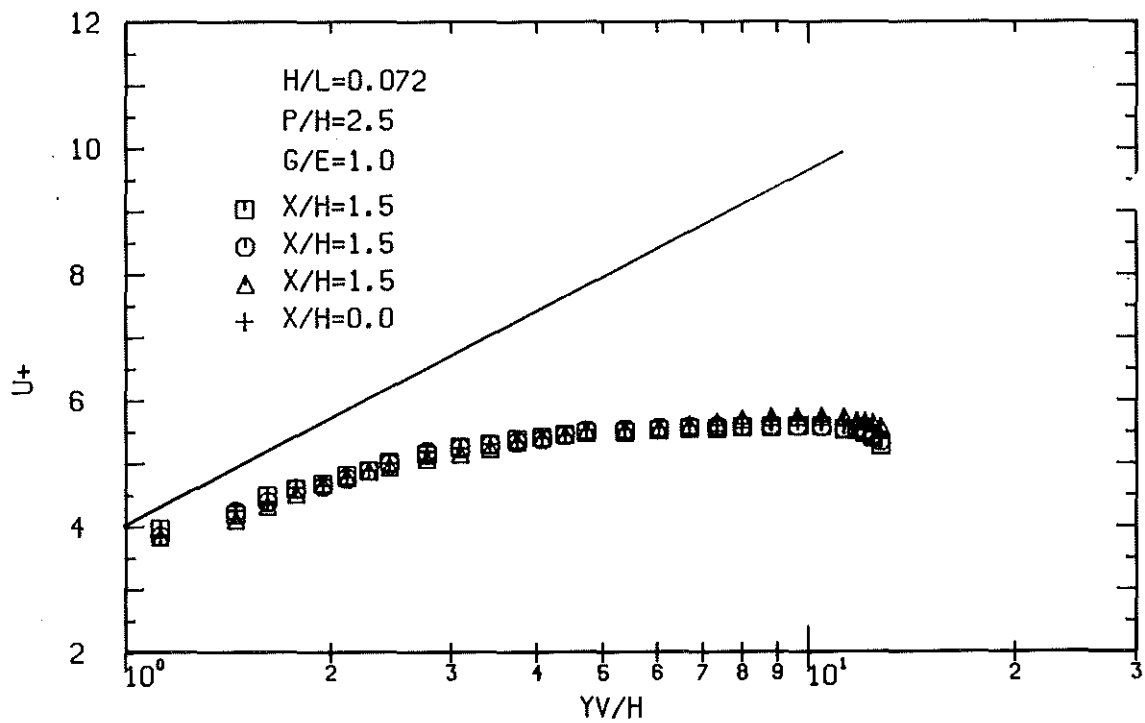
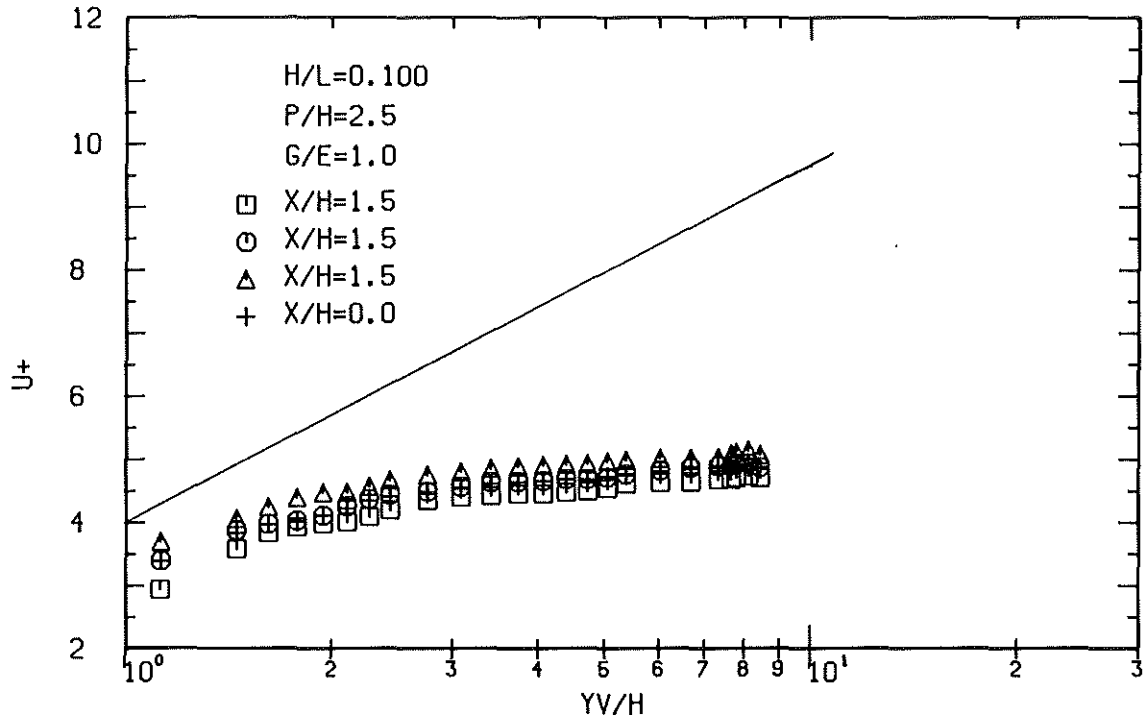


Fig.17.26-17.27: Non-dimensional velocity profiles at the rough wall (No.5, both walls rough)

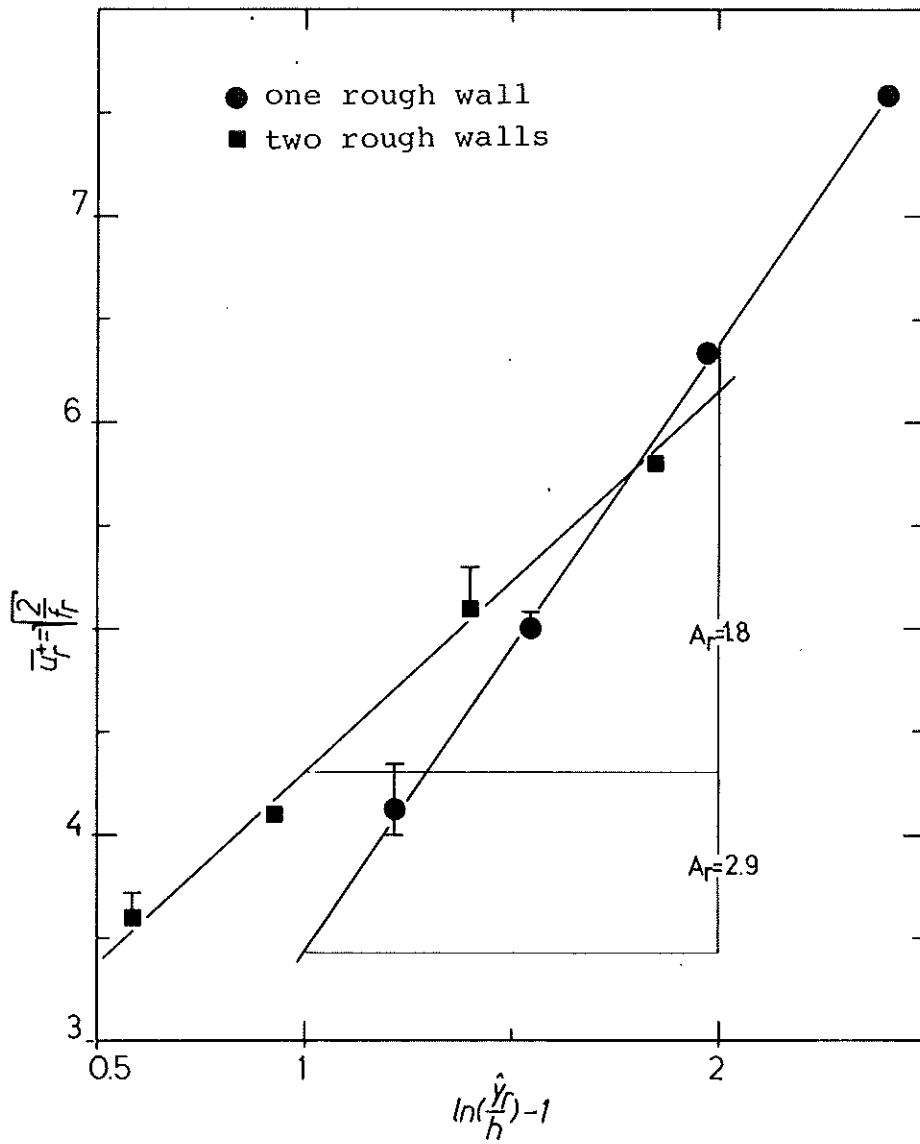


Fig.18: The average dimensionless velocity \bar{u}_r^+ of the rough zone versus $\ln(\hat{y}_r/h)-1$ for roughness No.5 in symmetrical and asymmetrical flow ($\epsilon=\epsilon_{vol}$)

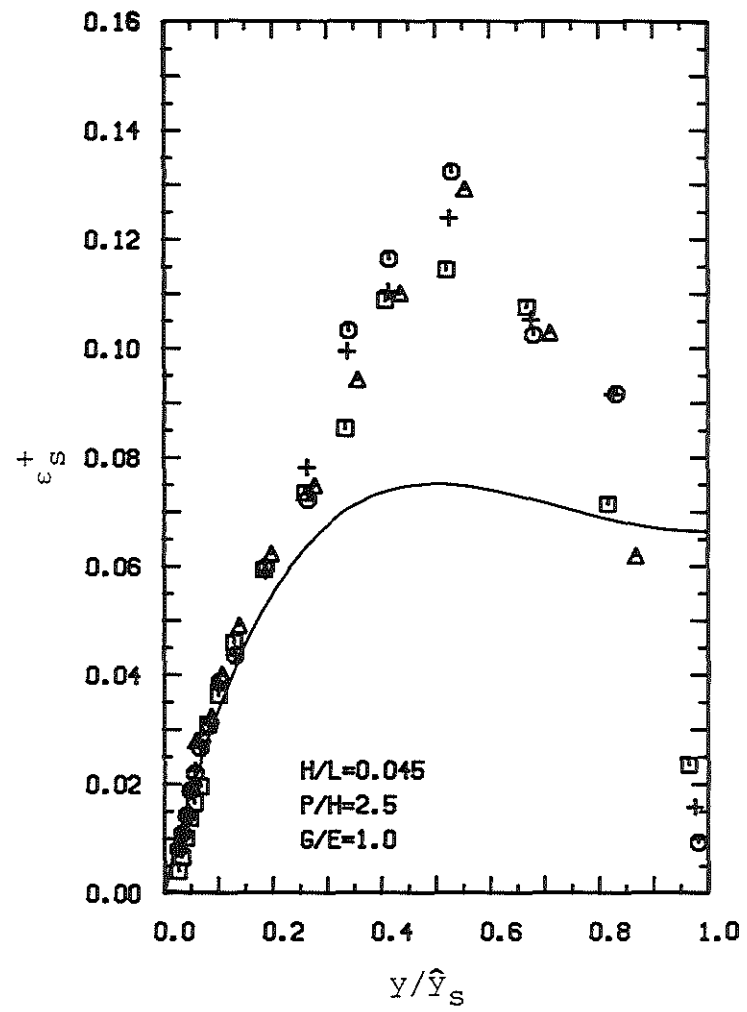
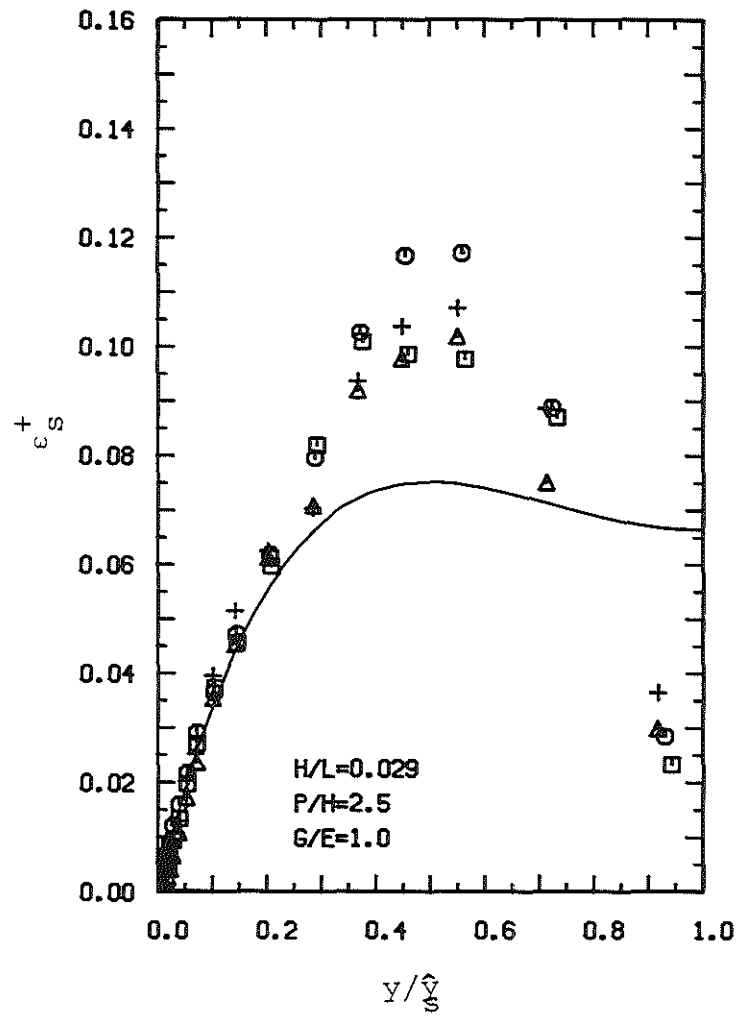


Fig.19.1-19.2: Eddy viscosity ϵ_s^+ in the smooth zone (No.5)

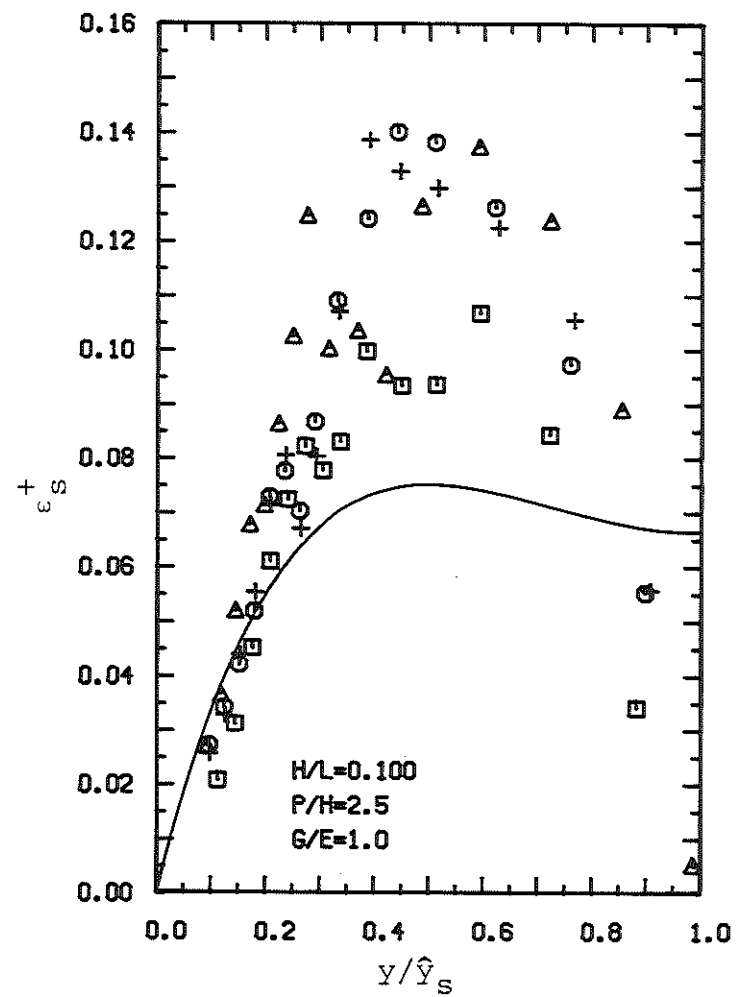
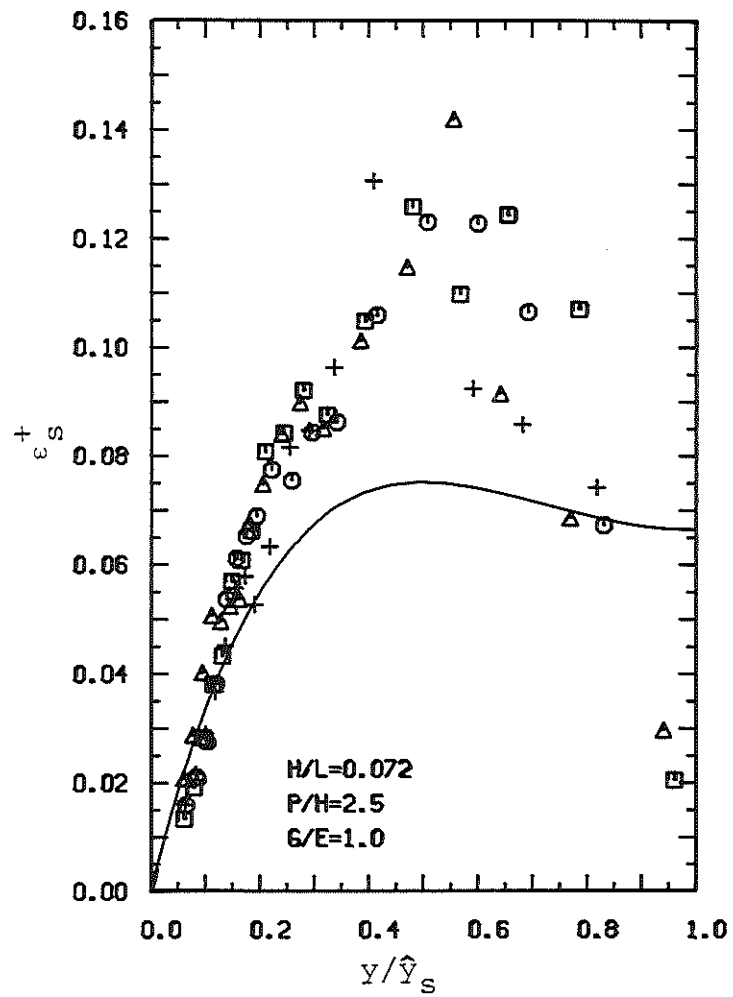


Fig.19.3-19.4: Eddy viscosity ϵ_s^+ in the smooth zone (No.5)

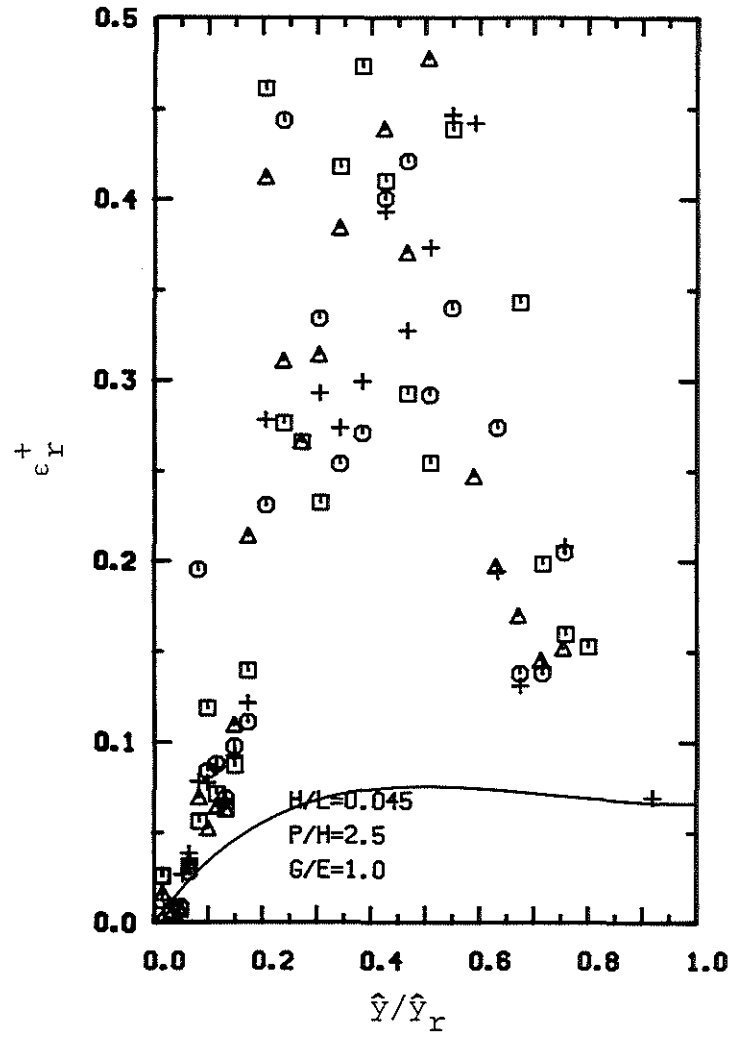
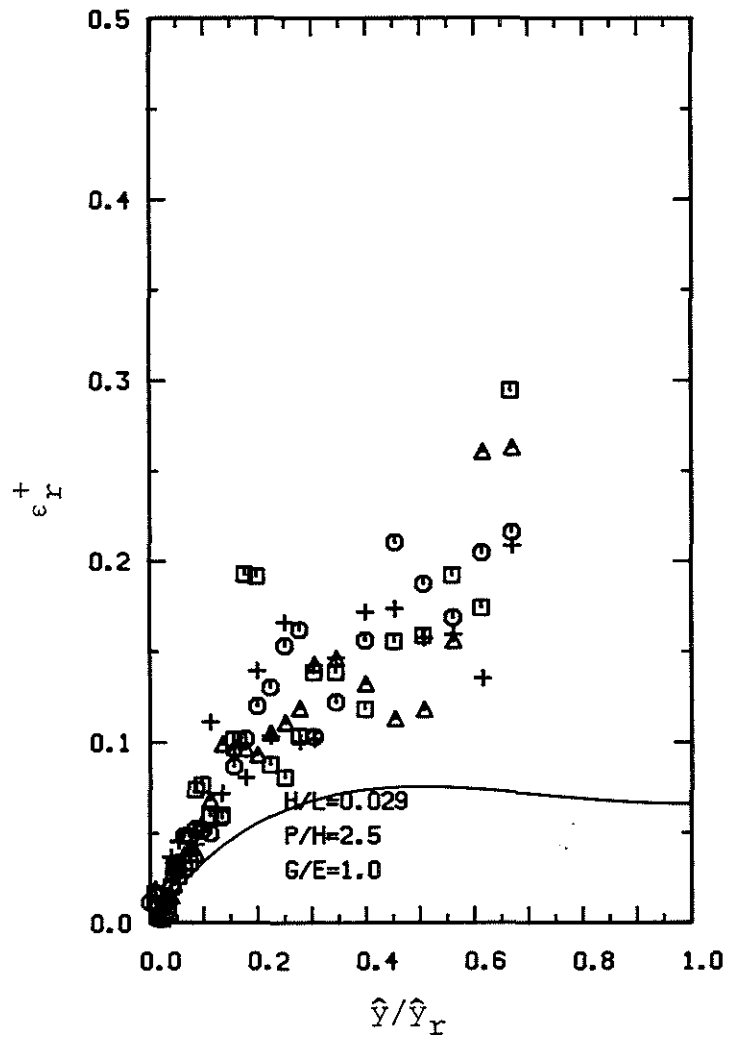


Fig.20.1-20.2: Eddy viscosity ϵ_r^+ in the rough zone (No.5)

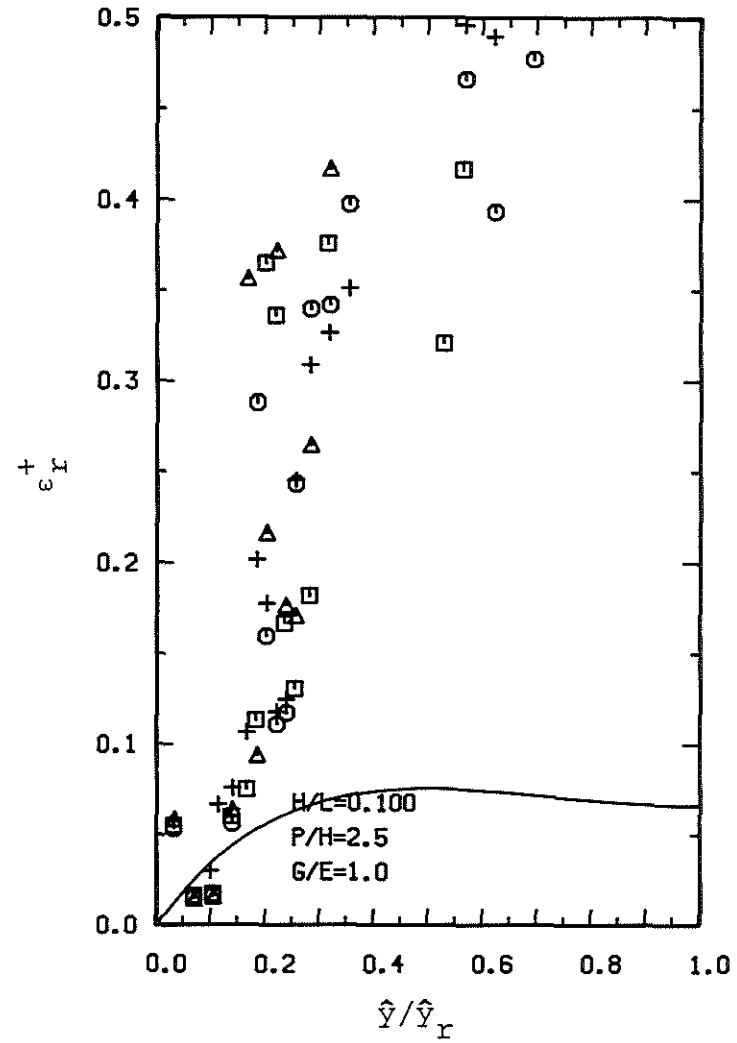
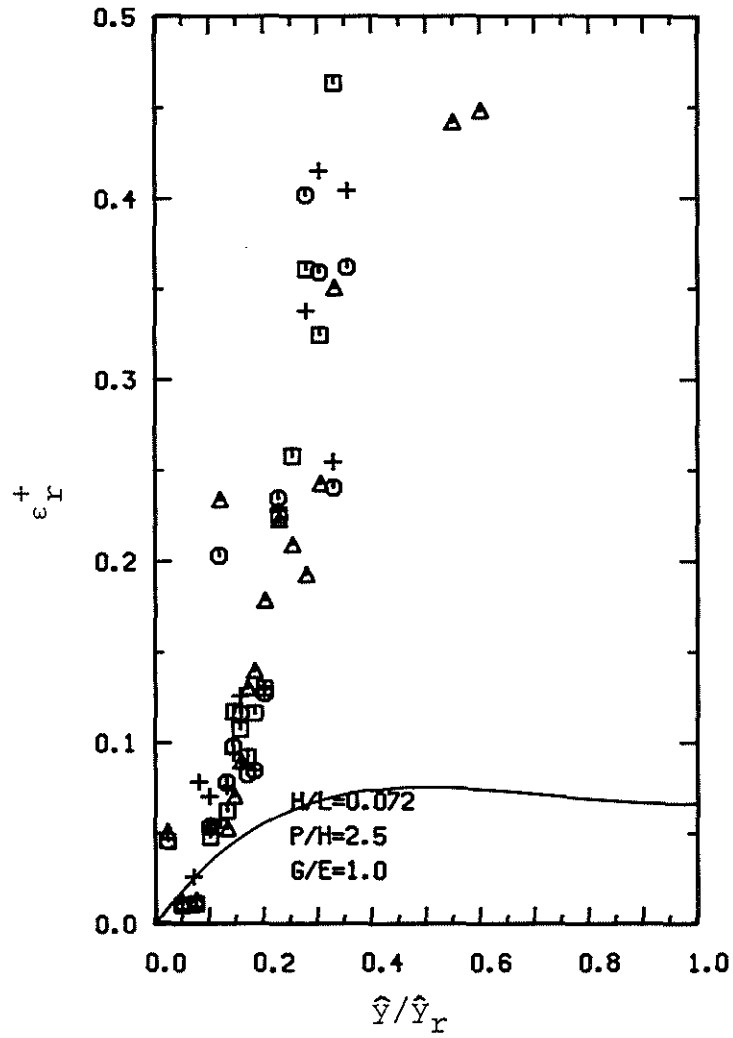


Fig.20.3-20.4: Eddy viscosity ϵ_r^+ in the rough zone (No.5)